

Title: Sub-Planckian Black Holes

Date: Mar 09, 2017 01:00 PM

URL: <http://pirsa.org/17030055>

Abstract: <p>The characteristics of black holes smaller than the Planck scale are addressed. These result from a modified metric that reproduces desirable aspects of a variety of disparate models in the sub-Planckian limit, while remaining Schwarzschild in the large mass limit. The self-dual nature of this solution has two interesting features: first, it naturally implies the Generalized Uncertainty Principle. Secondly, this metric exhibits an effective dimensional reduction feature, indicating that the gravitational physics of the sub-Planckian regime is effectively (1+1)-D.</p>

Sub-Planckian Black Holes



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09 March 2017 . Perimeter Institute . Waterloo, ON

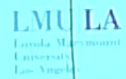
B. Carr, JM, P. Nicolini, *JHEP* **1507**:052 (2015) [arXiv:1504.07637]

R. Casadio, R. Cavalcanti, A. Guigno, JM, *PLB* **760**, 36 (2016) [arXiv:1509.09317]

L. Manfredi and JM, *AHEP* **2016**:1543741 (2016) [arXiv: 1610.00797]



Sub-Planckian Black Holes



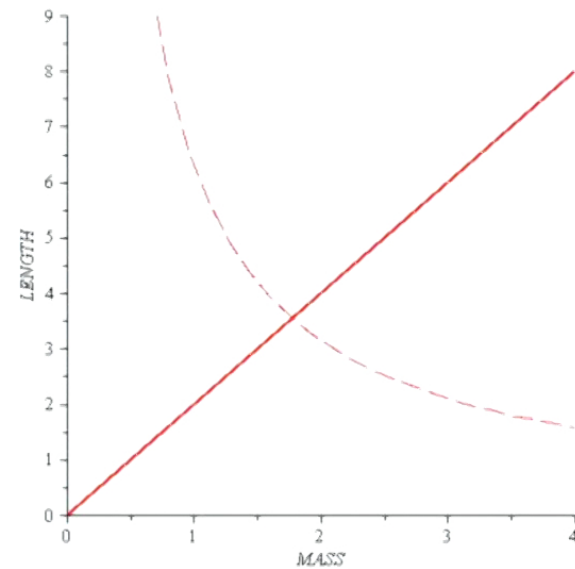
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Characteristic Scales of Nature

The large- and short-scale characteristics of Nature are defined by different and (apparently) disconnected **theories** and **length scales**

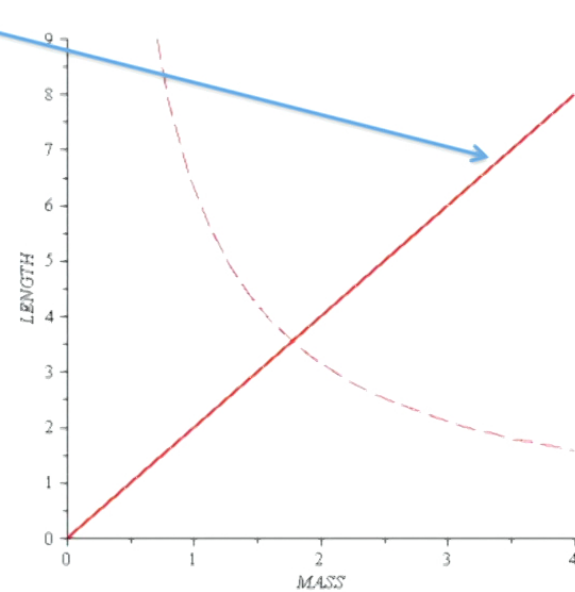


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Large: General Relativity

$$r_g = \frac{2GM}{c^2} \longrightarrow M$$



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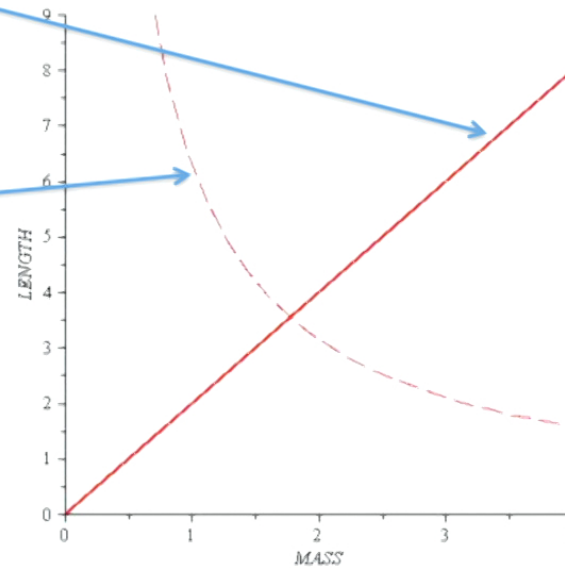
The large- and short-scale characteristics of Nature are defined by different and (apparently) disconnected **theories** and **length scales**

Large: General Relativity

$$r_g = \frac{2GM}{c^2} \longrightarrow M$$

Short: Quantum Mechanics

$$\lambda_C = \frac{\hbar}{Mc} \longrightarrow \frac{1}{M}$$

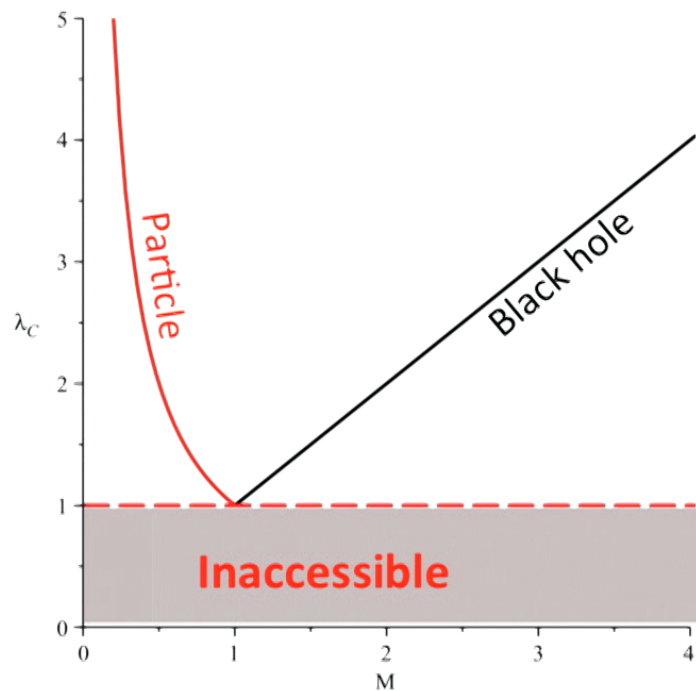


Planck Scale Criticality

The point $\lambda_C = r_g$ is a critical point

$$\lambda_C = r_g \implies M_{\min} = \frac{1}{\sqrt{G}} = M_{\text{Pl}}$$

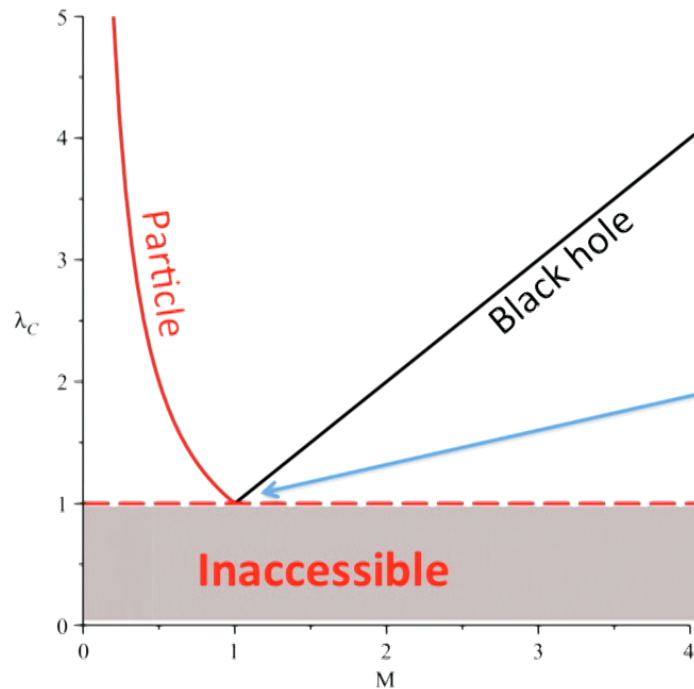
$$\lambda_{\min} = \ell_{\text{Pl}}$$



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$$\lambda_{\min} = \ell_{\text{Pl}}$$



Defines the smallest black hole, or alternatively the largest particle

The Generalized Uncertainty Principle

A momentum-dependent modification of the HUP

$$\Delta x \Delta p \geq \frac{\hbar}{2} (1 + \beta \Delta p^2)$$

[Veneziano, Europhys. Lett. **2**, 199 (1986); Amati, Ciafaloni, Veneziano, PLB **197**, 81 (1987); Amati, Ciafaloni, Veneziano, PLB **216**, 41 (1989)]

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Can be thought of as adding a *gravitational* correction to the standard quantum uncertainty that “connects” scale diagram:

$$\Delta x_Q \sim \frac{\hbar}{2\Delta p} \quad , \quad \Delta x_G \sim 2G\Delta p$$

$$\Delta x = \Delta x_Q + \Delta x_G \quad \Longrightarrow \quad \Delta x \sim \Delta p + \frac{1}{\Delta p}$$

Characterized by a *momentum / mass duality*

Black Holes and the GUP

QG effects become important for BHs at scales where GUP is relevant

1. Hot unstable remnants [Adler and Chen, Nucl.Phys.Proc.Sup. **124** (2003) 103]
 - Minimum sized BH with *maximum* temperature
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3. Cold stable remnants [Isi, JRM, Nicolini, JHEP **1311**:139 (2013)]

- Exact solution to deformed stress-energy tensor
- Minimum sized BH; thermodynamically-stable remnants

$$ds^2 = - \left(1 - \frac{2GM}{r} \gamma \left(2; \frac{r}{\sqrt{\beta}} \right) \right) dt^2 + \frac{dr^2}{1 - \frac{2GM}{r} \gamma \left(2; \frac{r}{\sqrt{\beta}} \right)} + r^2 d\Omega^2$$

GUP and Sub-Planckian Black Holes

[Carr, JM, Nicolini, JHEP 07:052 (2015)]

Can **black holes** exist *below* the Planck scale? $M_{\text{BH}} < M_{\text{Pl}}$

An approach to including the GUP in general relativity is to emphasize the **duality** in the black hole mass

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$$\Delta x \sim \frac{1}{\Delta p} + \Delta p$$


$$\Delta x_G \sim \frac{1}{M_{\text{bh}}} + M_{\text{bh}}$$

What Is the Mass in the GUP?

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Quantum: $M \sim \sqrt{s}$ ←

GR: $M = M_{\text{ADM}}$ →

M in classical Schwarzschild is ADM (Komar):

$$M \equiv \frac{1}{4\pi G} \int_{\partial\Sigma} d^2x \sqrt{\gamma^{(2)}} n_\mu \sigma_\nu \nabla^\mu K^\nu$$

$M < M_{pl}$ determined from approximate stress-energy interior to R_C

$$M \equiv \int_{\Sigma} d^3x \sqrt{\gamma} n_\mu K_\nu T^{\mu\nu} \simeq -4\pi \int_0^{R_C} dr r^2 T_0^0$$

GUP and Sub-Planckian Black Holes

[Carr, JM, Nicolini, JHEP 07:052 (2015)]

Proposal: $M = -4\pi \int_0^{\ell_P} dr r^2 T_0^0$ Unspecified quantum gravity stress-energy

Gives new M (duality): $M \longrightarrow M \left(1 + \frac{\beta}{2} \frac{M_{\text{Pl}}^2}{M^2} \right)$

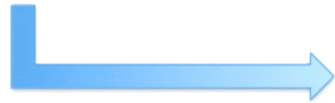
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Metric is:



$$ds^2 = F(r) dt^2 - F(r)^{-1} dr^2 - r^2 d\Omega^2$$
$$F(r) = 1 - \frac{2}{M_{\text{Pl}}^2} \frac{M}{r} \left(1 + \frac{\beta}{2} \frac{M_{\text{Pl}}^2}{M^2} \right)$$

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Planck mass is now critical point for which...



$$M \gg M_{\text{Pl}} \implies F(r) \sim 1 - \frac{M}{r}$$
$$M \ll M_{\text{Pl}} \implies F(r) \sim 1 - \frac{1}{Mr}$$

Black Hole Characteristics: Horizon

$$F(r_H) = 0$$



$$r_H = \frac{2}{M_{\text{Pl}}^2} \left(\frac{M^2 + \frac{\beta}{2} M_{\text{Pl}}^2}{M} \right)$$

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$$M \gg M_{\text{Pl}} \implies r_H \approx \frac{2M}{M_{\text{Pl}}^2}$$

$$M \sim M_{\text{Pl}} \implies r_H \sim \frac{2 + \beta}{M_{\text{Pl}}}$$

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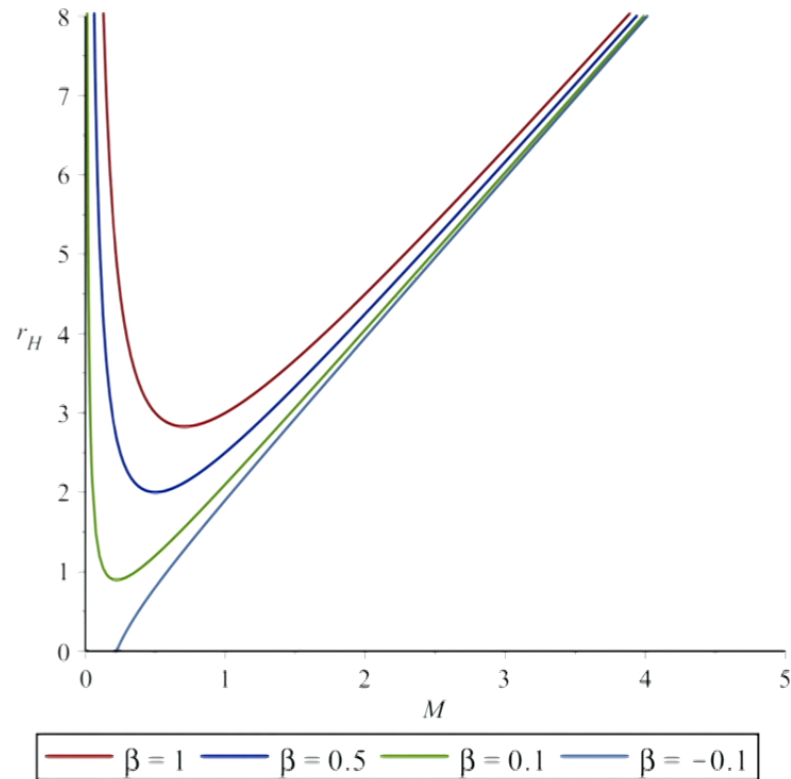


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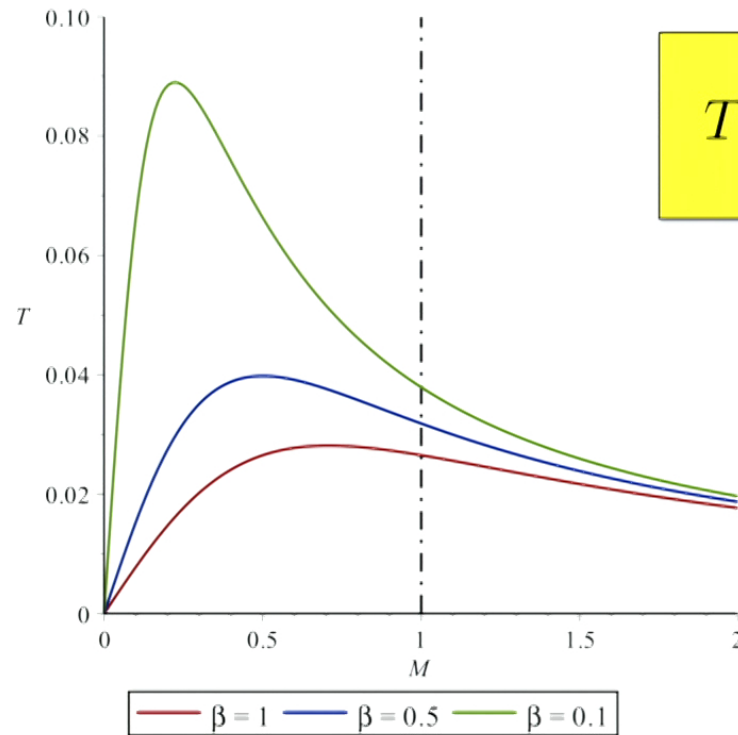
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Black Hole Characteristics: Temperature

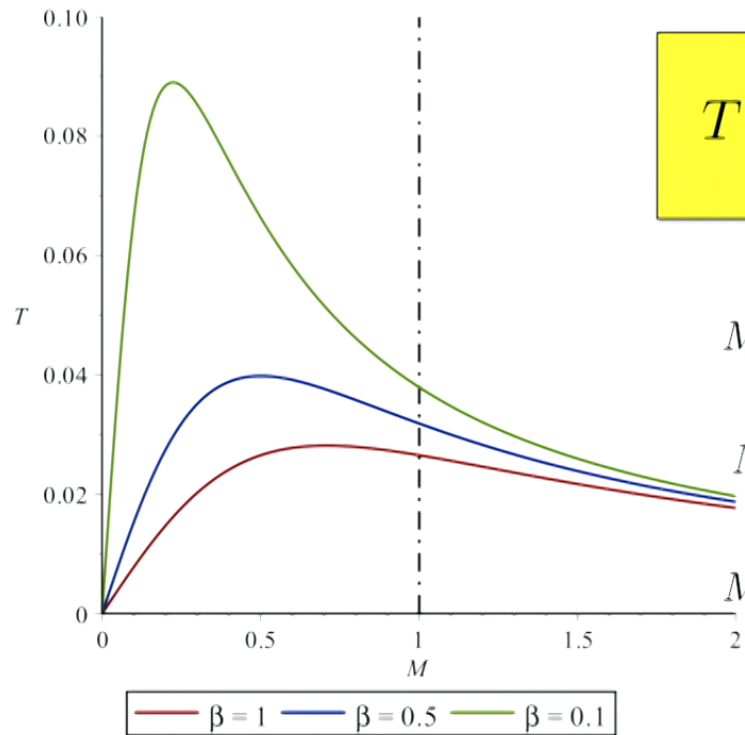
From surface gravity: $T = \frac{\kappa}{2\pi}$, $\kappa = \frac{1}{2} \frac{dF}{dr}(r = r_H)$



$$T = \frac{M_{\text{Pl}}^2}{8\pi M(1 + \beta M_{\text{Pl}}^2/2M^2)}$$

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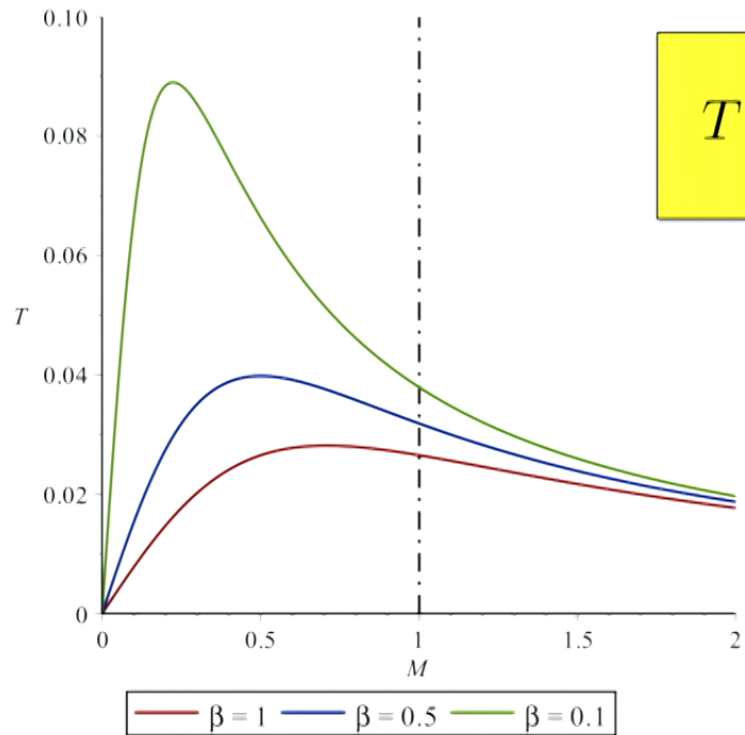
$$M \gg M_{\text{Pl}} \implies T \approx \frac{M_{\text{Pl}}^2}{8\pi M} \left[1 - \beta \left(\frac{M_{\text{Pl}}}{M} \right)^2 \right]$$

$$M \approx M_{\text{Pl}} \implies T \approx \frac{M_{\text{Pl}}}{8\pi(1 + \beta/2)}$$

$$M \ll M_{\text{Pl}} \implies T \approx \frac{M}{4\pi\beta} \left[1 - \frac{1}{\beta} \left(\frac{M}{M_{\text{Pl}}} \right)^2 \right]$$

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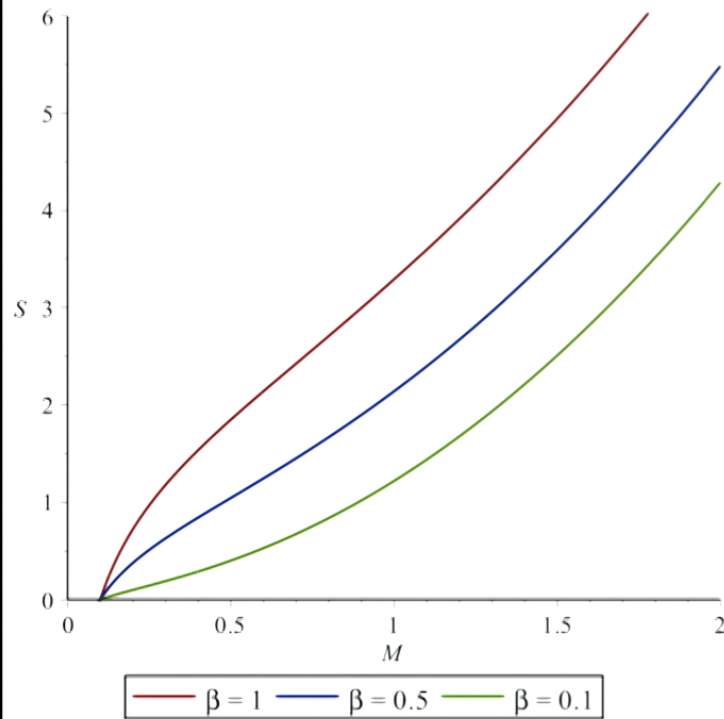
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Black Hole Characteristics: Entropy

$$S = \int_{M_0}^M \frac{dM'}{T(M')} \sim 4\pi \left(\frac{M^2}{M_{\text{Pl}}^2} - \frac{M_0^2}{M_{\text{Pl}}^2} + \beta \log \frac{M}{M_0} \right)$$

Black Hole Characteristics: Entropy

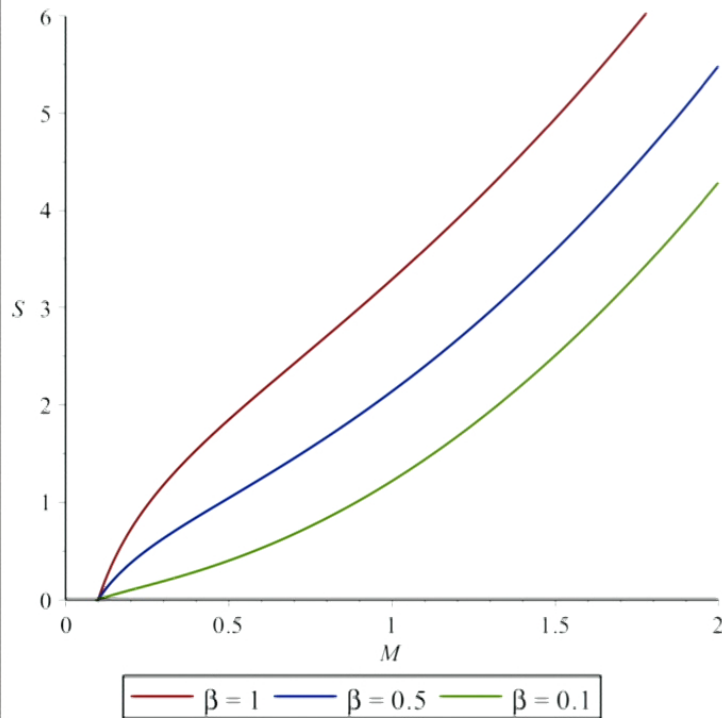
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$$\begin{aligned} M \gg M_{\text{Pl}} &\implies S \sim M^2 \\ M \ll M_{\text{Pl}} &\implies S \sim \log(M/M_0) \end{aligned}$$

I've seen these quantities before somewhere...

The (D+1)-dimensional Schwarzschild metric is...

$$ds^2 = \left(1 - \frac{R_D}{r^{D-2}}\right) dt^2 - \left(1 - \frac{R_D}{r^{D-2}}\right)^{-1} dr^2 - r^{D-1} d\Omega_{D-1}$$

where the horizon radius is....

$$R_D = \left(\frac{2G_D M}{D-2}\right)^{\frac{1}{D-2}} \quad \text{if } D > 2$$

and Gravitational constant is defined as

$$G_D = \frac{\ell_D^{D-2}}{m_D}$$

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and Gravitational constant is defined as

$$G_D = \frac{\ell_D^{D-2}}{m_D}$$



Matches an exact solution!
Can derive BH thermodynamics!

[Mann, Shiekh, Tarasov, NPB **341**, 134 (1990)]

(3+1)-D vs (1+1)-D Spacetime

(3+1)-D

$$g_{tt} = 1 - \frac{2G_N M}{r}$$

$$g_{rr} = -g_{tt}^{-1}$$

$$r_H \sim M$$

$$T \sim \frac{1}{M}$$

$$S \sim M^2$$

(1+1)-D

$$g_{tt} = 1 - G_1 M |x|$$

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Sub-Planckian regime

Dimensional reduction?

(1+1)-D

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The gravitational physics of the sub-Planckian regime is governed by an *effective* (1+1)-D

Dimensional Reduction, HEP, and QG

- Spectral (CDTs/EDTs, Hořava-Lifshitz, non-commutative)

[Ambjorn, Jurkiewicz, Loll, PRD72:064014 (2005);
Horava, PRL102, 161301 (2009);
Modesto and Nicolini, PRD81:104040 (2010)]

- Evolving (Vanishing) Dimensions

[Anchordoqui et al., PRD83,114046 (2011);
JRM and Stojkovic, PRL 106, 101101 (2011);
Anchordoqui et al., MPLA27,1250021 (2012)]

- Brane-world

[C. Callan, J. Maldacena, Nucl.Phys. B513, 198 (1998),
N. Constable, R. Myers, O. Tafjord, PRD61:106009 (2000)]

- Multifractal gravity

[Calcagni, PRL 104, 251301(2010);
Calcagni PLB697, 251 (2011);
Arzano *et al.*, PRD84:125002 (2011)]

Black Holes are not Particles

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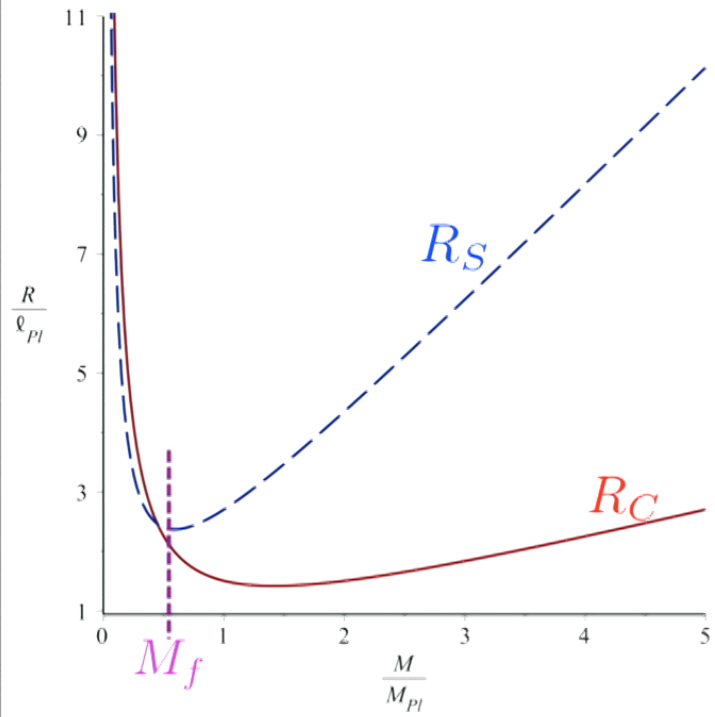
- Can ensure that the physical length scale of particles does not fall within their gravitational radius (*e.g.* neutrinos)
- Define two ways of approaching Planck scale [Carr 2013]:

- Generalized Compton:
$$R_C = \frac{1}{M} \left[1 + \alpha \left(\frac{M}{M_{\text{Pl}}} \right)^2 \right]$$

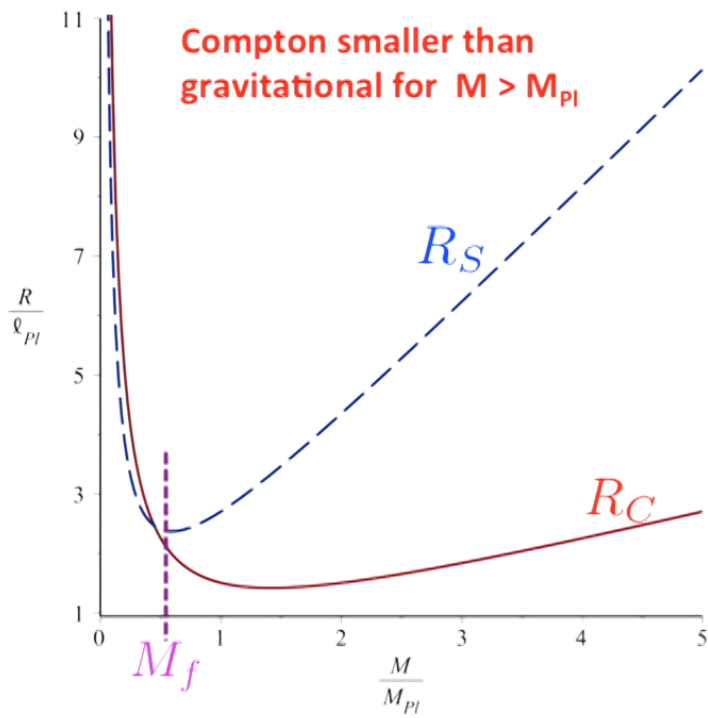
- Generalized event horizon:
$$R_S = 2M \left[1 + \frac{\beta}{2} \left(\frac{M_{\text{Pl}}}{M} \right)^2 \right]$$

- Coincide at the mass scale

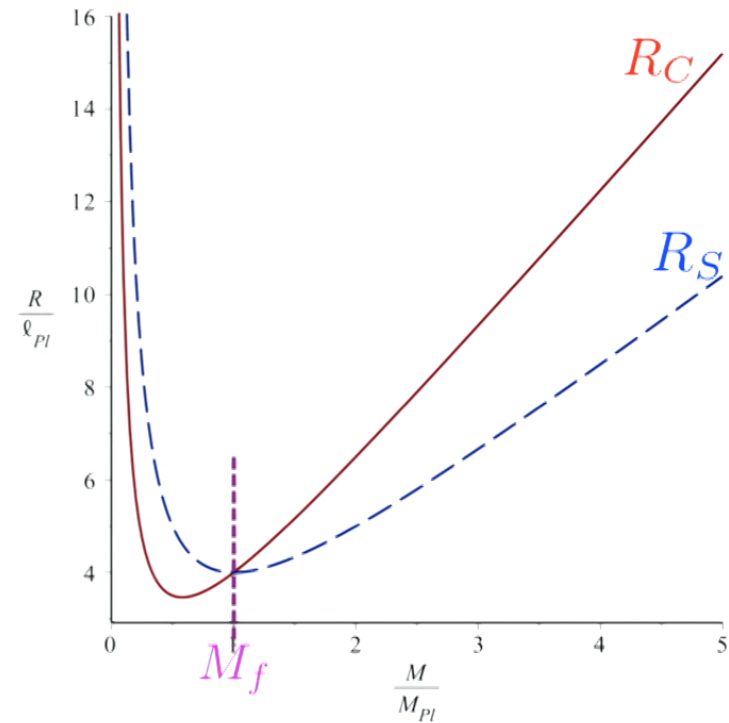
$$M_f = \sqrt{\frac{1 - \beta}{2 - \alpha}} M_{\text{Pl}},$$



$$\alpha = 0.5; \beta = 0.7$$



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$$\alpha = 3; \beta = 2$$

Production of Sub-Planckian BHs

- How can we understand production of (sub-Planckian) black holes in collisions?

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- Horizon Wavefunction Formalism

Given a particle created in e.g. a collision, what is the probability that...

1. ... the distribution is within a spherical region $R \leq r_H$
2. ... the particle's gravitational (horizon) radius is r_H
3. ... the particle lies completely within r_H

Production of Sub-Planckian BHs

- Given an eigenstate

$$\hat{H}|\psi_E\rangle = E|\psi_E\rangle \quad E = \frac{M_{\text{Pl}} R_H}{2\ell_{\text{Pl}}} \leftarrow \text{Hoop conjecture}$$

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$$|\psi_S\rangle = \sum_E C(E)|\psi_E\rangle \quad \psi_H(r_H) = C(M_{\text{Pl}} r_H / 2\ell_{\text{Pl}})$$

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- Define the spatial distribution of the particle:

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- Probability of being a black hole

- The particle lies completely within r_H , and r_H is its horizon radius!

Production of Sub-Planckian BHs

- Given an eigenstate

$$\hat{H}|\psi_E\rangle = E|\psi_E\rangle \quad E = \frac{M_{\text{Pl}} R_H}{2\ell_{\text{Pl}}} \quad \leftarrow \text{Hoop conjecture}$$

- Define the horizon wavefunction:

$$|\psi_S\rangle = \sum_E C(E)|\psi_E\rangle \quad \psi_H(r_H) = C(M_{\text{Pl}} r_H / 2\ell_{\text{Pl}})$$

- Define the spatial distribution of the particle:

$$\psi_S(r) = \text{Some (gaussian) distribution width } \ell \geq \lambda_C$$

- Probability of being a black hole

- The particle lies completely within r_H , and r_H is its horizon radius!

Production of Sub-Planckian BHs

[Casadio, Cavalcanti, Giugno, JM, PLB 760:36 (2016)]

- Probability that particle is inside a D-ball of radius :

$$P_S(r < r_H) = \Omega_{D-1} \int_0^{r_H} |\psi_S(r)|^2 r^{D-1} r$$

- Probability that radius of the horizon is r_H :

$$\mathcal{P}_H(r_H) = \Omega_{D-1} r^{D-1} |\psi_H(r_H)|^2$$

- Probability that particle lies within its horizon radius:

$$\mathcal{P}_{<}(r < r_H) = P_S(r < r + H) \mathcal{P}_H(r_H)$$

- **Probability that particle is a black hole:**

$$P_{\text{BH}} = \int_0^\infty \mathcal{P}_{<}(r < r_H) dr_H$$

The (D+1)-dimensional Schwarzschild metric is...

$$ds^2 = \left(1 - \frac{R_D}{r^{D-2}}\right) dt^2 - \left(1 - \frac{R_D}{r^{D-2}}\right)^{-1} dr^2 - r^{D-1} d\Omega_{D-1}$$

where the horizon radius is....

$$R_D = \left(\frac{2G_D M}{D-2}\right)^{\frac{1}{D-2}} \quad \text{if } D > 2 \quad = \quad \frac{1}{2G_1 M} \quad \text{if } D = 1$$

and Gravitational constant is defined as

$$G_D = \frac{\ell_D^{D-2}}{m_D}$$



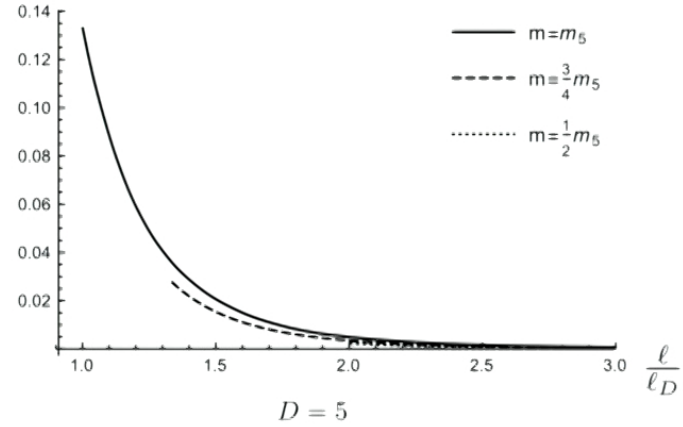
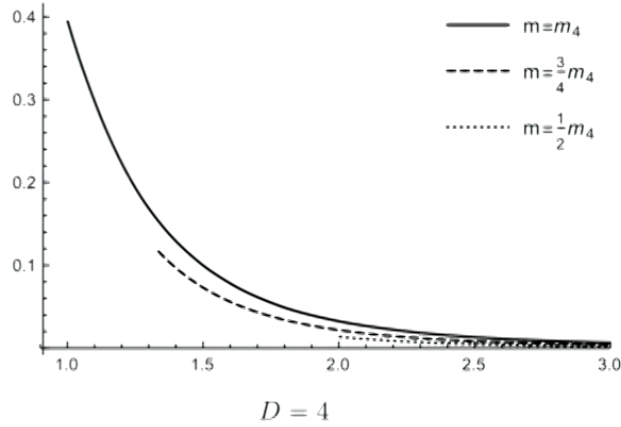
Matches an exact solution!

[Mann, Shiekh, Tarasov, NPB 341, 134 (1990)]

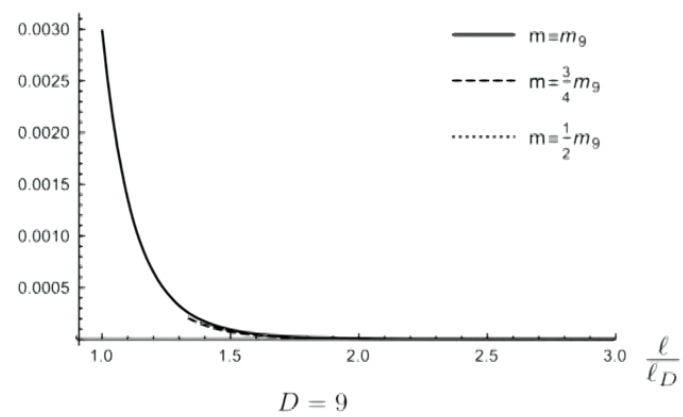
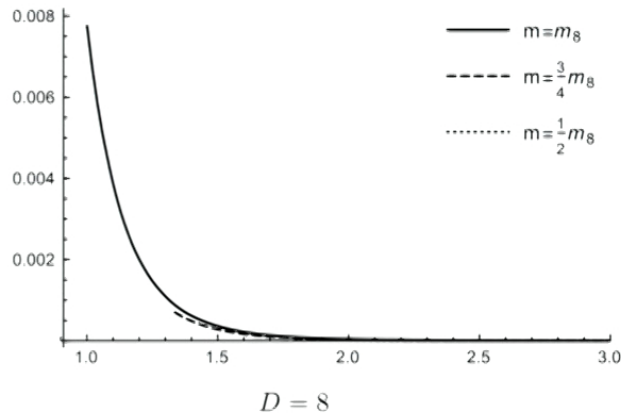
In (D+1)-dimensions...

$$P_{\text{BH}} = \frac{2}{\ell_D^D} \left(\frac{(D-2)m_D}{2\Delta} \right)^{\frac{D}{D-2}} \frac{D-2}{\Gamma\left(\frac{D}{2D-4}, \frac{m^2}{\Delta^2}\right) \Gamma\left(\frac{D}{2}\right)} \times \int_{R_D}^{\infty} \gamma\left(\frac{D}{2}, \frac{r_H^2}{\ell_D^2}\right) \exp\left\{-\left[\frac{(D-2)m_D}{2\Delta}\right]^2 \left(\frac{r_H}{\ell_D}\right)^{2(D-2)}\right\} r_H^{D-1} dr_H$$

P_{BH}



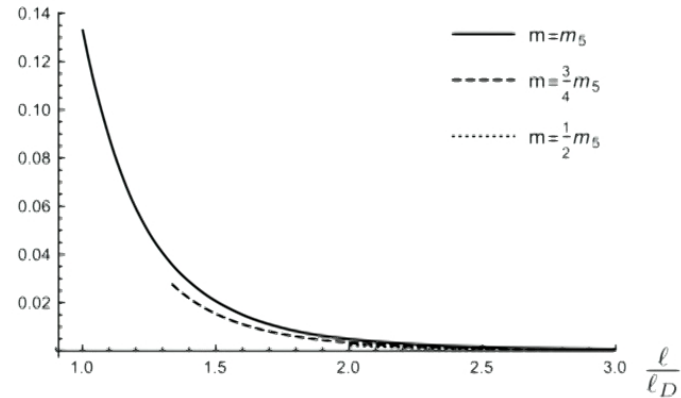
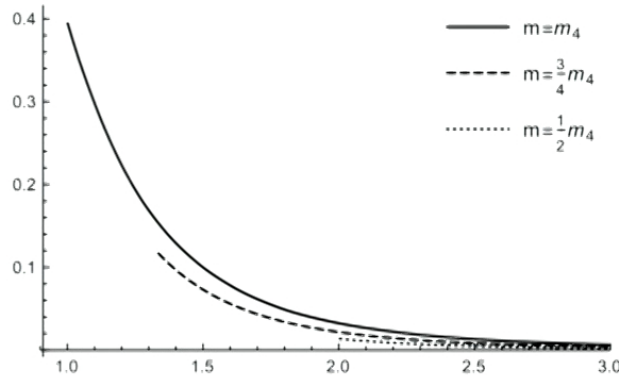
P_{BH}



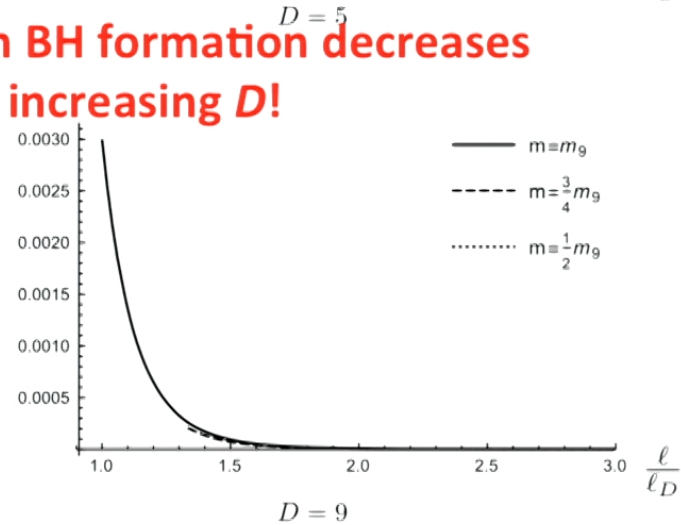
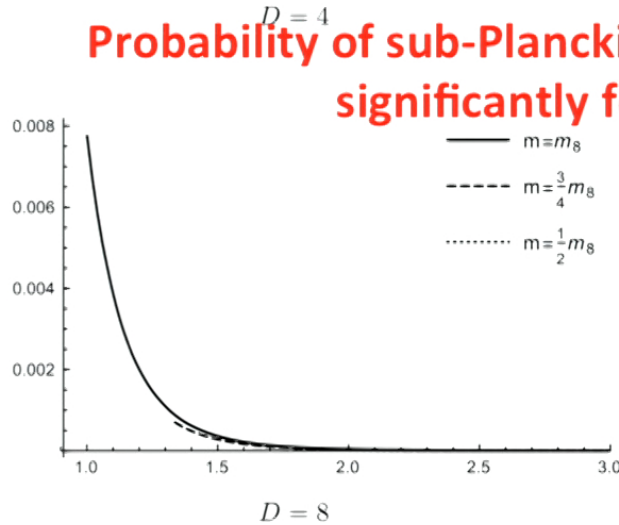
In (D+1)-dimensions...

$$P_{\text{BH}} = \frac{2}{\ell_D^D} \left(\frac{(D-2)m_D}{2\Delta} \right)^{\frac{D}{D-2}} \frac{D-2}{\Gamma\left(\frac{D}{2D-4}, \frac{m^2}{\Delta^2}\right) \Gamma\left(\frac{D}{2}\right)} \times \int_{R_D}^{\infty} \gamma\left(\frac{D}{2}, \frac{r_H^2}{\ell_D^2}\right) \exp\left\{-\left[\frac{(D-2)m_D}{2\Delta}\right]^2 \left(\frac{r_H}{\ell_D}\right)^{2(D-2)}\right\} r_H^{D-1} dr_H$$

P_{BH}



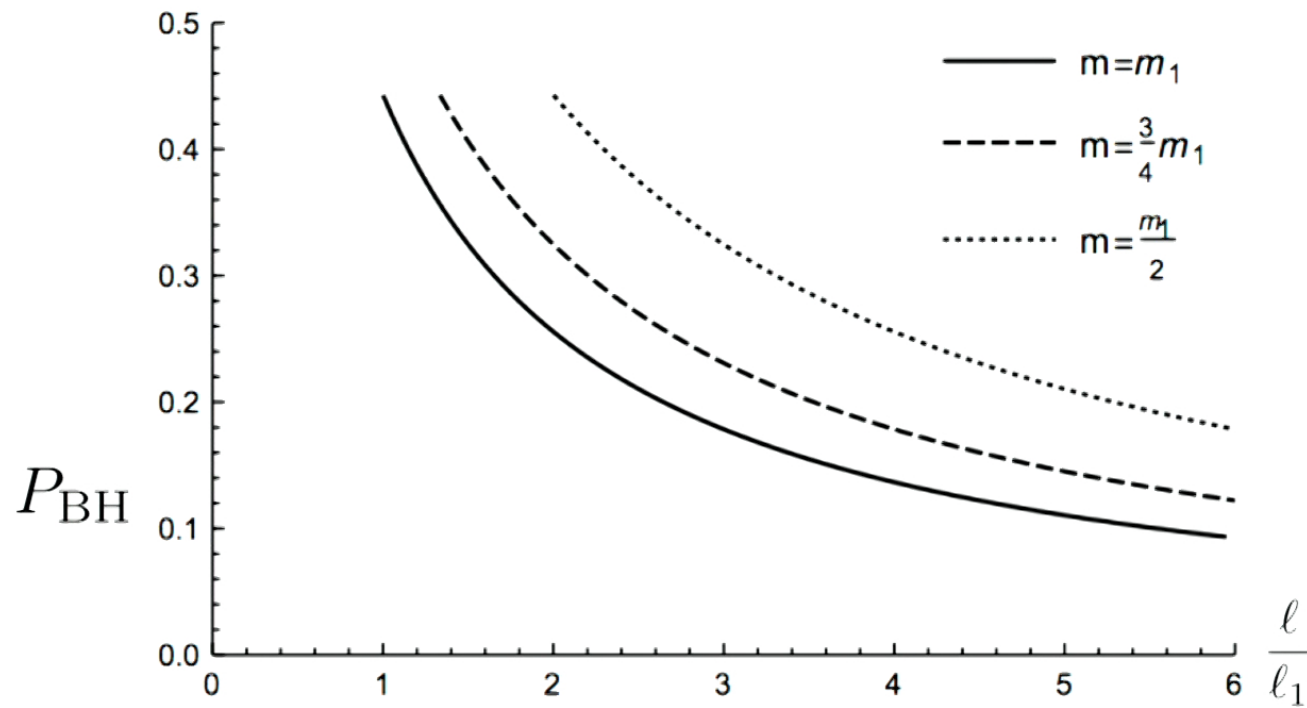
P_{BH}



Probability of sub-Planckian BH formation decreases significantly for increasing D !

In (1+1)-dimensions...

$$P_{\text{BH}} = \frac{4/\ell}{\Gamma\left(-\frac{1}{2}, \frac{m^2}{\Delta^2}\right)} \int_0^{R_1} \text{erf}\left(\frac{r_H}{\ell}\right) \exp\left\{-\frac{\ell^2}{4R_H^2}\right\} dr_H$$



HWF and GUP

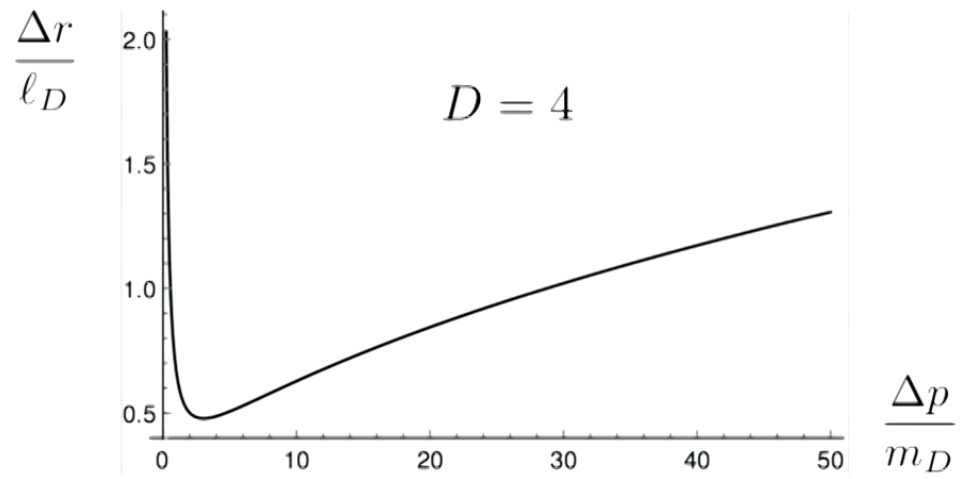
[Casadio, Cavalcanti, Giugno, JM, PLB 760:36 (2016)]

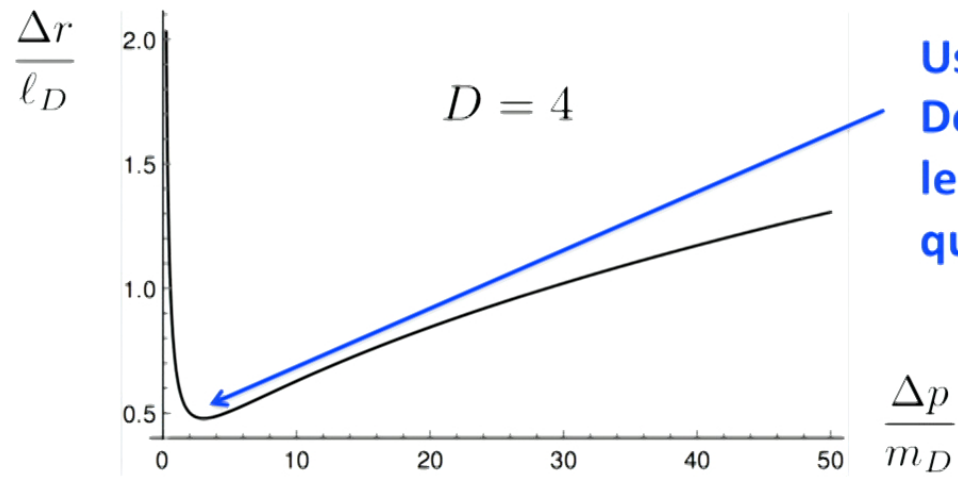
- *The Horizon Wavefunction formalism naturally implies the GUP*
- **Horizon and position expectations and uncertainties:**

$$\langle r_H^n \rangle = \Omega_{D-1} \int_0^\infty \tilde{\psi}_H^*(r) r_H^n \tilde{\psi}_H(r) r^{D-1} dr$$

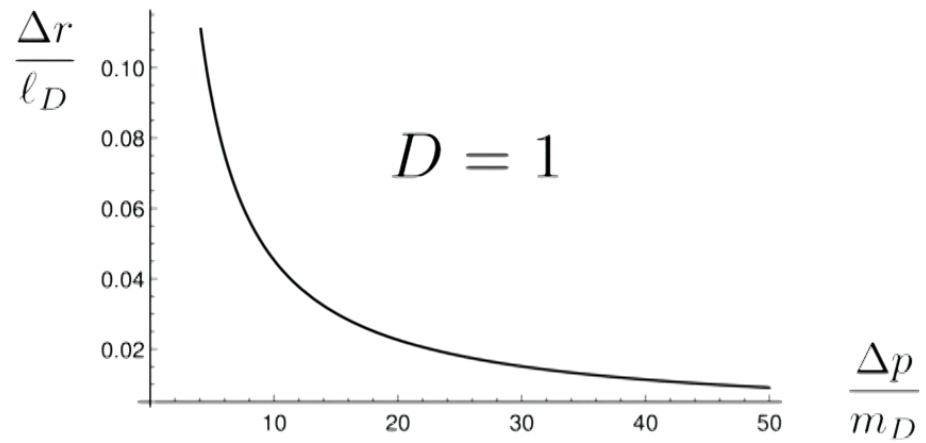
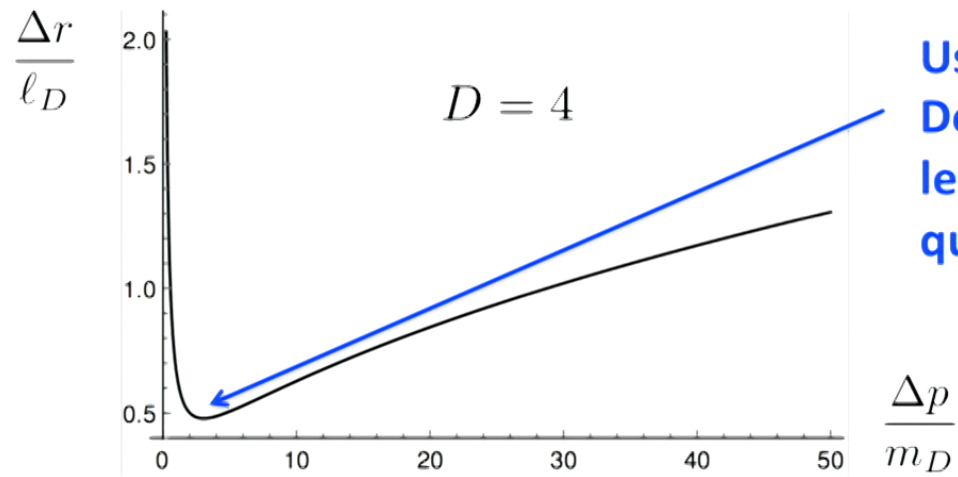
$$\langle r^n \rangle = \Omega_{D-1} \int_0^\infty \tilde{\psi}_S^*(r) r^n \tilde{\psi}_S(r) r^{D-1} dr$$

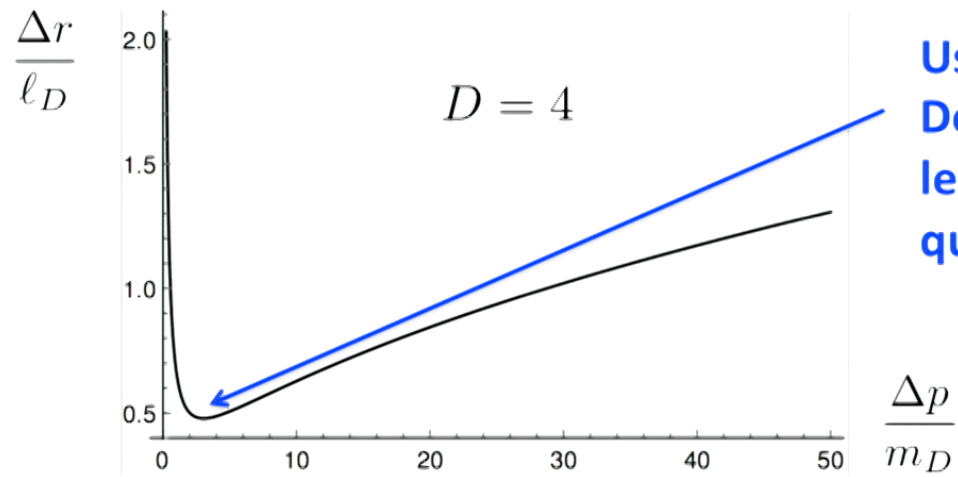
- **Total spatial uncertainty:** $\Delta r_{\text{tot}} = \Delta r + \alpha \Delta r_H$
- **Define momentum uncertainty similarly using** $\psi_S(p)$
- **Find Δr_{tot} in terms of Δp**



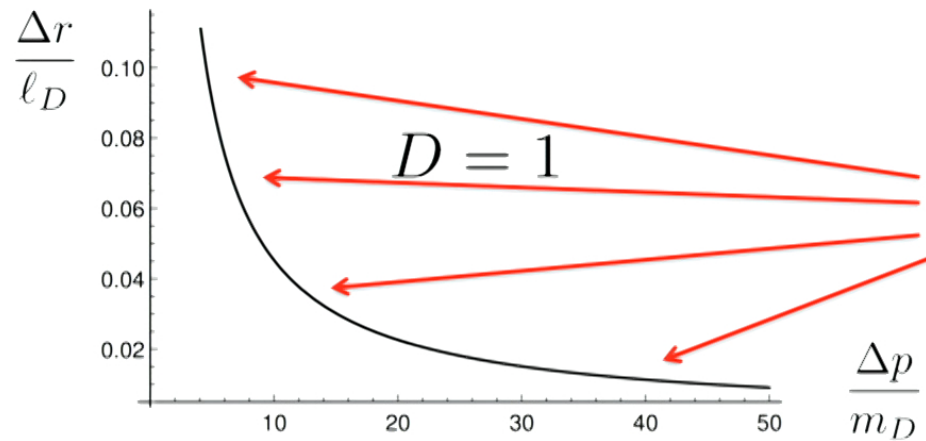


**Usual GUP;
 Defines minimum
 length scale (classical/
 quantum “boundary”)**





Usual GUP;
 Defines minimum
 length scale (classical/
 quantum "boundary")



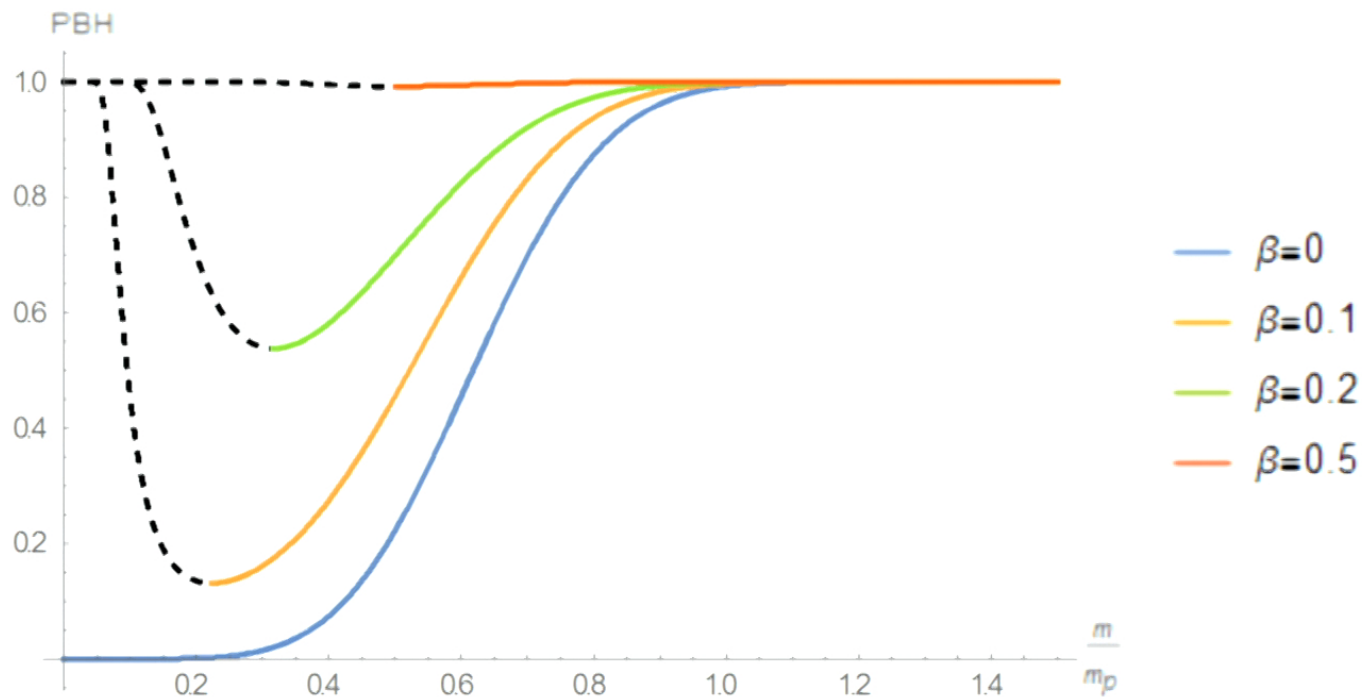
No minimum
 scale; theory
 is purely
 "quantum"

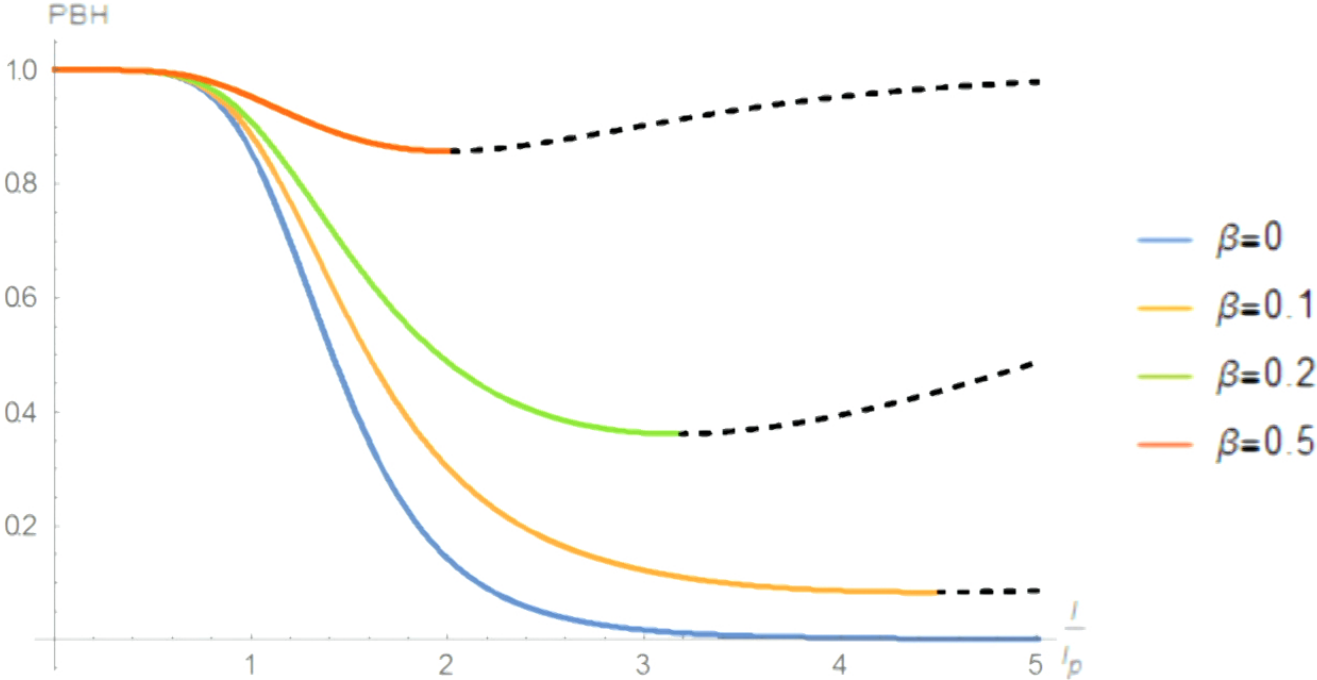
GUP Black Hole Production and HWF

Manfredi and JM, AHEP 2016:1543741 (2016)

- Given the GUP metric described earlier, the HWF is

$$\psi_H(r_H) = \frac{1}{4\ell_{\text{Pl}}^3} \sqrt{\frac{\ell^3}{\pi\Gamma\left(\frac{3}{2}, 1\right)}} \Theta(r_H - R_H) e^{-\frac{\ell^2 r_H^2}{8\ell^4}}$$





Conclusions and Future Musings

- “Encoding” the GUP duality in the mass gives a metric that exhibits dimensional reduction in the sub-Planckian regime
- Cures thermodynamic instability of evaporating Schwarzschild black holes
- Production of quantum black holes is “easier” in lower dimensions than higher (particularly sub-Planckian)
- GUP emerges again as a natural characteristic of gravitation
- Instead of a two regimes governed by **different theories** (GR and QM), we have a consistent theory (gravity) in two **different spacetime dimensions**



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Frankfurt Institute for Advanced Studies
24-28 July 2017

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Thank you!

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