Title: Quantum error-correction and black holes

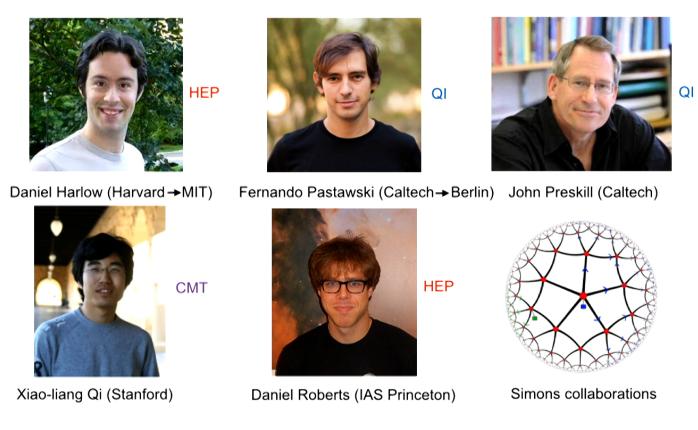
Date: Feb 22, 2017 04:00 PM

URL: http://pirsa.org/17020102

Abstract: It is commonly believed that quantum information is not lost in a black hole. Instead, it is encoded into non-local degrees of freedom in some clever way; like a quantum error-correcting code. In this talk, I will discuss recent attempts to resolve some paradoxes in quantum gravity by using the theory of quantum error-correction. First, I will introduce a simple toy model of the AdS/CFT correspondence based on tensor networks and demonstrate that the correspondence between the AdS gravity and CFT is indeed a realization of quantum codes. I will then show that the butterfly effect in black holes can be interpreted as non-local encoding of quantum information and can be quantitatively measured by out-of-time ordered correlations. Finally, I will discuss how out-of-time ordered correlations, measured outside the black hole horizon, may probe smoothness of the geometry across the horizon.

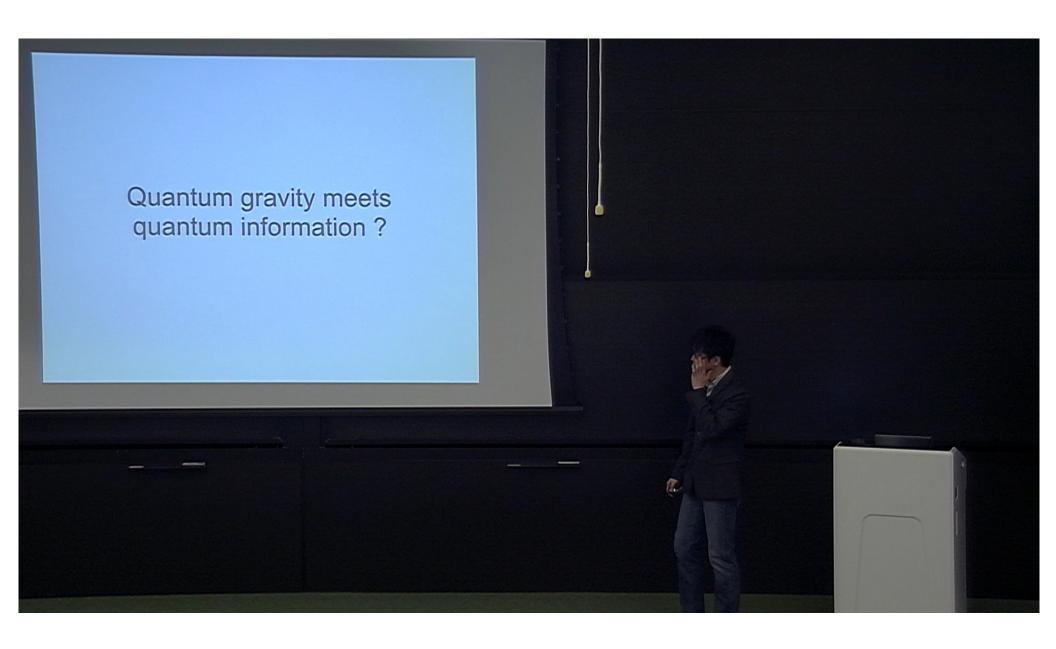
Pirsa: 17020102 Page 1/97

Collaborators



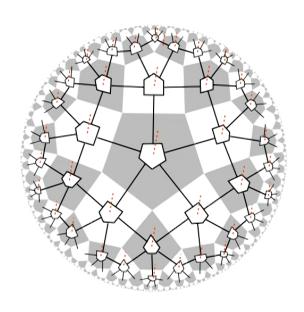
Thanks to Patrick Hayden, Pavan Hosur, Alexei Kitaev, Michael Walter and many others...

Pirsa: 17020102 Page 2/97



Pirsa: 17020102 Page 3/97

Quantum error-correction and black holes

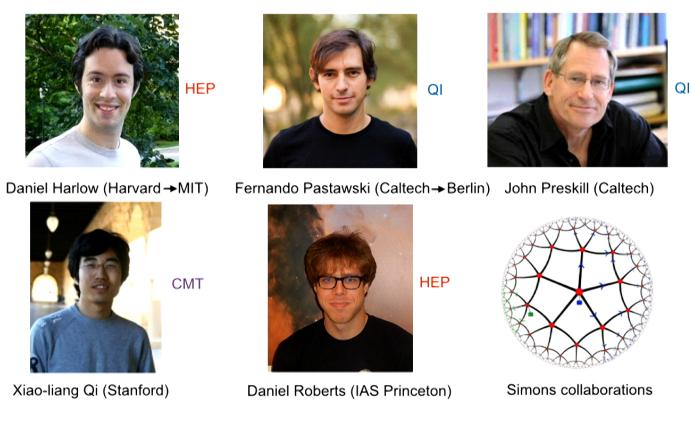


Beni Yoshida (Perimeter Institute)

@ Perimeter (Feb 2017)

Pirsa: 17020102 Page 4/97

Collaborators



Thanks to Patrick Hayden, Pavan Hosur, Alexei Kitaev, Michael Walter and many others...

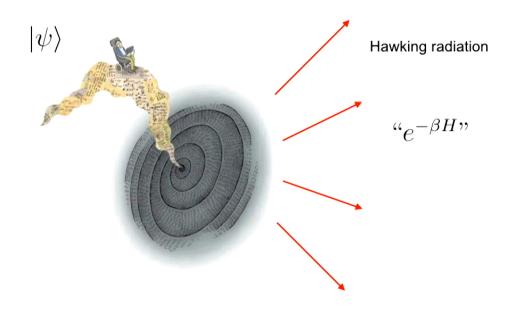
Pirsa: 17020102 Page 5/97

Information loss puzzle

- Quantum mechanics and general relativity are in serious conflicts!
 - (a) Quantum mechanics says that information is never lost.

$$|\psi(t)\rangle = e^{-iHt}|\psi(0)\rangle$$

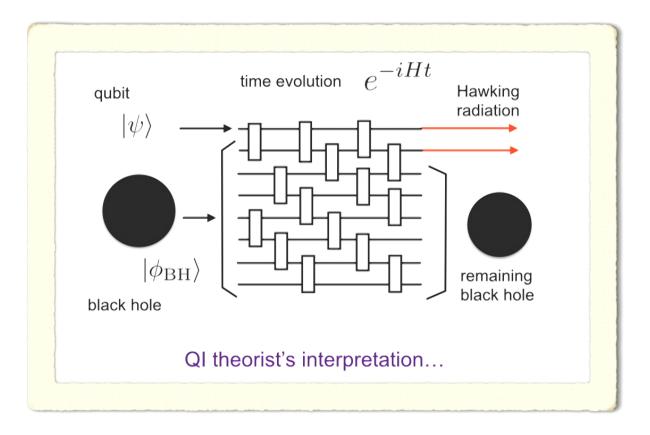
(b) General relativity says information is lost in black holes.



Pirsa: 17020102 Page 6/97

Why quantum information?

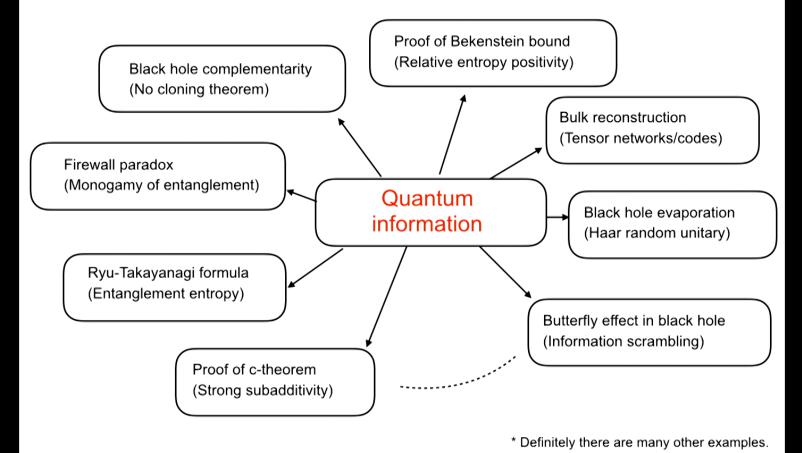
• Information loss puzzle = "Information problem in a quantum system"



Pirsa: 17020102 Page 7/97

Quantum information and HEP

• Many breakthrough ideas in quantum gravity come from quantum information theory



Pirsa: 17020102 Page 8/97

This talk:

A black hole is a quantum errorcorrecting code

Pirsa: 17020102

Part 1:

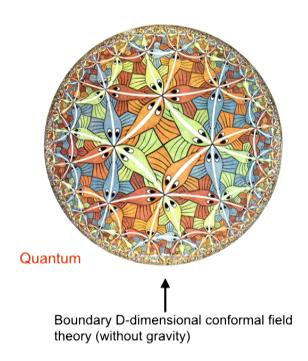
Simple toy model of the AdS/CFT correspondence

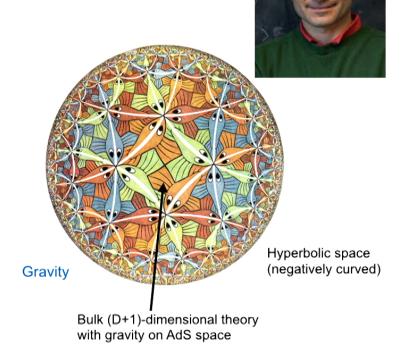
(joint with Harlow, Pastawski and Preskill 2015)

Pirsa: 17020102 Page 10/97

Anti-de Sitter/Conformal field theory (AdS/CFT) correspondence

• [Conjecture] Equivalence between string (gravity) theory in bulk and (certain types of) CFT on boundary (Maldacena 1997)



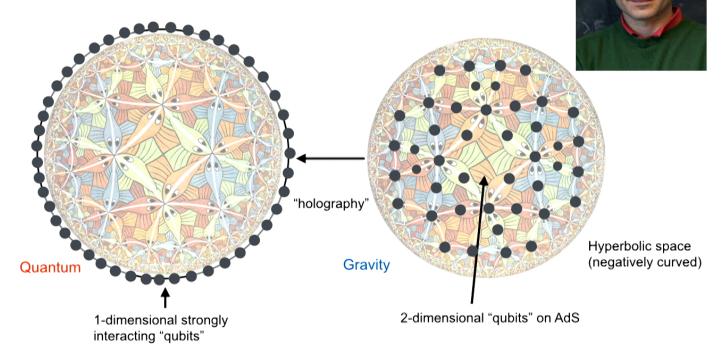


Pirsa: 17020102 Page 11/97

Anti-de Sitter/Conformal field theory (AdS/CFT) correspondence

• [Conjecture] Equivalence between string (gravity) theory in bulk and (certain types of) CFT on boundary (Maldacena 1997)

• [Holography] Bulk degrees of freedom are encoded in boundary, like a hologram.



* Finite-dimensional Hilbert space cartoon picture.

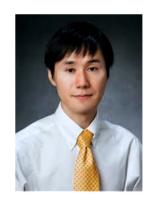
Pirsa: 17020102 Page 12/97

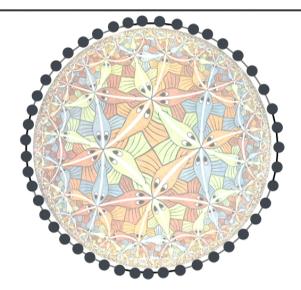
• [Ryu-Takayanagi formula 06]

Quantum $S(A) = \frac{1}{4G_N} \min_{\gamma_A} (\operatorname{area}(\gamma_A))$ Gravity

Minimize over spatial bulk surfaces homologous to A







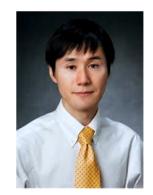
* Leading order in 1/G. * Casini-Huerta-Myers11, Lewkowycz-Maldacena13

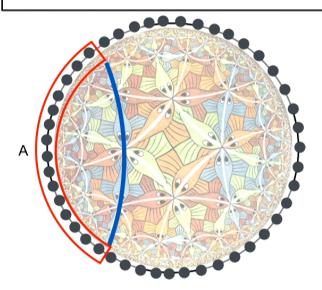
• [Ryu-Takayanagi formula 06]

Quantum $S(A) = \frac{1}{4G_N} \min_{\gamma_A} (\operatorname{area}(\gamma_A))$ Gravity

Minimize over spatial bulk surfaces homologous to A

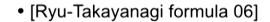






* Leading order in 1/G. * Casini-Huerta-Myers11, Lewkowycz-Maldacena13

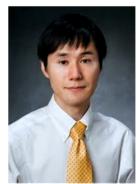
Gravity

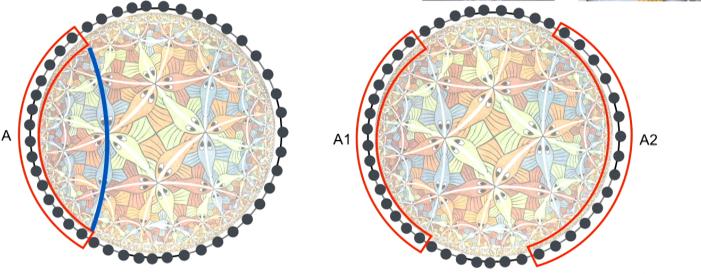


Quantum $S(A) = \frac{1}{4G_N} \min_{\gamma_A} (\operatorname{area}(\gamma_A))$

Minimize over spatial bulk surfaces homologous to A



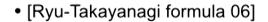




* Leading order in 1/G. * Casini-Huerta-Myers11, Lewkowycz-Maldacena13

Pirsa: 17020102 Page 15/97

Gravity

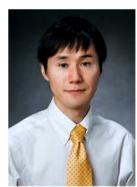


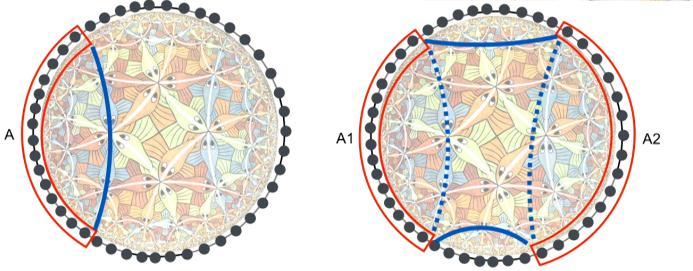
Quantum

 $S(A) = \frac{1}{4G_N} \min_{\gamma_A} (\operatorname{area}(\gamma_A))^{\bullet}$

Minimize over spatial bulk surfaces homologous to A





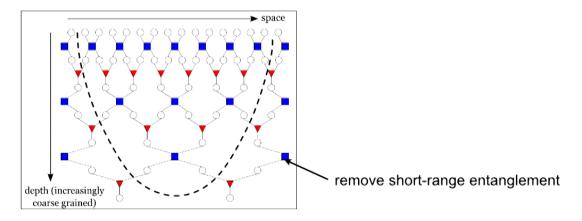


* Leading order in 1/G. * Casini-Huerta-Myers11, Lewkowycz-Maldacena13

Pirsa: 17020102 Page 16/97

MERA (Multiscale entanglement renormalization ansatz)

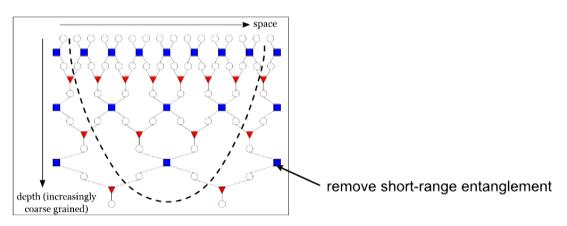
• Powerful numerical method to study strongly-correlated systems. (Vidal 07)



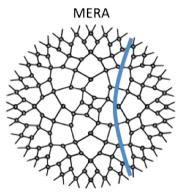
Pirsa: 17020102 Page 17/97

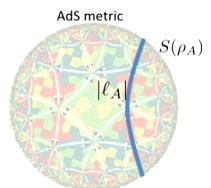
MERA (Multiscale entanglement renormalization ansatz)

• Powerful numerical method to study strongly-correlated systems. (Vidal 07)



 AdS/CFT correspondence can be explained by a tensor network ?
 (Swingle 09)

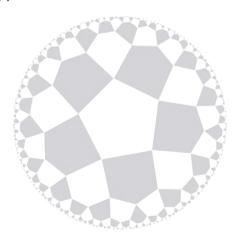




Pirsa: 17020102 Page 18/97

• [Entanglement wedge reconstruction]

A bulk operator ϕ can be represented by some integral of local boundary operators supported on A if ϕ is contained inside the entanglement wedge of A.

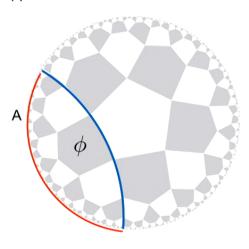


* HKLL for single intervals * Conjectured for multiple intervals by Aron Wall.

Pirsa: 17020102 Page 19/97

• [Entanglement wedge reconstruction]

A bulk operator ϕ can be represented by some integral of local boundary operators supported on A if ϕ is contained inside the entanglement wedge of A.

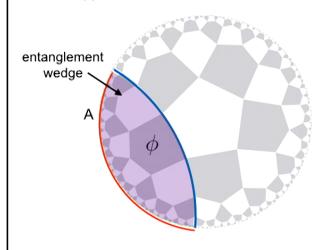


* HKLL for single intervals * Conjectured for multiple intervals by Aron Wall.

Pirsa: 17020102 Page 20/97

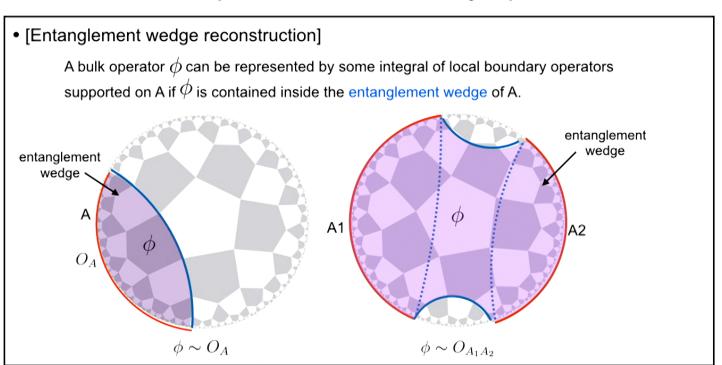
• [Entanglement wedge reconstruction]

A bulk operator ϕ can be represented by some integral of local boundary operators supported on A if ϕ is contained inside the entanglement wedge of A.



* HKLL for single intervals * Conjectured for multiple intervals by Aron Wall.

Pirsa: 17020102 Page 21/97



* HKLL for single intervals * Conjectured for multiple intervals by Aron Wall.

Pirsa: 17020102 Page 22/97

• [Entanglement wedge reconstruction] A bulk operator ϕ can be represented by some integral of local boundary operators supported on A if ϕ is contained inside the entanglement wedge of A. entanglement wedge $\phi \sim O_A \qquad \qquad \phi \sim O_{A_1A_2}$

Remarks for experts

- Entanglement wedge may go beyond black hole horizons (i.e. no firewall ?).
- "Proven" by using a "generalized" RT formula (Jefferis et al, Dong et al, Bao et al 2016)
- No explicit recipe is known for more than one intervals (AdS3) (For higher-dimensions, it's a bit more subtle).

* HKLL for single intervals * Conjectured for multiple intervals by Aron Wall.

Pirsa: 17020102 Page 23/97

Bulk locality paradox

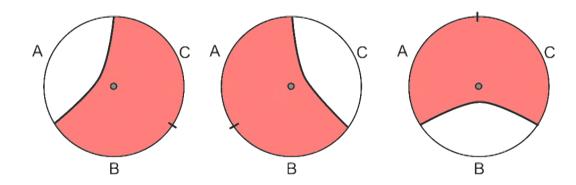
• The reconstruction leads to a <u>paradox</u>!

[Almheiri-Dong-Harlow 14]









* Uses Schur's lemma, assuming finite-dimensional factorizable Hilbert space

Pirsa: 17020102 Page 24/97

Bulk locality paradox

• The reconstruction leads to a <u>paradox</u>!

[Almheiri-Dong-Harlow 14]

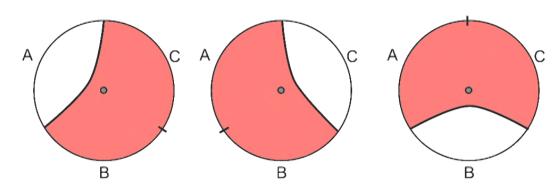






The bulk operator ϕ has non-trivial supports only on AB, BC, CA.

→ All the bulk operators must correspond to identity operators on the boundary?



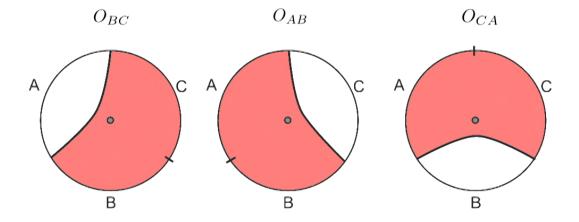
* Uses Schur's lemma, assuming finite-dimensional factorizable Hilbert space

Pirsa: 17020102 Page 25/97

Quantum error-correction in AdS/CFT?

• The AdS/CFT correspondence can be viewed as a quantum error-correcting code!

These operators may be different, but act in the same manner in a low energy subspace.



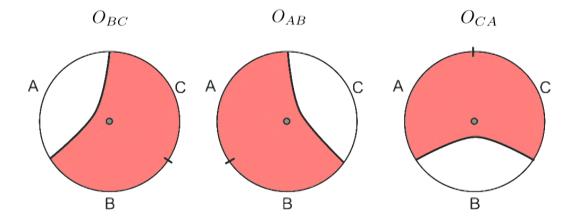
Pirsa: 17020102 Page 26/97

Quantum error-correction in AdS/CFT?

• The AdS/CFT correspondence can be viewed as a quantum error-correcting code!

These operators may be different, but act in the same manner in a low energy subspace.

Recall string operators in lattice gauge theory (or Z2 spin liquid).



Pirsa: 17020102 Page 27/97

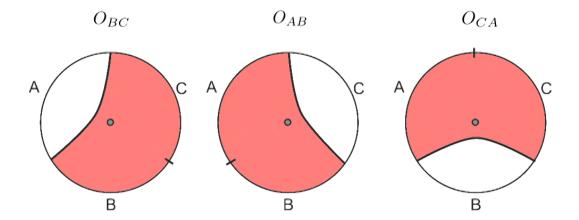
Quantum error-correction in AdS/CFT?

• The AdS/CFT correspondence can be viewed as a quantum error-correcting code!

These operators may be different, but act in the same manner in a low energy subspace.

Recall string operators in lattice gauge theory (or Z2 spin liquid).

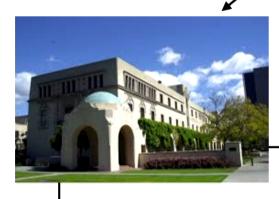
- Quantum secret-sharing code: Alice, Bob and Charlie share a quantum secret.
- Error-correction: Quantum information is protected against erasure of one party.



Pirsa: 17020102 Page 28/97



@ Caltech, 2014 November









Pirsa: 17020102 Page 29/97



@ Caltech, 2014 November



"Construct a toy model!"







Pirsa: 17020102 Page 30/97

A simple toy model

1 bulk qubit

in total, just 6 qubits

5 boundary qubits

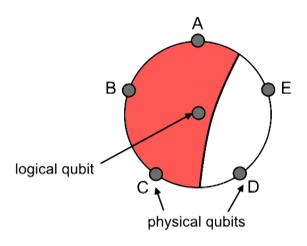
Pirsa: 17020102 Page 31/97

A simple toy model

1 bulk qubit

in total, just 6 qubits

5 boundary qubits



Pirsa: 17020102 Page 32/97

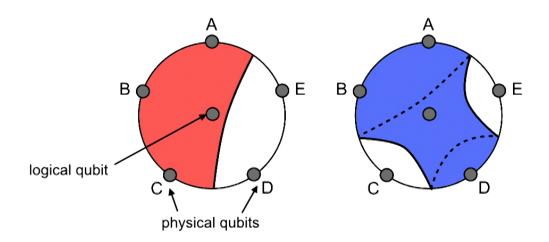
A simple toy model

1 bulk qubit

in total, just 6 qubits

5 boundary qubits

A bulk operator must have corresponding boundary operators on any region with three qubits.



Pirsa: 17020102 Page 33/97

Five-qubit quantum code

• The "simplest" quantum code which is made of qubits.

$$|\xi\rangle = \alpha|0\rangle + \beta|1\rangle \longrightarrow |\xi\rangle = \alpha|c_0\rangle + \beta|c_1\rangle$$

1-qubit input state

5-qubit output state

[DiVincenzo-Shor, Laflamme-Miquel-Paz-Zurek 1996]

Pirsa: 17020102 Page 34/97

Five-qubit quantum code

• The "simplest" quantum code which is made of qubits.

$$|\xi\rangle = \alpha|0\rangle + \beta|1\rangle \longrightarrow |\xi\rangle = \alpha|c_0\rangle + \beta|c_1\rangle$$

1-qubit input state

5-qubit output state

where

$$|c_{0}\rangle = |00000\rangle$$

$$+ |11000\rangle + |01100\rangle + |00110\rangle + |00011\rangle + |10001\rangle$$

$$- |10100\rangle - |01010\rangle - |00101\rangle - |10010\rangle - |01001\rangle$$

$$- |11110\rangle - |01111\rangle - |10111\rangle - |11011\rangle - |11101\rangle$$

$$|c_{1}\rangle = |11111\rangle$$

$$+ |00111\rangle + |10011\rangle + |11001\rangle + |11100\rangle + |01110\rangle$$

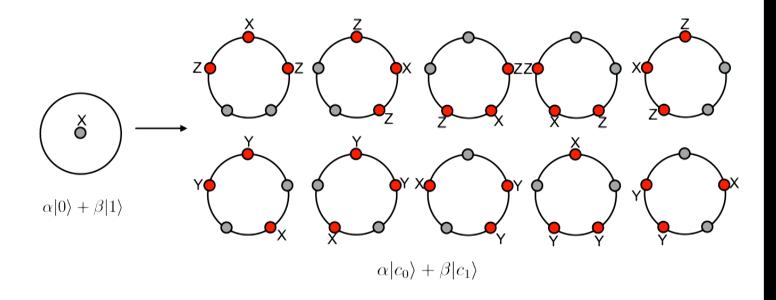
$$- |01011\rangle - |10101\rangle - |11010\rangle - |01101\rangle - |10110\rangle$$

$$- |00001\rangle - |10000\rangle - |01000\rangle - |00100\rangle - |00010\rangle$$

[DiVincenzo-Shor, Laflamme-Miquel-Paz-Zurek 1996]

Operator correspondence in five-qubit code

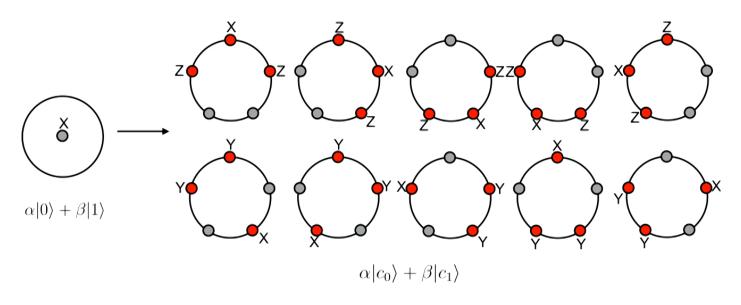
• Pauli X operator in the input (bulk) corresponds to following 3-body operators.



Pirsa: 17020102 Page 36/97

Operator correspondence in five-qubit code

• Pauli X operator in the input (bulk) corresponds to following 3-body operators.

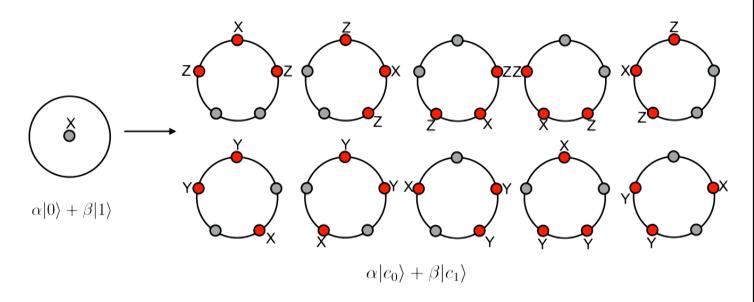


• So, the five-qubit code is a very small toy quantum gravity!

Pirsa: 17020102 Page 37/97

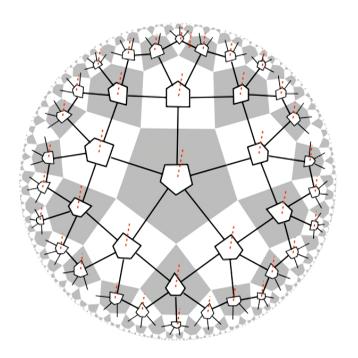
Operator correspondence in five-qubit code

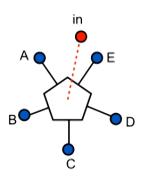
• Pauli X operator in the input (bulk) corresponds to following 3-body operators.



- So, the five-qubit code is a very small toy quantum gravity!
- Error-correction : losing 2 qubits is OK.

• A tiling of the five qubit code via tensor network technique



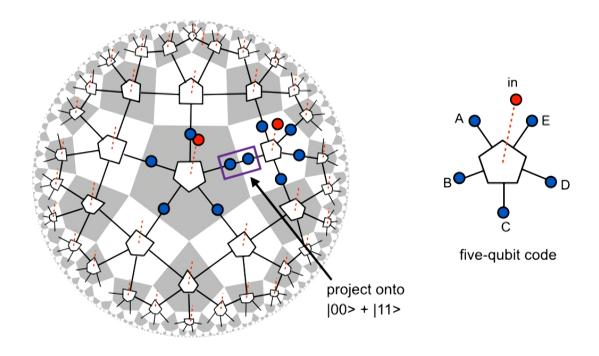


five-qubit code

* Tensor network and AdS/CFT: Vidal07, Swingle12, Qi13, Czech et al15 ...

Pirsa: 17020102 Page 39/97

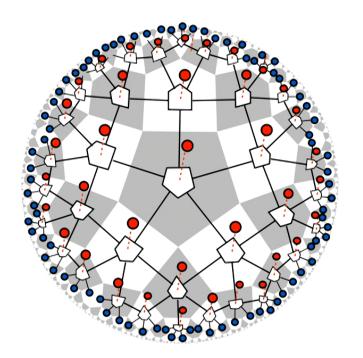
• A tiling of the five qubit code via tensor network technique

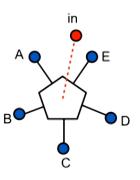


 * Tensor network and AdS/CFT : Vidal07, Swingle12, Qi13, Czech et al15 \dots

Pirsa: 17020102 Page 40/97

• A tiling of the five qubit code via tensor network technique



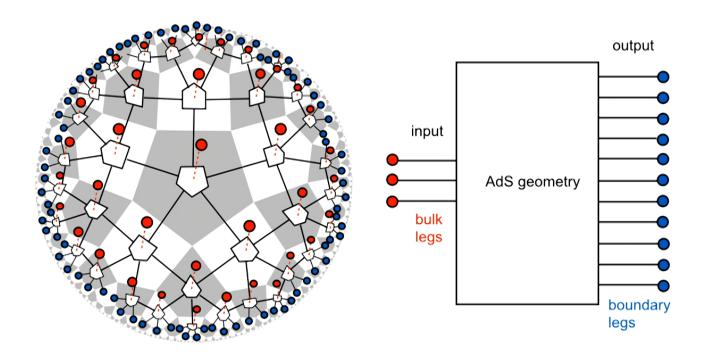


five-qubit code

* Tensor network and AdS/CFT: Vidal07, Swingle12, Qi13, Czech et al15 ...

Pirsa: 17020102 Page 41/97

• A tiling of the five qubit code via tensor network technique

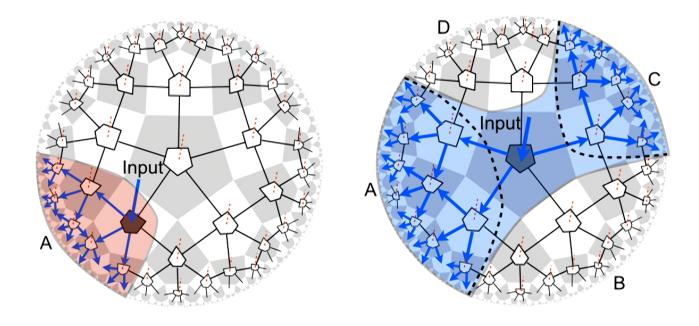


* Tensor network and AdS/CFT: Vidal07, Swingle12, Qi13, Czech et al15 ...

Pirsa: 17020102 Page 42/97

Entanglement wedge reconstruction

• 1 in & 3 out (operator pushing)

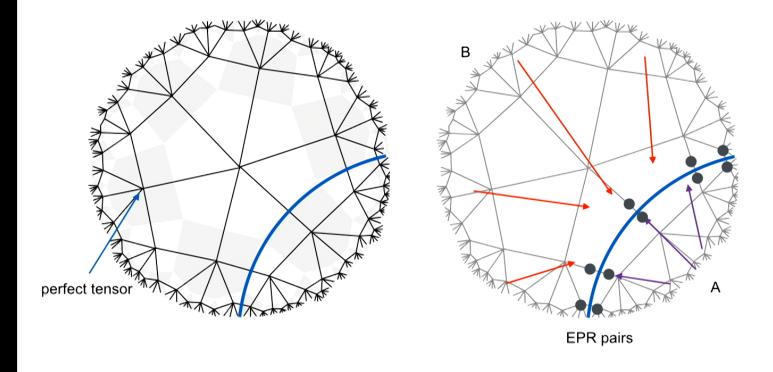


Pirsa: 17020102 Page 43/97

Holographic state

• The Ryu-Takayanagi formula holds (tiling without bulk legs)

Coarse-graining (RG transformation) = Distillation of EPR pairs along the geodesic

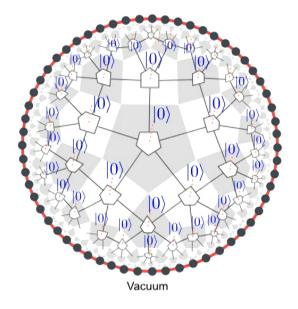


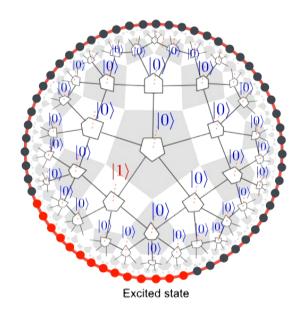
Pirsa: 17020102 Page 44/97

• Qubits on the bulk represent matter fields, coupled to the geometry (via 5-qubit code tensors).

Pirsa: 17020102 Page 45/97

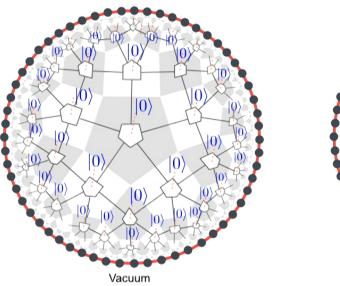
• Qubits on the bulk represent matter fields, coupled to the geometry (via 5-qubit code tensors).

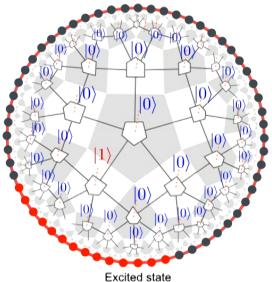




Pirsa: 17020102 Page 46/97

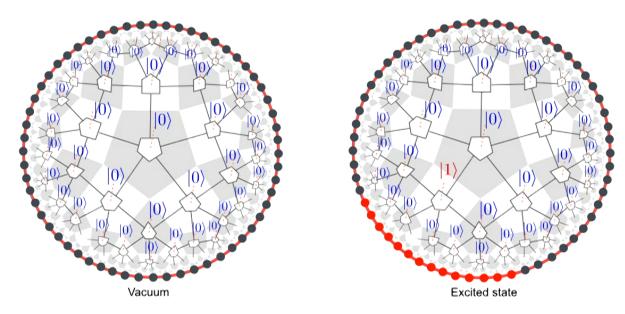
- Qubits on the bulk represent matter fields, coupled to the geometry (via 5-qubit code tensors).
- Bulk Hilbert space $\mathcal{H}_{\mathrm{bulk}}$ is much smaller than boundary Hilbert space $\mathcal{H}_{\mathrm{bdy}}$.
 - The model captures "perturbations" around a fixed geometry.



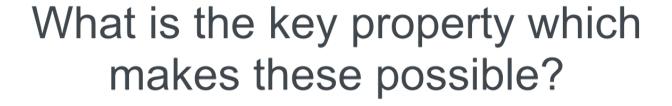


Pirsa: 17020102 Page 47/97

- Qubits on the bulk represent matter fields, coupled to the geometry (via 5-qubit code tensors).
- Bulk Hilbert space $\mathcal{H}_{\mathrm{bulk}}$ is much smaller than boundary Hilbert space $\mathcal{H}_{\mathrm{bdy}}$.
 - → The model captures "perturbations" around a fixed geometry.
- Going <u>outside the codeword space</u> = changing geometries (eg. <u>micro black holes</u>)



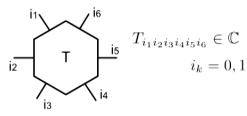
Pirsa: 17020102 Page 48/97



Pirsa: 17020102 Page 49/97

• A tensor with 6 legs

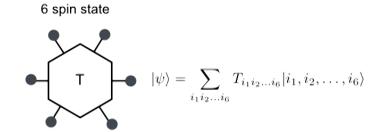




[Choi-Jamilkowski isomorphism]

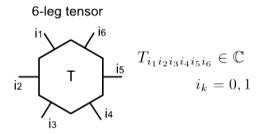
Pirsa: 17020102 Page 50/97

- A tensor with 6 legs
 - 6-leg tensor $\overrightarrow{\mathbf{I}}_{\mathbf{i}2} \xrightarrow{\mathbf{i}_{\mathbf{i}}} \overrightarrow{\mathbf{I}}_{i_{1}i_{2}i_{3}i_{4}i_{5}i_{6}} \in \mathbb{C}$ $\overrightarrow{\mathbf{I}}_{i_{1}i_{2}i_{3}i_{4}i_{5}i_{6}} \in \mathbb{C}$ $\overrightarrow{\mathbf{I}}_{k} = 0, 1$
- A state with 6 qubits

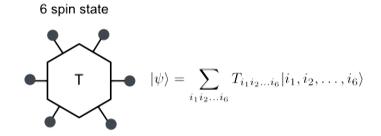


[Choi-Jamilkowski isomorphism]

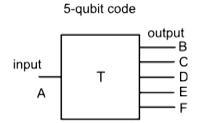
• A tensor with 6 legs



A state with 6 qubits



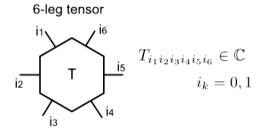
• A quantum code with 1 input & 5 output



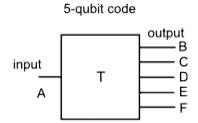
A and B are not entangled.

[Choi-Jamilkowski isomorphism]

• A tensor with 6 legs

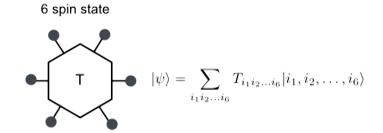


• A quantum code with 1 input & 5 output



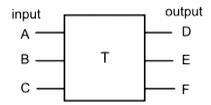
A and B are not entangled.

A state with 6 qubits



• A linear operator with 3 input & 3 output

3-qubit quantum operation



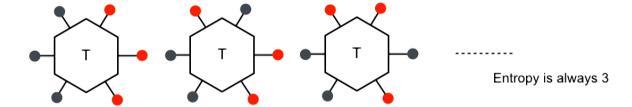
unitary if ABC is maximally entangled with DEF

[Choi-Jamilkowski isomorphism]

5-qubit code is "perfect"

• The 6-qubit state (or 5-qubit code) is "perfect".

[Def] The state is maximally entangled in any bipartition.

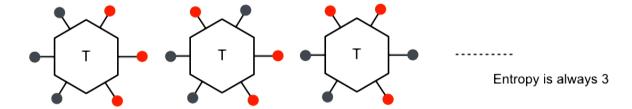


Pirsa: 17020102 Page 54/97

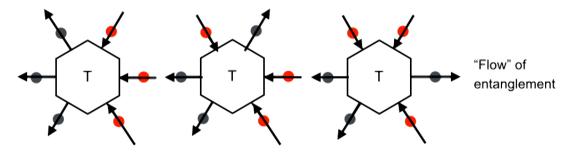
5-qubit code is "perfect"

• The 6-qubit state (or 5-qubit code) is "perfect".

[Def] The state is maximally entangled in any bipartition.



• Any "3-in to 3-out" defines unitary operator.



• Perfect tensor networks generate quantum codes with similar AdS/CFT properties.

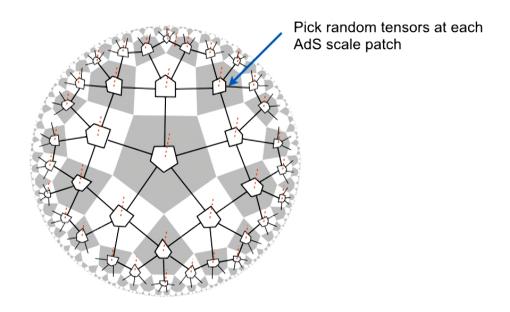
["Bit threading" Freedman-Headrick]

Pirsa: 17020102 Page 55/97

Space-time from "something random"?

• Perfect tensors are easy to create!

A random state is nearly maximally entangled along any bipartition (eg: Page's argument, Canonical typicality).



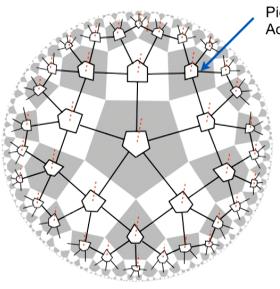
• Analytical solution of random tensor network [Hastings15, Hayden et al 16]

Pirsa: 17020102 Page 56/97

Space-time from "something random"?

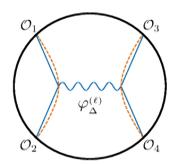
• Perfect tensors are easy to create!

A random state is nearly maximally entangled along any bipartition (eg: Page's argument, Canonical typicality).



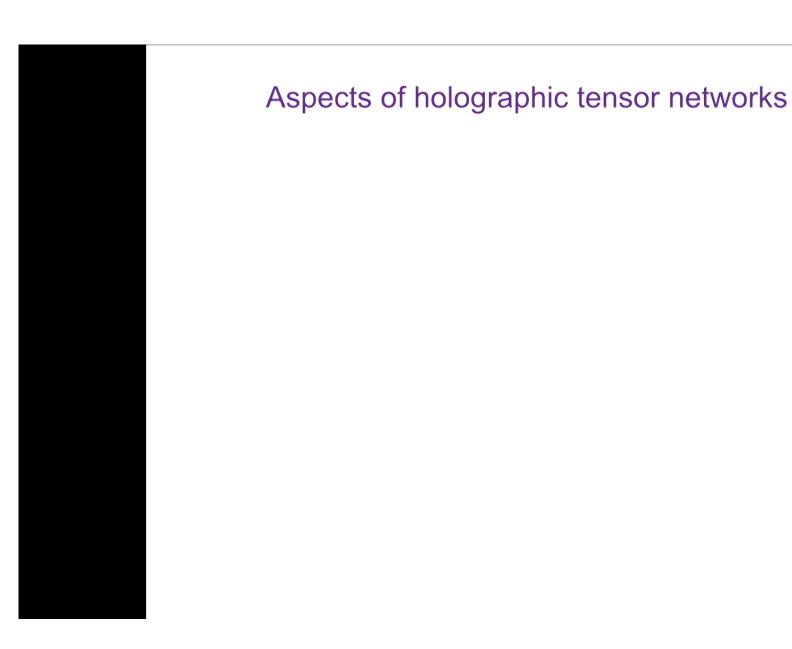
Pick random tensors at each AdS scale patch

Interesting "1/S" behaviours (S: spin dimensions)



• Analytical solution of random tensor network [Hastings15, Hayden et al 16]

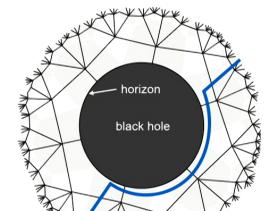
Pirsa: 17020102 Page 57/97



Pirsa: 17020102 Page 58/97

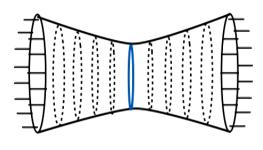
Aspects of holographic tensor networks

* Toy models of black holes (smoothness of the horizon? interior of black holes?)



One-sided black hole

Two-sided black hole

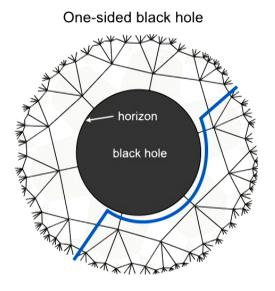


Einstein-Rosen bridge

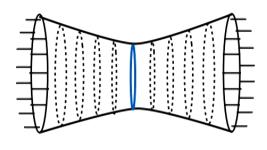
Pirsa: 17020102 Page 59/97

Aspects of holographic tensor networks

- * Toy models of black holes (smoothness of the horizon? interior of black holes?)
- * Complexity = Volume = Action? [Susskind and his friends] (number of tensors = complexity?)
- * Generalizations: symmetries, "matrix QM"-like network, sub-AdS locality, kinematic space
- * Tensor network in de Sitter space ? (Dark matter problem, as mentioned by Eric Verlinde)



Two-sided black hole

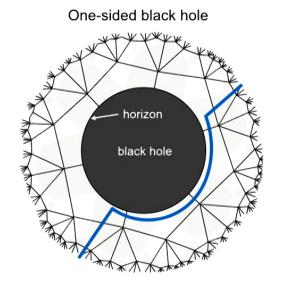


Einstein-Rosen bridge

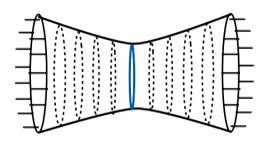
Pirsa: 17020102 Page 60/97

Aspects of holographic tensor networks

- * Toy models of black holes (smoothness of the horizon? interior of black holes?)
- * Complexity = Volume = Action? [Susskind and his friends] (number of tensors = complexity?)
- * Generalizations: symmetries, "matrix QM"-like network, sub-AdS locality, kinematic space
- * Tensor network in de Sitter space ? (Dark matter problem, as mentioned by Eric Verlinde)
- * Difficulty: No Hamiltonian, No dynamics. Need more inputs to be more realistic...



Two-sided black hole



Einstein-Rosen bridge

Pirsa: 17020102 Page 61/97

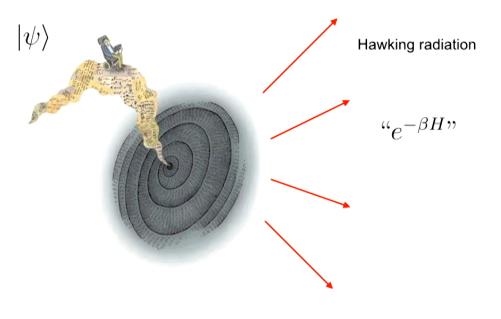
Part 2:

Black hole is a quantum error-correcting code

Pirsa: 17020102 Page 62/97

Information loss puzzle

• Is quantum information lost ?

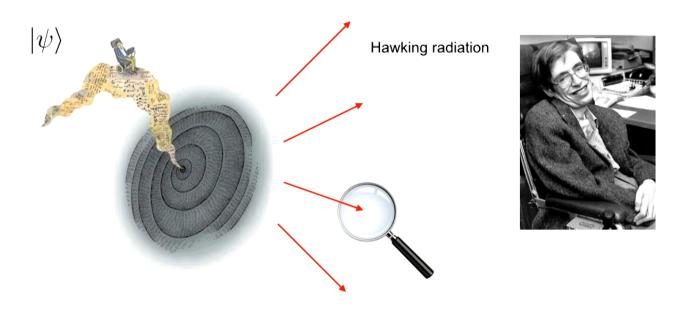




Pirsa: 17020102 Page 63/97

Information loss puzzle

• Or hidden into some non-local degrees of freedom ? (Scrambling!)

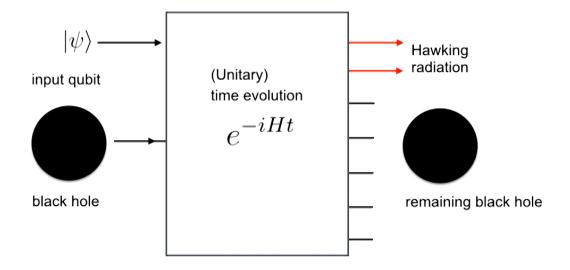


Locally it looks like " $e^{-\beta H}$ ", but globally it is not .

Pirsa: 17020102 Page 64/97

Black hole is a quantum error-correcting code?

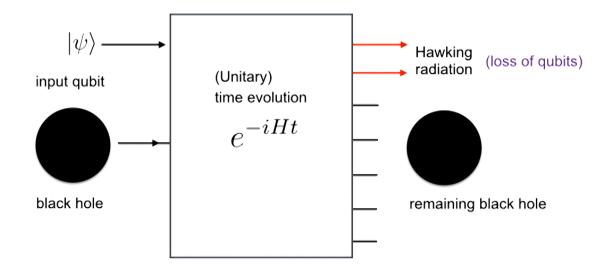
• Scrambling is very similar to how quantum error-correcting codes work.



Pirsa: 17020102 Page 65/97

Black hole is a quantum error-correcting code?

• Scrambling is very similar to how quantum error-correcting codes work.

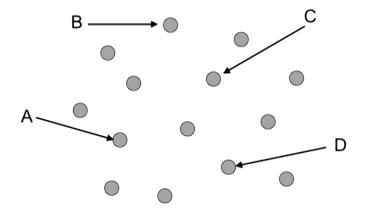


• Input quantum state is protected from loss of qubits (Hawking radiation)

Pirsa: 17020102 Page 66/97

Out-of-time ordered correlation functions

• Black holes have some "hidden" correlations (Kitaev14, Shenker-Stanford13)



• Previously considered by Larkin and Ovchinikov in 1960s

Pirsa: 17020102 Page 67/97

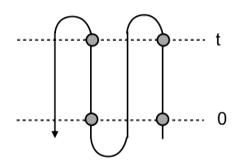
Out-of-time ordered correlation functions

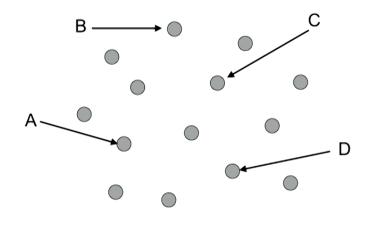
• Black holes have some "hidden" correlations (Kitaev14, Shenker-Stanford13)

$$OTOC = \langle O_A(0)O_B(t)O_C(0)O_D(t) \rangle$$

$$O_B(t) = e^{-iHt}O_B(0)e^{iHt}$$

$$O_D(t) = e^{-iHt}O_D(0)e^{iHt}$$





• Previously considered by Larkin and Ovchinikov in 1960s

Out-of-time ordered correlation functions detect scrambling of quantum information.

(joint with Hosur, Qi and Roberts 2015)

Pirsa: 17020102 Page 69/97

State-Tensor duality, again

• Unitary operator acting on n qubits can be viewed as a state on 2n qubits.

[Choi-Jamilkowski 1960s, Hayden-Preskill, Hartman-Maldacena]

Pirsa: 17020102 Page 70/97

State-Tensor duality, again

• Unitary operator acting on n qubits can be viewed as a state on 2n qubits.

$$U = \sum_{i,j} U_{ij} |i\rangle\langle j|$$
 out $|U\rangle = \sum_{i,j} U_{ij} |i\rangle \otimes |j\rangle$ $|U\rangle$

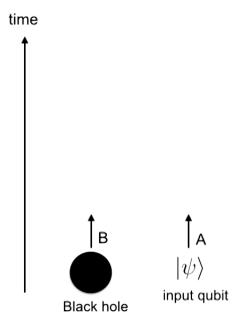
[Choi-Jamilkowski 1960s, Hayden-Preskill, Hartman-Maldacena]

State-Tensor duality, again

• Unitary operator acting on n qubits can be viewed as a state on 2n qubits.

$$U = \sum_{i,j} U_{ij} |i\rangle\langle j|$$
 in U out $|U\rangle = \sum_{i,j} U_{ij} |i\rangle\otimes|j\rangle$ $|U\rangle$

• Viewing the black hole dynamics as a quantum state.



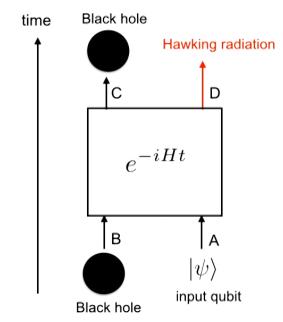
[Choi-Jamilkowski 1960s, Hayden-Preskill, Hartman-Maldacena]

State-Tensor duality, again

• Unitary operator acting on n qubits can be viewed as a state on 2n qubits.

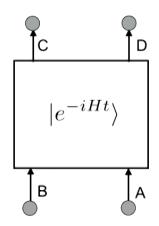
$$U = \sum_{i,j} U_{ij} |i\rangle\langle j|$$
 out $|U\rangle = \sum_{i,j} U_{ij} |i\rangle\otimes|j\rangle$ $|U\rangle$

• Viewing the black hole dynamics as a quantum state.



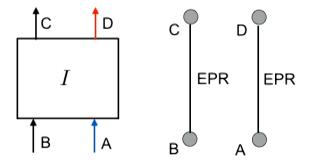
Entanglement between input and output

eg. between A and D = radiation and input



[Choi-Jamilkowski 1960s, Hayden-Preskill, Hartman-Maldacena]

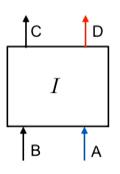
• An identity operator

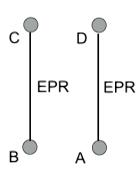


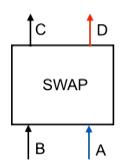
Pirsa: 17020102 Page 74/97

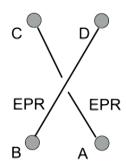
• An identity operator

SWAP operator



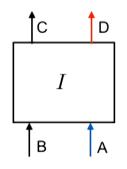


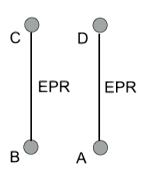




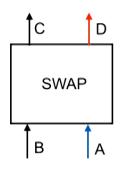
Pirsa: 17020102 Page 75/97

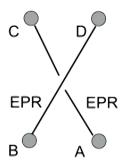
• An identity operator



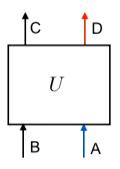


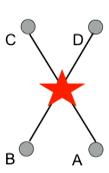
SWAP operator





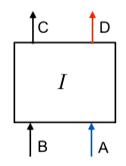
• Scrambling

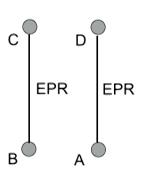




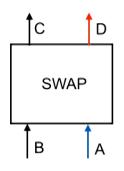
Pirsa: 17020102 Page 76/97

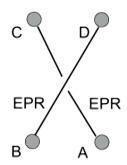
• An identity operator



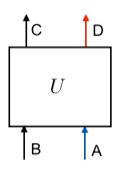


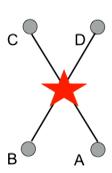
• SWAP operator





Scrambling





Quantum error-correcting code postulate

No correlation between local inputs and local outputs

"mutual information"

$$I(A,D) = 0$$
 $I(A,C) = 0$

• Average of OTO over local operators A and D at T=infty [Hosur-Qi-Roberts-BY15]

$$|\langle O_A(0)O_D(t)O_A(0)O_D(t)\rangle|_{\text{ave}}$$



Pirsa: 17020102 Page 78/97

• Average of OTO over local operators A and D at T=infty [Hosur-Qi-Roberts-BY15]

$$|\langle O_A(0)O_D(t)O_A(0)O_D(t)
angle|_{
m ave}$$



Pirsa: 17020102 Page 79/97

Average of OTO over local operators A and D at T=infty [Hosur-Qi-Roberts-BY15]

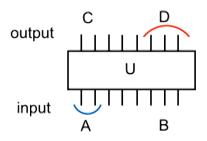
$$|\langle O_A(0)O_D(t)O_A(0)O_D(t)\rangle|_{\rm ave} = \frac{1}{4^{a+d}} \sum_{O_A,O_D} \langle O_A(0)O_D(t)O_A(0)O_D(t)\rangle$$
 Pauli operators (unitary 1-design) average over OA, OD



Pirsa: 17020102 Page 80/97

Average of OTO over local operators A and D at T=infty [Hosur-Qi-Roberts-BY15]

$$|\langle O_A(0)O_D(t)O_A(0)O_D(t)\rangle|_{\rm ave} = \frac{1}{4^{a+d}} \sum_{O_A,O_D} \langle O_A(0)O_D(t)O_A(0)O_D(t)\rangle$$
 Pauli operators (unitary 1-design) average over OA, OD





Pirsa: 17020102 Page 81/97

Average of OTO over local operators A and D at T=infty [Hosur-Qi-Roberts-BY15]

$$|\langle O_A(0)O_D(t)O_A(0)O_D(t)\rangle|_{\text{ave}} = \frac{1}{4^{a+d}} \sum_{\substack{O_A,O_D\\\text{eventure}}} \langle O_A(0)O_D(t)O_A(0)O_D(t)\rangle$$

$$= 2^{n-a-d-S_{BD}^{(2)}}$$
Renyi-2 entropy
$$A \qquad D$$

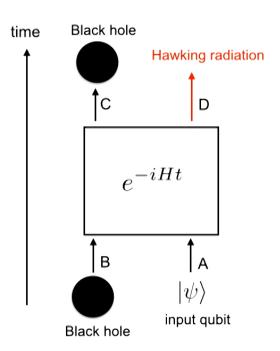
$$= D$$

$$\text{input} \qquad A \qquad B$$

• If $OTO \simeq 0$ then, $S_{BD}^{(2)}$ is large

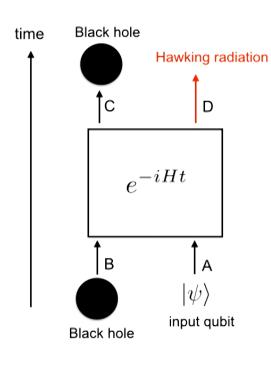
This implies the mutual information $I^{(2)}(B,D)=S_B^{(2)}+S_D^{(2)}-S_{BD}^{(2)}$ is small

• If $OTO \simeq 0$ then, $I^{(2)}(A,C) = S_A^{(2)} + S_C^{(2)} - S_{AC}^{(2)}$ is also small.



[Hosur-Qi-Roberts-BY15, Roberts-BY16, BY in prep]

Pirsa: 17020102 Page 83/97



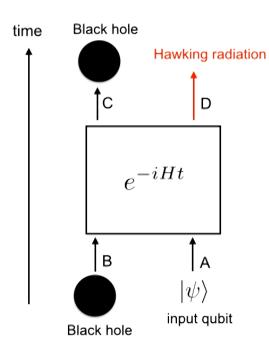
• Hawking-Unruh semiclassical calculation

$$\langle O_A(0)O_A(t)\rangle \to 0$$
 $\Rightarrow I(A,D) \to 0$

• Kitaev-Shenker-Stanford shockwave calculation

$$\langle O_A(0)O_D(t)O_A(0)O_D(t)\rangle \to 0$$

[Hosur-Qi-Roberts-BY15, Roberts-BY16, BY in prep]



• Hawking-Unruh semiclassical calculation

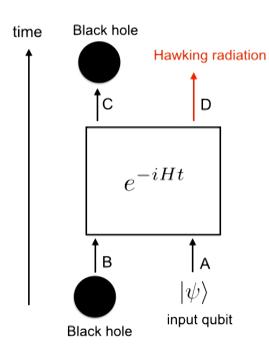
$$\langle O_A(0)O_A(t)\rangle \to 0$$
 $\Rightarrow I(A,D) \to 0$

• Kitaev-Shenker-Stanford shockwave calculation

$$\langle O_A(0)O_D(t)O_A(0)O_D(t)\rangle \to 0$$

$$\Rightarrow I(A,C) \to 0$$

[Hosur-Qi-Roberts-BY15, Roberts-BY16, BY in prep]



• Hawking-Unruh semiclassical calculation

$$\langle O_A(0)O_A(t)\rangle \to 0$$
 $\Rightarrow I(A,D) \to 0$

• Kitaev-Shenker-Stanford shockwave calculation

$$\langle O_A(0)O_D(t)O_A(0)O_D(t)\rangle \to 0$$

$$\Rightarrow I(A,C) \to 0$$

So, a black hole is a quantum error-correcting code.

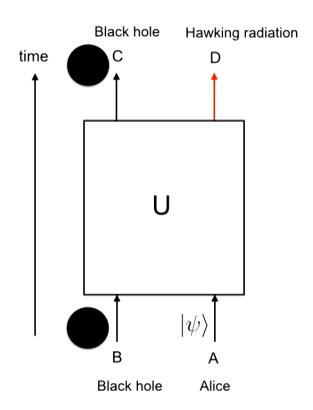
[Hosur-Qi-Roberts-BY15, Roberts-BY16, BY in prep]

Any new insight into information loss puzzle?

Pirsa: 17020102 Page 87/97

Hayden-Preskill thought experiment (2007)

• Can Bob reconstruct Alice's quantum state?



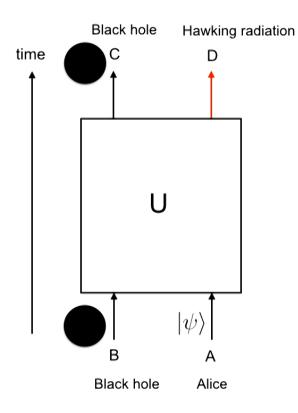




Pirsa: 17020102 Page 88/97

Hayden-Preskill thought experiment (2007)

• Can Bob reconstruct Alice's quantum state?







Bob has an access to

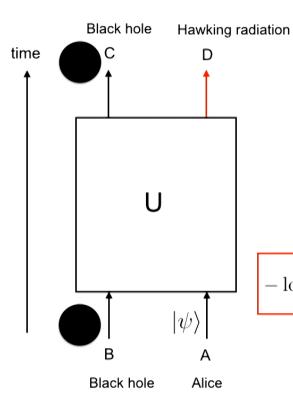
B: Black hole initial state

D: Hawking radiation

Pirsa: 17020102 Page 89/97

Hayden-Preskill thought experiment (2007)

• Can Bob reconstruct Alice's quantum state?







Bob has an access to

B: Black hole initial state

D : Hawking radiation

$$-\log_2 |\langle O_A(0)O_D(t)O_A(0)O_D(t)\rangle| = I^{(2)}(A, BD)$$

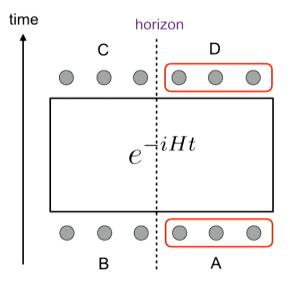
[Roberts-BY16]

Small OTOC ——— Bob's success

Pirsa: 17020102 Page 90/97

Seeing "behind the horizon" with OTOCs?

• OTOCs provide a way of seeing behind the black hole horizon.



Pirsa: 17020102 Page 91/97

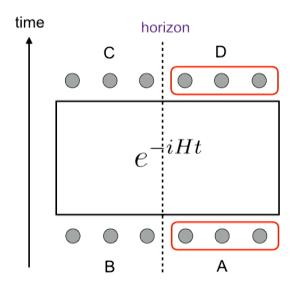
Seeing "behind the horizon" with OTOCs?

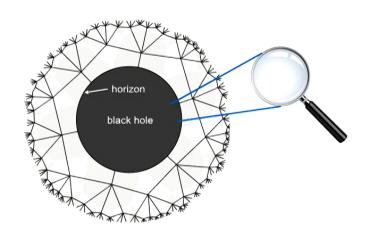
• OTOCs provide a way of seeing behind the black hole horizon.

Assume that we have an access only to A

$$|\langle O_A(0)O_D(t)O_A(0)O_D(t)\rangle|_{\text{ave}} \longrightarrow I(A,C)$$

We can study cross-horizon correlations.





Pirsa: 17020102 Page 92/97

Seeing "behind the horizon" with OTOCs?

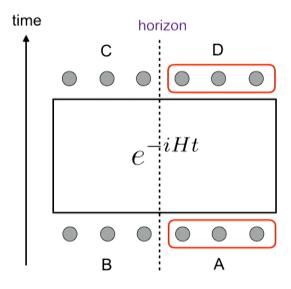
• OTOCs provide a way of seeing behind the black hole horizon.

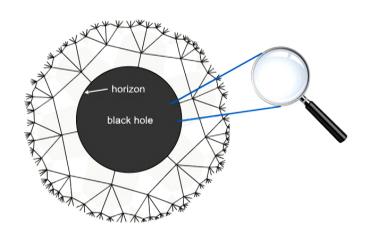
Assume that we have an access only to A

$$|\langle O_A(0)O_D(t)O_A(0)O_D(t)\rangle|_{\text{ave}} \longrightarrow I(A,C)$$

We can study cross-horizon correlations.

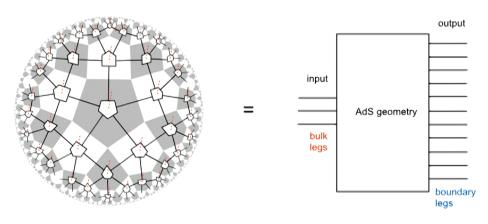
• But, we need a finite-temperature generalization... (ongoing work with Mozgunov and Kitaev)



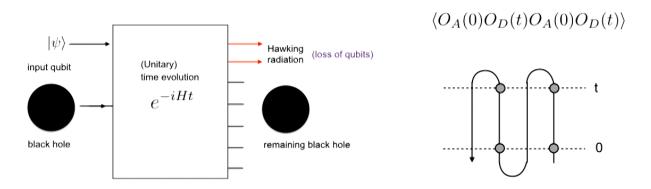


Pirsa: 17020102 Page 93/97

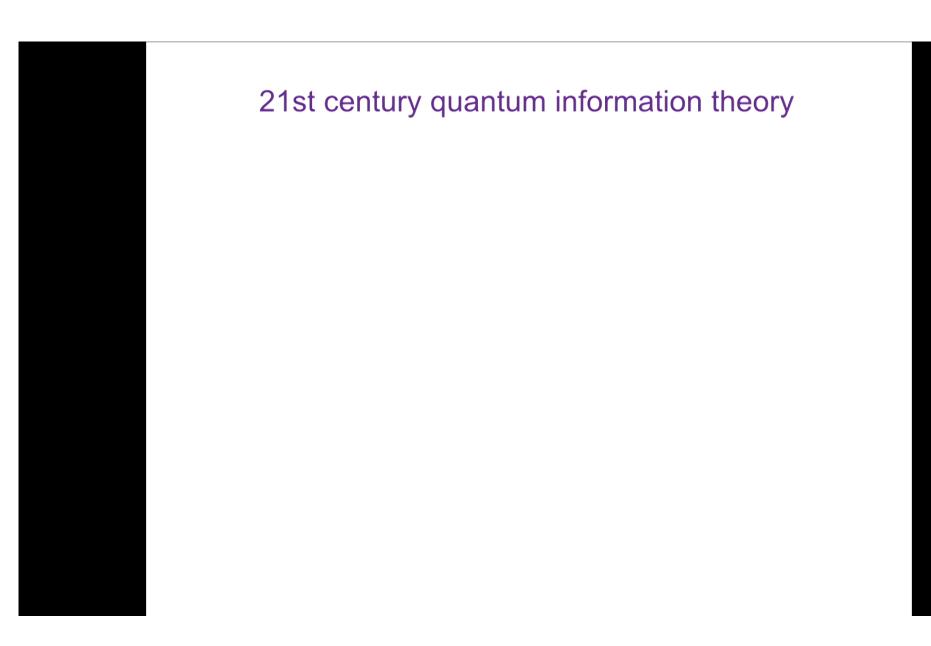
Part 1: AdS/CFT correspondence is a quantum error-correcting code.



Part 2: Black hole is a quantum error-correcting code.



Pirsa: 17020102 Page 94/97



Pirsa: 17020102 Page 95/97

21st century quantum information theory

Entanglement entropy (1932)



von Neumann

Pirsa: 17020102 Page 96/97

21st century quantum information theory

Entanglement entropy (1932)

Factoring algorithm (1994)



von Neumann



Peter Shor

Pirsa: 17020102 Page 97/97