

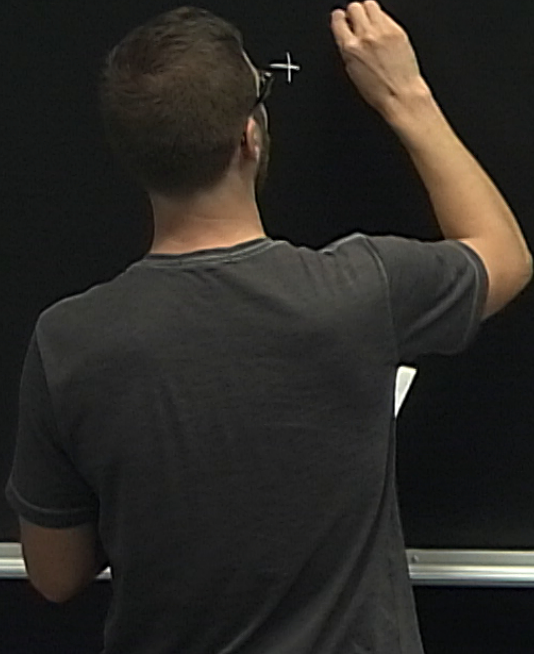
Title: 2016/2017 Statistical Mechanics 2 - Roger Melko - Lecture 17

Date: Mar 01, 2017 10:30 AM

URL: <http://pirsa.org/17030001>

Abstract:

$$\begin{aligned}
 S'[\psi_k] &= S_0[\psi_k] + \langle S_{int} \rangle_{0, \gamma} - \frac{1}{2} \left[\langle S_{int}^2 \rangle_{0, \gamma} - \langle S_{int} \rangle_{0, \gamma}^2 \right] \\
 &= \frac{1}{2} \int_0^{1/b} \frac{d^d k}{(2\pi)^d} (k^2 + r) |\psi_k(k)|^2
 \end{aligned}$$



$$\delta'[\Psi_L] = S_0[\Psi_L] + \langle S_{int} \rangle_{0,1} - \frac{1}{2} \left[\langle S_{int}^2 \rangle_{0,1} - \langle S_{int} \rangle_{0,1}^2 \right]$$

$$= \frac{1}{2} \int_0^{M_b} \frac{d^d k}{(2\pi)^d} (k^2 + r) |\Psi_L(k)|^2$$

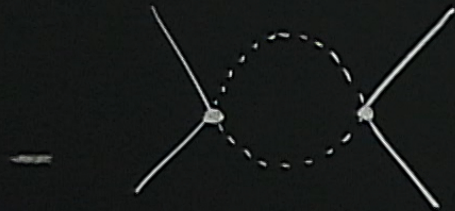
$$+ \frac{32\pi}{2} I_1 \int_0^1 \frac{d^d k}{(2\pi)^d} |\Psi_L(k)|^2 \quad \text{"D term"}$$

$$- \frac{1}{2} \left\langle \begin{array}{c} \text{fast} \\ \text{slow} \end{array} \right\rangle + \begin{array}{c} \text{X} \\ \text{X} \end{array} + \dots \right\rangle_{0,1} \quad (25 \text{ terms})$$

$$+ \frac{1}{2} \left\langle \begin{array}{c} \text{X} \\ \text{X} \end{array} + \dots \right\rangle_{0,1} \left\langle \begin{array}{c} \text{X} \\ \text{X} \end{array} + \dots \right\rangle_{0,1} \quad (5 \text{ in each})$$

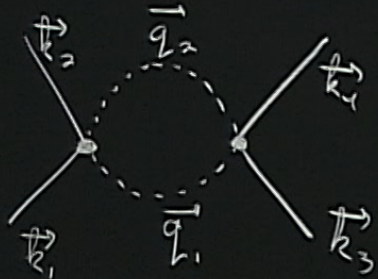
each vertex is $\frac{24}{4!} (2\pi)^d \delta(\vec{k}_1 + \vec{k}_2 + \vec{k}_3 + \vec{k}_4)$

and \diagdown slow fields & \cdots fast fields



each vertex $\rightarrow \frac{2i}{4!} (2\pi)^d \delta(\vec{k}_1 + \vec{k}_2 + \vec{k}_3 + \vec{k}_4)$

and \swarrow slow fields & \searrow fast fields



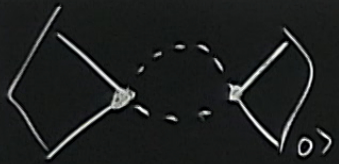
combinatorial factor of 36

δ -functions
say

$$\left. \begin{aligned} \vec{k}_1 + \vec{k}_2 + \vec{q}_1 + \vec{q}_2 &= 0 \\ \vec{k}_3 + \vec{k}_4 - \vec{q}_1 - \vec{q}_2 &= 0 \end{aligned} \right\}$$

$$\vec{k}_1 + \vec{k}_2 + \vec{k}_3 + \vec{k}_4 = 0$$

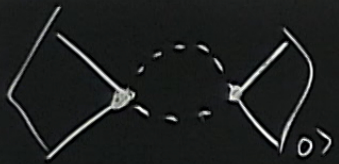
let: $\vec{q}_1 = \vec{q}$
 $\vec{q}_2 = -\vec{q} - \vec{k}_1 - \vec{k}_2$



$$\begin{aligned}
 &= \left(\frac{24}{4}\right)^2 \int_{k_1, k_4} (2\pi)^4 \delta(k_1 + k_2 + k_3 - k_4) \cdot \\
 &\quad \cdot \varphi_L(k_1) \varphi_L(k_2) \varphi_L(k_3) \varphi_L(k_4) \\
 &\quad \cdot \int_{\Lambda/b}^{\Lambda} \frac{d^d q}{(2\pi)^d} G_0(\vec{q}) G_0(-k_1 - k_2 - q)
 \end{aligned}$$

$\underbrace{\hspace{15em}}_{I_2}$

Note:



$$= \left(\frac{2i}{4}\right)^2 \int_{k_1, k_4} (2\pi)^d \delta(k_1 + k_2 + k_3 - k_4) \cdot$$

$$\cdot \varphi_L(k_1) \varphi_L(k_2) \varphi_L(k_3) \varphi_L(k_4)$$

$$\cdot \int_{\Lambda/b}^{\Lambda} \frac{d^d q}{(2\pi)^d} G_0(\vec{q}) G_0(-k_1 - k_2 - q)$$

I_2

Note:

$$I_2 = \int_{\Lambda/b}^{\Lambda} \frac{d^d q}{(2\pi)^d} \frac{1}{(r+q^2)} \frac{1}{[r+(-k_1-k_2-q)^2]}$$

$k_1 \dots k_n$

$$\cdot \varphi_L(k_1) \varphi_L(k_2) \varphi_L(k_3) \varphi_L(k_4)$$

$$\cdot \int_{\Lambda/b}^{\Lambda} \frac{d^d q}{(2\pi)^d} G_0(\vec{q}) G_0(-k_1 - k_2 - q)$$

I_2

Note:

$$I_2 = \int_{\Lambda/b}^{\Lambda} \frac{d^d q}{(2\pi)^d} \frac{1}{(r+q^2)} \frac{1}{[r+(-k_1-k_2-q)^2]}$$

Note:

$$I_2 = \int_{\Lambda_{1/b}}^{\Lambda} \frac{d^d q}{(2\pi)^d} \frac{1}{(r+q^2)} \frac{1}{[r + (-\vec{k}_1 - \vec{k}_2 - \vec{q})^2]}$$

$$I_2 = \int_{\Lambda_{1/b}}^{\Lambda} \frac{d^d q}{(2\pi)^d} \frac{1}{(r+q^2)^2} \left[1 - \frac{2\vec{q} \cdot (-\vec{k}_1 - \vec{k}_2) + (-\vec{k}_1 - \vec{k}_2)^2}{(r+q^2)} + \dots \right]$$

$$I_2 = \int_{\Lambda/b}^{\Lambda} \frac{d^d q}{(2\pi)^d} \frac{1}{(r+q^2)^2} \left[1 - \frac{2\vec{q} \cdot (-\vec{k}_1 - \vec{k}_2) + (-\vec{k}_1 - \vec{k}_2)^2}{(r+q^2)} + \dots \right]$$

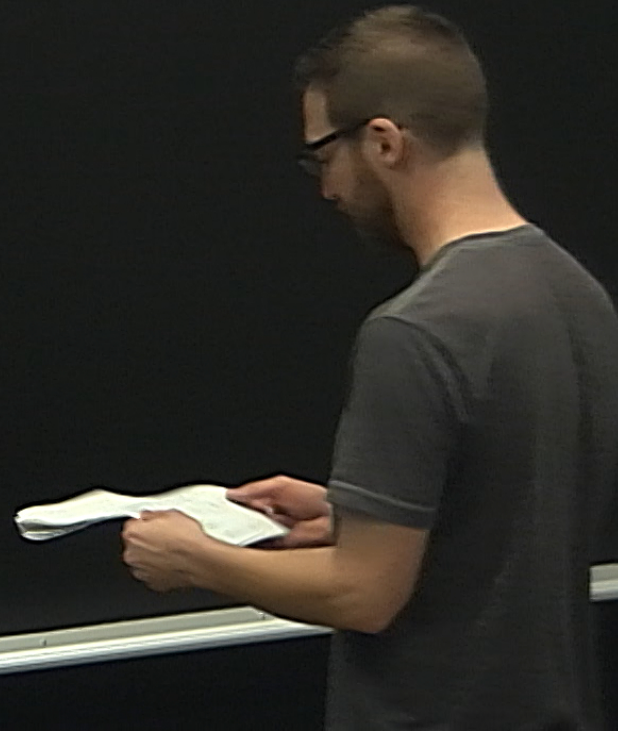
after F.T. back into real space
these give you additional terms
past ψ^4 : $\psi^2 (\nabla\psi)^2$, $\psi^2 \nabla^2 \psi^2$
etc.

so call

$$I_2 = \int_{\Lambda/b}^{\Lambda} \frac{d^d q}{(2\pi)^d} \frac{1}{(r+q^2)^2}$$

and ignore for now the
fact that parameter space
isn't closed at ψ^4

$$I_2 = \int_{\frac{1}{b}}^1 \frac{d^2 q}{(2\pi)^d} \frac{1}{(r+q^2)^2} = \frac{S_d}{(2\pi)^d} \int_{\frac{1}{b}}^1 dq \frac{q^{d-1}}{(r+q^2)^2}$$



$$\begin{aligned}
 I_2 &= \int_{\frac{1}{b}}^1 \frac{d^2 q}{(2\pi)^d} \frac{1}{(r+q^2)^2} = \frac{S_d}{(2\pi)^d} \int_{\frac{1}{b}}^1 dq \frac{q^{d-1}}{(r+q^2)^2} \\
 &= \frac{S_d}{(2\pi)^d} \int_{\frac{1}{b}}^1 dq \frac{q^{d-1}}{q^4} \left[1 - \frac{2r}{q^2} + \dots \right]
 \end{aligned}$$

$$\begin{aligned}
&= \frac{S_d}{(2\pi)^d} \int_{\frac{r}{b}}^1 dq \frac{r}{q^4} \left[1 - \frac{r}{q^2} + \dots \right] \\
&= \frac{S_d}{(2\pi)^d} \int_{\frac{r}{b}}^1 dq \left[q^{d-5} - 2r q^{d-7} + \dots \right] \\
&= \frac{S_d}{(2\pi)^d} \left[\frac{1}{d-4} \left(1 - \left(\frac{1}{b}\right)^{d-4} \right) - 2r \frac{1}{d-6} \left(1 - \left(\frac{1}{b}\right)^{d-6} \right) + \dots \right]
\end{aligned}$$

use again $1 - \left(\frac{1}{b}\right)^n = 1 - e^{-n \ln b} \cong n \ln b$

$$I_2 = \frac{S_d}{(2\pi)^d} \left[\Lambda^{d-4} - 2r \Lambda^{d-6} + \dots \right] \Delta l$$

$$= \frac{S_d}{(2\pi)^d} \frac{\Lambda^d \Delta l}{(r + \Lambda^2)^2}$$

CAUTION

use again $1 - \left(\frac{1}{b}\right)^n = 1 - e^{-nsl} \cong nsl$

$$I_2 = \frac{S_d}{(2\pi)^d} \left[\Lambda^{d-4} - 2r\Lambda^{d-6} + \dots \right] \Delta l$$

$$= \frac{S_d}{(2\pi)^d} \frac{\Lambda^d \Delta l}{(r + \Lambda^2)^2}$$

$$\langle \dots \rangle = \left(\frac{y}{4}\right)^2 I_2 \int_{h_1, \dots, h_4} (2\pi)^d \delta(\vec{k}_1 + \dots + \vec{k}_4) \Psi_{\zeta}(\vec{k}_1) \Psi_{\zeta}(\vec{k}_2) \Psi_{\zeta}(\vec{k}_3) \Psi_{\zeta}(\vec{k}_4)$$

collect all terms up to second order in y to get S'

Rescale: $\vec{k}' = \vec{k} b$ and $\varphi'(\vec{k}_i) = b^{-\frac{d+2}{2}} \varphi_{<}(\frac{\vec{k}}{b})$

$$\left\{ \begin{array}{l} r' = (r + 3u I_1) b^3 \\ z' = (z - 9u^2 I_2) b^{4-d} \end{array} \right.$$

$$\left. \begin{array}{l} \vec{k}_1 + \vec{k}_2 + \vec{q}_1 - \vec{q}_2 = 0 \\ \vec{k}_3 + \vec{k}_4 - \vec{q}_1 - \vec{q}_2 = 0 \end{array} \right\}$$

$$\vec{k}_1 + \vec{k}_2 + \vec{k}_3 + \vec{k}_4 = 0$$

let: $\vec{q}_1 = \vec{q}$
 $\vec{q}_2 = -\vec{q} - \vec{k}_1 - \vec{k}_2$

In differential form

$$\begin{cases} \frac{dr}{dl} = 2r + 3u \frac{S_d}{(2\pi)^d} \frac{\Lambda^d}{(r+\Lambda^2)} \\ \frac{du}{dl} = (4-d)u - 9u^2 \frac{S_d}{(2\pi)^d} \frac{\Lambda^d}{(r+\Lambda^2)^2} \end{cases}$$

In differential form

$$\begin{cases} \frac{dr}{dl} = 2r + 3u \frac{S_d}{(2\pi)^d} \frac{\Lambda^d}{(r+\Lambda^2)} \\ \frac{du}{dl} = (4-d)u - 9u^2 \frac{S_d}{(2\pi)^d} \frac{\Lambda^d}{(r+\Lambda^2)^2} \end{cases}$$

solving $\frac{dr}{dl} = \frac{du}{dl} = 0$

$$u^* \approx \frac{(4-d)}{9} \frac{(2\pi)^d}{S_d} \Lambda^{4-d} \quad \text{near } d=4$$

In differential form

$$\begin{cases} \frac{dr}{dl} = 2r + 3u \frac{S_d}{(2\pi)^d} \frac{\Lambda^d}{(r+\Lambda^2)} \\ \frac{du}{dl} = (4-d)u - 9u^2 \frac{S_d}{(2\pi)^d} \frac{\Lambda^d}{(r+\Lambda^2)^2} \end{cases}$$

Solving $\frac{dr}{dl} = \frac{du}{dl} = 0$

$$u^* \approx \frac{(4-d)}{9} \frac{(2\pi)^d}{S_d} \Lambda^{4-d} \approx \frac{\epsilon}{9} \frac{(2\pi)^4}{S_4} \quad \begin{matrix} \text{near} \\ d=4 \end{matrix}$$

In differential form

$$\begin{cases} \frac{dr}{dl} = 2r + 3u \frac{S_d}{(2\pi)^d} \frac{\Lambda^d}{(r+\Lambda^2)} \\ \frac{du}{dl} = (4-d)u - 9u^2 \frac{S_d}{(2\pi)^d} \frac{\Lambda^d}{(r+\Lambda^2)^2} \end{cases}$$

Solving $\frac{dr}{dl} = \frac{du}{dl} = 0$

$$u^* \approx \frac{(4-d)}{9} \frac{(2\pi)^d}{S_d} \Lambda^{4-d} \approx \frac{\epsilon}{9} \frac{(2\pi)^4}{S_4} \quad \text{near } d=4$$

also $S_4 = \frac{2\pi^{4/2}}{\Gamma(4/2)} = \frac{2\pi^2}{\Gamma(2)} = 2\pi^2$

$$u^* = \frac{\epsilon}{9} \frac{16\pi^4}{2\pi^2} = \frac{8}{9} \epsilon \pi^2$$

Note:

$$r^* = -\frac{\epsilon}{6} \Lambda^2$$

after F.T. back into real space
these give you additional terms
past φ^4 :

$$\varphi^2 (\nabla \varphi)^2, \quad \varphi^2 \nabla^2 \varphi^2$$

etc.

so call

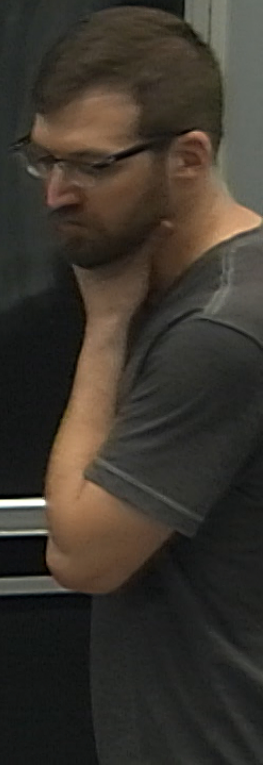
$$I_2 = \int_{1/b}^1 \frac{d^d q}{(2\pi)^d} \frac{1}{(r+q^2)^2}$$

and ignore for now the
fact that parameter space
isn't closed at φ^4

Define: $K_d \equiv \frac{S_d \Lambda^d}{(2\pi)^d}$ $K_4 \approx \frac{S_4}{(2\pi)^4} \Lambda^4 = \frac{\Lambda^4}{8\pi^2}$

$$\frac{dn}{dl} = 2n + 3u \frac{K_d}{r + \Lambda^2}$$

$$\frac{du}{dl} = (4-d)u - 9u^2 K_d \frac{1}{(r + \Lambda^2)^2}$$



CAUTION

Define: $K_d \equiv \frac{3_d \Lambda^d}{(2\pi)^d}$ $K_4 \equiv \frac{3_4}{(2\pi)^4} \Lambda^4 = \frac{1}{8\pi^2}$

$$\begin{cases} \frac{dr}{dl} = 2r + 3u \frac{K_d}{r + \Lambda^2} \\ \frac{du}{dl} = (4-d)u - 9u^2 K_d \frac{1}{(r + \Lambda^2)^2} \end{cases}$$

$$SK' = b' SK \quad \text{using } b' = e^{y \Delta l}$$

$$SK' = e^{y \Delta l} SK \approx (1 + y \Delta l) SK$$

$$\frac{SK' - SK}{\Delta l} = y SK \Rightarrow$$

$$\boxed{\frac{d}{dl} SK = y SK}$$

So γ is the eigenvalue of a "M" matrix

$$\frac{d}{dl} \begin{pmatrix} s_r \\ s_u \end{pmatrix} = M \begin{pmatrix} s_r \\ s_u \end{pmatrix}$$

where: define $\frac{dr}{dl} = \beta_r$ and $\frac{du}{dl} = \beta_u$

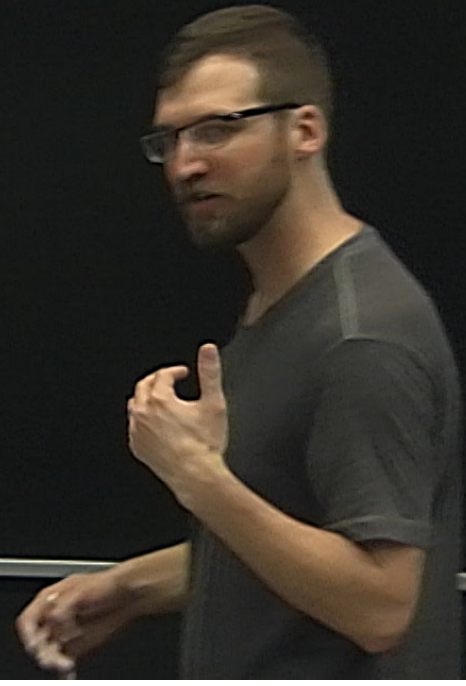
and $M = \begin{pmatrix} \frac{\partial \beta_r}{\partial r} & \frac{\partial \beta_r}{\partial u} \\ \frac{\partial \beta_u}{\partial r} & \frac{\partial \beta_u}{\partial u} \end{pmatrix}_{r^*, u^*}$

$\alpha r | x$

$$M = \begin{pmatrix} 2 - \frac{\omega}{3} & \frac{3}{8\pi^2} \lambda^2 \\ 0 & -\varepsilon \end{pmatrix}$$

$$y_t = y_r = 2 - \frac{1}{3} \varepsilon$$

$$y_u = -\varepsilon$$



CAUTION

$$M = \begin{pmatrix} 2 - \frac{\epsilon}{3} & \frac{3}{8\pi^2} \Lambda^2 \\ 0 & -\epsilon \end{pmatrix}$$

$$y_t = y_r = 2 - \frac{1}{3} \epsilon$$

$$y_u = -\epsilon$$

check: $\vec{e}^t = (1, 0)$, $\vec{e}^u = \left(-\frac{3\Lambda^2}{16\pi^2}, 1 \right)$

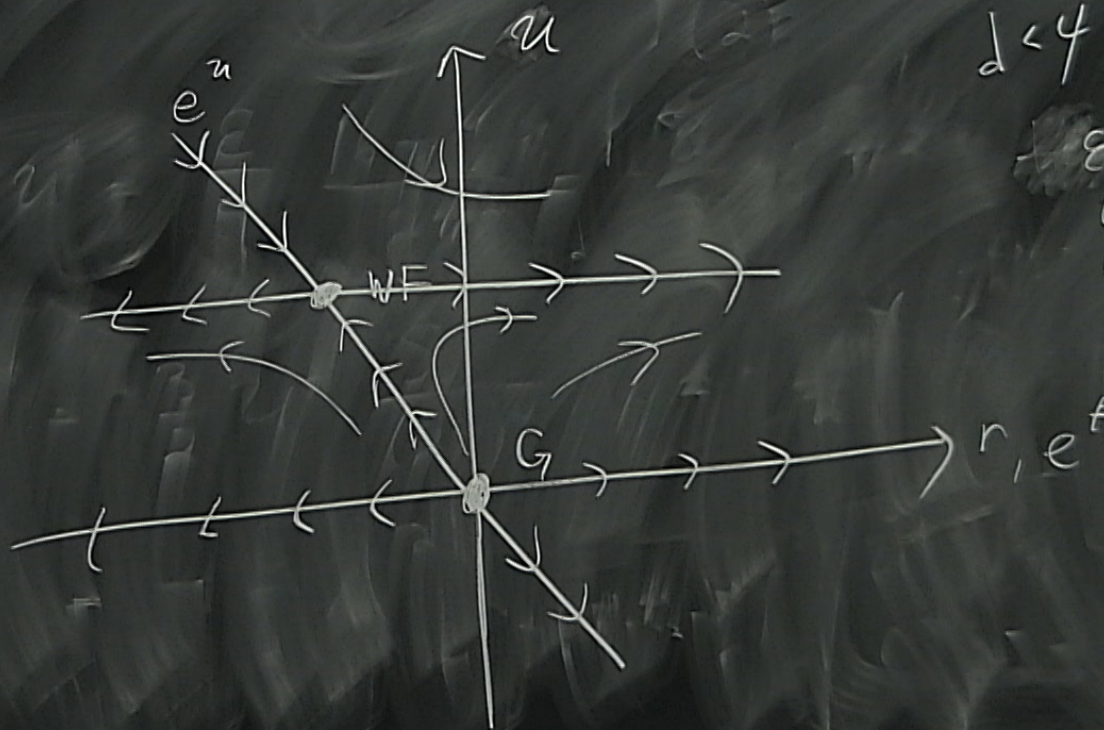
these should be the same eigenvectors as the 1st cumulant case.

Let's look at $d < 4$: Gaussian F.P. at $u^* = r^* = 0$

W.F. at $(r^*, u^*) = \left(-\frac{1}{6} \epsilon \Lambda^2, \frac{8}{9} \epsilon \pi^2 \right)$

Define: $K = S_d \Lambda^d$ $\quad \quad \quad S_4 \quad \Lambda^4 \quad \Lambda$

$$(r^*, z^*) = \frac{8}{9} \epsilon \pi^2 \left(-\frac{3}{16} \frac{\Lambda^2}{\pi^2}, 1 \right) = \frac{8}{9} \epsilon \pi^2 e^{\vec{u}}$$



$d < 4$

$$\epsilon = 4 - d > 0$$

y_t : relevant

y_z : irrelevant

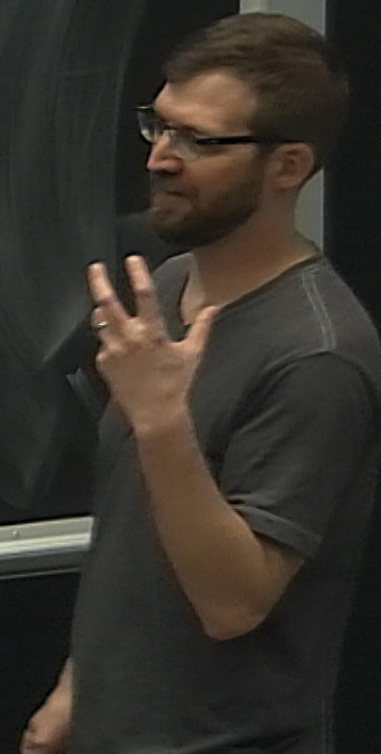
$$= 2 - \frac{d}{2} - \frac{\epsilon}{6} \frac{d}{2} = 2 - \frac{d}{2} + \frac{4}{2} - 2 \frac{d}{2} \frac{\epsilon}{6}$$

$$= \frac{1}{2} \epsilon - \frac{d}{12} \epsilon$$

set $d=4$

$$= \frac{\epsilon}{6} - \frac{\epsilon}{6} =$$

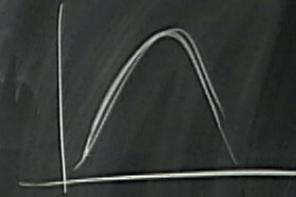
$$\boxed{\frac{\epsilon}{6} + \mathcal{O}(\epsilon^2)}$$



$$= 2 - \frac{d}{12} - \frac{\epsilon}{6} \frac{d}{2} = 2 - \frac{d}{2} + \frac{4}{2} - 2 \frac{d}{2} \frac{\epsilon}{6}$$

$$= \frac{1}{2} \epsilon - \frac{d}{12} \epsilon \quad \text{set } \underline{d=4}$$

$$= \frac{\epsilon}{12} - \frac{\epsilon}{3} = \boxed{\frac{\epsilon}{6} + \mathcal{O}(\epsilon^2)}$$



For many other exponents you need a magnetic field

$$S[\psi] = \int d^d x \left[\frac{1}{2} (\nabla \psi)^2 + \frac{r}{2} \psi^2 + \frac{u}{4} \psi^4 - B \psi \right]$$

$$h' = hb^{1 + \frac{d}{2}} \quad (\text{exercise})$$

$$y_h = 1 + \frac{d}{2} = 1 + \frac{d}{2} - \frac{4}{2} + \frac{4}{2} = 3 - \frac{1}{2} \epsilon$$

$$\gamma = \frac{2y_h - d}{y_t} = \frac{6 + \frac{d}{2} - d}{2 - \frac{\epsilon}{3}} = \left(6 - \frac{d}{2}\right) \frac{1}{2} \left(1 + \frac{\epsilon}{6}\right) = 1 + \frac{\epsilon}{6}$$

$$\beta = \frac{d - y_h}{y_t} = \frac{1}{2} - \frac{\epsilon}{6} \quad (\text{check})$$

$$\delta = \frac{y_h}{d - y_h} = 3 + \epsilon$$

CAUTION

$$h' = hb^{1 + \frac{d}{2}} \quad (\text{exercise})$$

$$y_h = 1 + \frac{d}{2} = 1 + \frac{d}{2} - \frac{4}{2} + \frac{4}{2} = 3 - \frac{1}{2} \epsilon$$

$$\gamma = \frac{2y_h - d}{y_t} = \frac{6 + \frac{d}{2} - d}{2 - \frac{\epsilon}{3}} = \left(6 - \frac{d}{2}\right) \frac{1}{2} \left(1 + \frac{\epsilon}{6}\right) = 1 + \frac{\epsilon}{6}$$

$$\beta = \frac{d - y_h}{y_t} = \frac{1}{2} - \frac{\epsilon}{6} \quad (\text{check})$$

$$\delta = \frac{y_h}{d - y_h} = 3 + \epsilon$$

CAUTION