

Title: The asymptotic safety paradigm for quantum gravity and matter

Date: Mar 02, 2017 02:30 PM

URL: <http://pirsa.org/17030050>

Abstract: <p>In this talk, I will discuss the asymptotic safety paradigm, and will highlight that it can provide a framework for a predictive ultraviolet completion for gravity and matter. Specifically, I will discuss compelling hints that exist for the realization of asymptotic safety in pure gravity, and will then present recent progress on the case of gravity coupled to Standard Model matter. In particular, I will highlight results that show how to forge a link between physics at the Planck scale and physics at the electroweak scale, in order to impose observational constraints on the microscopic quantum gravity dynamics.</p>

The asymptotic safety paradigm for gravity & matter

Astrid Eichhorn
University of Heidelberg

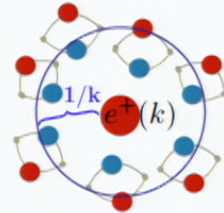
based on work with
N. Christiansen, A. Held,
S. Lippoldt, J. Pawłowski



March 2, 2017
Quantum Gravity Seminar, Perimeter Institute for Theoretical Physics

The asymptotic safety paradigm

$$\int \mathcal{D}\varphi e^{-S[\varphi]}$$

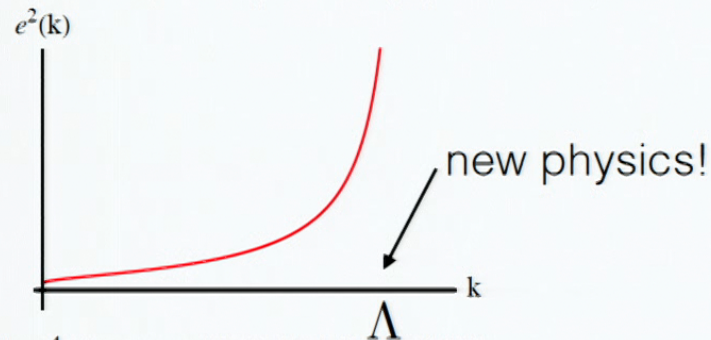


quantum fluctuations generate scale dependent couplings

UV behavior:

effective theory:

Landau pole (triviality problem)

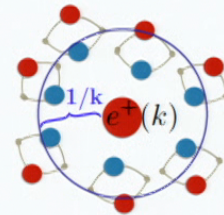


$\lambda \varphi^4$ [Aizenmann '81; Frohlich '82, Wolff '09]

U(1) gauge theory [Gell-Mann, Low '54; Gockeler et al. '97; Gies, Jaeckel '04]

The asymptotic safety paradigm

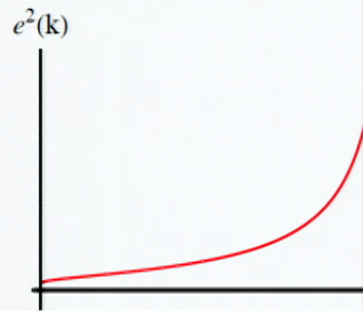
$$\int \mathcal{D}\varphi e^{-S[\varphi]}$$



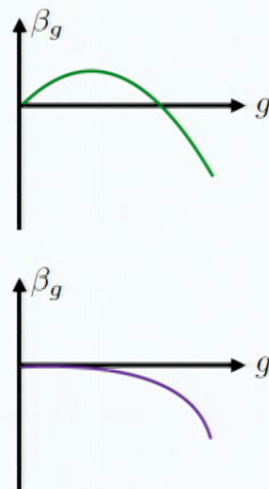
quantum fluctuations generate scale dependent couplings

UV behavior:

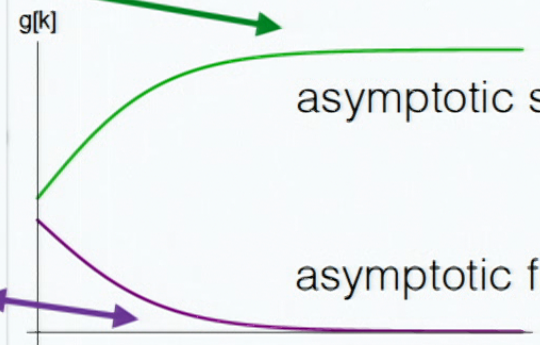
effective theory:
Landau pole (trivial



$\lambda \varphi^4$ [Aizenmann '81; Frohlich '82]



fundamental theory:
scale-invariance (RG fixed point)

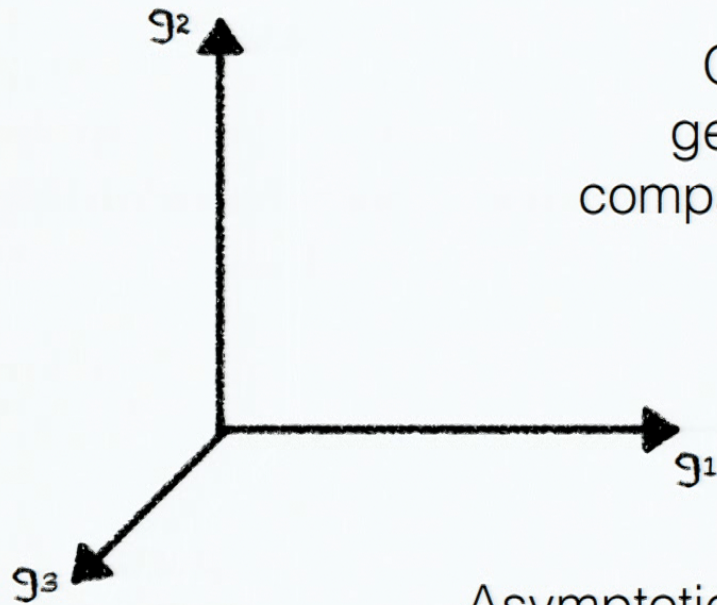


asymptotic safety

asymptotic freedom

U(1) gauge theory [Gell-Mann, Low '54; Gockeler et al. '97; Gies, Jaeckel '04]

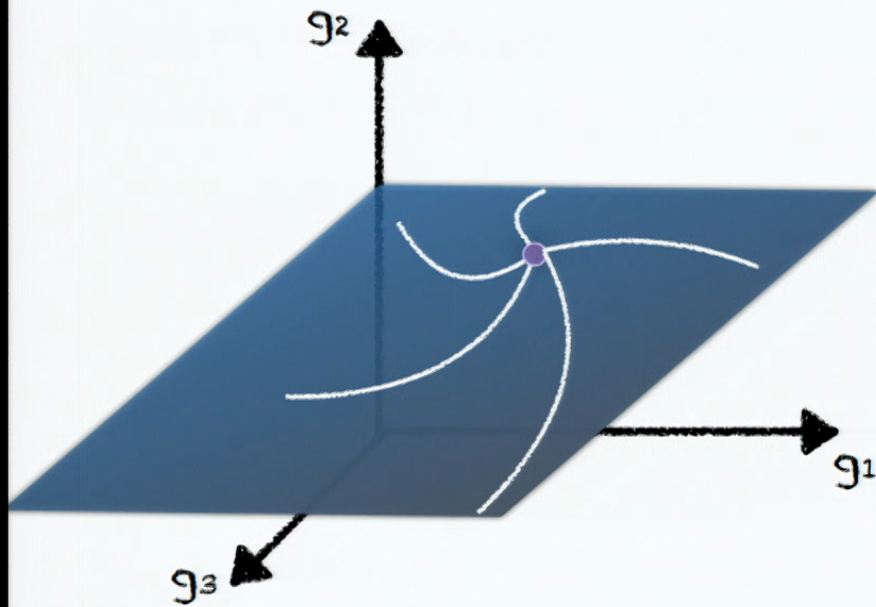
Theory space



Quantum fluctuations
generate all interactions
compatible with the symmetries

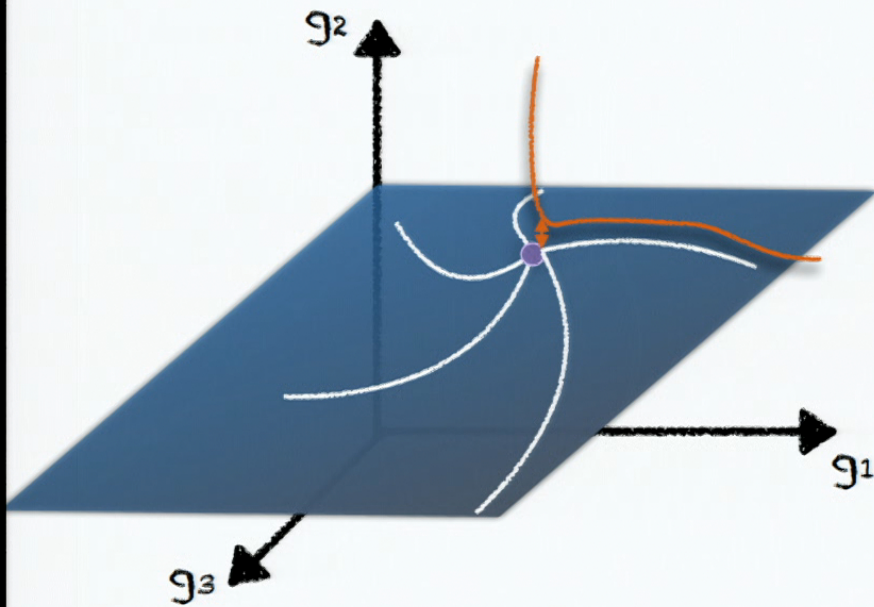
Asymptotic safety/freedom:
Fixed point in infinite-dimensional theory space

Irrelevant directions: Predictions from asymptotic safety



UV-critical surface:
UV-attractive directions
 \leftrightarrow free parameters

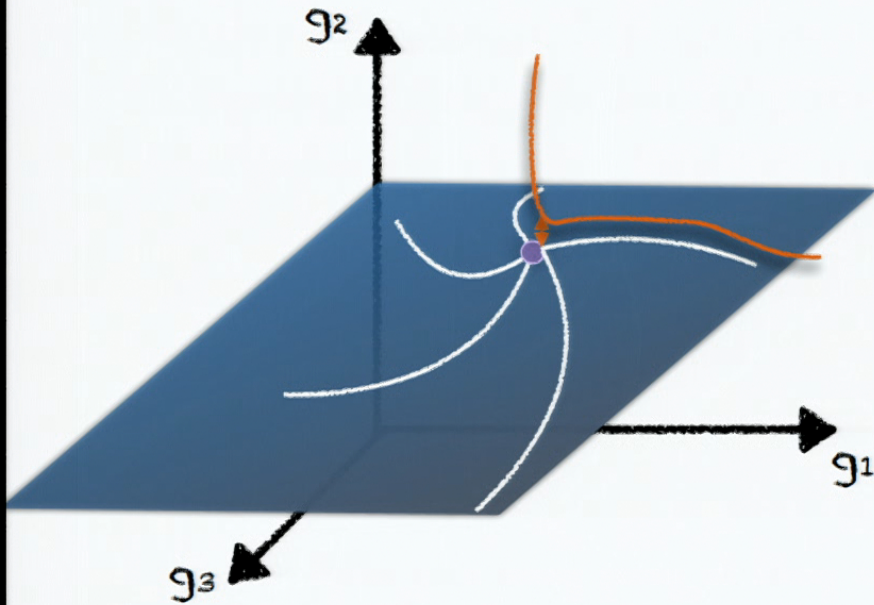
Irrelevant directions: Predictions from asymptotic safety



UV-critical surface:
UV-attractive directions
 \leftrightarrow free parameters

UV-repulsive directions
(irrelevant couplings)
 \leftrightarrow predictions

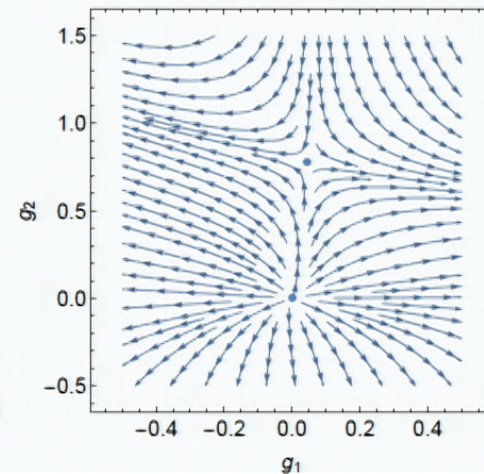
Irrelevant directions: Predictions from asymptotic safety



UV-critical surface:
UV-attractive directions
 \leftrightarrow free parameters

UV-repulsive directions
(irrelevant couplings)
 \leftrightarrow predictions

towards infrared:
irrelevant couplings determined
(QG fluctuations force irrelevant couplings to specific values)
asymptotically free models:
couplings with negative canonical dimension

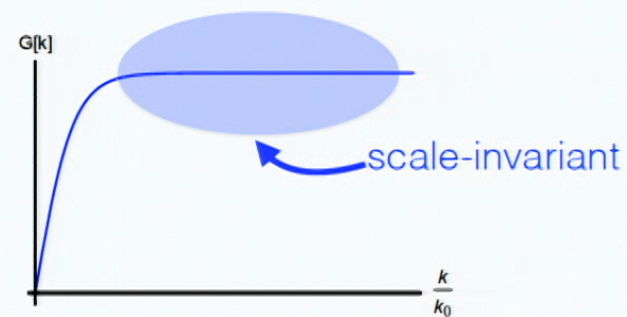
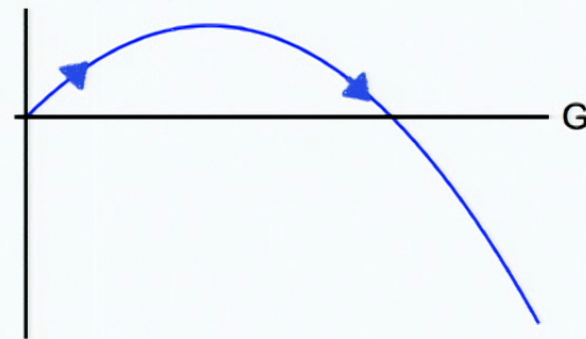


Mechanism for asymptotic safety

Balancing canonical
and quantum scaling:

$$[G_N] = 2 - d \rightarrow \text{dim'less: } G = G_N k^{d-2}$$

$$\beta_G = (d-2)G \pm G^2$$



Mechanism for asymptotic safety

Balancing canonical and quantum scaling:

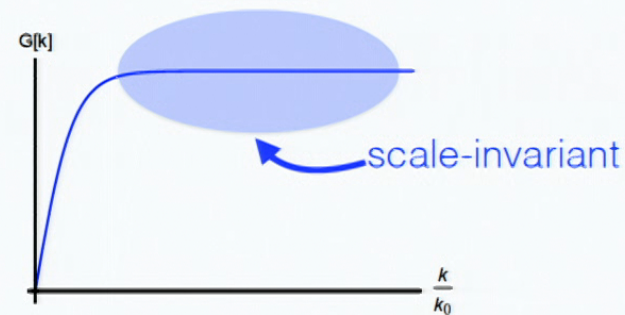
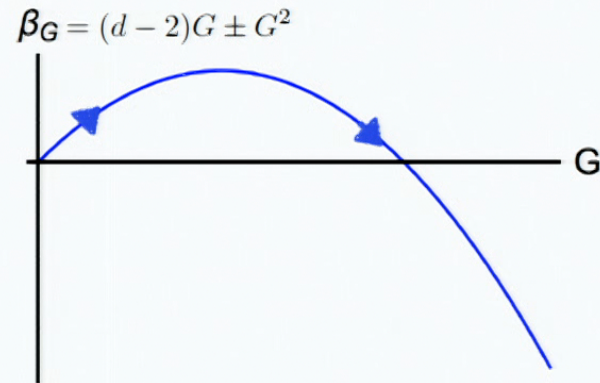
ϵ expansion in $d = 2 + \epsilon$

$$\beta_G = \epsilon G - \frac{38}{3} G^2$$

Weinberg '76; Christensen, Duff '78
Gastmans, Kallosh, Truffin '78

Goal: extension to $d=4$?

$$[G_N] = 2 - d \rightarrow \text{dim'less: } G = G_N k^{d-2}$$



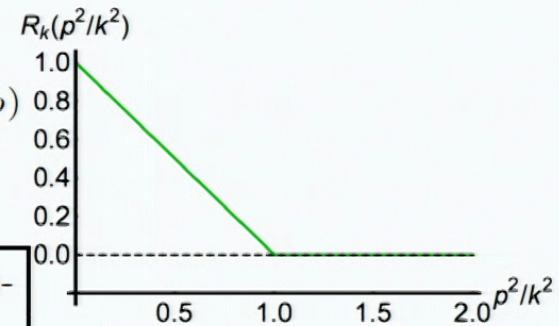
Functional Renormalization Group

probe scale dependence of QFT

$$e^{-\Gamma_k[\phi]} = \int \mathcal{D}\varphi e^{-S[\varphi] - \frac{1}{2} \int \varphi(-p) R_k(p) \varphi(p)}$$

Wetterich '93; Morris, '94

scale- and momentum-dependent "mass"



S

Γ_k contains effect of quantum fluctuations above k

$$\Gamma_k = \sum_i g_i(k) \int d^d x \mathcal{O}^i$$



$\Gamma_{k \rightarrow 0}$

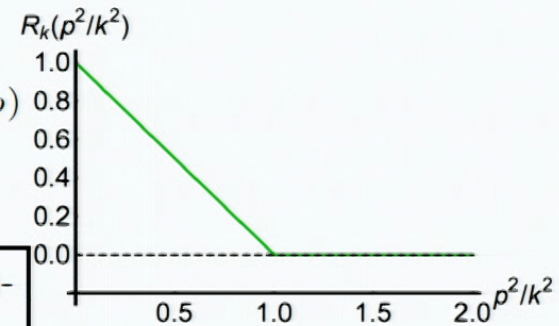
Functional Renormalization Group

probe scale dependence of QFT

$$e^{-\Gamma_k[\phi]} = \int \mathcal{D}\varphi e^{-S[\varphi] - \frac{1}{2} \int \varphi(-p) R_k(p) \varphi(p)}$$

Wetterich '93; Morris, '94

scale- and momentum-dependent "mass"



S

prediction of asymptotic safety

Γ_k contains effect of quantum fluctuations above k

$$\Gamma_k = \sum_i g_i(k) \int d^d x \mathcal{O}^i$$

$$\rightarrow k \partial_k \Gamma_k = \sum_i \beta_{g_i} \int d^d x \mathcal{O}^i$$



$\Gamma_{k \rightarrow 0}$

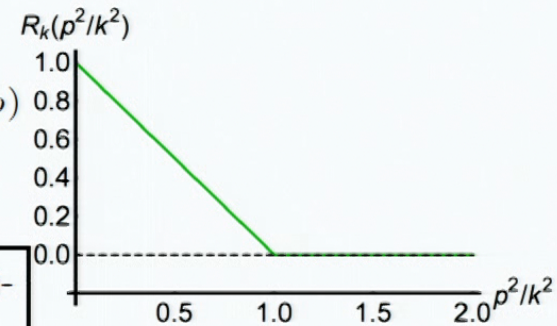
Functional Renormalization Group

probe scale dependence of QFT

$$e^{-\Gamma_k[\phi]} = \int \mathcal{D}\varphi e^{-S[\varphi] - \frac{1}{2} \int \varphi(-p) R_k(p) \varphi(p)}$$

Wetterich '93; Morris, '94

scale- and momentum-dependent "mass"



Wetterich equation:

$$\partial_k \Gamma_k = \frac{1}{2} \text{STr} \left(\Gamma_k^{(2)} + R_k \right)^{-1} \partial_k R_k = \text{Diagram}$$

$$\rightarrow k \partial_k \Gamma_k = \sum_i \beta_{g_i} \int d^d x \mathcal{O}^i$$

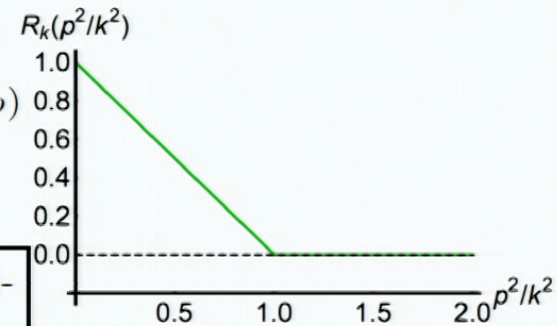
Functional Renormalization Group

probe scale dependence of QFT

$$e^{-\Gamma_k[\phi]} = \int \mathcal{D}\varphi e^{-S[\varphi] - \frac{1}{2} \int \varphi(-p) R_k(p) \varphi(p)}$$

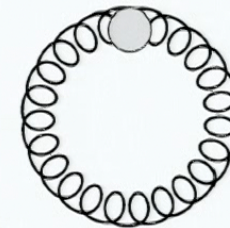
Wetterich '93; Morris, '94

scale- and momentum-dependent "mass"



Wetterich equation:

$$\partial_k \Gamma_k = \frac{1}{2} \text{STr} \left(\Gamma_k^{(2)} + R_k \right)^{-1} \partial_k R_k =$$



practical calculations: truncation

$$\Gamma_k = \sum_i g_i(k) \int d^d x \mathcal{O}^i$$

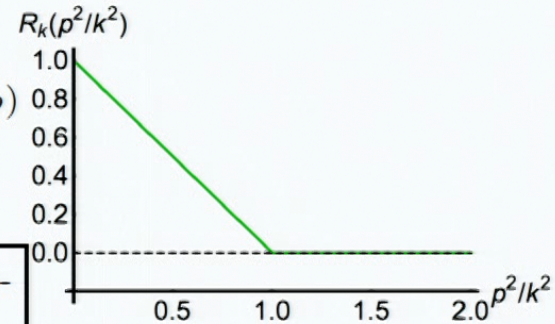
Functional Renormalization Group

probe scale dependence of QFT

$$e^{-\Gamma_k[\phi]} = \int \mathcal{D}\varphi e^{-S[\varphi] - \frac{1}{2} \int \varphi(-p) R_k(p) \varphi(p)}$$

Wetterich '93; Morris, '94

scale- and momentum-dependent "mass"



Wetterich equation:

$$\partial_k \Gamma_k = \frac{1}{2} \text{STr} \left(\Gamma_k^{(2)} + R_k \right)^{-1} \partial_k R_k =$$



Asymptotically Safe Gravity: Reuter, '96

background field method: $g_{\mu\nu} = \bar{g}_{\mu\nu} + h_{\mu\nu}$

→ shift symmetry: $(\bar{g}_{\mu\nu} + \epsilon_{\mu\nu}) + (h_{\mu\nu} - \epsilon_{\mu\nu}) = g_{\mu\nu}$

→ distinction between background & fluctuation couplings

$$\int \mathcal{D}g_{\mu\nu} \rightarrow \int \mathcal{D}h_{\mu\nu}$$

$$h_{\mu\nu} R_k (\bar{D}^2)^{\mu\nu\kappa\lambda} h_{\kappa\lambda}$$

Manrique, Reuter '09; Manrique, Reuter, Saueressig '10, '11; Christiansen, Litim, Pawłowski, Rodigast '12; Becker, Reuter '14; Christiansen, Knorr, Meibohm, Pawłowski, Reichert '15; Denz, Pawłowski, Reichert '16

shift Ward identity Dietz, Morris '15; Morris '16; Percacci, Vacca '16

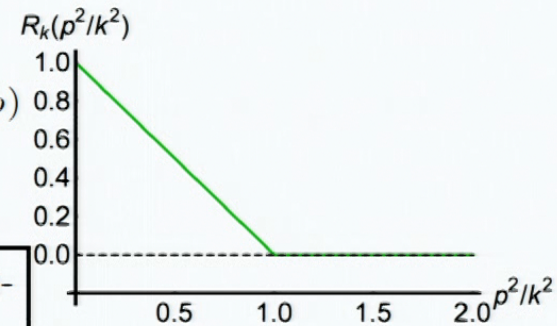
Functional Renormalization Group

probe scale dependence of QFT

$$e^{-\Gamma_k[\phi]} = \int \mathcal{D}\varphi e^{-S[\varphi] - \frac{1}{2} \int \varphi(-p) R_k(p) \varphi(p)}$$

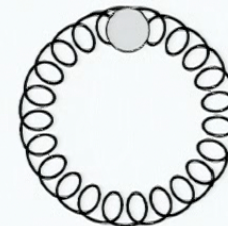
Wetterich '93; Morris, '94

scale- and momentum-dependent "mass"



Wetterich equation:

$$\partial_k \Gamma_k = \frac{1}{2} \text{STr} \left(\Gamma_k^{(2)} + R_k \right)^{-1} \partial_k R_k =$$



matrix/tensor models: A.E., Koslowski '13, '14, '17

group field theories:

Benedetti, Ben Geloun, Oriti '14;
Benedetti, Lahoche '15; Ben Geloun, Martini, Oriti '15, '16;
Ben Geloun, Koslowski '16; Carrozza, Lahoche '16

Functional Renormalization Group & ϵ - expansion

perturbative evidence for asymptotic safety:

expansion in $d = 2 + \epsilon$

$$\beta_G = \epsilon G - \frac{38}{3} G^2$$

Weinberg '76; Christensen, Duff '78
Gastmans, Kallosh, Truffin '78

Goal: Match onto $d > 2$ results
and extend to $d=4$

FRG suitable tool:

extend pert. controlled FP away from crit. dim.

→ see non-gravity models

Functional Renormalization Group & ϵ - expansion

perturbative evidence for asymptotic safety:

expansion in $d = 2 + \epsilon$

$$\beta_G = \epsilon G - \frac{38}{3} G^2$$

Weinberg '76; Christensen, Duff '78
Gastmans, Kallosh, Truffin '78

PHYSICAL REVIEW D 90, 025018 (2014)

\mathfrak{S}

Critical $O(N)$ models in $6 - \epsilon$ dimensions

Lin Fei,¹ Simone Giombi,¹ and Igor R. Klebanov^{1,2}

- unitary UV completion of $O(N)$ models in $d > 4$

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \phi_i)^2 + \frac{1}{2}(\partial_\mu z)^2 + \frac{1}{2}gz\phi_i\phi_i + \frac{1}{6}\lambda z^3$$

Goal: Match onto $d > 2$ results
and extend to $d=4$

FRG suitable tool:

extend pert. controlled FP away from crit. dim.

→ see non-gravity models

Functional Renormalization Group & ϵ - expansion

perturbative evidence for asymptotic safety:

expansion in $d = 2 + \epsilon$

$$\beta_G = \epsilon G - \frac{38}{3} G^2$$

PHYSICAL REVIEW D 90, 025018 (2014)

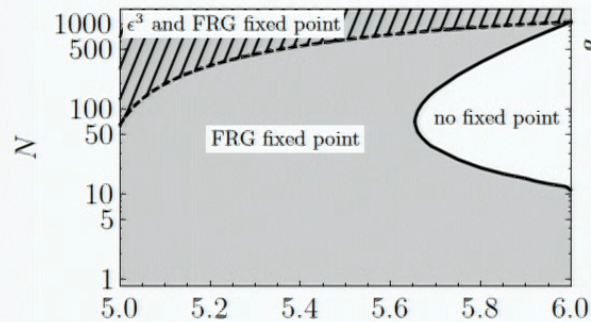
\mathfrak{G}
Critical $O(N)$ models in $6 - \epsilon$ dimensions

Lin Fei,¹ Simone Giombi,¹ and Igor R. Klebanov^{1,2}

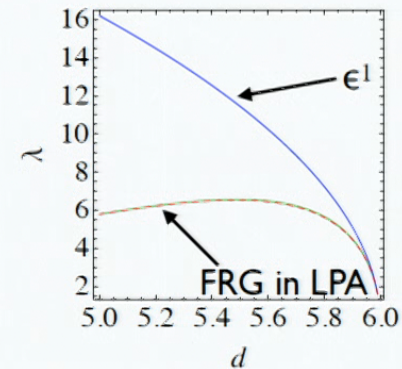
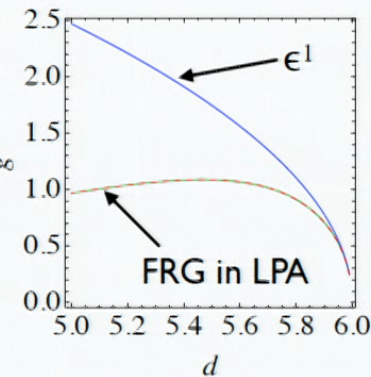
- unitary UV completion of $O(N)$ models in $d > 4$

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \phi_i)^2 + \frac{1}{2}(\partial_\mu z)^2 + \frac{1}{2}gz\phi_i\phi_i + \frac{1}{6}\lambda z^3$$

Weinberg '76; Christensen, Duff '78
Gastmans, Kallosh, Truffin '78



d [A.E., Janssen, Scherer '16]



Functional Renormalization Group & ϵ - expansion

perturbative evidence for asymptotic safety:

FRG in $d = 2 + \epsilon$

expansion in $d = 2 + \epsilon$

$$\beta_G = \epsilon G - \frac{38}{3} G^2$$

$$\beta_G = \epsilon G - \frac{38}{3} G^2$$

Nink '15, Falls '15, '17

(regulator- independent)

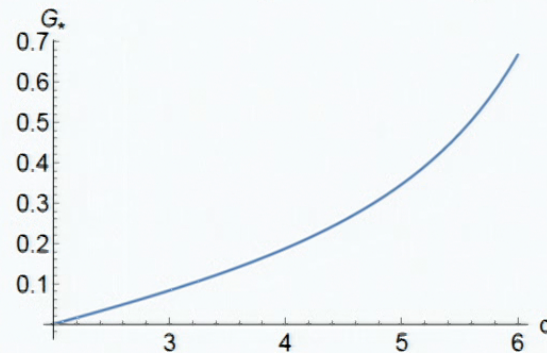
Weinberg '76; Christensen, Duff '78
Gastmans, Kallosh, Truffin '78

continuation of fixed point towards $d=4$?

$$\beta_G = (d - 2)G - \frac{2}{3} \left(18 - \frac{d(d - 3)}{2} \right) G^2$$

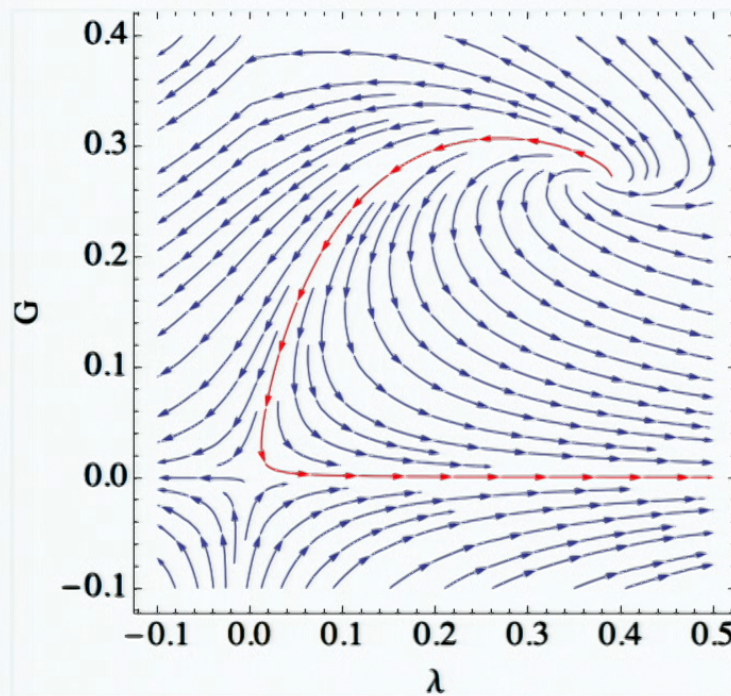
Falls '15

interacting fixed point
for gravity:
well-controlled limit for $d \rightarrow 2$



The case for asymptotic safety in gravity

$$\Gamma_k = -\frac{1}{16\pi G_N(k)} \int d^4x \sqrt{g} (R - 2\bar{\lambda}(k)) \quad G(k) = G_N(k)k^2, \quad \lambda(k) = \bar{\lambda}(k)/k^2$$

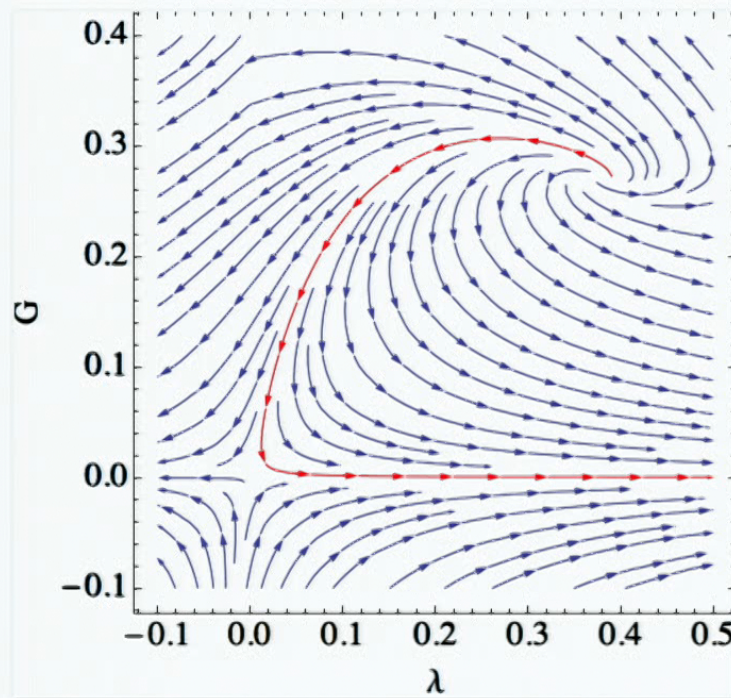


interacting fixed point
with ~ 3 (?) relevant directions

Reuter, '96; Reuter, Saueressig '01, Litim, '03

The case for asymptotic safety in gravity

$$\Gamma_k = -\frac{1}{16\pi G_N(k)} \int d^4x \sqrt{g} (R - 2\bar{\lambda}(k)) \quad G(k) = G_N(k)k^2, \quad \lambda(k) = \bar{\lambda}(k)/k^2$$

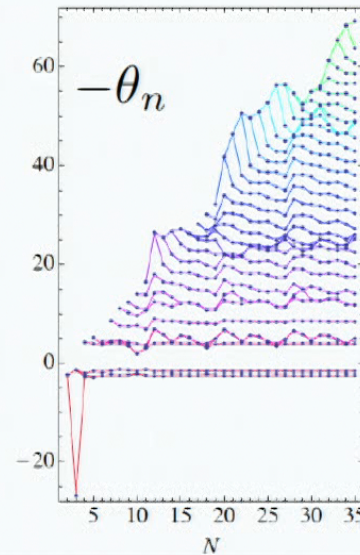


Reuter, '96; Reuter, Saueressig '01, Litim, '03

interacting fixed point
with ~ 3 (?) relevant directions

$$f(R) = \sum_{n=0}^N a_n R^n$$

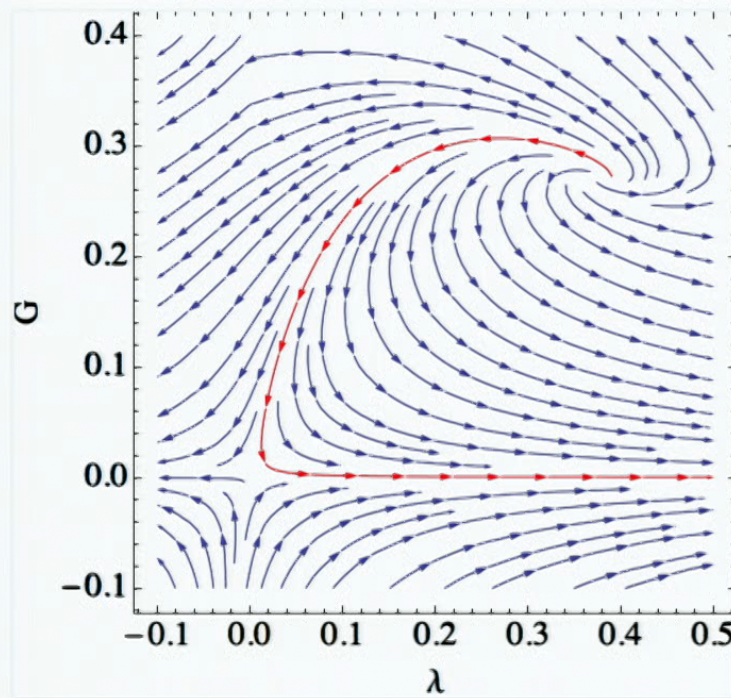
Reuter, Lauscher, '02; Codello, Percacci, Rahmede, '09;
Benedetti, Caravelli, '12; Dietz, Morris, '12;
Falls, Litim, Nikolakopoulos, Rahmede, '13, '14
Demmel, Saueressig, Zanusso, '15; Eichhorn '15



- towards convergence in extensions of the truncation
- near-canonical scaling: guidance for well-controlled truncations

The case for asymptotic safety in gravity

$$\Gamma_k = -\frac{1}{16\pi G_N(k)} \int d^4x \sqrt{g} (R - 2\bar{\lambda}(k)) \quad G(k) = G_N(k)k^2, \quad \lambda(k) = \bar{\lambda}(k)/k^2$$

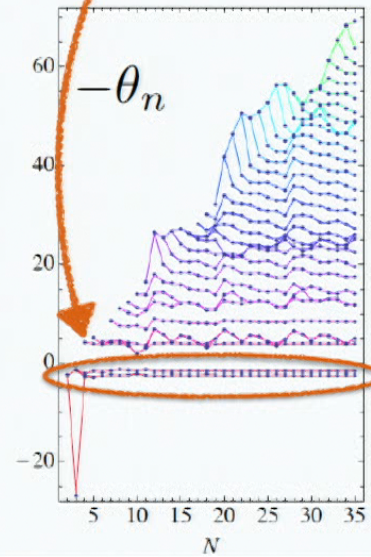


Reuter, '96; Reuter, Saueressig '01, Litim, '03

interacting fixed point
with ~ 3 (?) relevant directions

$$f(R) = \sum_{n=0}^N a_n R^n$$

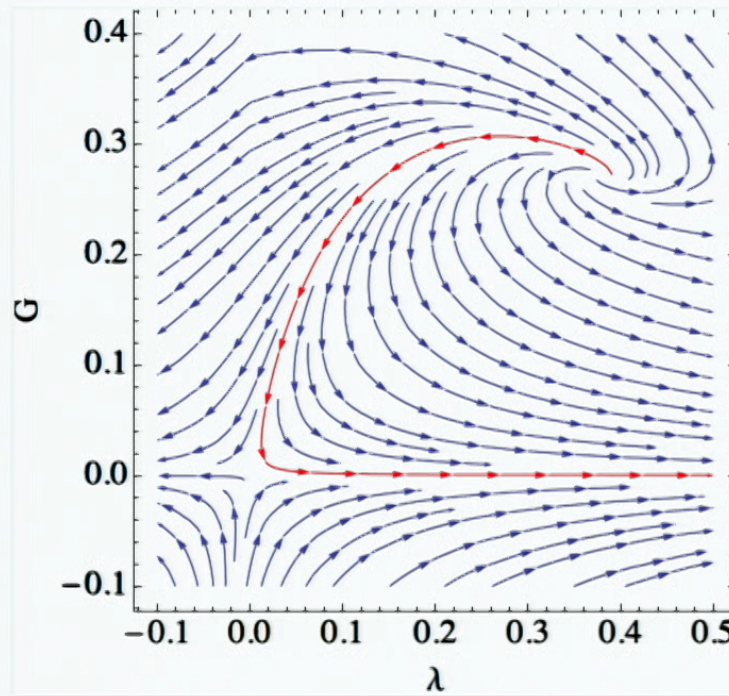
Reuter, Lauscher, '02; Codello, Percacci, Rahmede, '09;
Benedetti, Caravelli, '12; Dietz, Morris, '12;
Falls, Litim, Nikolakopoulos, Rahmede, '13, '14
Demmel, Saueressig, Zanusso, '15; Eichhorn '15



- towards convergence in extensions of the truncation
- near-canonical scaling: guidance for well-controlled truncations

The case for asymptotic safety in gravity

$$\Gamma_k = -\frac{1}{16\pi G_N(k)} \int d^4x \sqrt{g} (R - 2\bar{\lambda}(k)) \quad G(k) = G_N(k)k^2, \quad \lambda(k) = \bar{\lambda}(k)/k^2$$



interacting fixed point
with ~ 3 (?) relevant directions

$$f(R) = \sum_{n=0}^N a_n R^n$$

Reuter, Lauscher, '02; Codello, Percacci, Rahmede, '09;
Benedetti, Caravelli, '12; Dietz, Morris, '12;
Falls, Litim, Nikolakopoulos, Rahmede, '13, '14
Demmel, Saueressig, Zanusso, '15; Eichhorn '15

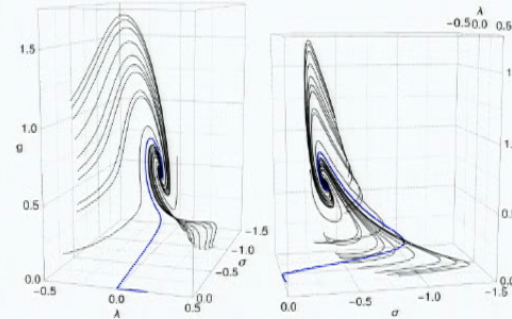
$$R^2, \quad R_{\mu\nu} R^{\mu\nu}$$

Benedetti, Machado, Saueressig, '09

Christiansen, '16

$$C_{\mu\nu}{}^{\kappa\lambda} C_{\kappa\lambda}{}^{\rho\sigma} C_{\rho\sigma}{}^{\mu\nu}$$

Gies, Knorr, Lippoldt, Saueressig '16



perturbative counterterms: asymptotic safety

A link that matters

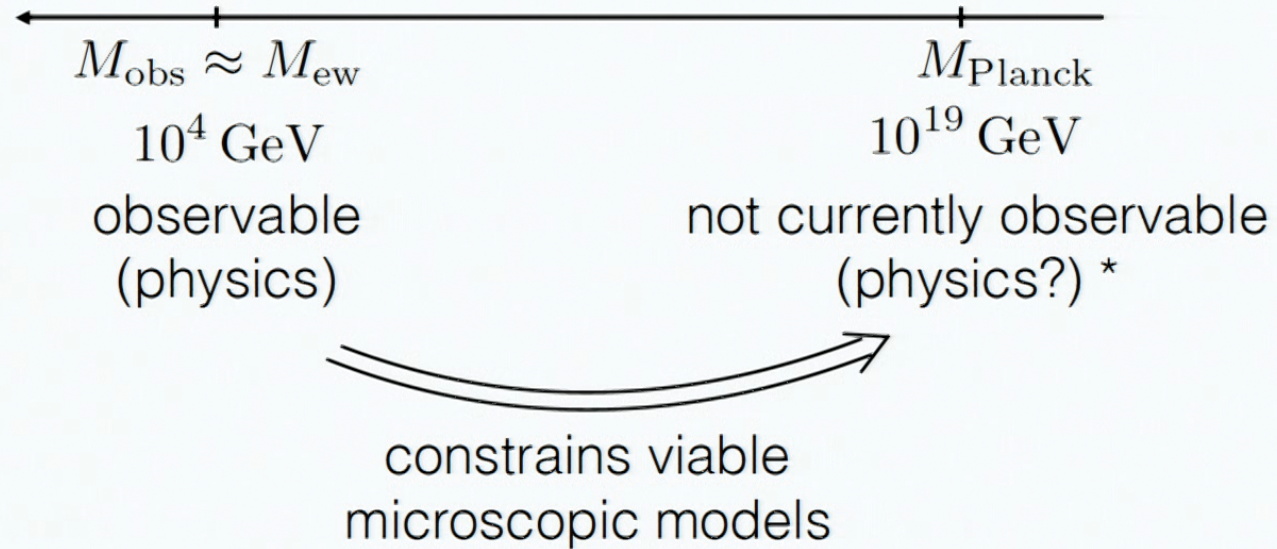


A link that matters



* “In any region of physics where very little is known, one must keep to the experimental basis if one is not to indulge in wild speculation that is almost certain to be wrong. “
(P.A.M. Dirac)

A link that matters



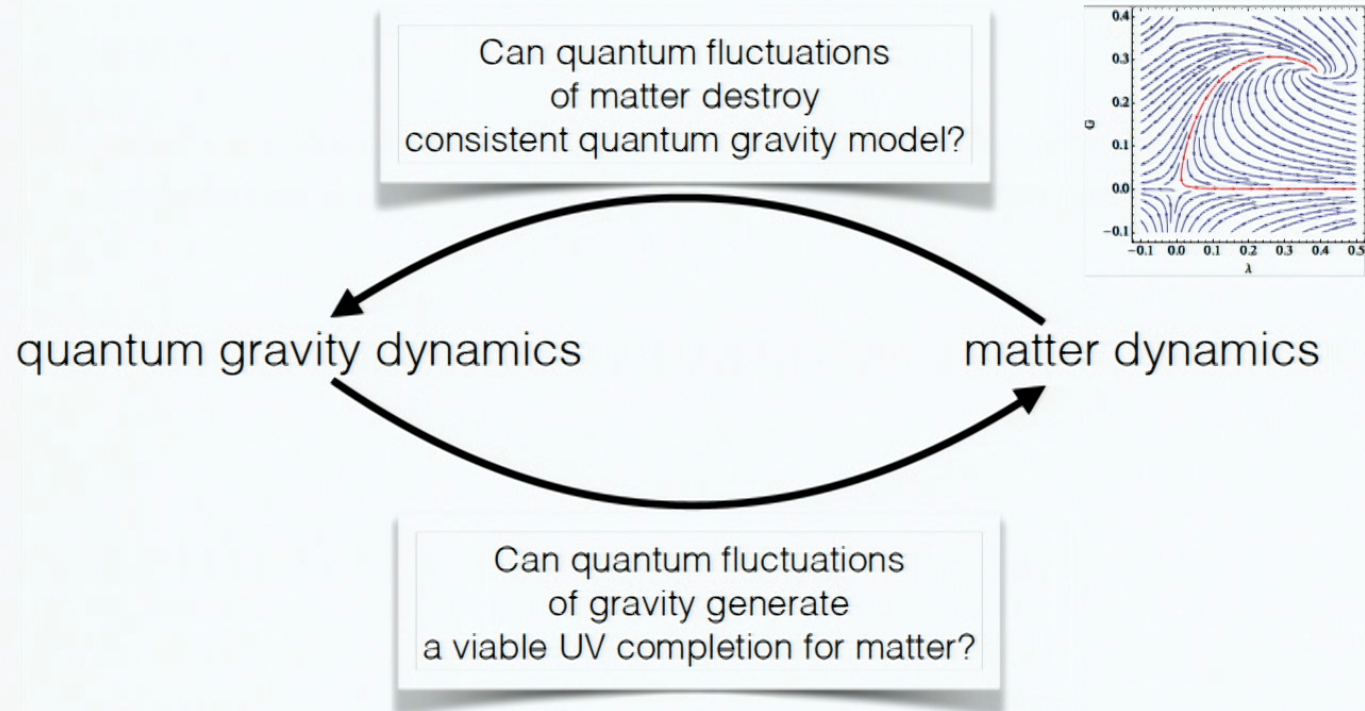
* "In any region of physics where very little is known, one must keep to the experimental basis if one is not to indulge in wild speculation that is almost certain to be wrong."
(P.A.M. Dirac)

Asymptotic safety for quantum gravity and matter

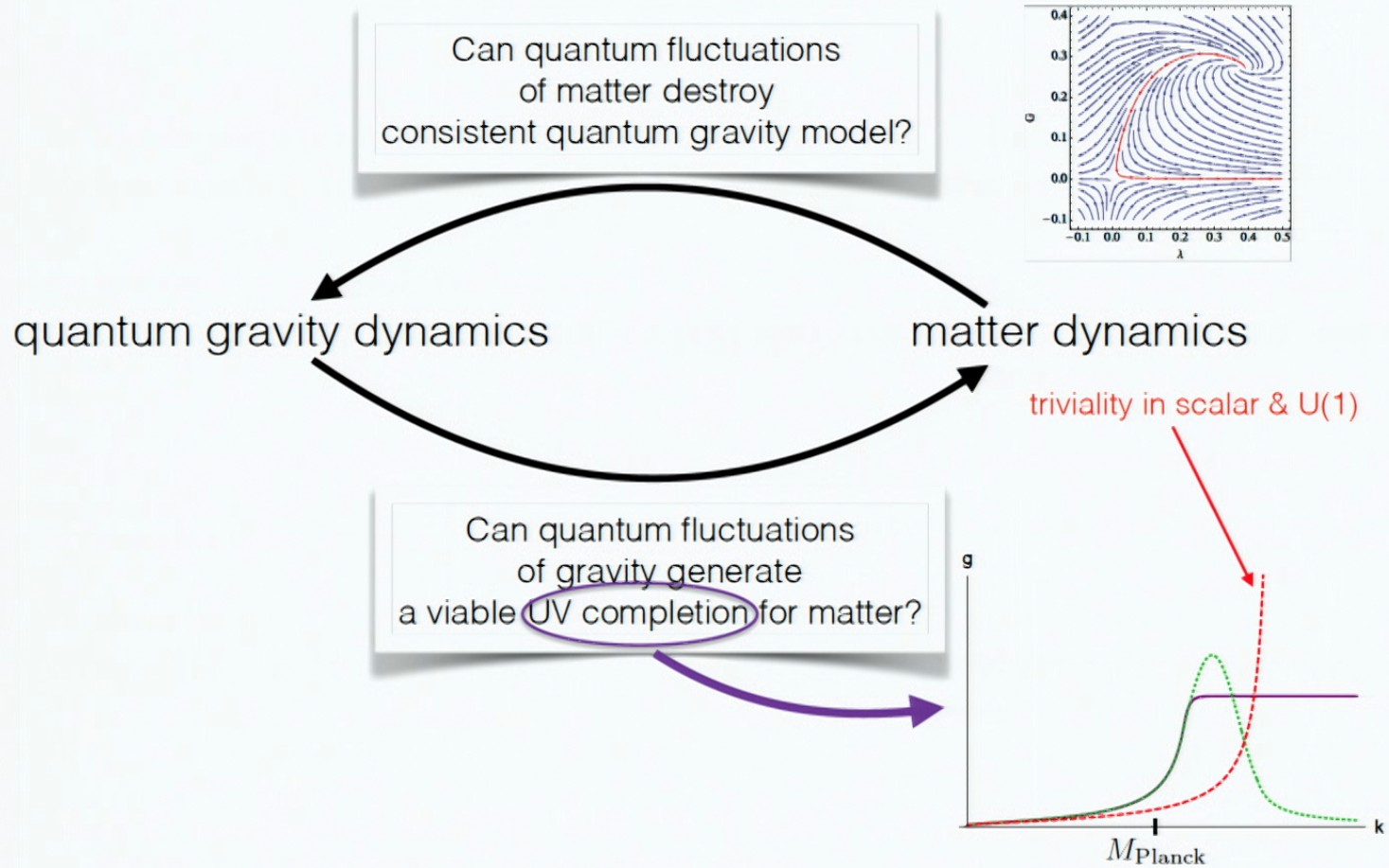
quantum gravity dynamics

matter dynamics

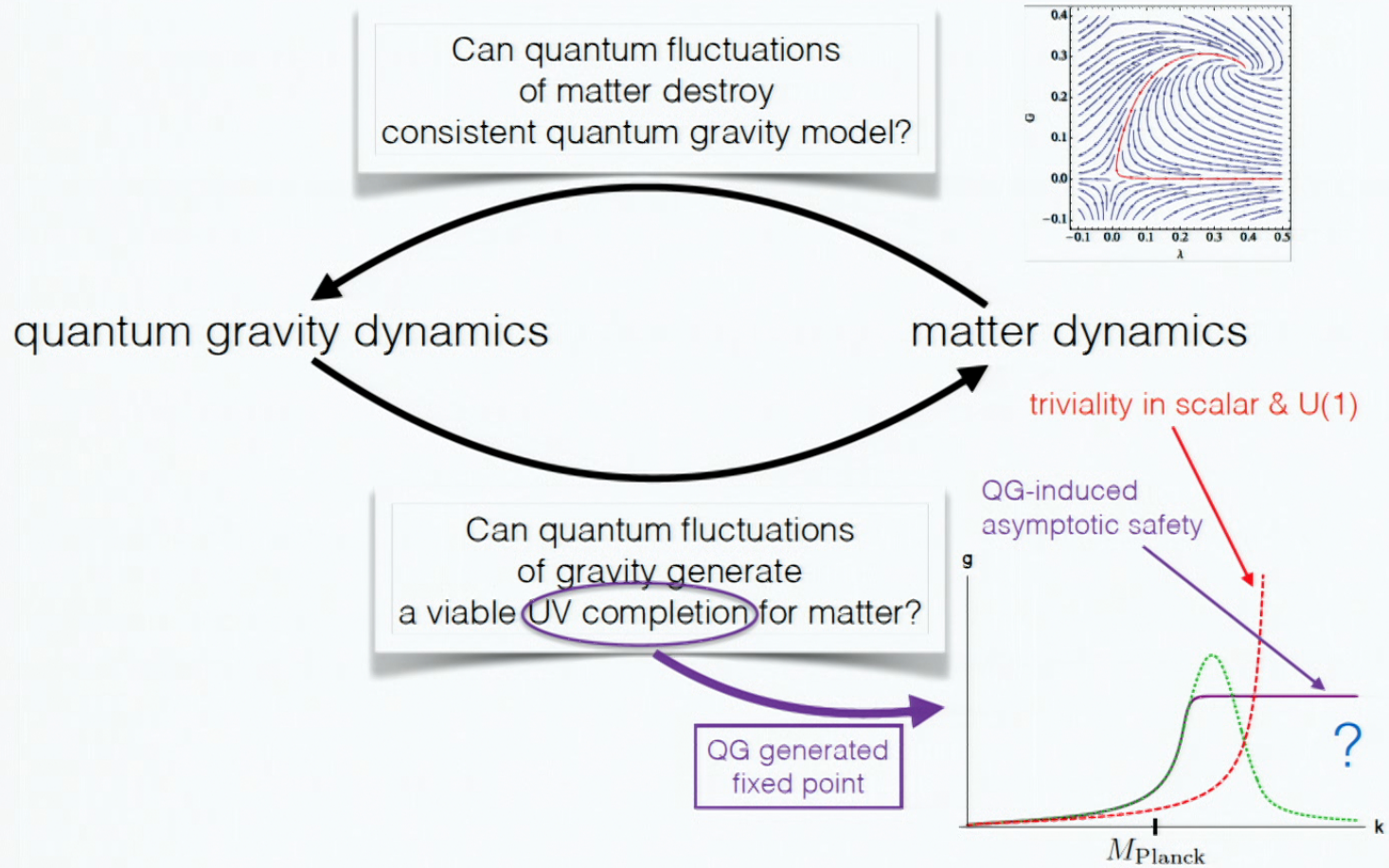
Asymptotic safety for quantum gravity and matter



Asymptotic safety for quantum gravity and matter

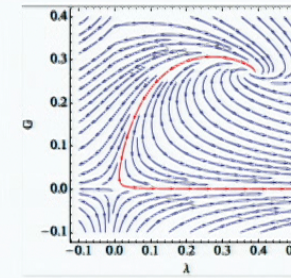


Asymptotic safety for quantum gravity and matter



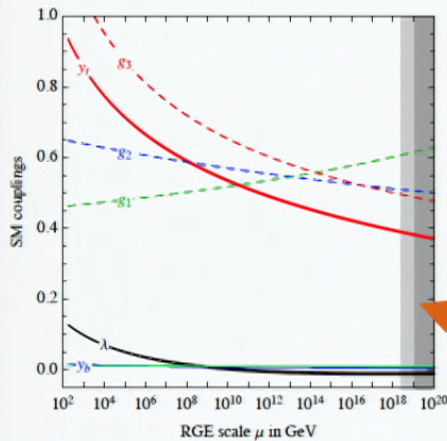
Asymptotic safety for quantum gravity and matter

Can quantum fluctuations of matter destroy consistent quantum gravity model?



quantum gravity dynamics

matter dynamics

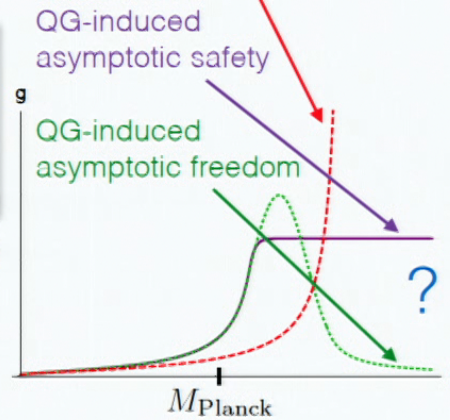


Can quantum fluctuations of gravity generate a viable UV completion for matter?

match onto SM at Planck scale

QG generated fixed point

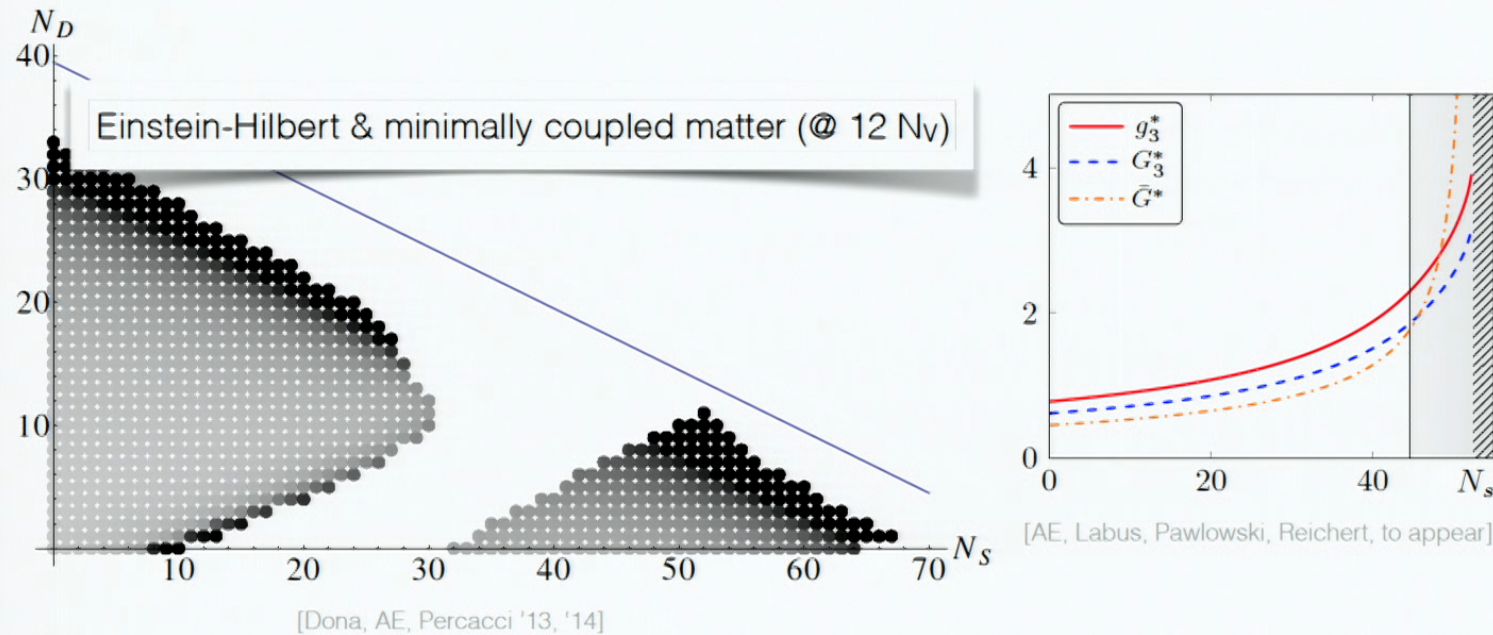
triviality in scalar & U(1)



Buttazzo et al. '13

Matter matters

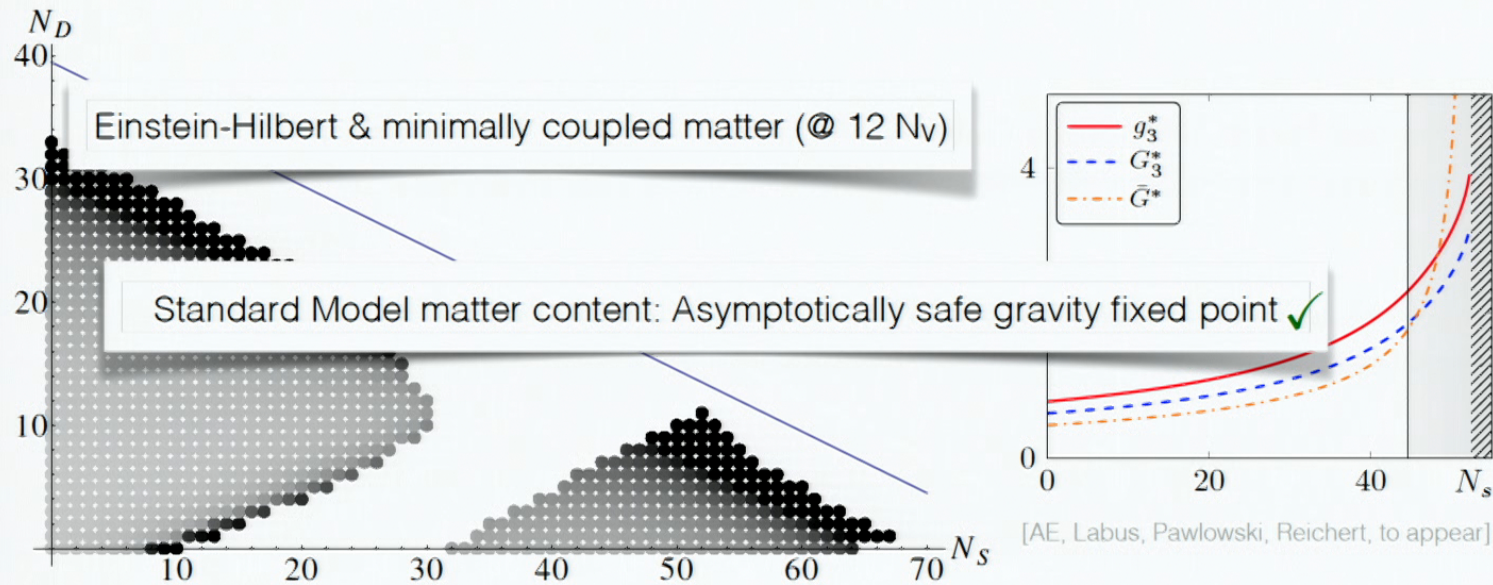
Quantum fluctuations of matter alter the UV dynamics for gravity



$$\beta_G = 2G - G^2 (\#_{\text{grav}} - N_S \#_{\text{scalar}}), \quad \#_{\text{grav}} > 0, \quad \#_{\text{scalar}} > 0$$

Matter matters

Quantum fluctuations of matter alter the UV dynamics for gravity



[Dona, AE, Percacci '13, '14]

[AE, Labus, Pawłowski, Reichert, to appear]

$$\beta_G = 2G - G^2 (\#_{\text{grav}} - N_S \#_{\text{scalar}}), \quad \#_{\text{grav}} > 0, \quad \#_{\text{scalar}} > 0$$

Asymptotically safe solution to the U(1) triviality problem

Asymptotically safe solution to the U(1) triviality problem

$$\Gamma_k = \Gamma_{k \text{ EH}} + \frac{1}{4g_{U(1)}^2} \int d^4x \sqrt{g} F^2$$

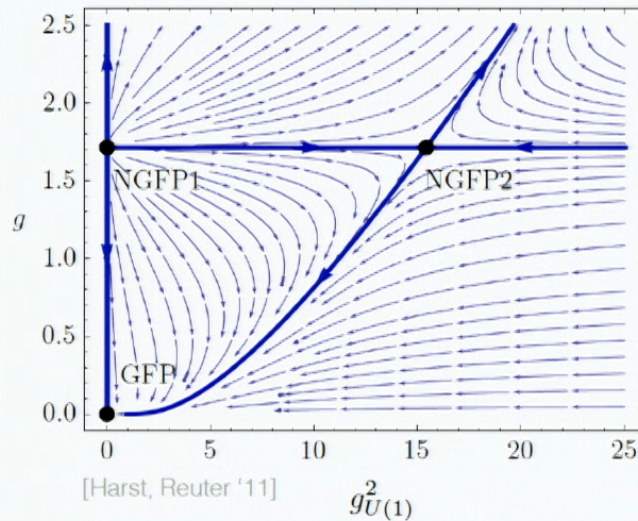
Newton coupling

$$\partial_t g_{U(1)}^2 = \left(\frac{2}{3\pi} N_f g_{U(1)}^2 - \frac{6}{\pi} g \Phi_1(0) \right) g_{U(1)}^2$$

[Harst, Reuter '11]

non-Abelian case:

[Daum, Harst, Reuter '10; Folkerts, Litim, Pawłowski '11]



QG-induced asymptotic freedom?

Asymptotically safe solution to the U(1) triviality problem

$$\Gamma_k = \Gamma_{k \text{ EH}} + \frac{1}{4g_{U(1)}^2} \int d^4x \sqrt{g} F^2$$

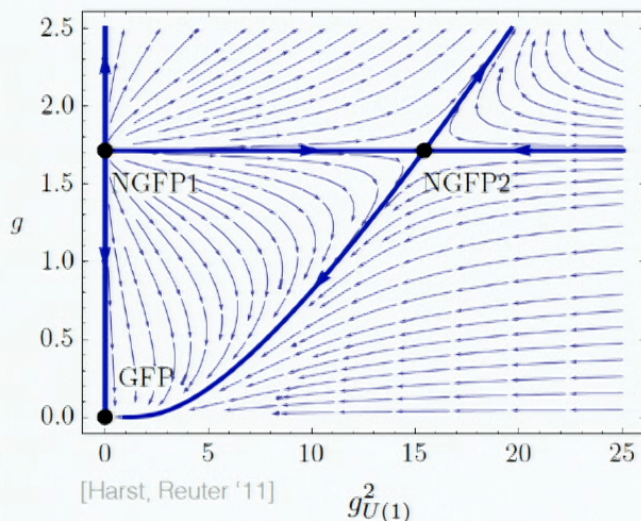
Newton coupling

$$\partial_t g_{U(1)}^2 = \left(\frac{2}{3\pi} N_f g_{U(1)}^2 - \frac{6}{\pi} g \Phi_1(0) \right) g_{U(1)}^2$$

[Harst, Reuter '11]

non-Abelian case:

[Daum, Harst, Reuter '10; Folkerts, Litim, Pawłowski '11]

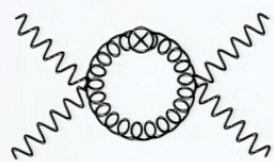


[Harst, Reuter '11]

QG-induced asymptotic freedom?

$$\sqrt{g} g^{\mu\nu} g^{\kappa\lambda} F_{\mu\kappa} F_{\nu\lambda}$$

QG-induced photon-interactions



$$w_2 (F^2)^2$$

$$\beta_{w_2} = 4w_2 + 40g^2 - \frac{7}{2\pi} g w_2 + \frac{1}{8\pi^2} w_2^2$$

shift free to interacting fixed point

[Christiansen, AE, 2017]

Asymptotically safe solution to the U(1) triviality problem

$$\Gamma_k = \Gamma_{k \text{ EH}} + \frac{1}{4g_{U(1)}^2} \int d^4x \sqrt{g} F^2$$

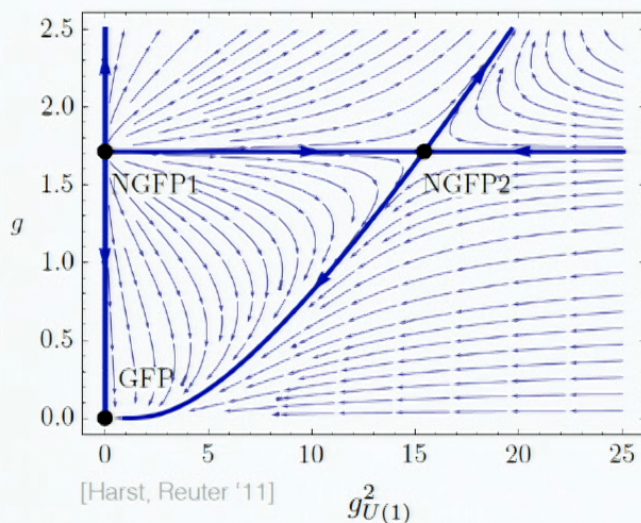
Newton coupling

$$\partial_t g_{U(1)}^2 = \left(\frac{2}{3\pi} N_f g_{U(1)}^2 - \frac{6}{\pi} g \Phi_1(0) \right) g_{U(1)}^2$$

[Harst, Reuter '11]

non-Abelian case:

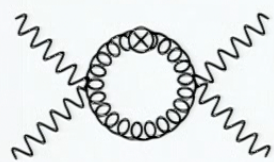
[Daum, Harst, Reuter '10; Folkerts, Litim, Pawłowski '11]



[Harst, Reuter '11]

QG-induced asymptotic freedom?

$$\sqrt{g} g^{\mu\nu} g^{\kappa\lambda} F_{\mu\kappa} F_{\nu\lambda}$$



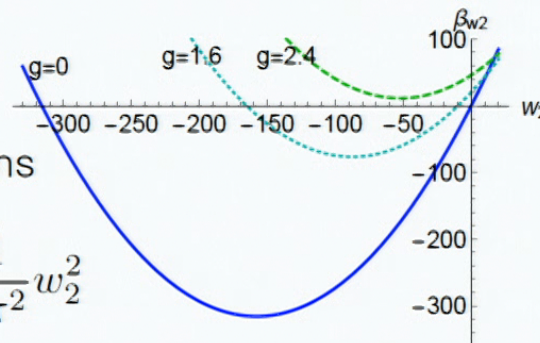
$$w_2 (F^2)^2$$

QG-induced photon-interactions

$$\beta_{w_2} = 4w_2 + 40g^2 - \frac{7}{2\pi} g w_2 + \frac{1}{8\pi^2} w_2^2$$

shift free to interacting fixed point

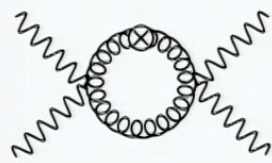
[Christiansen, AE, 2017]



Asymptotically safe solution to the U(1) triviality problem

gravity must remain “weak”
for a real fixed point in
U(1) gauge theory

$$\sqrt{g} g^{\mu\nu} g^{\kappa\lambda} F_{\mu\kappa} F_{\nu\lambda}$$

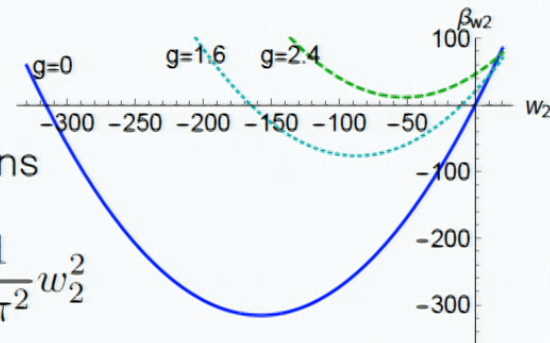


$$w_2 (F^2)^2$$

QG-induced photon-interactions

$$\beta_{w_2} = 4w_2 + 40g^2 - \frac{7}{2\pi}g w_2 + \frac{1}{8\pi^2}w_2^2$$

shift free to
interacting fixed point



[Christiansen, AE, 2017]

Asymptotically safe solution to the U(1) triviality problem
interaction structure in gauge sector:

$$\frac{1}{4g_{U(1)}^2} F^2 + w_2 F^4$$

interacting fixed point

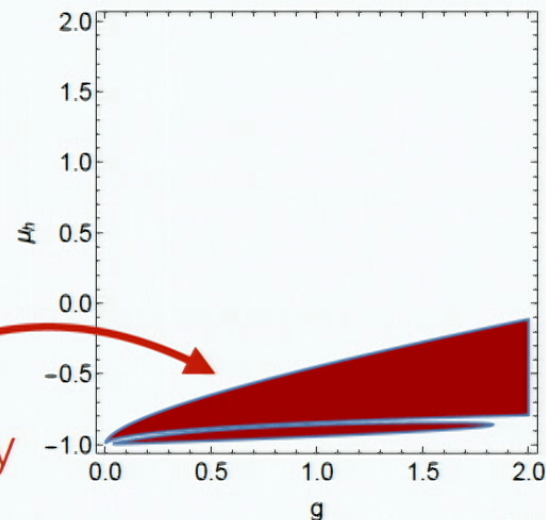
asymptotically free fixed point?

Asymptotically safe solution to the U(1) triviality problem

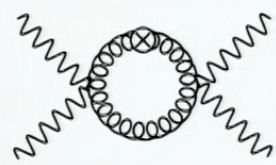
gravity must remain "weak"
for a real fixed point in
U(1) gauge theory

effective
gravity coupling $\frac{g}{1 + \mu_h}$

strong gravity
excluded



$$\sqrt{g} g^{\mu\nu} g^{\kappa\lambda} F_{\mu\kappa} F_{\nu\lambda}$$

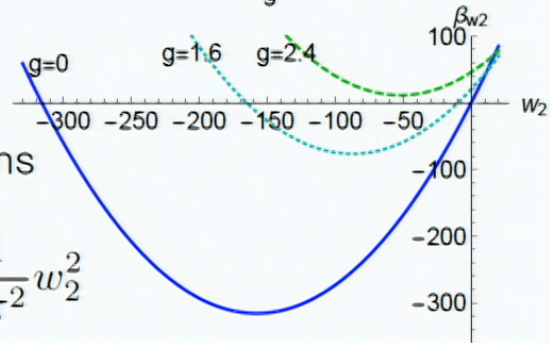


$$w_2 (F^2)^2$$

QG-induced photon-interactions

$$\beta_{w_2} = 4w_2 + 40g^2 - \frac{7}{2\pi} g w_2 + \frac{1}{8\pi^2} w_2^2$$

shift free to
interacting fixed point



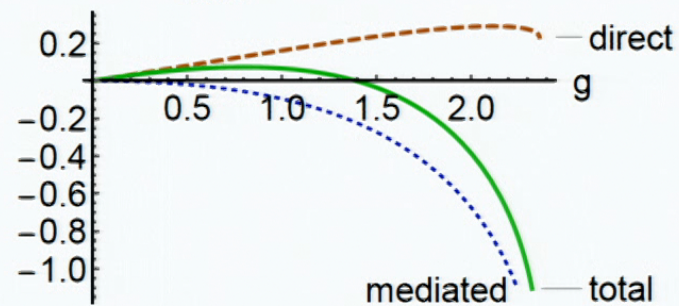
[Christiansen, AE, 2017]

Asymptotically safe solution to the U(1) triviality problem
 interaction structure in gauge sector:

$$\frac{1}{4g_{U(1)}^2} F^2 + w_2 F^4$$

↙ asymptotically free fixed point?
 ↘ interacting fixed point

asymptotic freedom in $g_{U(1)}$: $\theta_{g_{U(1)}^2} = \theta_{g_{U(1)}^2}|_{\text{grav}} + \theta_{g_{U(1)}^2}|_{w_2} > 0$
 contributions to $\theta_{g_{U(1)}^2}$



Asymptotically safe solution to the U(1) triviality problem

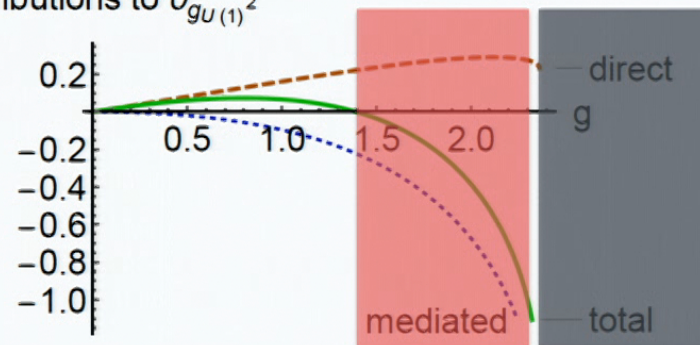
interaction structure in gauge sector:

$$\frac{1}{4g_{U(1)}^2} F^2 + w_2 F^4$$

$\frac{1}{4g_{U(1)}^2}$ → asymptotically free fixed point?
 $w_2 F^4$ → interacting fixed point

asymptotic freedom in $g_{U(1)}$: $\theta_{g_{U(1)}^2} = \theta_{g_{U(1)}^2}|_{\text{grav}} + \theta_{g_{U(1)}^2}|_{w_2} > 0$

contributions to $\theta_{g_{U(1)}^2}$



gravity-induced interactions destroy asymptotic freedom
 gravity too strong to support fixed point in gauge sector

Asymptotically safe solution to the U(1) triviality problem

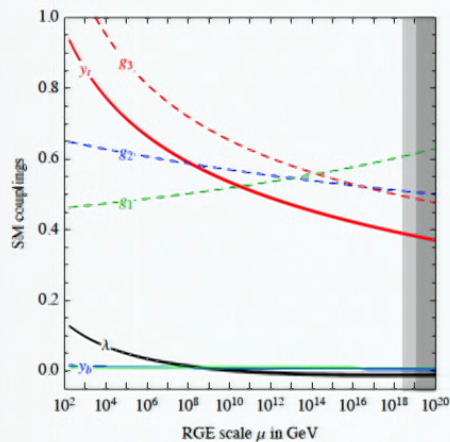
interaction structure in gauge sector:

$$\frac{1}{4g_{U(1)}^2} F^2 + w_2 F^4$$

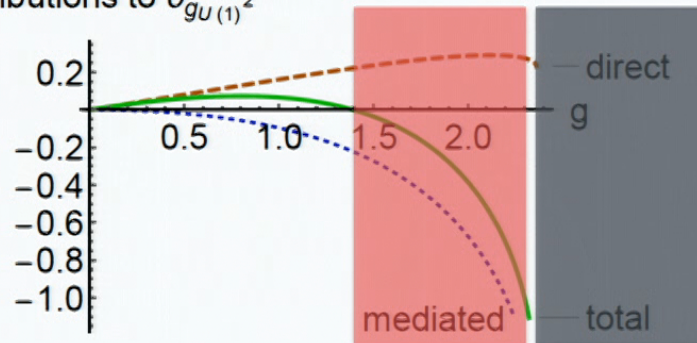
$\frac{1}{4g_{U(1)}^2}$ → asymptotically free fixed point?
 $w_2 F^4$ → interacting fixed point

$$g_{U(1)}^2(k) = \left(\frac{k}{k_0}\right)^{-\theta_{g_{U(1)}^2}} \Rightarrow g_{U(1)}^2(M_{\text{Planck}}) \approx 0$$

asymptotic freedom in $g_{U(1)}$: $\theta_{g_{U(1)}^2} = \theta_{g_{U(1)}^2}|_{\text{grav}} + \theta_{g_{U(1)}^2}|_{w_2} > 0$



contributions to $\theta_{g_{U(1)}^2}$



gravity-induced interactions destroy asymptotic freedom
 gravity too strong to support fixed point in gauge sector

Asymptotically safe solution to the U(1) triviality problem

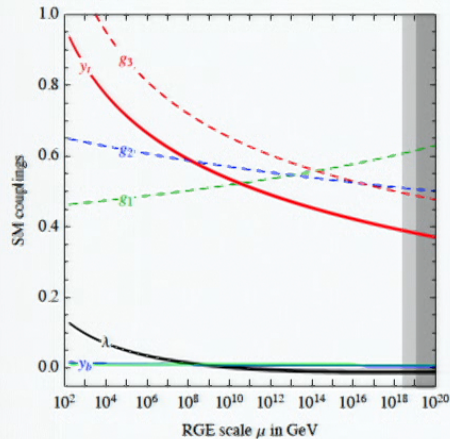
interaction structure in gauge sector:

$$\frac{1}{4g_{U(1)}^2} F^2 + w_2 F^4$$

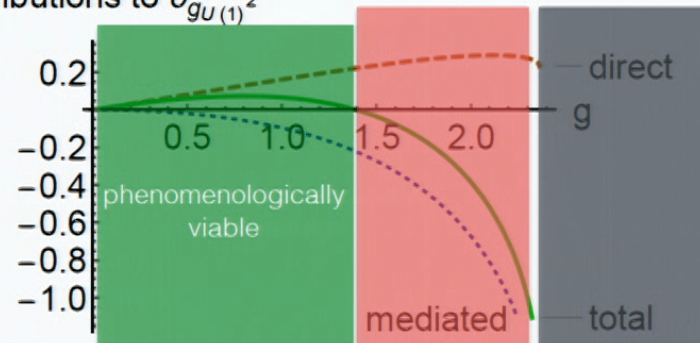
$\frac{1}{4g_{U(1)}^2}$ → asymptotically free fixed point?
 $w_2 F^4$ → interacting fixed point

$$g_{U(1)}^2(k) = \left(\frac{k}{k_0}\right)^{-\theta_{g_{U(1)}^2}} \Rightarrow g_{U(1)}^2(M_{\text{Planck}}) \approx 0$$

asymptotic freedom in $g_{U(1)}$: $\theta_{g_{U(1)}^2} = \theta_{g_{U(1)}^2}|_{\text{grav}} + \theta_{g_{U(1)}^2}|_{w_2} > 0$



contributions to $\theta_{g_{U(1)}^2}$



gravity-induced interactions
destroy asymptotic freedom

gravity too strong to support
fixed point in gauge sector

QG-induced fixed-point structure

interaction structure in gauge sector:

$$\frac{1}{4g_{U(1)}^2} F^2 + w_2 F^4$$

free fixed point

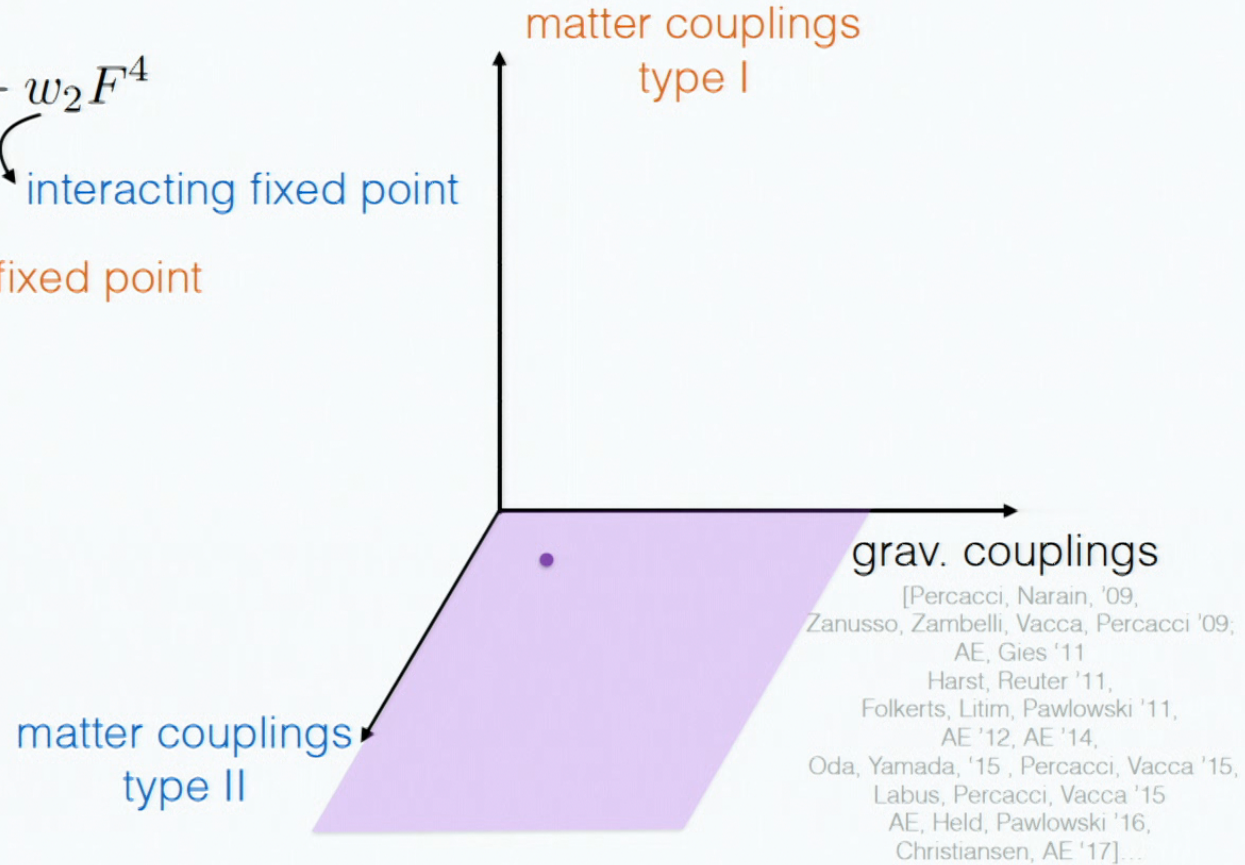
interacting fixed point

QG-induced fixed-point structure

interaction structure in gauge sector:

$$\frac{1}{4g_{U(1)}^2} F^2 + w_2 F^4$$

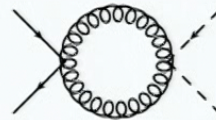
free fixed point (pointing to $\frac{1}{4g_{U(1)}^2}$)
interacting fixed point (pointing to $w_2 F^4$)



Symmetries dictate fixed-point structure

$$\begin{aligned}
 iZ_\psi \int d^4x \sqrt{g} \bar{\psi} \not{\nabla} \psi & \quad \frac{Z_\phi}{2} \int d^4x \sqrt{g} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi \\
 \psi \rightarrow e^{i\alpha\gamma_5} \psi & \quad \phi \rightarrow -\phi \\
 \bar{\psi} \rightarrow \bar{\psi} e^{i\alpha\gamma_5} & \quad \text{shift symm. } \phi \rightarrow \phi + a
 \end{aligned}$$

induced interactions (non-free)



fermions & scalars $\chi (\partial_\mu \phi \partial^\mu \phi) \bar{\psi} \not{\nabla} \psi \rightarrow \chi_* \neq 0$ [AE, Held, Pawłowski '16]

→ fixed point in matter sector cannot be completely free
 nonzero interactions: symmetry of kinetic term

Symmetries dictate fixed-point structure

$$iZ_\psi \int d^4x \sqrt{g} \bar{\psi} \not{\nabla} \psi$$

$$\psi \rightarrow e^{i\alpha\gamma_5} \psi$$

$$\bar{\psi} \rightarrow \bar{\psi} e^{i\alpha\gamma_5}$$

$$\frac{Z_\phi}{2} \int d^4x \sqrt{g} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi$$

$$\phi \rightarrow -\phi$$

$$\text{shift symm. } \phi \rightarrow \phi + a$$

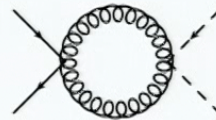
$$iy \int d^4x \sqrt{g} \bar{\psi} \psi \phi$$

$$\phi \rightarrow -\phi$$

$$\psi \rightarrow e^{i\frac{\pi}{2}\gamma_5} \psi$$

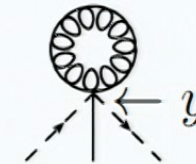
$$\bar{\psi} \rightarrow \bar{\psi} e^{i\frac{\pi}{2}\gamma_5}$$

induced interactions (non-free)



QG flucs: global symmetries ✓

fermions & scalars $\chi (\partial_\mu \phi \partial^\mu \phi) \bar{\psi} \not{\nabla} \psi \rightarrow \chi_* \neq 0$ [AE, Held, Pawłowski '16]



$$\beta_y|_{\text{grav}} = \# y g$$

→ fixed point in matter sector cannot be completely free
nonzero interactions: symmetry of kinetic term

Symmetries dictate fixed-point structure

$$iZ_\psi \int d^4x \sqrt{g} \bar{\psi} \not{\nabla} \psi$$

$$\psi \rightarrow e^{i\alpha\gamma_5} \psi$$

$$\bar{\psi} \rightarrow \bar{\psi} e^{i\alpha\gamma_5}$$

$$\frac{Z_\phi}{2} \int d^4x \sqrt{g} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi$$

$$\phi \rightarrow -\phi$$

$$\text{shift symm. } \phi \rightarrow \phi + a$$

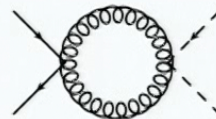
$$iy \int d^4x \sqrt{g} \bar{\psi} \psi \phi$$

$$\phi \rightarrow -\phi$$

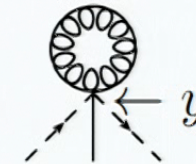
$$\psi \rightarrow e^{i\frac{\pi}{2}\gamma_5} \psi$$

$$\bar{\psi} \rightarrow \bar{\psi} e^{i\frac{\pi}{2}\gamma_5}$$

induced interactions (non-free)



QG flucs: global symmetries ✓



fermions & scalars $\chi (\partial_\mu \phi \partial^\mu \phi) \bar{\psi} \not{\nabla} \psi \rightarrow \chi_* \neq 0$ [AE, Held, Pawłowski '16]

$$\beta_y|_{\text{grav}} = \# y g$$

minimally interacting FP:
free in couplings with
reduced symmetry,
interacting in couplings with
full symmetry of kin. terms

Symmetries dictate fixed-point structure

$$iZ_\psi \int d^4x \sqrt{g} \bar{\psi} \not{\nabla} \psi$$

$$\psi \rightarrow e^{i\alpha\gamma_5} \psi$$

$$\bar{\psi} \rightarrow \bar{\psi} e^{i\alpha\gamma_5}$$

$$\frac{Z_\phi}{2} \int d^4x \sqrt{g} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi$$

$$\phi \rightarrow -\phi$$

$$\text{shift symm. } \phi \rightarrow \phi + a$$

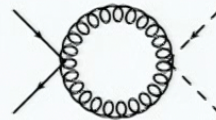
$$iy \int d^4x \sqrt{g} \bar{\psi} \psi \phi$$

$$\phi \rightarrow -\phi$$

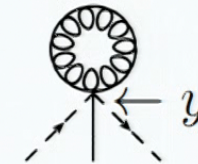
$$\psi \rightarrow e^{i\frac{\pi}{2}\gamma_5} \psi$$

$$\bar{\psi} \rightarrow \bar{\psi} e^{i\frac{\pi}{2}\gamma_5}$$

induced interactions (non-free)



QG flucs: global symmetries ✓



fermions & scalars $\chi (\partial_\mu \phi \partial^\mu \phi) \bar{\psi} \not{\nabla} \psi \rightarrow \chi_* \neq 0$ [AE, Held, Pawłowski '16]

$$\beta_y|_{\text{grav}} = \# y g$$

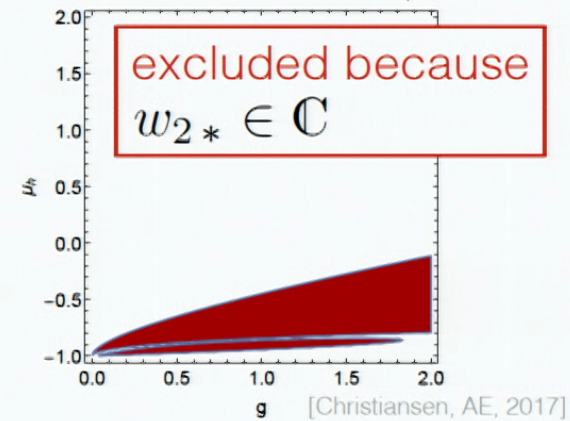
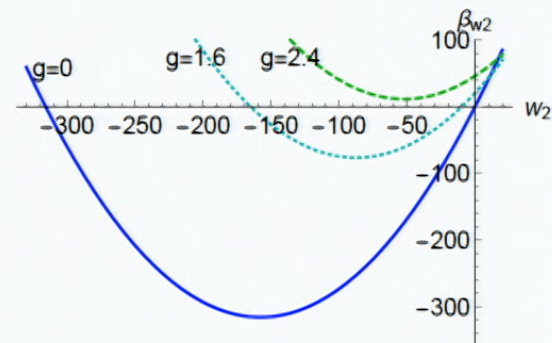
- Real couplings at induced fixed point?
- (Ir)relevant interactions?

minimally interacting FP:
free in couplings with
reduced symmetry,
interacting in couplings with
full symmetry of kin. terms

Asymptotic safety for matter requires “weakly interacting” gravity

gauge fields

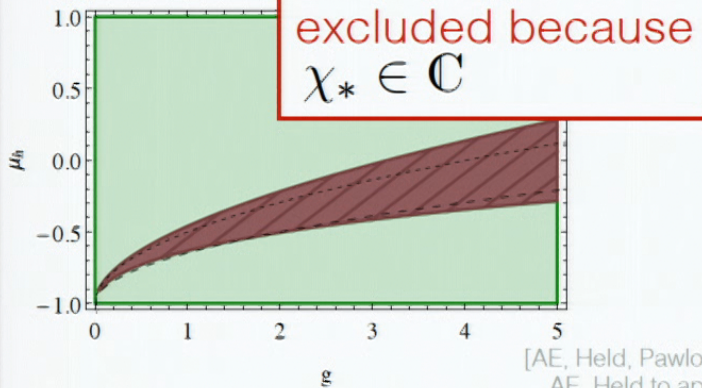
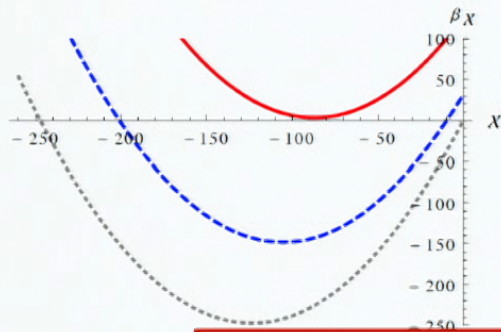
induced interactions $w_2 (F^2)^2 \rightarrow w_{2*} \neq 0$



Asymptotic safety for matter requires “weakly interacting” gravity

fermions & scalars
induced interactions

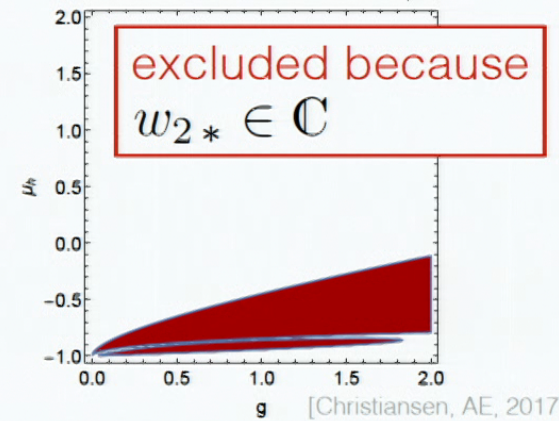
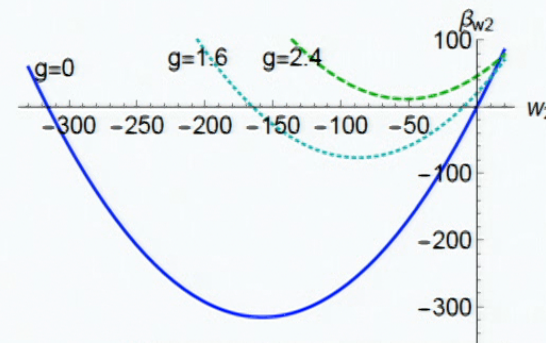
$$\chi (\partial_\mu \phi \partial^\mu \phi) \bar{\psi} \not{\nabla} \psi \rightarrow \chi_* \neq 0$$



[AE, Held, Pawłowski '16
AE, Held to appear]

gauge fields

induced interactions $w_2 (F^2)^2 \rightarrow w_{2*} \neq 0$



[Christiansen, AE, 2017]

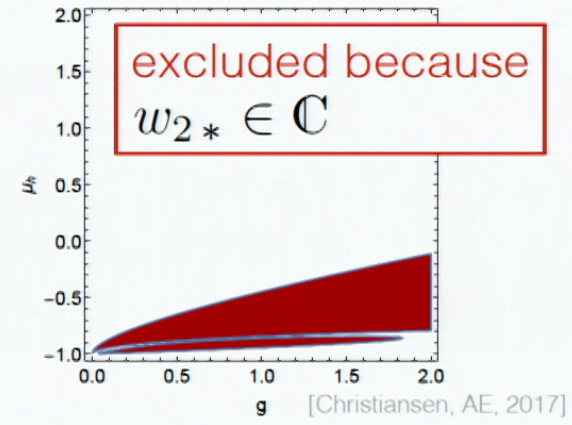
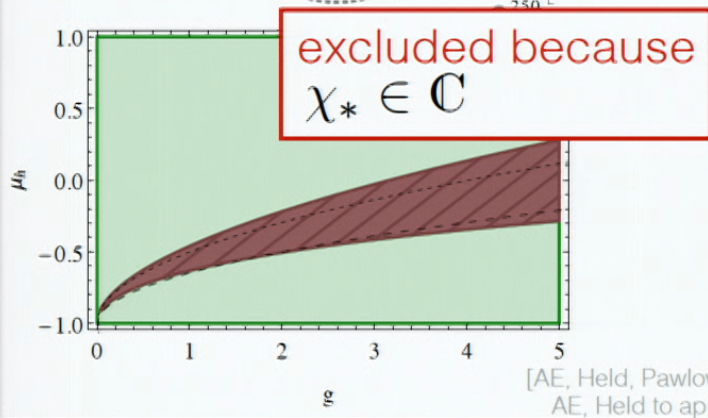
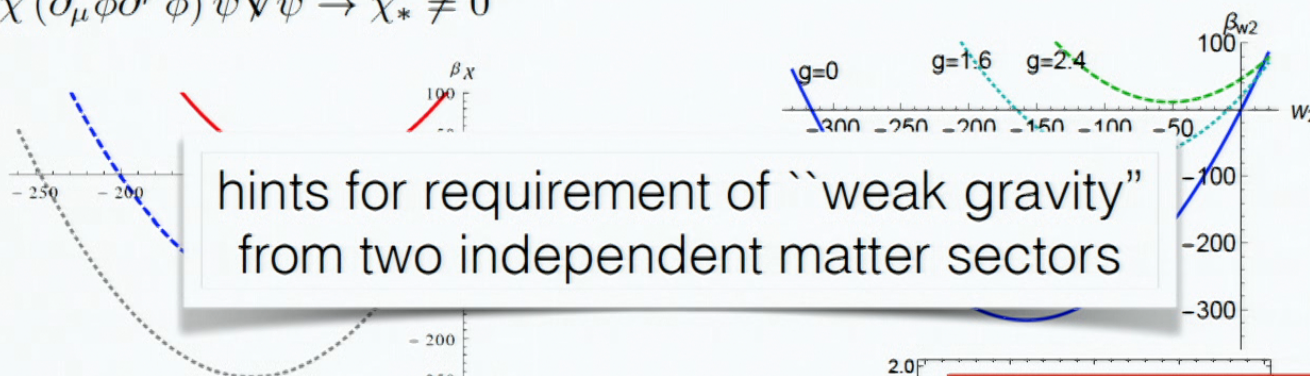
Asymptotic safety for matter requires “weakly interacting” gravity

fermions & scalars
induced interactions

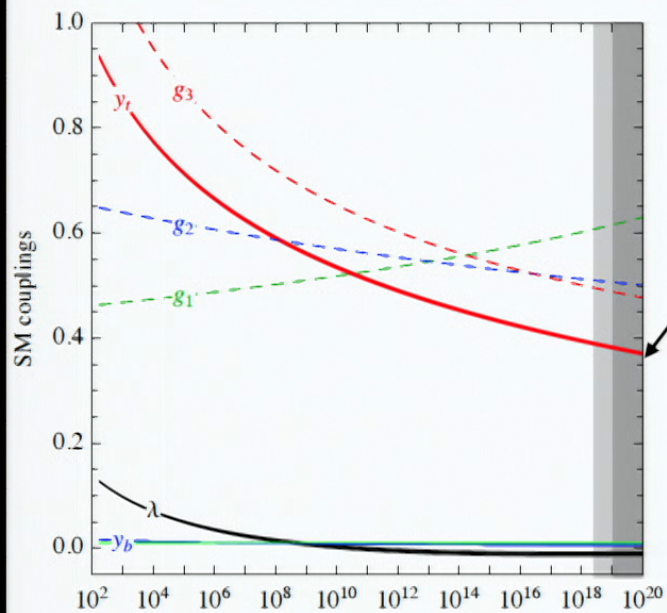
$$\chi (\partial_\mu \phi \partial^\mu \phi) \bar{\psi} \not{\nabla} \psi \rightarrow \chi_* \neq 0$$

gauge fields

induced interactions $w_2 (F^2)^2 \rightarrow w_{2*} \neq 0$



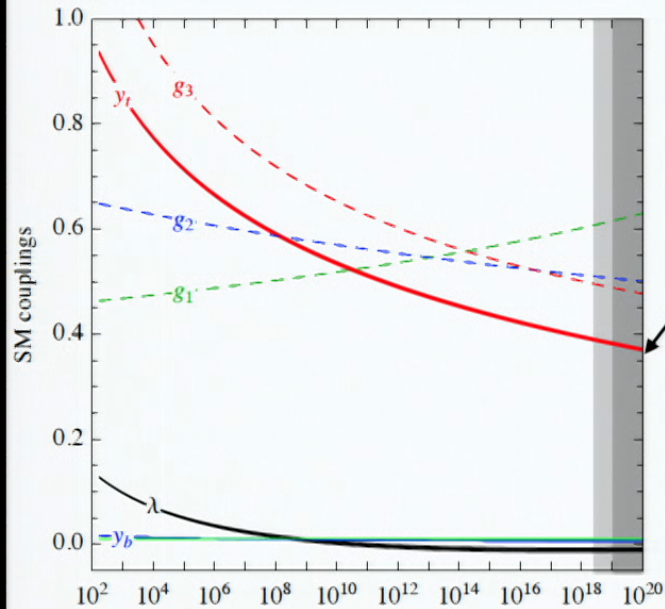
Linking ASQG to the electroweak scale



Connection to viable
UV fixed point?

[Buttazzo et al. '13] RGE scale μ in GeV

Linking ASQG to the electroweak scale



Connection to viable
UV fixed point?

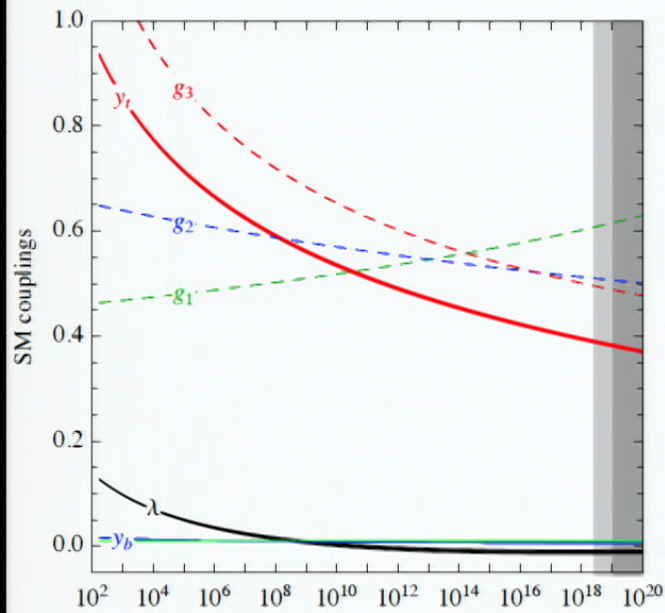
[Buttazzo et al. '13] RGE scale μ in GeV

toy model: simple Higgs-Yukawa model:

$$\Gamma_k = \frac{Z_\phi}{2} \int d^4x \sqrt{g} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi + i Z_\psi \int d^4x \sqrt{g} \bar{\psi} \not{\nabla} \psi + i y \int d^4x \sqrt{g} \phi \bar{\psi} \psi$$

[AE, Held, Pawłowski '16; AE, Held, to appear]

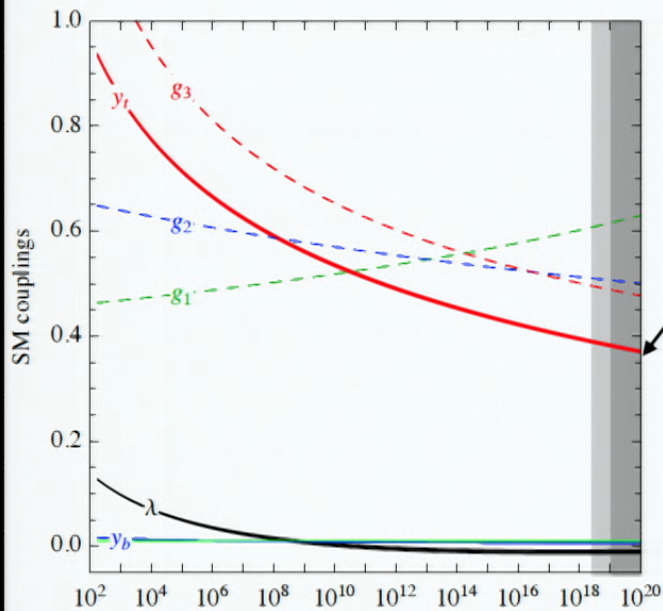
Linking ASQG to the electroweak scale



[Buttazzo et al. '13] RGE scale μ in GeV

$$\beta_y = -\theta_y \Big|_{\text{grav}} y g + \frac{y^3}{8\pi^2} \rightarrow y_* = 0$$

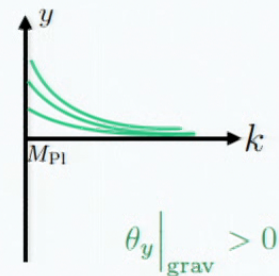
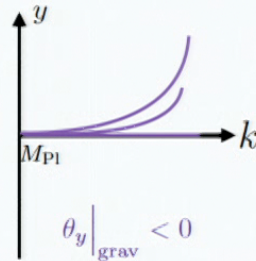
Linking ASQG to the electroweak scale



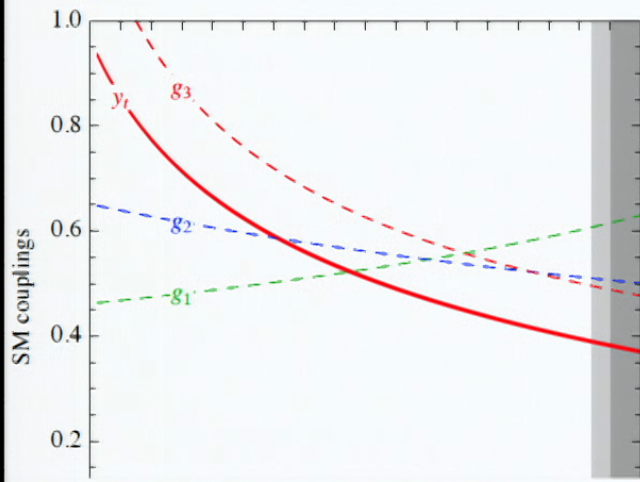
[Buttazzo et al. '13] RGE scale μ in GeV

$$\beta_y = -\theta_y \Big|_{\text{grav}} y g + \frac{y^3}{8\pi^2} \rightarrow y_* = 0$$

When can this fixed point be connected to viable Planck-scale physics?

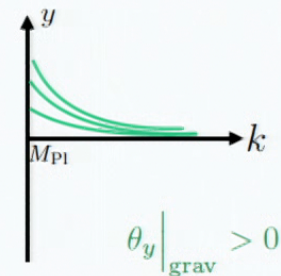
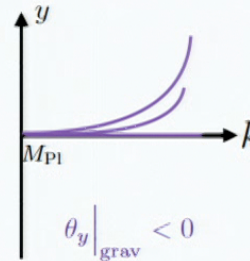


Linking ASQG to the electroweak scale



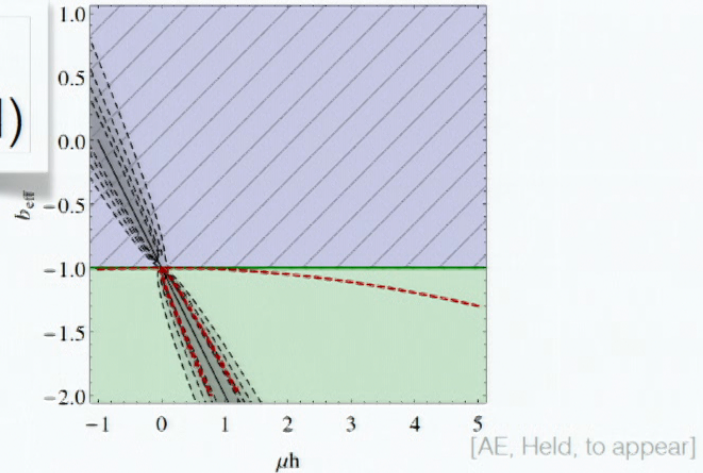
$$\beta_y = -\theta_y \Big|_{\text{grav}} y g + \frac{y^3}{8\pi^2} \rightarrow y_* = 0$$

When can this fixed point be connected to viable Planck-scale physics?

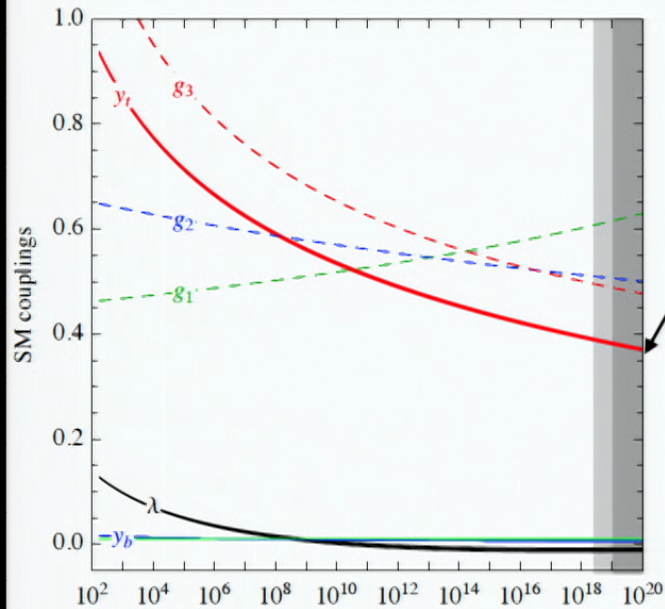


constraint on fixed-point values
(outlook: Higgs-top-bottom model)

$$\Gamma_{k \text{ grav}} = \frac{-1}{16\pi G} Z_h \int d^4x \sqrt{g} (R - 2\Lambda) + \frac{1}{16\pi} Z_h \int d^4x \sqrt{g} (aR^2 + bR_{\mu\nu}R^{\mu\nu} + cR\Box R + dR_{\mu\nu}\Box R^{\mu\nu})$$



Linking QG to the electroweak scale



[Buttazzo et al. '13] RGE scale μ in GeV

determines entire
top mass at e/w scale

$$y \bar{t} t \phi \rightarrow (y \cdot \langle \phi \rangle) \bar{t} t$$

chiral symmetry:

no microscopic fermion masses

$$m_\psi \bar{\psi} \psi = \frac{m_\psi}{2} (\bar{\psi}_R \psi_L + \bar{\psi}_L \psi_R)$$

in the Standard Model

→ guaranteed automatically below M_{Pl} or nontrivial constraint on fundamental model?

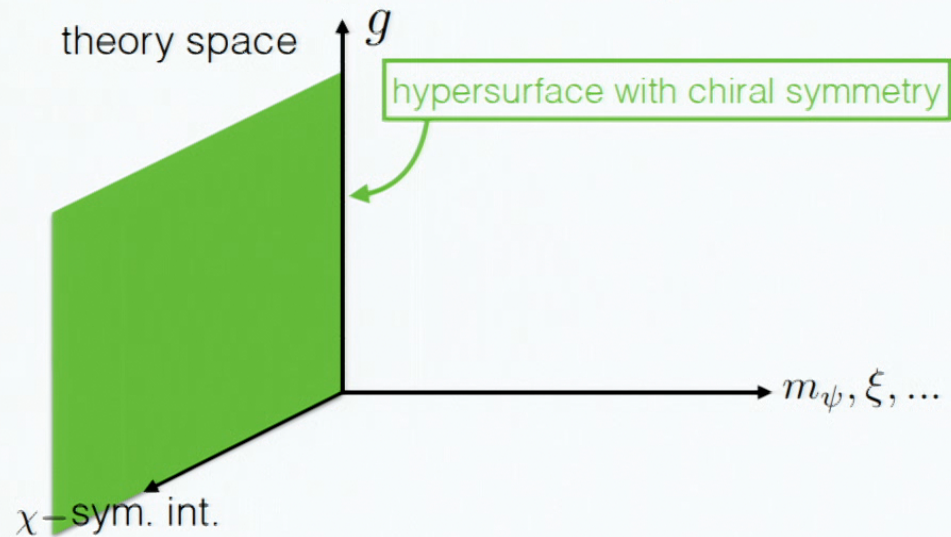
Chiral fermions

chiral symmetry: no microscopic fermion masses $m_\psi \bar{\psi}\psi$ in the Standard Model
 → guaranteed automatically below M_{Pl} or nontrivial constraint on fundamental model?

$$\Gamma_k = \Gamma_{\text{grav}} + iZ_\psi \int d^4x \sqrt{g} \bar{\psi}^i \nabla \psi^i + i\bar{m}_\psi \int d^4x \sqrt{g} \bar{\psi}^i \psi^i$$

$$+ i\bar{\xi} \int d^4x \sqrt{g} R \bar{\psi}^i \psi^i + i\bar{\zeta} \int d^4x \sqrt{g} \bar{\psi}^i \nabla^2 \psi^i$$

[AE, Lippoldt '16]

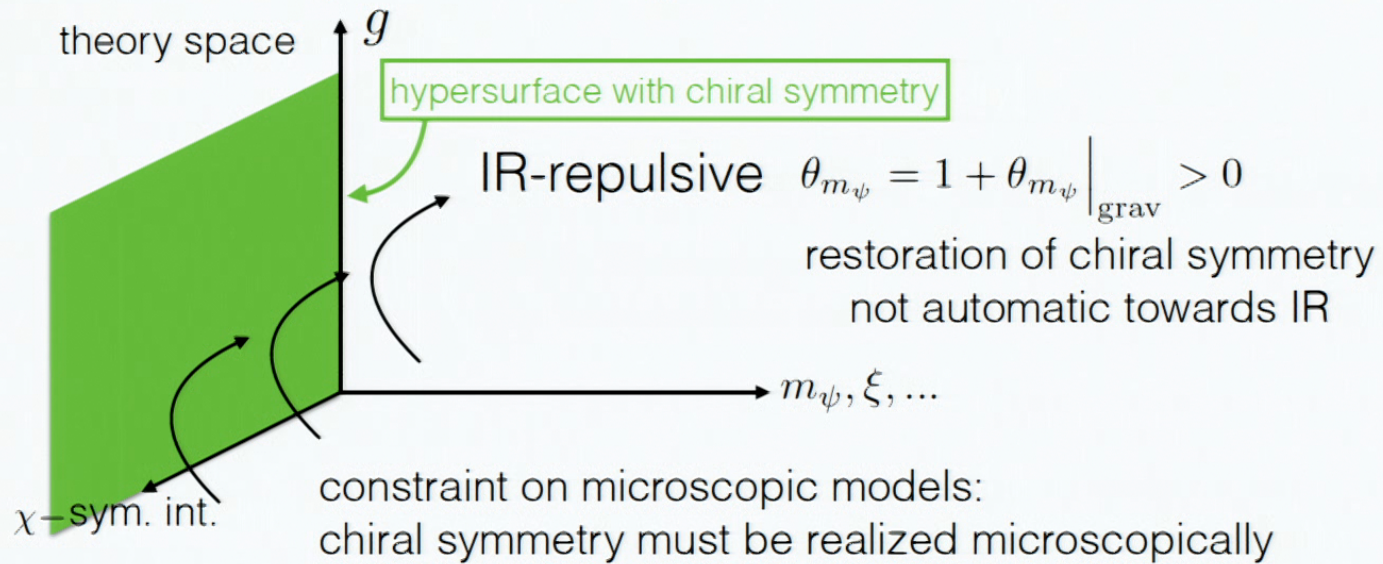


Chiral fermions

chiral symmetry: no microscopic fermion masses $m_\psi \bar{\psi}\psi$ in the Standard Model
 → guaranteed automatically below M_{Pl} or nontrivial constraint on fundamental model?

$$\Gamma_k = \Gamma_{\text{grav}} + iZ_\psi \int d^4x \sqrt{g} \bar{\psi}^i \nabla \psi^i + i\bar{m}_\psi \int d^4x \sqrt{g} \bar{\psi}^i \psi^i \quad [\text{AE, Lippoldt '16}]$$

$$+ i\bar{\xi} \int d^4x \sqrt{g} R \bar{\psi}^i \psi^i + i\bar{\zeta} \int d^4x \sqrt{g} \bar{\psi}^i \nabla^2 \psi^i$$



Chiral fermions in asymptotic safety

$$\Gamma_k = \Gamma_{\text{grav}} + iZ_\psi \int d^4x \sqrt{g} \bar{\psi}^i \not{\nabla} \psi^i + i\bar{m}_\psi \int d^4x \sqrt{g} \bar{\psi}^i \psi^i$$

$$+ i\bar{\xi} \int d^4x \sqrt{g} R \bar{\psi}^i \psi^i + i\bar{\zeta} \int d^4x \sqrt{g} \bar{\psi}^i \nabla^2 \psi^i$$

[AE, Lippoldt '16]

| fixed point | symmetry for fermions | g_* | λ_* | m_{ψ^*} | ξ_* | ζ_* | η_ψ | θ_1 | θ_2 | θ_3 | θ_4 | θ_5 |
|---------------------|-----------------------|-------|-------------|--------------|---------|-----------|-------------|------------|-------------------|------------|--------------------|------------|
| chiral non-Gaussian | chiral | 2.52 | -0.42 | 0 | 0 | 0 | -0.17 | 3.54 | 1.34 | 0.84 | -0.73 | -1.27 |
| chiral non-Gaussian | chiral | 2.52 | -0.42 | 0 | - | - | -0.17 | 3.54 | 1.34 | 0.84 | - | - |
| chiral non-Gaussian | chiral | 2.52 | -0.42 | - | 0 | - | -0.17 | 3.54 | 1.34 | - | -0.69 | - |
| chiral non-Gaussian | chiral | 2.52 | -0.42 | - | - | 0 | -0.17 | 3.54 | 1.34 | - | - | -1.30 |
| non-Gaussian | none | 1.00 | -0.27 | 1.01 | 1.10 | -2.49 | -0.56 | 3.65 | 1.66 | 0.59 | $-2.50 \pm i 1.60$ | |
| non-Gaussian | none | 2.52 | -0.41 | - | 0.74 | - | -0.15 | 3.54 | $1.37 \pm i 0.04$ | | - | - |

chirally symmetric interacting fixed point:
realization of chiral symmetry on all scales

Conclusions & Outlook

quantum gravity might provide a UV completion for the Standard Model

- asymptotically safe solution to the U(1) triviality problem
- existence of matter fixed point requires weak gravity
- (phenomenologically) viable matter fixed point constrains gravitational parameter space
- existence of chiral fermions accommodated in ASQG (chiral symmetry not IR attractive)

Outlook:

