

Title: PSI 2016/2017 Quantum Information (Review) - Lecture 9 (Eduardo Martin-Martinez)

Date: Mar 06, 2017 09:00 AM

URL: <http://pirsa.org/17030033>

Abstract:

March 5-6th 2017 PSI

# Relativistic Quantum Information

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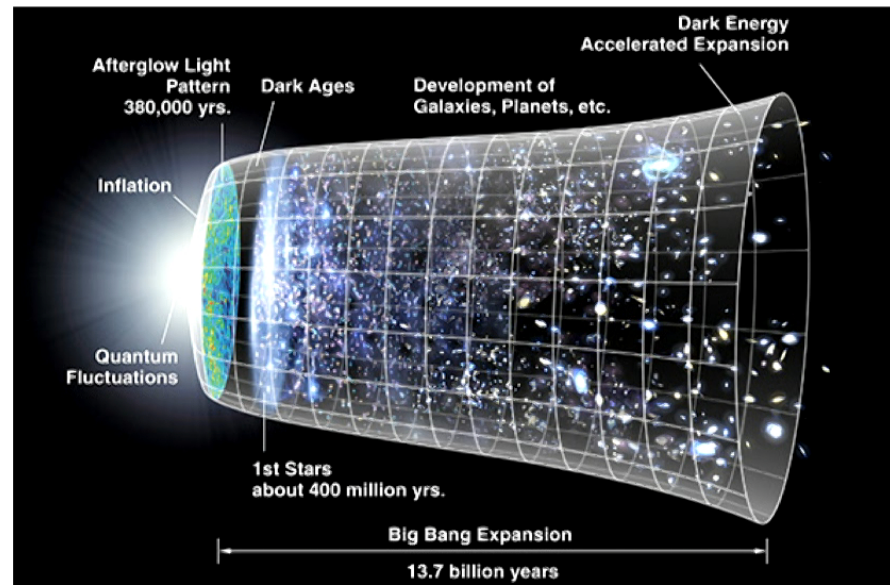


Eduardo Martín-Martínez

Professor of Applied Mathematics (University of Waterloo)  
Institute for Quantum Computing  
Perimeter Institute for Theoretical Physics



# Fundamental Topics: Cosmology



How much can we learn about the Early Universe nowadays?

# **Fundamental Topics: Black Hole Information Loss Problem**

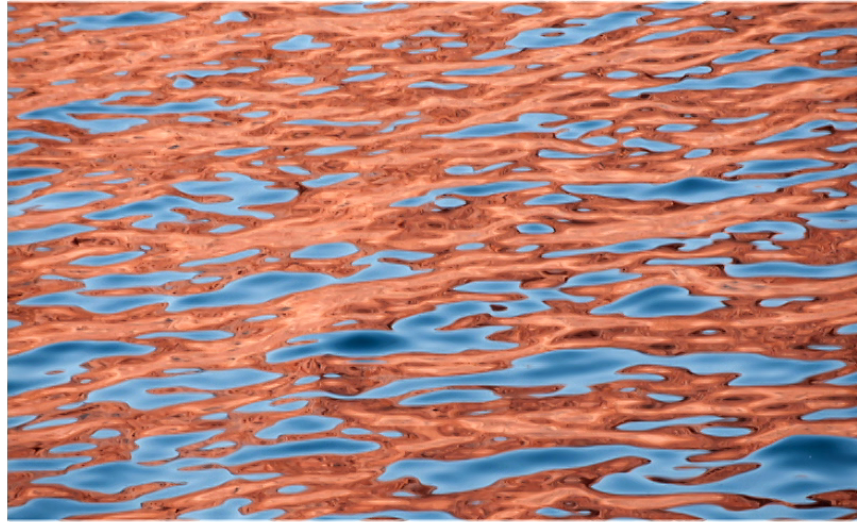


**Quantum Mechanics preserves information.**

**Black Holes: Does Nature destroy information?**

**Or does the information escape in the form of Hawking Radiation?**

# Fundamental Topics: Vacuum Fluctuations

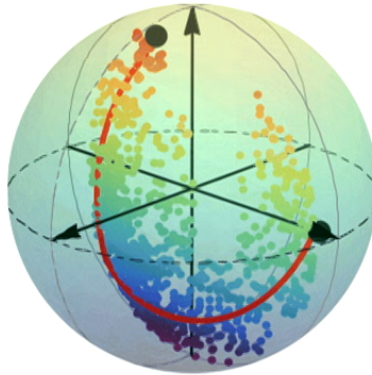


**In Nature, the vacuum is not empty. Only on average.**

**Vacuum Fluctuations contain Information about curvature of spacetime.**

**Quantum noise is special: It can assist communication!**

# **Fundamental Topics: Quantum Measurements and Localization**



**Quantum Theory is a probabilistic theory.**

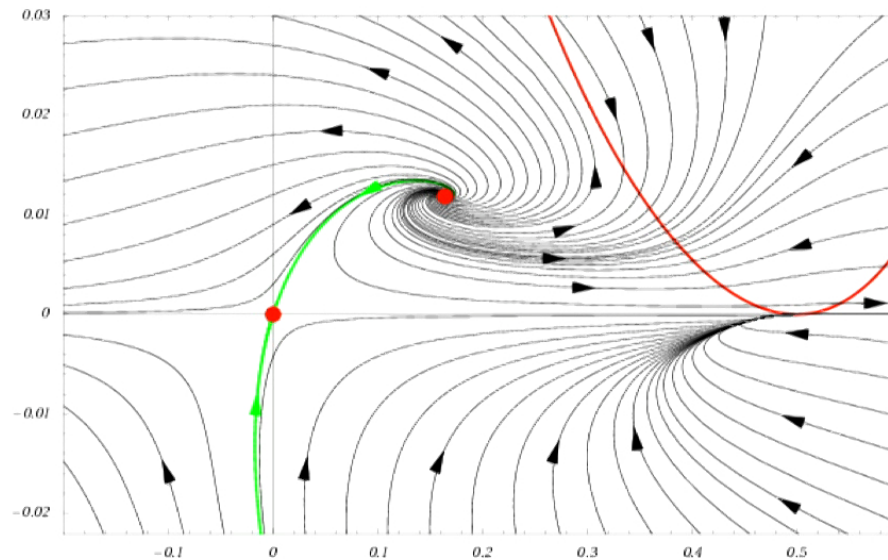
**The outcome of measurements is not deterministically predictable.**

**Quantum-to-Classical transition.**

**Relativistic considerations in the localization of Information**



# Fundamental Topics: Fixed points in Quantum evolution



The problem of equilibration in Quantum Theory and in Gravity.

Quantum Thermodynamics

Quantum Control

# Fundamental Topics: Quantum Gravity

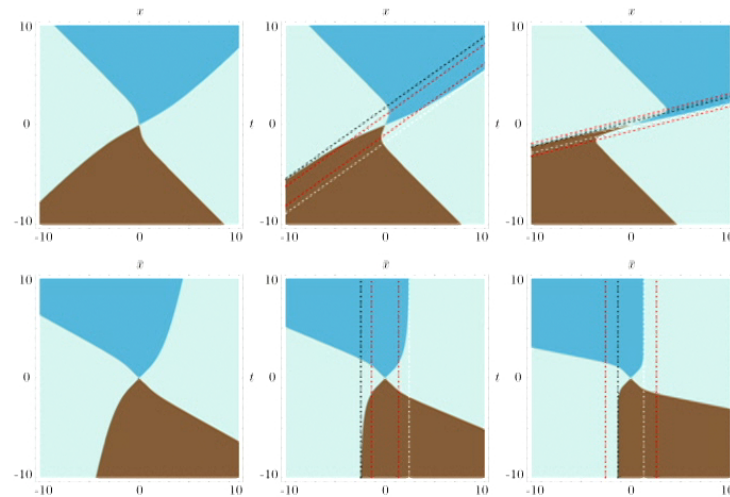


One of the most important challenges of modern Mathematical Physics:  
Quantum Theory for Gravitation

# Fundamental Topics: “Spacetime Engineering”

Violate energy conditions:

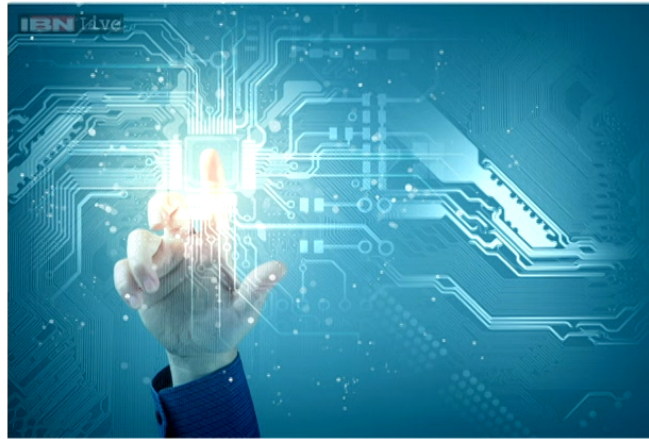
- Warp drives?
- Wormholes?



For more info: N.Funai, E. martin-Martinez: [arXiv:1701.03805](https://arxiv.org/abs/1701.03805)

# Technological Applications

Development of Mathematical tools that can be applied to experiments and technologies:



- Quantum Entanglement and Quantum Resources
- Communication
- Metrology
- Quantum Control and Simulations



# Getting Familiar with Spacetime

Relativity Matters for Quantum Information

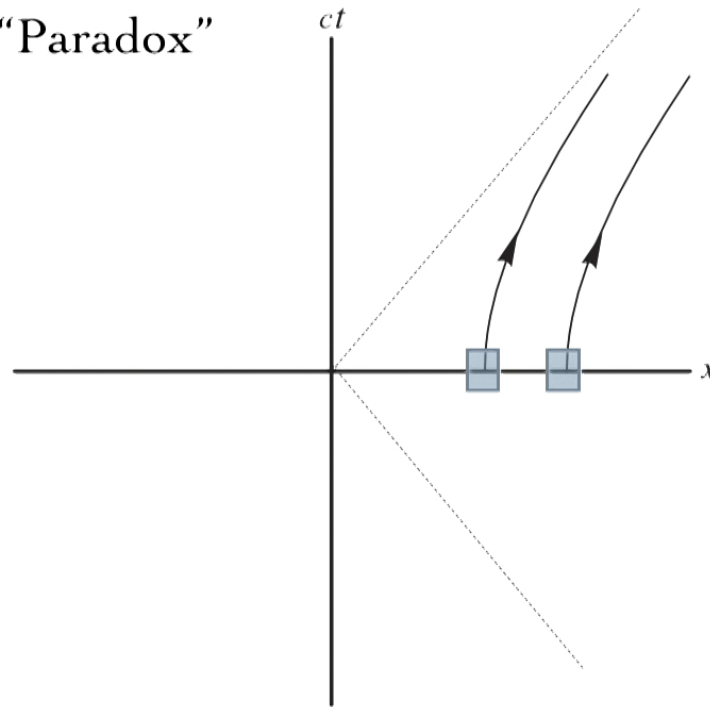
# Same Physics, Different Descriptions

Bell Rocket “Paradox”



# Same Physics, Different Descriptions

Bell Rocket "Paradox"

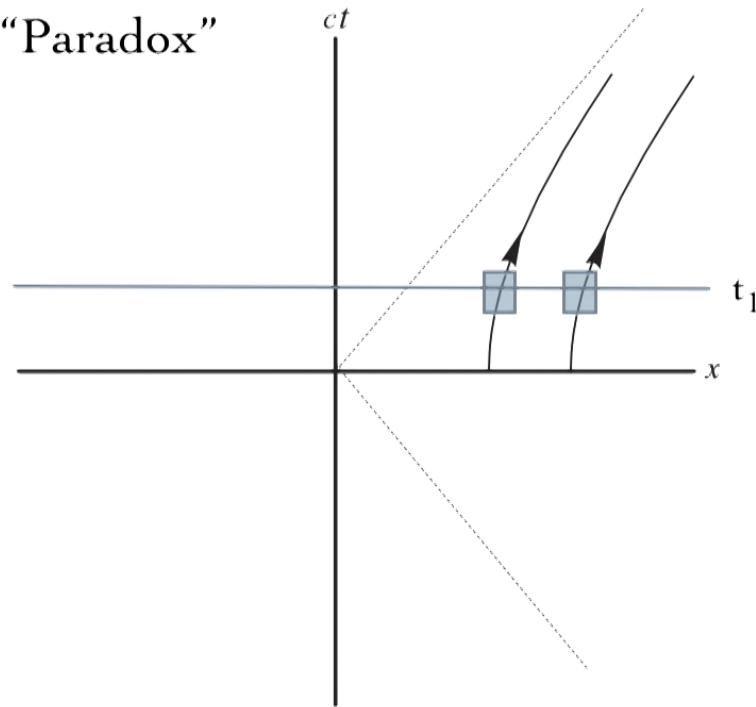


Does the rope break or not??

Why??

# Same Physics, Different Descriptions

Bell Rocket "Paradox"



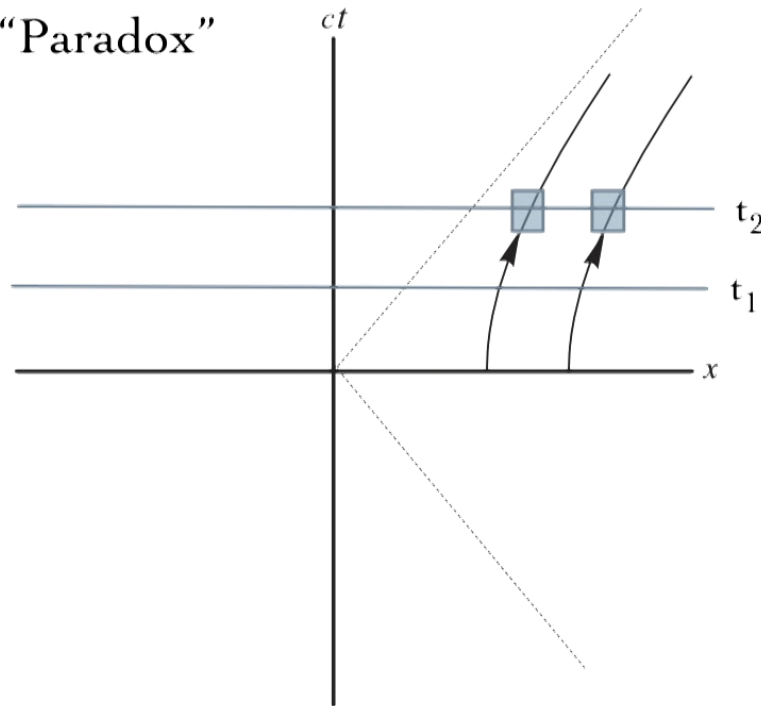
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# Same Physics, Different Descriptions

Bell Rocket "Paradox"

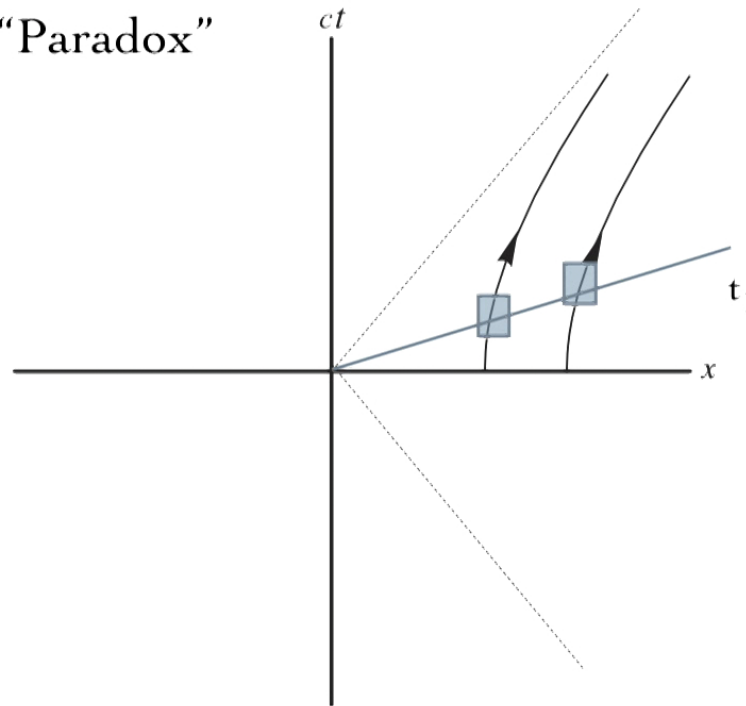


Does the rope break or not??

Why??

# Same Physics, Different Descriptions

Bell Rocket "Paradox"

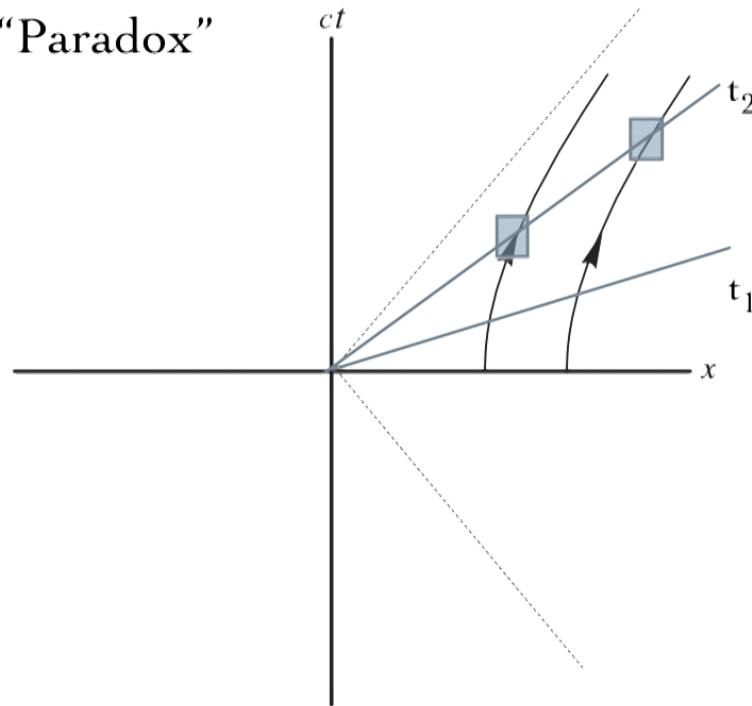


Does the rope break or not??

Why??

# Same Physics, Different Descriptions

Bell Rocket "Paradox"



Does the rope break or not??

Why??

# Same Physics, Different Descriptions

PHYSICS:

The rope breaks, all right!

PHENOMENOLOGY:

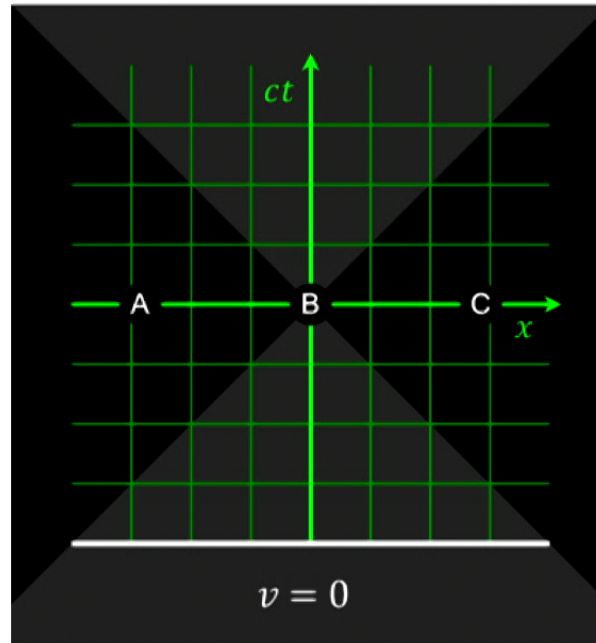
**For the accelerated observer A:** Because rocket B is faster than us!

**For the observer on the ground:** Because both rockets go equally faster and faster, the length of the rope Lorentz-contracts!



# Getting Familiar with Spacetime

Simultaneity is Relative!



Two observers in different states of motion would not agree about what happens first

# Getting Familiar with Quantum Mechanics

Schrödinger's cat

$$\frac{1}{\sqrt{2}}|\text{cat sitting}\rangle + \frac{1}{\sqrt{2}}|\text{cat lying}\rangle$$

# Quantum Entanglement



## EINSTEIN ATTACKS QUANTUM THEORY

Scientist and Two Colleagues  
Find It Is Not 'Complete'  
Even Though 'Correct.'

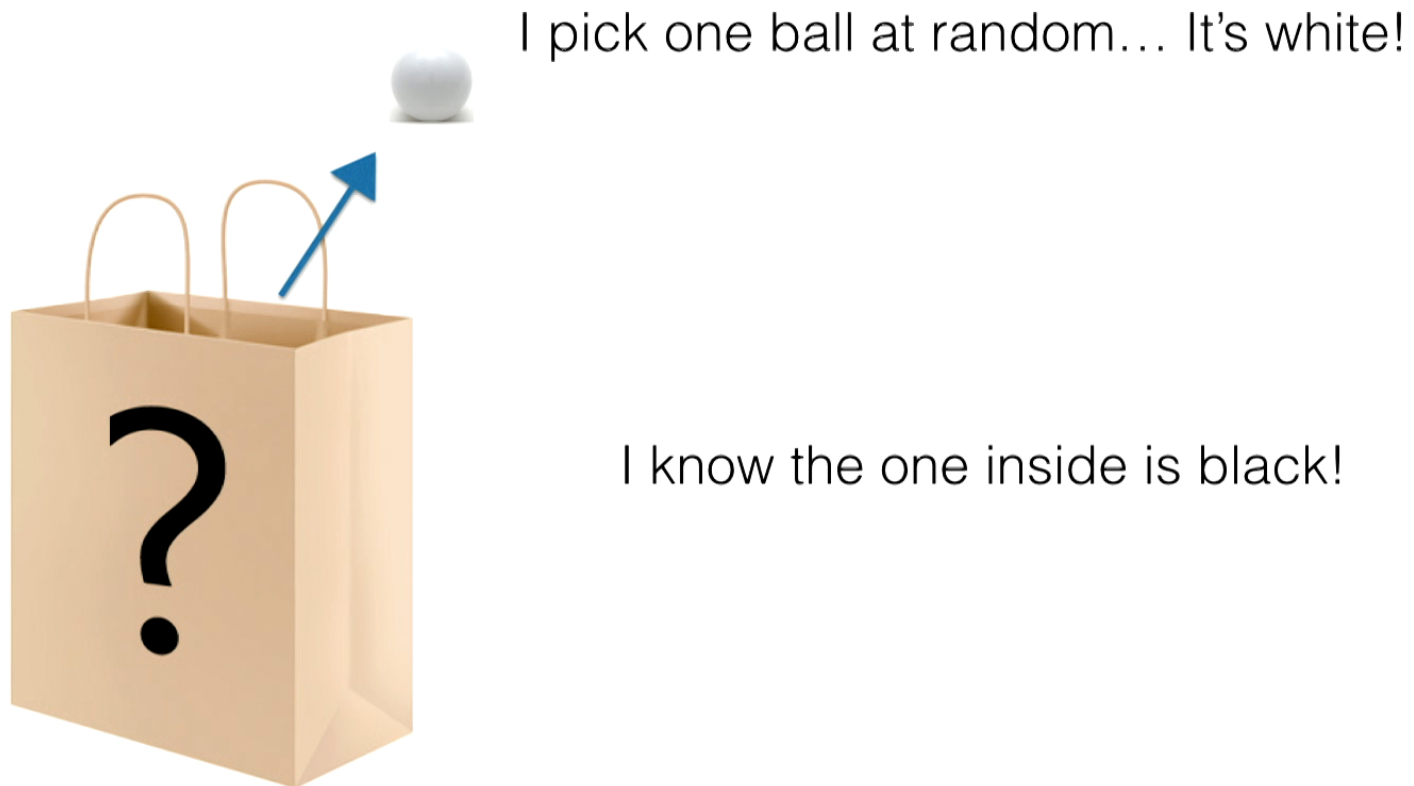
SEE FULLER ONE POSSIBLE

Believe a Whole Description of  
'the Physical Reality' Can Be  
Provided Eventually.

# Quantum Entanglement Vs Classical Correlations



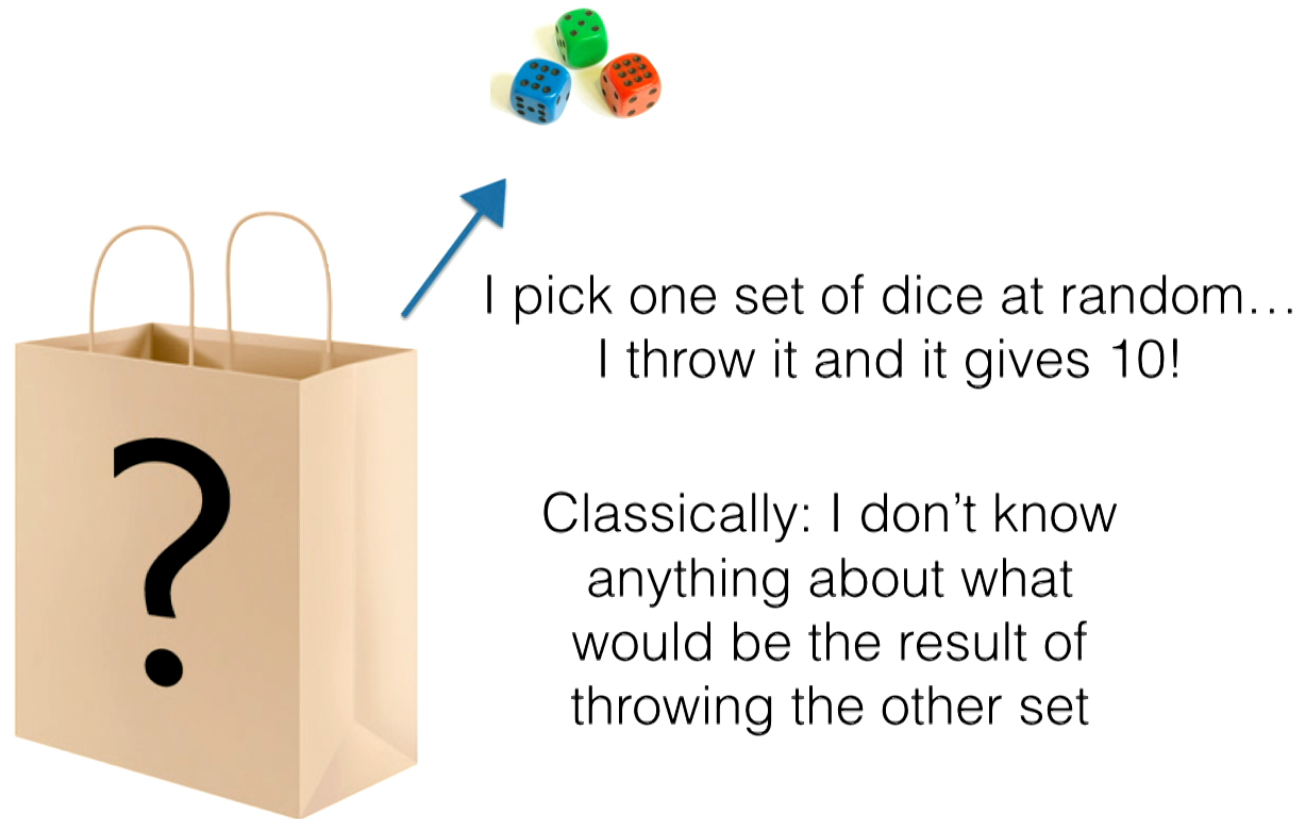
# Quantum Entanglement Vs Classical Correlations



# Quantum Entanglement Vs Classical Correlations



# Quantum Entanglement Vs Classical Correlations



# What entanglement is not

POPULAR  
SCIENCE SUBSCRIBE

## The Race To Prove 'Spooky' Quantum Connection May Have a Winner

Entanglement breakthrough could lead to unhackable Internet

By Devin Powell August 29, 2015

Particles don't obey the same rules as people. Poke a particle, and another one far away can instantly respond the touch -- without any messages passing through the space between, as if the two particles were one. "Entanglement" is what quantum physics calls the intimate connection.

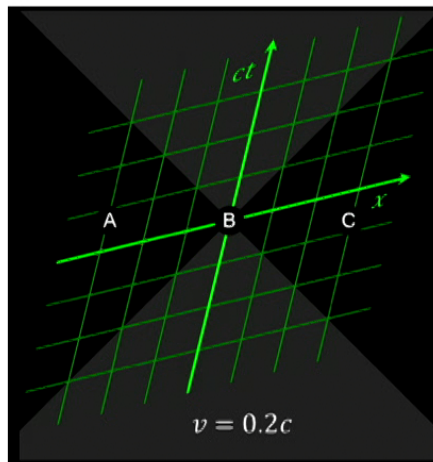
Einstein called it "spooky." To his dying day, he refused to believe that nature could be so unreasonable.

From <http://www.popsci.com>

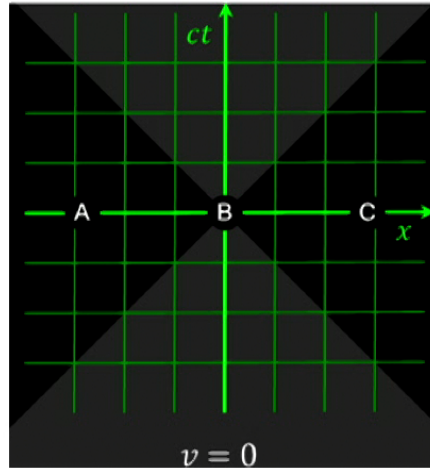


# Getting Familiar with Spacetime

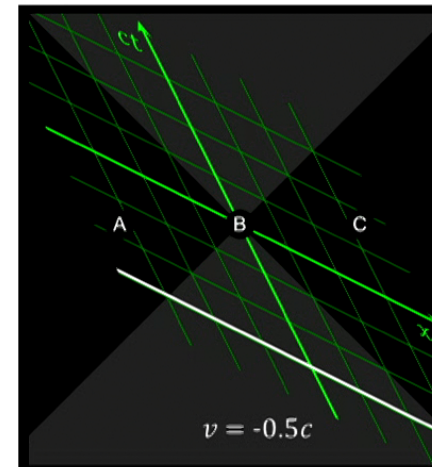
Simultaneity is Relative!  
Who collapses what??  
Who pokes what??



A happens after C

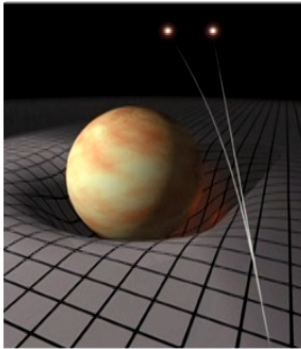


A and C are simultaneous

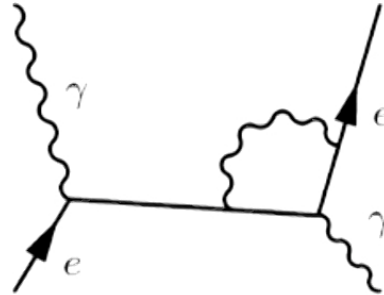


A happens before C

# Relativistic Quantum Information



General relativity



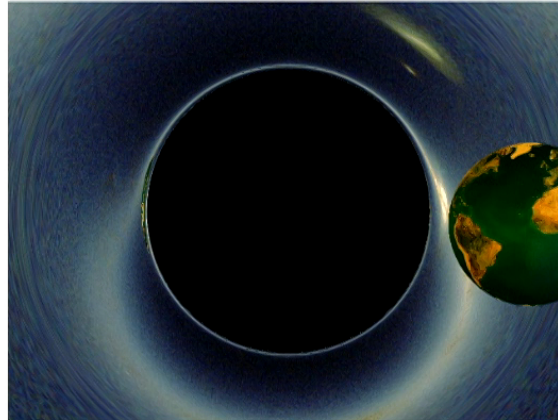
Quantum field theory



Quantum information

- Harness relativistic approaches to “do more” in quantum information processing.
- Study the structure of spacetime and the quantum nature of gravity via quantum informational tools

# THE BLACK HOLE INFORMATION PARADOX



# Entanglement in a Stellar Collapse

Once upon a time...

There was... NADA  $\Psi_0 = |0\rangle$

But then...

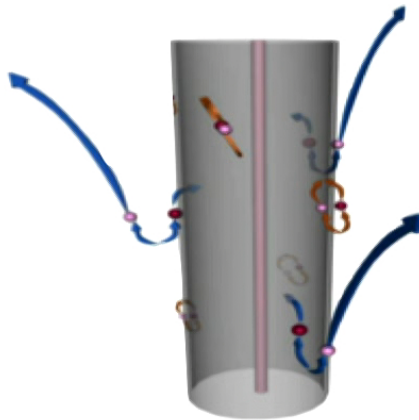
What happened to the field!?



# Entanglement in a Stellar Collapse

Vacuum in the far past evolves into two mode squeezed state between infalling and outgoing modes in the far future

$$\text{Tr}_{\text{hor}} (|0\rangle \langle 0|) = \bigotimes_{\omega} \frac{1}{\cosh^2 r} \sum \tanh^{2n} r_{\omega} |n_{\omega}\rangle_{\text{out}} \langle n_{\omega}|_{\text{out}}$$



We only see the modes that reach the future!



# Black hole Information loss problem

If we believe in quantum theory, information cannot be lost...

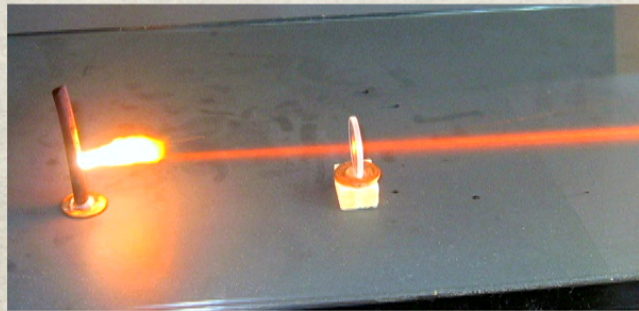


# Black hole Information loss problem

If we believe in quantum theory, information cannot be lost...

After corrections, the outflow may not be entirely thermal...

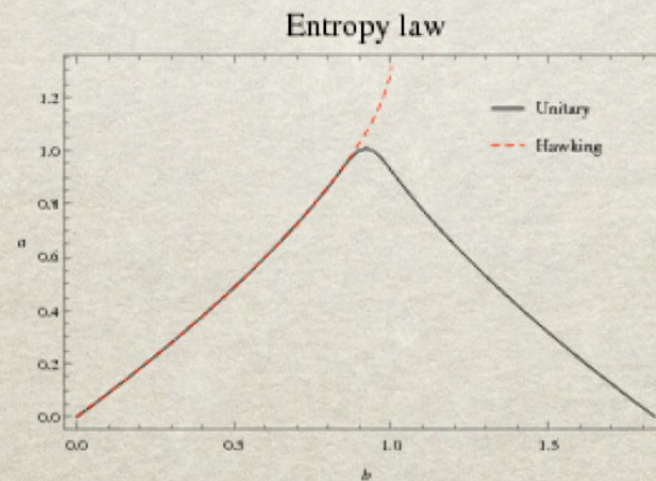
Like when a piece of charcoal burns



# Black holes Information loss problem

Page Hypothesis:

Entanglement between radiation emitted at different times  
in the black hole life!



Page time



# Black holes Information loss problem

So... The outflow is not entirely thermal...

Hold on!! that's potentially even worse!!

# Black hole Information Paradox

A: Radiation emitted after Page Time

B: Infalling Radiation

C: Radiation emitted before Page Time

Entropy subadditivity:

$$S(\rho_{ABC}) + S(\rho_A) \leq S(\rho_{AB}) + S(\rho_{AC})$$

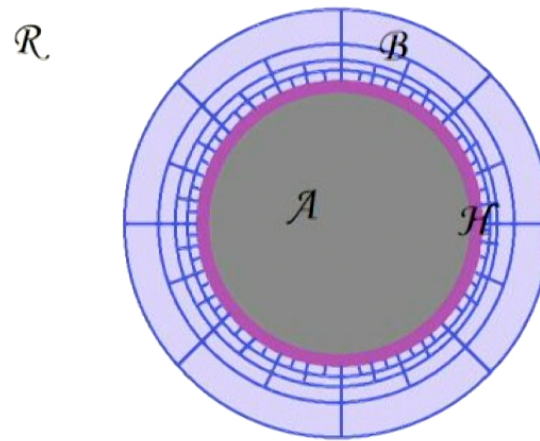
Entanglement subadditivity:

$$\mathcal{E}(A, B) + \mathcal{E}(A, C) \leq \mathcal{E}(A, BC)$$



# HAWKING RADIATION

What about monogamy?



Susskind's Black Hole complementary  
The Harlow-Hayden Conjecture

# Black hole Information Paradox

## Possible Solution: Firewalls!

Almheiri, Ahmed; Marolf, Donald; Polchinski, Joseph; Sully, James. Journal of High Energy Physics 2013 (2).




# Black hole Information Paradox (Firewalls)

What-if scenario:

Somehow dynamics is such that it destroys the correlations between “in” and “out” regions

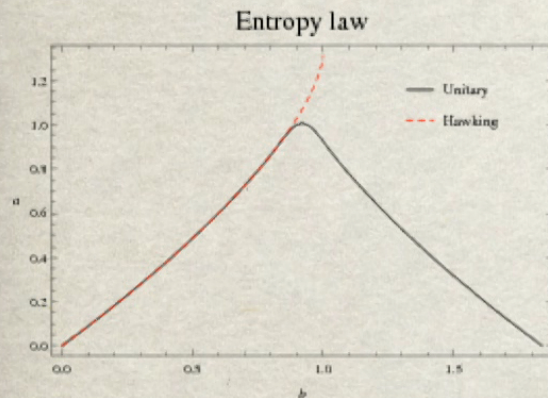
Entanglement subadditivity:

$$\mathcal{E}(A, B) + \mathcal{E}(A, C) \leq \mathcal{E}(A, BC)$$

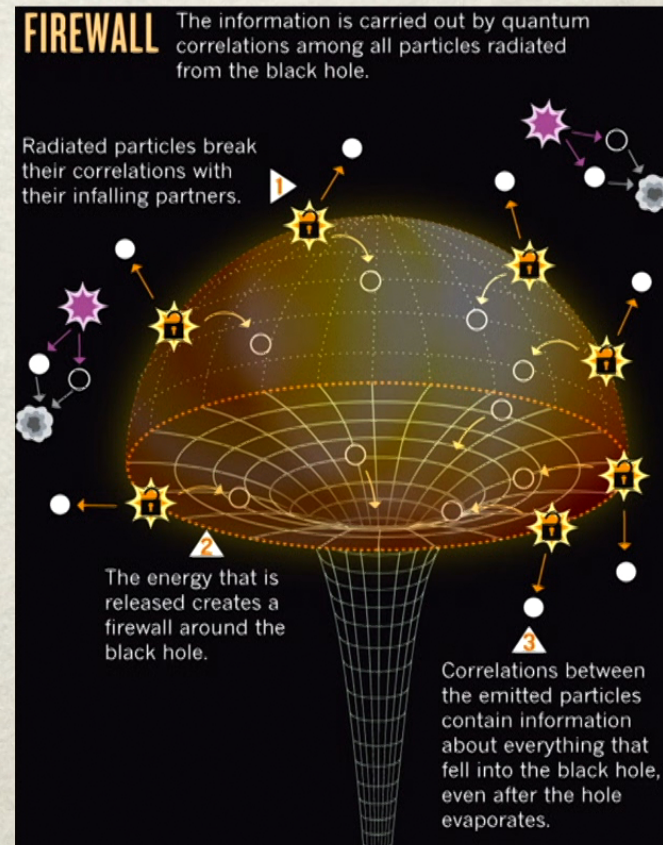
 Make this zero



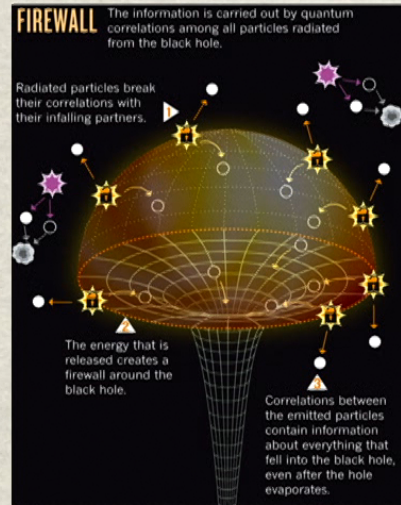
# Black hole Information Paradox (Firewalls)



“Charcoalization” of the BH



# Are Firewalls really ‘Monsters’?



Divergences in the stress-energy tensor: Violence at the horizon

PRL 115, 031301 (2015)

PHYSICAL REVIEW LETTERS

week ending  
17 JULY 2015

## $(1+1)$ D Calculation Provides Evidence that Quantum Entanglement Survives a Firewall

Eduardo Martín-Martínez<sup>1,2,3</sup> and Jorma Louko<sup>4</sup>

<sup>1</sup>*Institute for Quantum Computing, University of Waterloo, Waterloo, Ontario N2L 3G1, Canada*

<sup>2</sup>*Department of Applied Mathematics, University of Waterloo, Waterloo, Ontario N2L 3G1, Canada*

<sup>3</sup>*Perimeter Institute for Theoretical Physics, Waterloo, Ontario N2L 2Y5, Canada*

<sup>4</sup>*School of Mathematical Sciences, University of Nottingham, Nottingham NG7 2RD, United Kingdom*

(Received 26 February 2015; published 14 July 2015)

We analyze how preexisting entanglement between two Unruh-DeWitt particle detectors evolves when one of the detectors falls through a Rindler firewall in  $(1+1)$ -dimensional Minkowski space. The firewall effect is minor and does not wash out the detector-detector entanglement, in some regimes even preserving the entanglement better than Minkowski vacuum. The absence of cataclysmic events should continue to hold for young black hole firewalls. A firewall's prospective ability to resolve the information paradox must hence hinge on its detailed gravitational structure, presently poorly understood.

DOI: 10.1103/PhysRevLett.115.031301

PACS numbers: 04.70.Dy, 04.60.-m, 04.62.+v

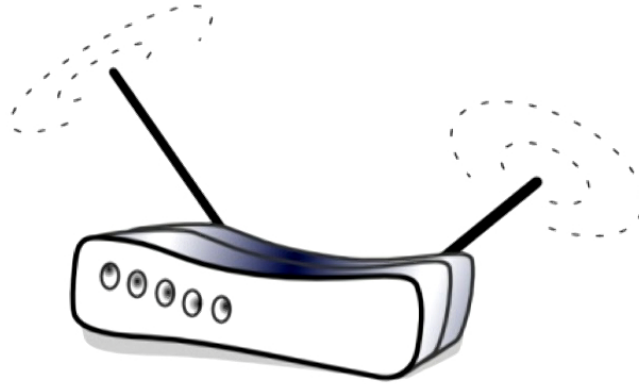
# The two results we will talk about

Vacuum Entanglement Harvesting and Farming

Quantum Collect Calling



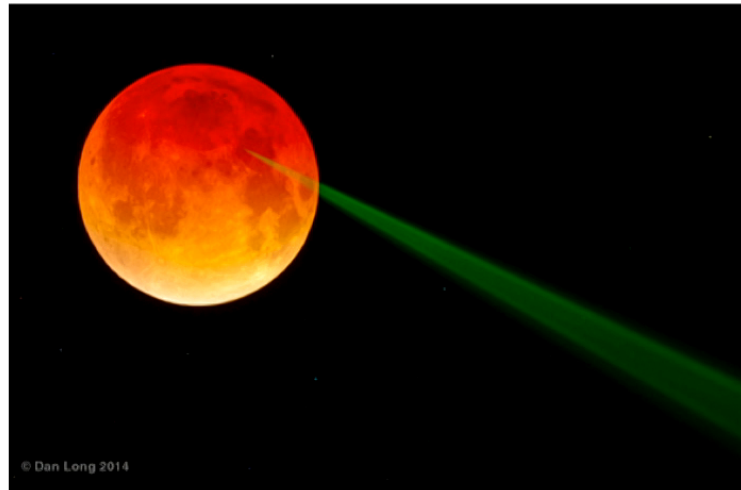
# Communication through massless fields



**General properties of wireless communication**

# Communication through the EM field

Communication mediated by 'real' energy-carrying quanta



An emitter emits photons. A receiver captures photons.

# Communication through the EM field

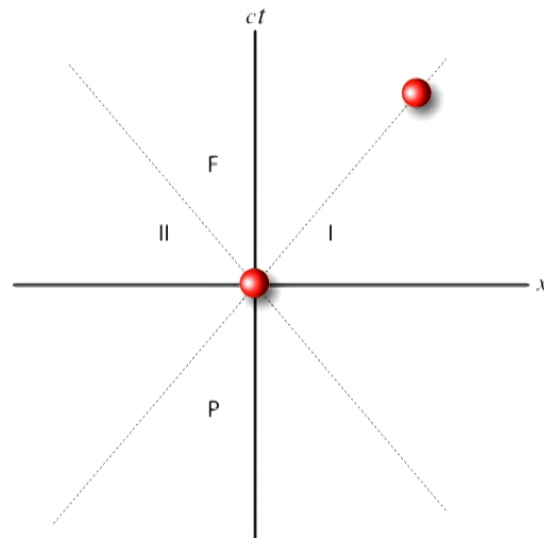
Information flow carried by (an average) energy flow



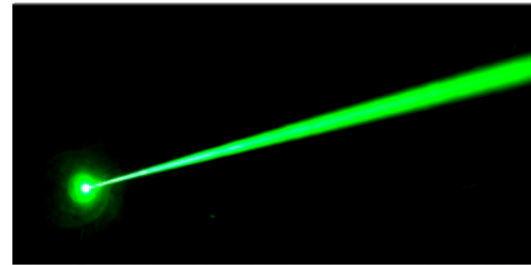
Information reaches you when energy reaches you

# Communication through the EM field

Communication is **only** possible at the speed of light (in vacuum)



If you miss the beam, you miss the message

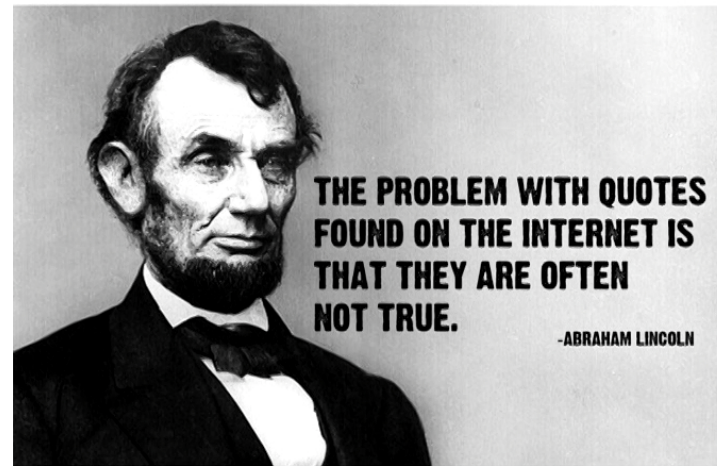




# Communication through massless fields

## Communication through a masses fields in the vacuum

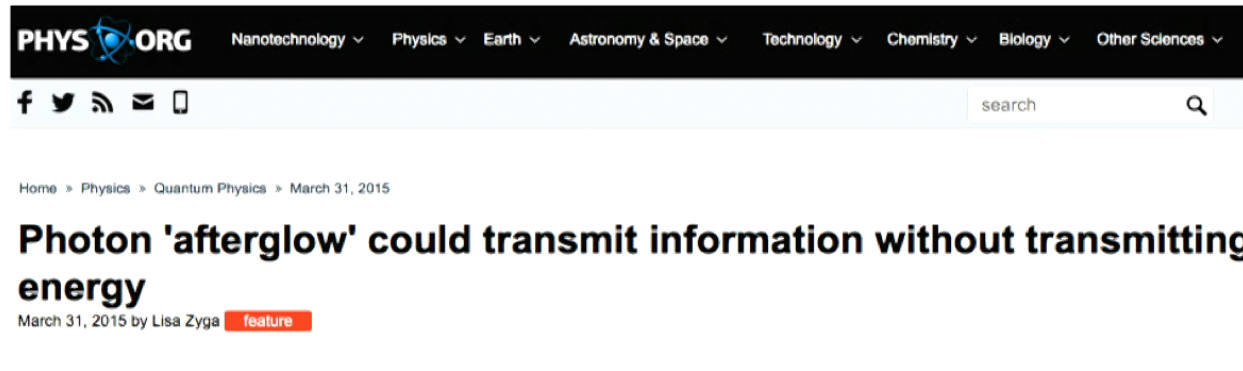
- Only At the speed of light.
- Through the exchange of real quanta
- Information flow carried by energy flow.
- Miss the beam, miss the message



# Communication through massless fields

## Communication through a massless fields in vacuum

- Information propagates arbitrarily slow even for massless field.
- Recover the message even if the beam is missed.
- Information flow not supported by real quanta (photons) flow.
- Information flow in absence of energy flow.



# Mathematical Methods: Beyond the Strong Huygens Principle

Subtleties in the behaviour of the solutions of certain PDEs:  
The strong Huygens principle

*The Green's function of the (massless) wave equation in 3+1D Minkowski space has support only on the light cone. Hence, any disturbances propagate strictly along null geodesics (at the speed of light)*

# Mathematical Methods: Beyond the Strong Huygens Principle

Subtleties in the behaviour of the solutions of certain PDEs:  
The strong Huygens principle

*The Green's function of the (massless) wave equation in 3+1D Minkowski space has support only on the light cone. Hence, any disturbances propagate strictly along null geodesics (at the speed of light)*

Exploitable when emitters are quantum!

# TECHNICAL DETAILS

R. H. Jonsson, E. Martin-Martinez, A. Kempf, Phys. Rev. Lett. 114, 110505 (2015)

A. Blasco, L. J. Garay, M. Martin-Benito, E. Martin-Martinez, Phys. Rev. Lett. 114, 141103 (2015)

A. Blasco, L. J. Garay, M. Martin-Benito, E. Martin-Martinez, Phys. Rev. D 93, 024055 (2016)

P. Simidzija, E. Martin-Martinez, Phys. Rev. D 95, 025002 (2017)

## See also:

R. H. Jonsson, J. of Phys. A, 44, 445402 (2016)

## STRONG HUYGENS PRINCIPLE

The radiation Green's function (or equivalently the commutator) of a massless field has support only on the light-cone

$$\square G(x, x') = -4\pi\delta_4(x, x') \qquad [\Phi(x), \Phi(x')] = \frac{i}{4\pi}G(x, x')$$

————→ **Communication** has support only on the **light-cone**

True in 3+1 Flat spacetime

## BEYOND THE STRONG HUYGENS PRINCIPLE

In general: if there is curvature (unless there is conformal invariance)

————→ In curved spacetimes, **communication through massless fields** is not confined to the light-cone, but there can be a leakage of information towards the **inside of the light-cone decoupled from energy propagation**.

SPATIALLY **FLAT**, **OPEN FRW** SPACETIME 3+1D:

$$ds^2 = a(\eta)^2(-d\eta^2 + dr^2 + r^2 d\Omega^2)$$

$\eta$  : conformal time  
 $a(\eta)$  : scale factor  
 $t$  : cosmological time,  
 $dt = a(\eta)d\eta$   
 units:  $\hbar = c = 1$

This geometry will be generated by:

a **perfect fluid** with a constant density-to-pressure ratio  $p = w\rho$

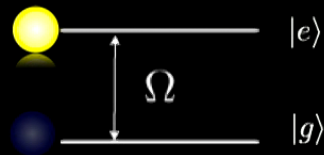
→ the **scale factor** evolves as  $a \propto \eta^{\alpha + \frac{1}{2}} \propto t^{\frac{2\alpha + 1}{2\alpha + 3}}$  with  $\alpha = \frac{3 - 3w}{6w + 2}$

A TEST **SCALAR FIELD** QUANTIZED IN THE **BUNCH-DAVIS VACUUM**  
WILL BE COUPLED TO THE BACKGROUND GEOMETRY.



## Unruh-DeWitt DETECTOR

-Two-level system



-Energy gap ground-excited states:

$\Omega$

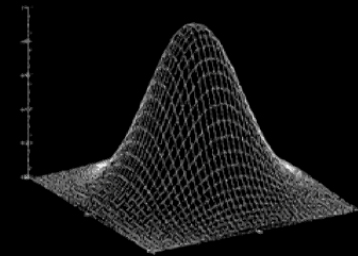
-Monopole moment operator:

$$\nu = \{A, B\}$$

$$\mu_\nu(t) = |e_\nu\rangle\langle g_\nu|e^{i\Omega_\nu t} + |g_\nu\rangle\langle e_\nu|e^{-i\Omega_\nu t}$$

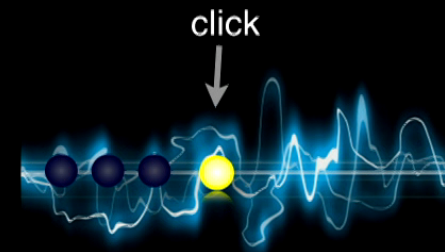
-Spatially smeared:

$$F(\vec{x}, t) = \frac{1}{\sigma^3 \sqrt{\pi^3}} e^{-a(t)^2 \vec{x}^2 / \sigma^2}$$



Detectors:  $|\psi_\nu\rangle = \alpha_\nu |e_\nu\rangle + \beta_\nu |g_\nu\rangle$

## ALICE & BOB's DETECTOR MODEL



DETECTOR-FIELD  
**INTERACTION**  
HAMILTONIAN

$$H_{I,\nu} = \lambda_\nu \chi_\nu(t) \mu_\nu(t) \int d^3\mathbf{x} a(t)^3 F[\mathbf{x} - \mathbf{x}_\nu(t), t] \Phi[\mathbf{x}, \eta(t)]$$

Coupling strength

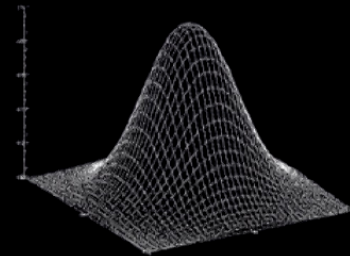
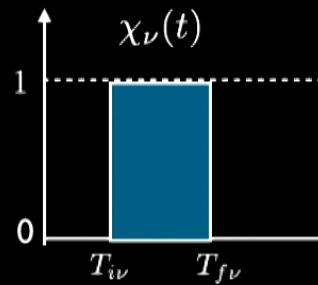
Monopole  
moment

Scale  
factor

Smearing function

Detector's trajectory

Switching function



Total Interaction  
Hamiltonian:

$$H_I = H_{I,A} + H_{I,B}$$

## TRANSMISSION OF INFORMATION

Influence of the presence of  $A$  on  $B$   $\longrightarrow$  **SIGNALING ESTIMATOR,  $S$**

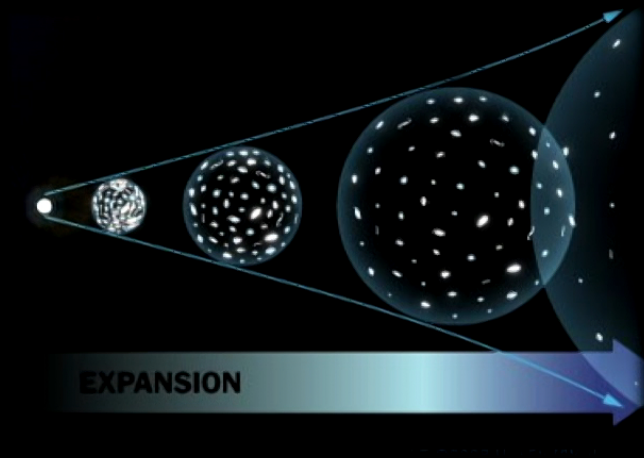
how much information can be sent?  $\longrightarrow$  **CHANNEL CAPACITY,  $C$**

## TRANSMISSION OF INFORMATION

Influence of the presence of A on B  $\longrightarrow$   $\left\{ \begin{array}{l} \text{SIGNALING ESTIMATOR, } \mathbf{S} \\ P_e(t) = |\alpha_B|^2 + P_{vac}(t) + \mathbf{S}(t). \end{array} \right.$

how much information can be sent?  $\longrightarrow$  **CHANNEL CAPACITY, C**

## THE **BIG BANG** Setting



BIG BANG CASE, ST. COSMOLOGICAL MODEL: GENERAL RELATIVITY

## SCALAR FIELD: COUPLING TO GRAVITY

### KLEIN-GORDON EQUATION

$$(\square - m^2 + \xi R)\phi = 0 \quad \square = \frac{1}{\sqrt{|g|}} \partial_\mu \left( \sqrt{|g|} g^{\mu\nu} \partial_\nu \right)$$

### CONFORMAL COUPLING

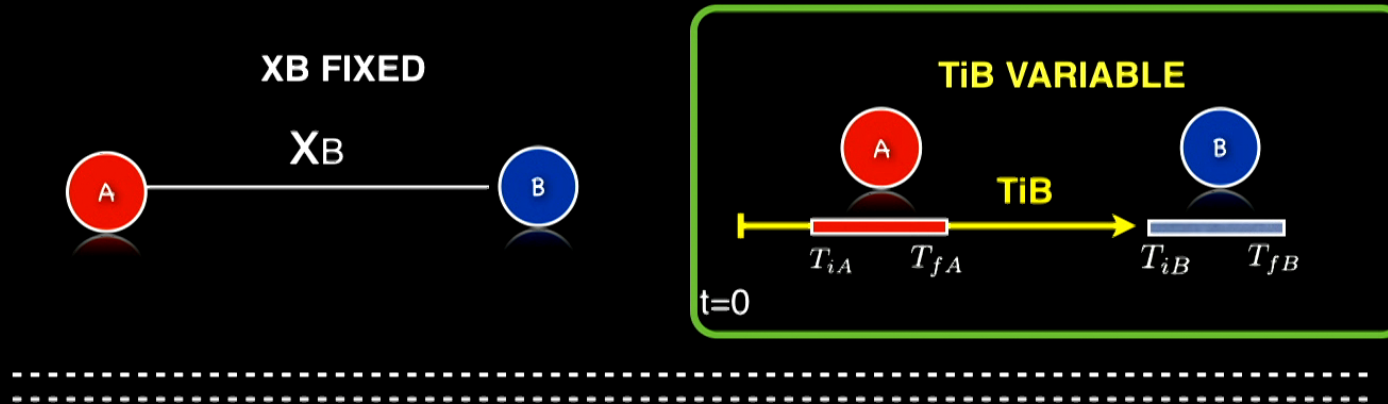
$$\xi = \frac{1}{6} \quad \text{Yields Conformally Invariant Action}$$

### MINIMAL COUPLING

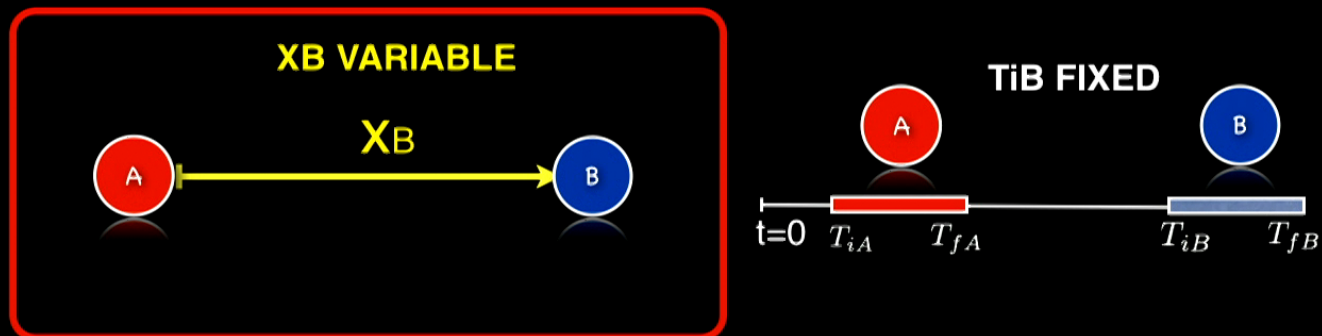
$$\xi = 0 \quad \text{Gives good predictions (Cosmology, etc..)}$$



## CASE : Variation of **temporal** separation



## CASE: Variation of **spatial** separation



## IS INFORMATION TRANSMITED?

Influence of the presence of A on B  $\longrightarrow$   $\left\{ \begin{array}{l} \text{SIGNALING ESTIMATOR, } \mathbf{S} \\ P_e(t) = |\alpha_B|^2 + P_{vac}(t) + \mathbf{S}(t). \end{array} \right.$

CONFORMAL COUPLING

$$S = \lambda_A \lambda_B S_2 + \mathcal{O}(\lambda_\nu^4)$$

$$S_2 = 4 \int a(t)^3 d^3 \mathbf{x} dt \int a(t')^3 d^3 \mathbf{x}' dt' \chi_A(t) \chi_B(t') \text{Re}(\alpha_A^* \beta_A) F(\mathbf{x} - \mathbf{x}_A, t) \\ \times F(\mathbf{x}' - \mathbf{x}_B, t) \text{Re}(\alpha_B^* \beta_B [\phi(\mathbf{x}, t), \phi(\mathbf{x}', t')])$$

$$[\phi(\mathbf{x}, t), \phi(\mathbf{x}', t')] = \frac{i}{4\pi} \left[ \frac{\delta(\Delta\eta + |\mathbf{x} - \mathbf{x}'|) - \delta(\Delta\eta - |\mathbf{x} - \mathbf{x}'|)}{a(t)a(t')|\mathbf{x} - \mathbf{x}'|} \right]$$

$$\Delta\eta = \eta(t) - \eta(t')$$

$$|\psi_{0,\nu}\rangle = \alpha_\nu |e_\nu\rangle + \beta_\nu |g_\nu\rangle$$

## CHANNEL CAPACITY

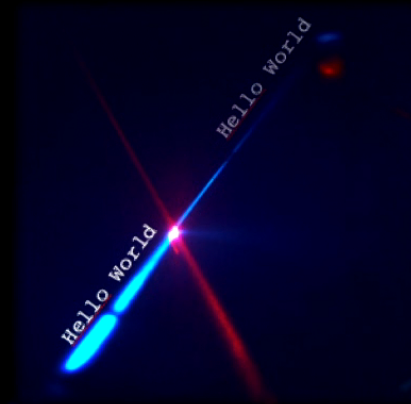
To obtain a lower bound to the channel capacity, we use a simple  
**COMMUNICATION PROTOCOL:**

- **Alice** encodes “1” by coupling her detector A to the field, and “0” by not coupling it.
- Later **Bob** switches on B and measures its energy. If B is excited, Bob interprets a “1”, and a “0” otherwise.

$$C \simeq \lambda_A^2 \lambda_B^2 \frac{2}{\ln 2} \left( \frac{S_2}{4|\alpha_B||\beta_B|} \right)^2 + \mathcal{O}(\lambda_\nu^6)$$

(noisy asymmetric binary channel)

Robert H. Jonsson, Eduardo Martín-Martínez, and Achim Kempf.  
Quantum Collect Calling.  
Phys. Rev. Lett. 114, 110505 (2015).



## CONFORMAL COUPLING

$$[\phi(\mathbf{x}, t), \phi(\mathbf{x}', t')] = \frac{i}{4\pi} \left[ \frac{\delta(\Delta\eta + |\mathbf{x} - \mathbf{x}'|) - \delta(\Delta\eta - |\mathbf{x} - \mathbf{x}'|)}{a(t)a(t')|\mathbf{x} - \mathbf{x}'|} \right]$$

support on the  
light cone

Decay with Spatial  
separation

Not surprising:

Conformal Symmetry makes it  
too similar to flat spacetime

## SIGNALING ESTIMATOR, S

$$[\phi(\mathbf{x}_A, t_A), \phi(\mathbf{x}_B, t_B)] = i \frac{\theta(\eta(t_B) - \eta(t_A)) - \theta(\eta(t_A) - \eta(t_B))}{(2\pi)^3 |\mathbf{x} - \mathbf{x}'| a(\eta(t_A)) a(\eta(t_B))} \int_0^\infty dk k \sin(k|\mathbf{x} - \mathbf{x}'|) \hat{g}(\eta(t_A), \eta(t_B), k)$$

$$\hat{g}(\eta, \eta', k) = \frac{8\pi}{k} \sqrt{\left| \frac{\eta}{\eta'} \right|} \frac{\text{sgn}(\eta') [J_{\alpha-1/2}(k|\eta|) Y_{\alpha-1/2}(k|\eta'|) - Y_{\alpha-1/2}(k|\eta|) J_{\alpha-1/2}(k|\eta'|)]}{Y_{\alpha-1/2}(k|\eta'|) [J_{\alpha-3/2}(k|\eta'|) - J_{\alpha+1/2}(k|\eta'|)] - J_{\alpha-1/2}(k|\eta'|) [Y_{\alpha-3/2}(k|\eta'|) - Y_{\alpha+1/2}(k|\eta')]$$

$J_\alpha$   $Y_\alpha$  BESSEL FUNCTIONS

**MATTER DOMINATED  
UNIVERSE**  $\longrightarrow \alpha = 2 \longrightarrow a \propto \eta^2 \propto t^{2/3}$

$$J_{2-1/2}(k|\eta|) = \sqrt{2/\pi} \frac{1}{\sqrt{k|\eta|}} \left[ -\cos(k\eta) + \frac{\sin(k\eta)}{k\eta} \right]$$

$$Y_{2-1/2}(k|\eta|) = \sqrt{2/\pi} \frac{\text{sgn}(\eta)}{\sqrt{k|\eta|}} \left[ -\sin(k\eta) + \frac{\cos(k\eta)}{k\eta} \right]$$



MINIMAL COUPLING

SIGNALING  
ESTIMATOR, S

MATTER DOMINATED  
UNIVERSE  $\longrightarrow \alpha = 2 \longrightarrow a \propto \eta^2 \propto t^{2/3}$

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## CONFORMAL COUPLING

**SIGNALING  
ESTIMATOR, S**

$$[\phi(\mathbf{x}, t), \phi(\mathbf{x}', t')] = \frac{i}{4\pi} \left[ \frac{\delta(\Delta\eta + |\mathbf{x} - \mathbf{x}'|) - \delta(\Delta\eta - |\mathbf{x} - \mathbf{x}'|)}{a(t)a(t')|\mathbf{x} - \mathbf{x}'|} \right]$$

support on the  
light cone

Decay with Spatial  
separation

**VIOLATION OF  
STRONG HUYGENS  
PRINCIPLE !!!!**

## MINIMAL COUPLING

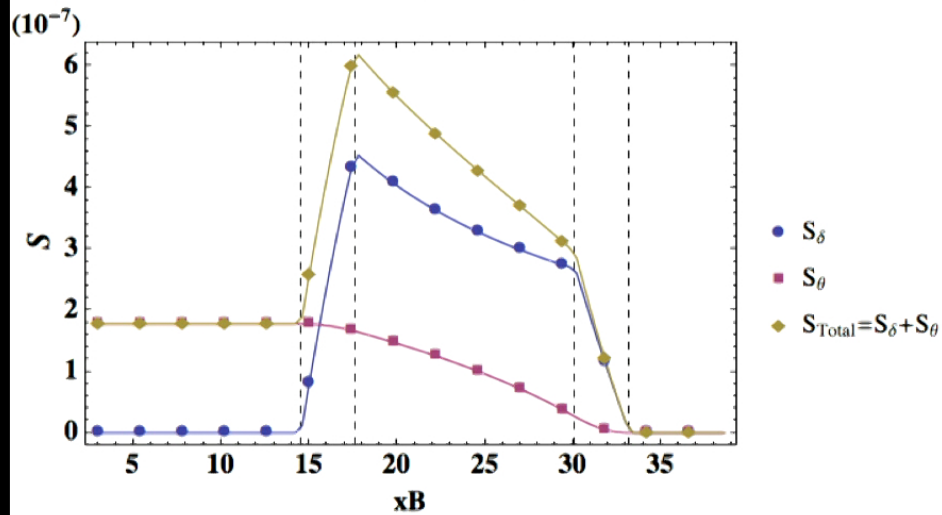
$$[\phi(\mathbf{x}, t), \phi(\mathbf{x}', t')] = \frac{i}{4\pi} \left[ \frac{\delta(\Delta\eta + |\mathbf{x} - \mathbf{x}'|) - \delta(\Delta\eta - |\mathbf{x} - \mathbf{x}'|)}{a(t)a(t')|\mathbf{x} - \mathbf{x}'|} + \frac{\theta(-\Delta\eta - |\mathbf{x} - \mathbf{x}'|) - \theta(\Delta\eta - |\mathbf{x} - \mathbf{x}'|)}{a(t)a(t')\eta(t)\eta(t')} \right]$$

Does NOT decay  
with Spatial  
separation

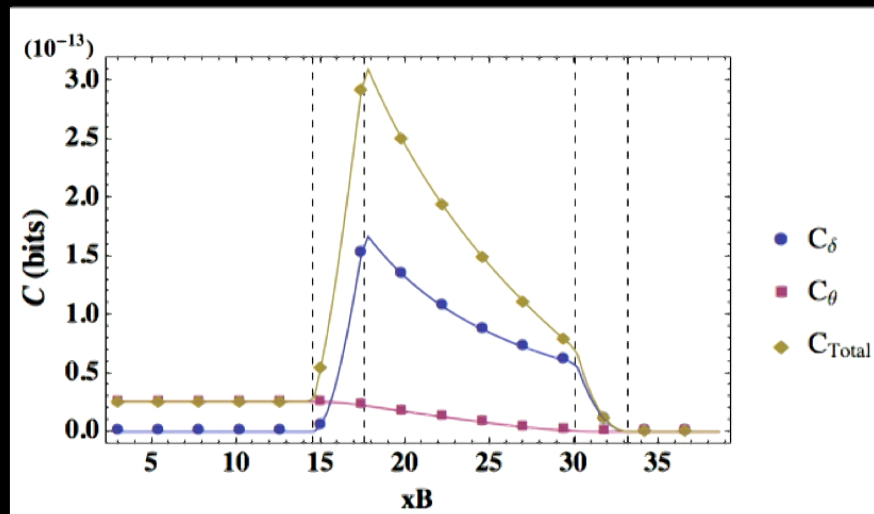
Timelike-  
leakage

Case:  
Variation of  
**spatial**  
separation

## MINIMAL COUPLING



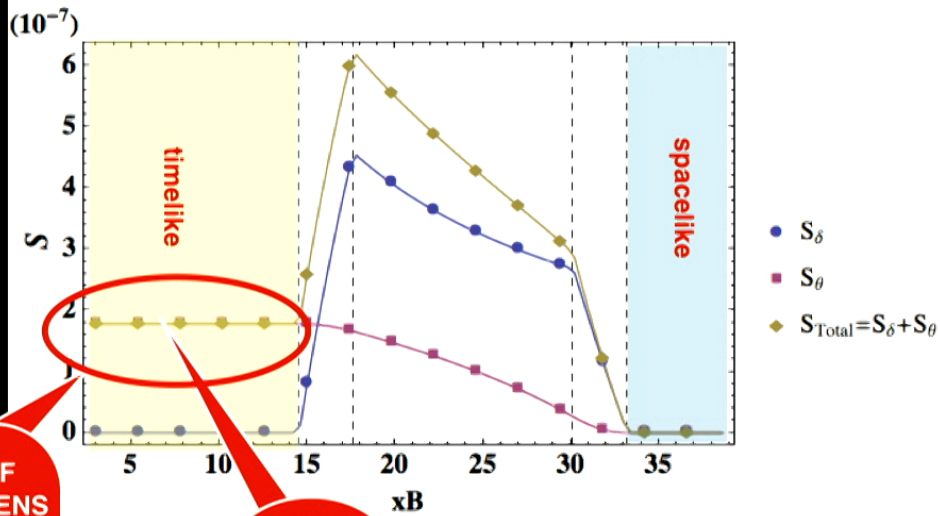
**SIGNALING**  
ESTIMATOR,  $S$



**CHANNEL**  
**CAPACITY**

Case:  
Variation of  
**spatial**  
separation

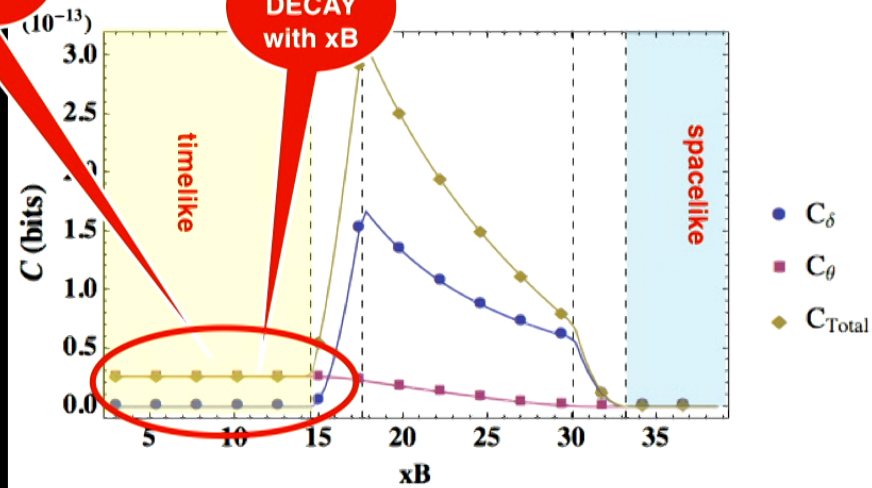
## MINIMAL COUPLING



**SIGNALING  
ESTIMATOR,  $S$**

**VIOLATION OF  
STRONG HUYGENS  
PRINCIPLE !!!!**

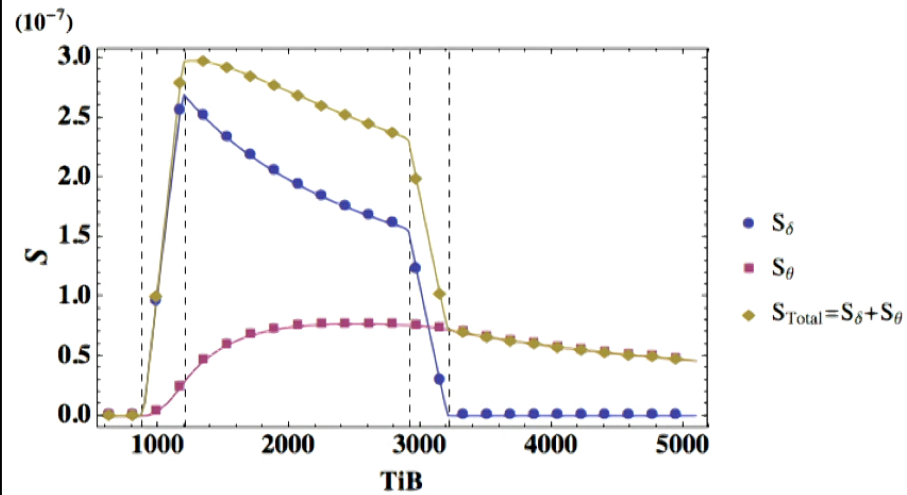
**NO  
DECAY  
with  $xB$**



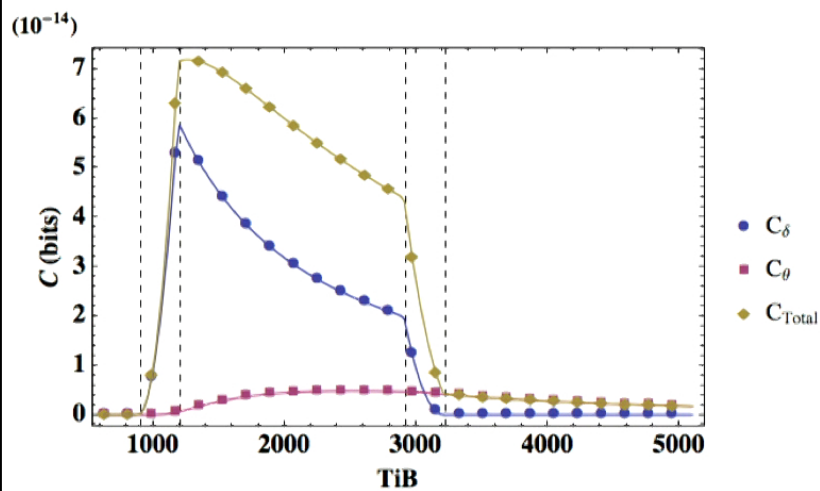
**CHANNEL  
CAPACITY**

Case:  
Variation of  
temporal  
separation

## MINIMAL COUPLING



**SIGNALING  
ESTIMATOR, S**

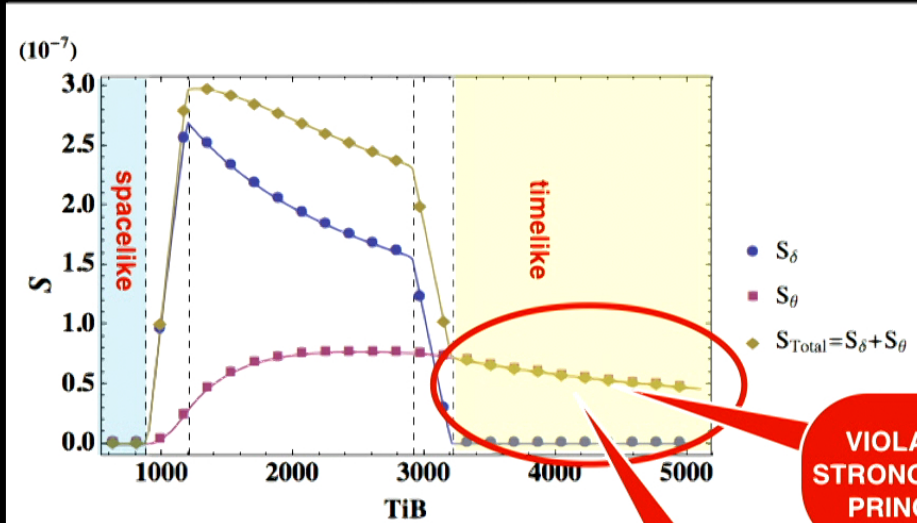


**CHANNEL  
CAPACITY**

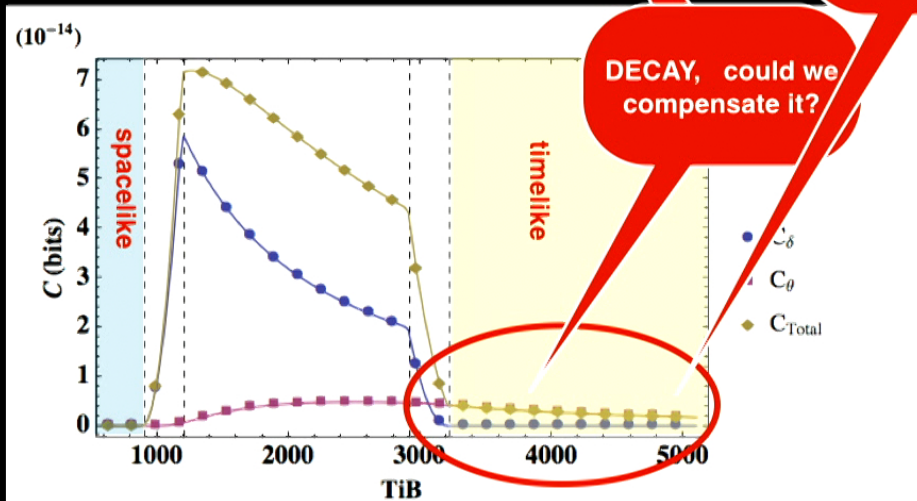


Case:  
Variation of  
**temporal**  
separation

## MINIMAL COUPLING



**SIGNALING  
ESTIMATOR,  $S$**



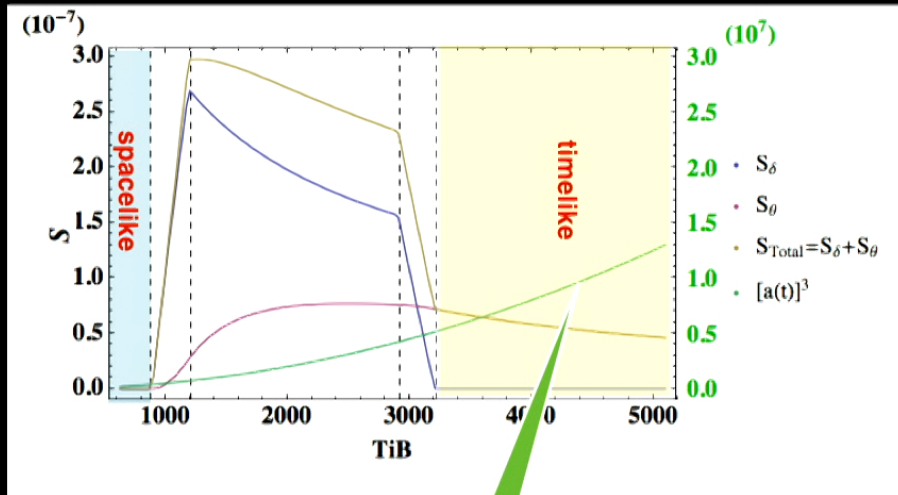
**CHANNEL  
CAPACITY**

**VIOLATION OF  
STRONG HUYGENS  
PRINCIPLE !!!!**

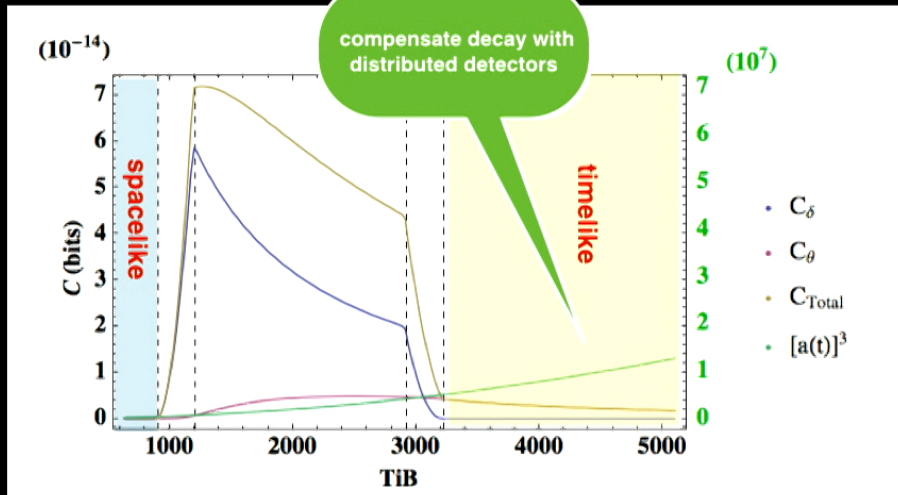
**DECAY, could we  
compensate it?**

Case:  
Variation of  
**temporal**  
separation

## MINIMAL COUPLING



**SIGNALING**  
ESTIMATOR, S



**CHANNEL**  
**CAPACITY**

# Exponential Expansion (deSitter): No decay in time!

P. Simidzija, E. Martin-Martinez, Phys. Rev. D 95, 025002 (2017)

Cosmological  
cataclysm!!!!

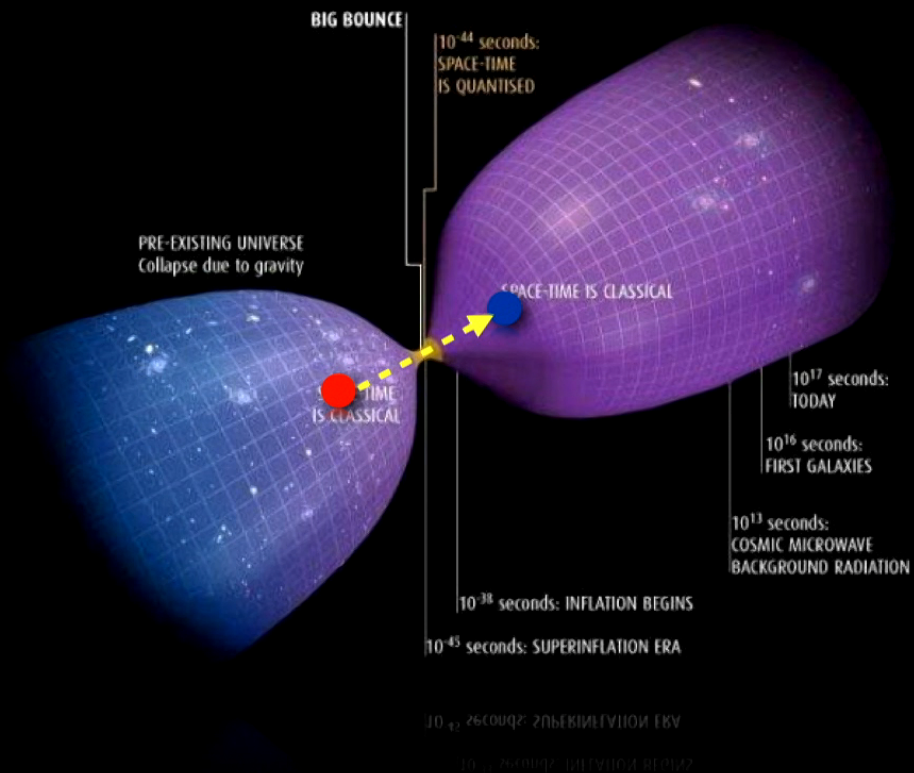


information

## Setting QUANTUM BOUNCE

### EXAMPLE: LOOP QUANTUM COSMOLOGY

- Replaces the Big Bang by a Big Bounce
- Bridge between two large classical universes: a contracting and an expanding cosmological phase



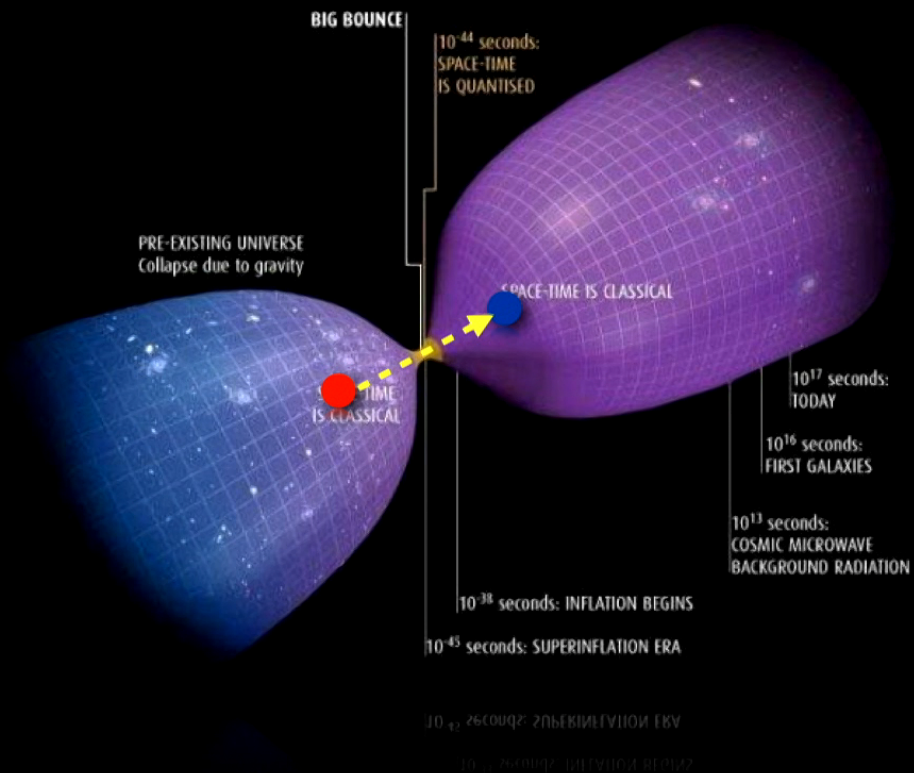
Martin Bojowald. Loop quantum cosmology. Living Rev.Rel., 11:4, 2008.



## Setting QUANTUM BOUNCE

### EXAMPLE: LOOP QUANTUM COSMOLOGY

- Replaces the Big Bang by a Big Bounce
- Bridge between two large classical universes: a contracting and an expanding cosmological phase



Martin Bojowald. Loop quantum cosmology. Living Rev.Rel., 11:4, 2008.

Atoms (or any complex system) will not survive a quantum bounce

Imagine an ancient (pre-bounce) and very advanced civilization

What would you do if you wanted your legacy to survive?



Encode the information in the quantum field:  
detectors and field get entangled.