

Title: Harnessing Topological Quantum Matter

Date: Mar 10, 2017 04:00 PM

URL: <http://pirsa.org/17020099>

Abstract: <p>Topological quantum computing requires phases of matter which host fractionalized excitations that are neither bosons nor fermions. I will present a new route toward realizing such fractionalized phases of matter by literally building on existing topological phases. I will first discuss how existing topological phases, when interfaced with other systems, can exhibit a “topological proximity effect” in which nontrivial topology of a different nature is induced in the neighboring system. Then, I will show how this enables a new entanglement based technique (the “topological bootstrap”) for upgrading topological phases from the integer into the fractional quantum Hall variety. Finally, I will highlight the rich phenomenology of systems with interacting Majorana modes. Such systems can exhibit physics ranging from black hole scrambling to supersymmetry and from alternative surface code architectures to topological phases in three dimensions with completely immobile excitations. I will discuss my plans for understanding both general properties and specific models of such fascinating systems.</p>



# Harnessing Topological Quantum Matter

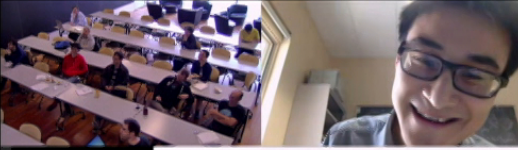
Tim Hsieh

KITP, UCSB

Perimeter Institute Seminar

3/10/17





skype™

# The Maple Bacon Donut





# Outline

## I. Topological Materials

Nature Communications 3:982 (2012)  
Nature Materials 13:178 (2014)

w/ Liang Fu (MIT)  
Junwei Liu (MIT)  
Arun Bansil (Northeastern)  
Wenhui Duan (Tsinghua)  
Jagadeesh Moodera (MIT)  
Peng Wei (UCR)

## II. Topological Proximity Effect

PRL 116, 086802 (2016)  
PRB Rapid 94, 161112 (R) (2016)

w/ Hiro Ishizuka (U. Tokyo)  
Leon Balents (KITP)  
Taylor Hughes (UIUC)

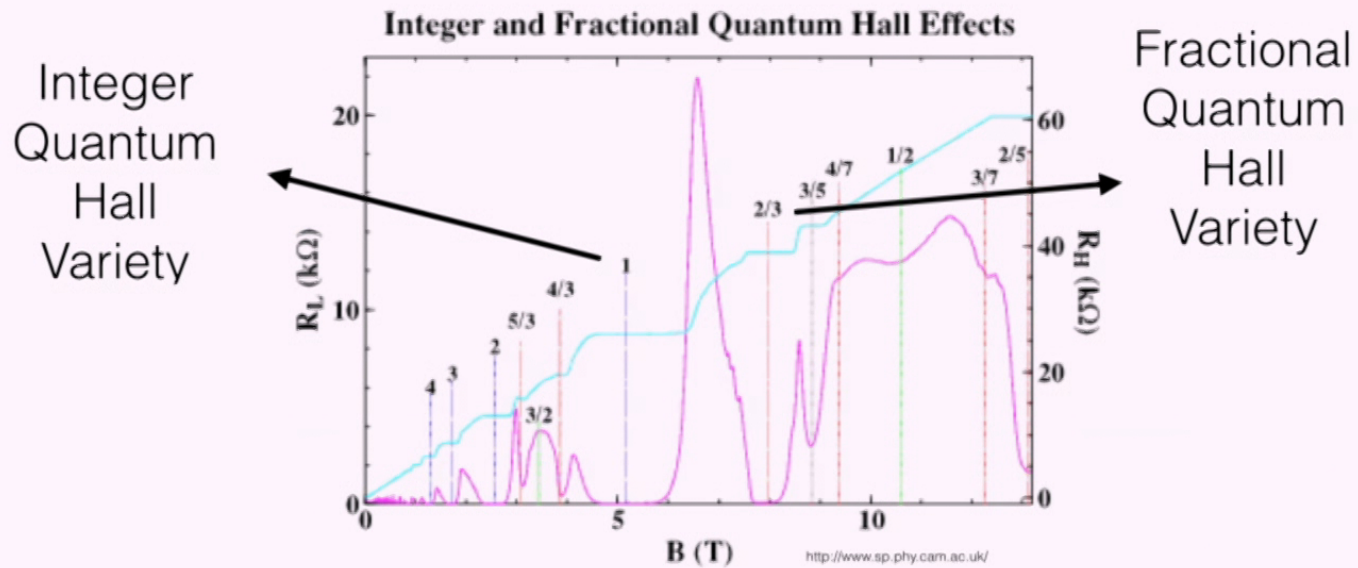
## III. Topological Bootstrap

arXiv:1610.04614 (2016)

w/ Yuan-Ming Lu (Ohio State U.)  
Andreas Ludwig (UCSB)



# A Tale of Two Topologies



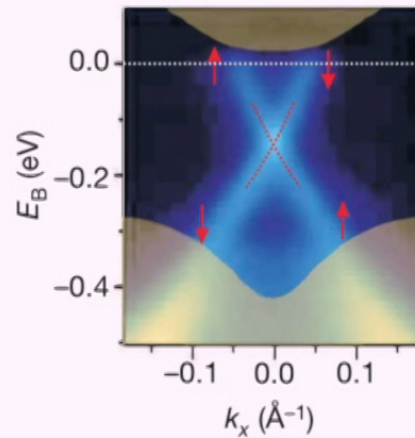


# Topological Materials



$\text{Bi}_2\text{Se}_3$

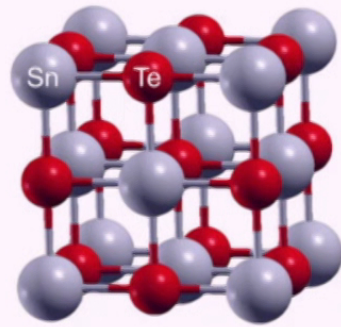
**time reversal symmetry** protects gapless surface states



D. Hsieh et.al. Nature 2009



# Topological Materials

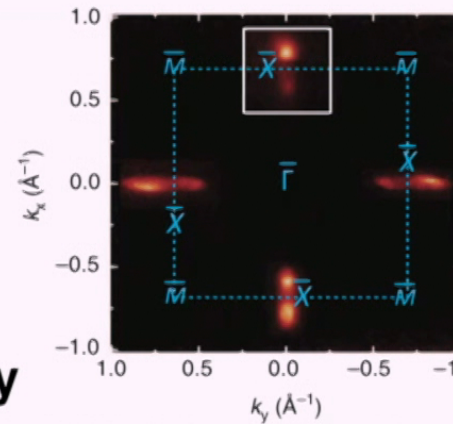


SnTe:  
Topological Crystalline Insulator (TCI)

TH, et.al. Nature Communications 3:982 (2012)

**Mirror reflection symmetry**  
protects gapless surface states

**No need for time reversal symmetry**

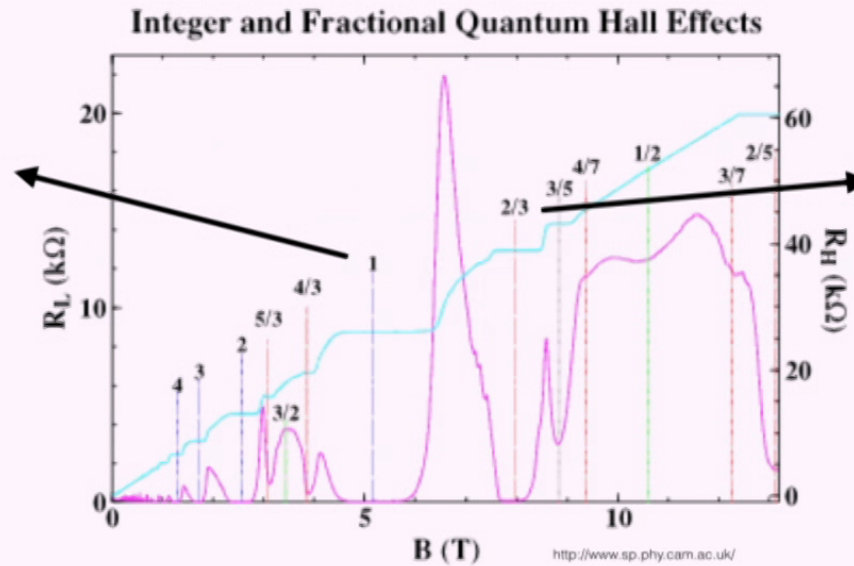


Pb<sub>1-x</sub>Sn<sub>x</sub>Te ARPES (S. Xu et.al. Nature Materials, 2012)



# A Tale of Two Topologies

Integer  
Quantum  
Hall  
Variety



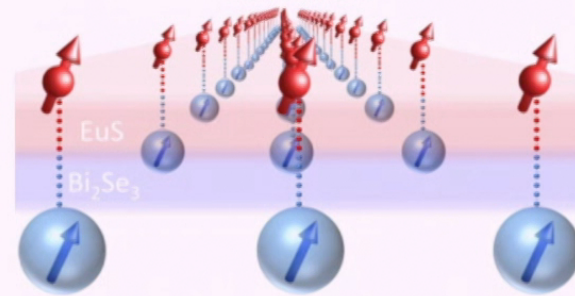
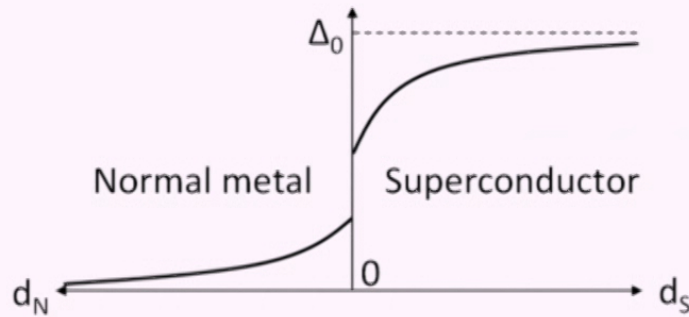
Fractional  
Quantum  
Hall  
Variety





# What is the effect of topological phases on proximate systems?

Existing proximity effects all have local order:



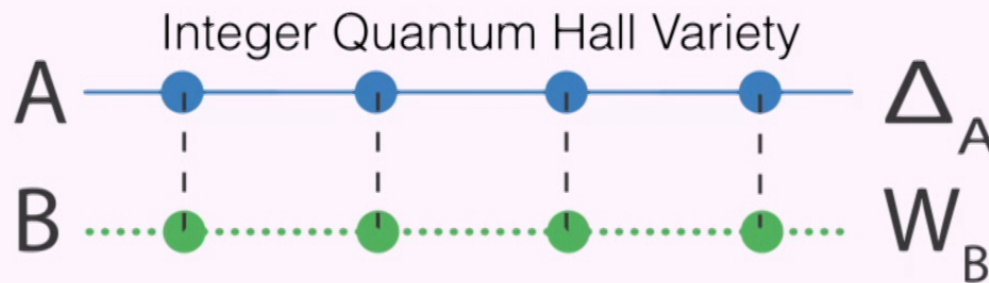
Katmis, et.al. Nature 2016

Is there a topological proximity effect?



# General Setup I

$$H = H_A + \cancel{H_B} + H_{AB}$$



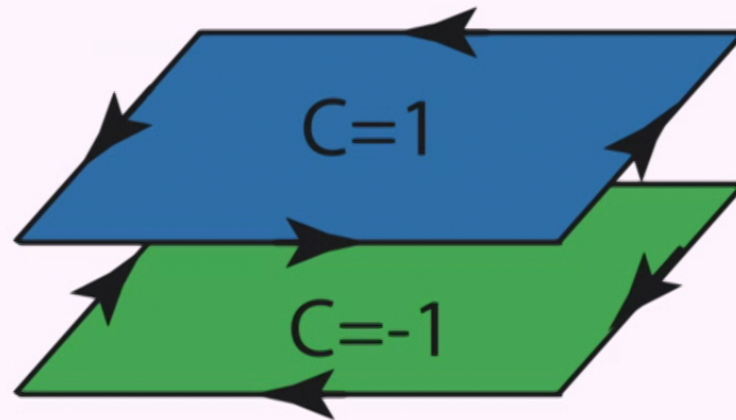
$$W_B \ll g_{AB} \ll \Delta_A$$



# Integer Quantum Hall Variety

Chern insulator: Integer Chern number

Dictates number of chiral gapless edge modes





# Integer Quantum Hall Variety

Chern insulator: Integer Chern number

Dictates number of chiral gapless edge modes





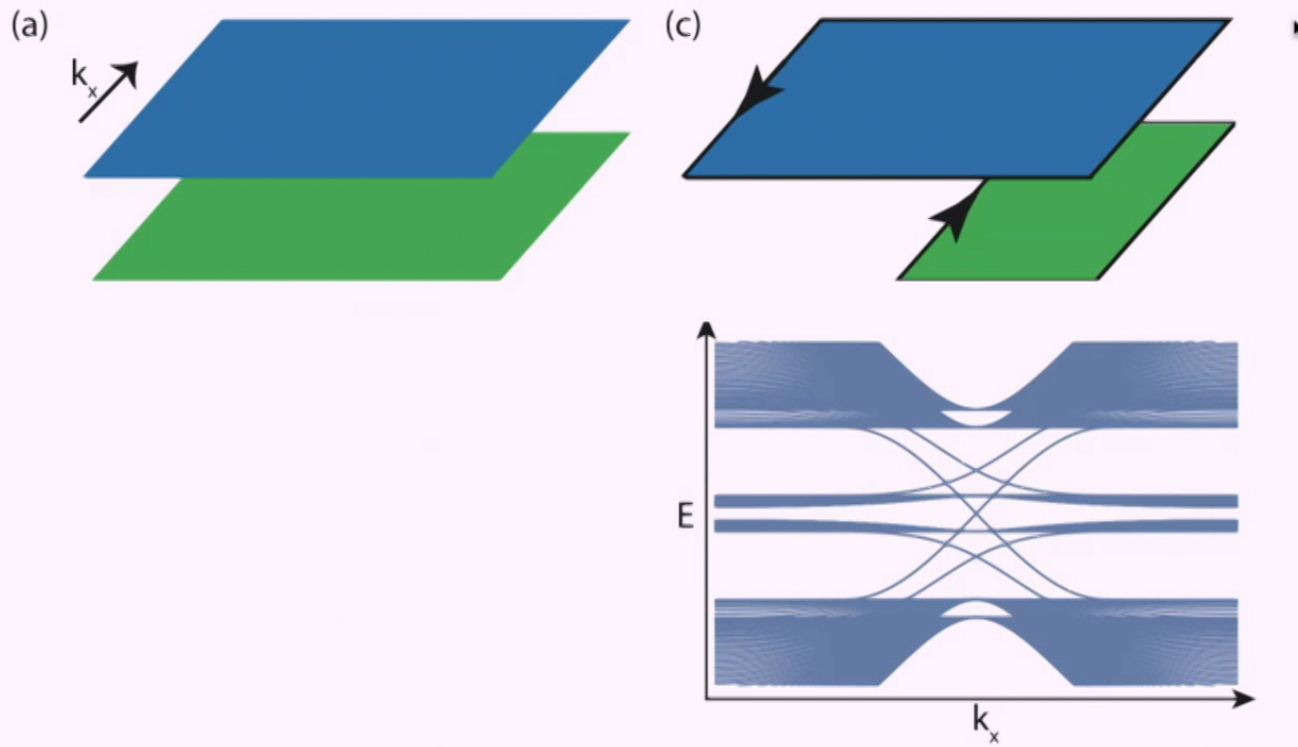
# Chern Insulator, Proximitized



$$h_B^{\text{eff}}(k) = -\frac{4g^2}{\Delta(k)} |\psi_k^{\text{ex}}\rangle \langle \psi_k^{\text{ex}}|$$



# Phenomenology





# Generalizes to **All** Free Fermion Topological Phases

$$\begin{pmatrix} E1 & 0 & g & 0 \\ 0 & E2 & 0 & g \\ g & 0 & 0 & 0 \\ 0 & g & 0 & 0 \end{pmatrix}$$

$$H = \begin{pmatrix} H_A^d & g \\ g & 0 \end{pmatrix}$$

$$H_A^d = UH_AU^{-1}$$

$E_n^0$ : eigenvalues of  $H_A$



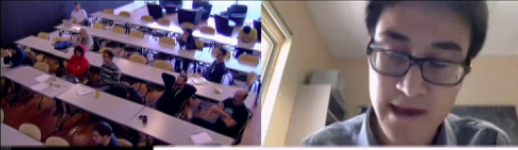
# Generalizes to **All** Free Fermion Topological Phases

$$H = \bigoplus_n \begin{pmatrix} E_n^0 & g \\ g & 0 \end{pmatrix}$$

Gap increases monotonically

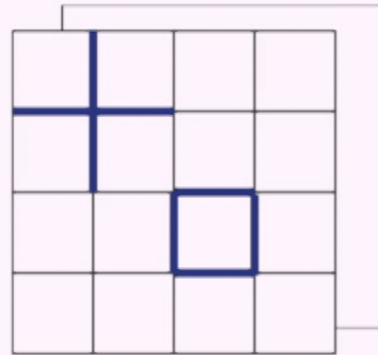
No phase transition from weak to strong coupling





# Fractionalized Phases

Kitaev's Toric Code



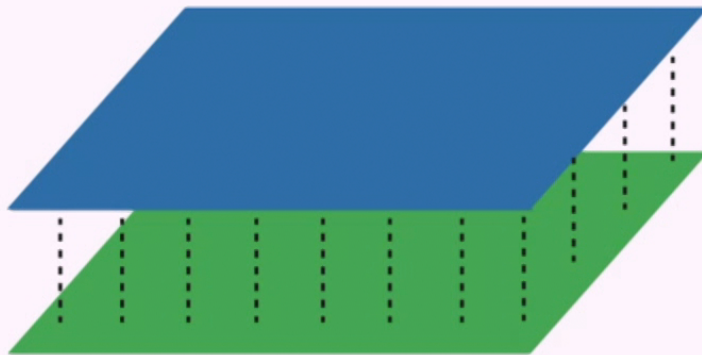
$$H_A = - \sum_s \prod_{l \in s} \sigma_l^z - \sum_p \prod_{l \in p} \sigma_l^x,$$

Toric code  
duplicates itself!

$$H_{AB} = g \sum_i (\sigma_i^z \tilde{\sigma}_i^z + \sigma_i^x \tilde{\sigma}_i^x)$$

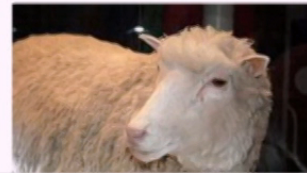
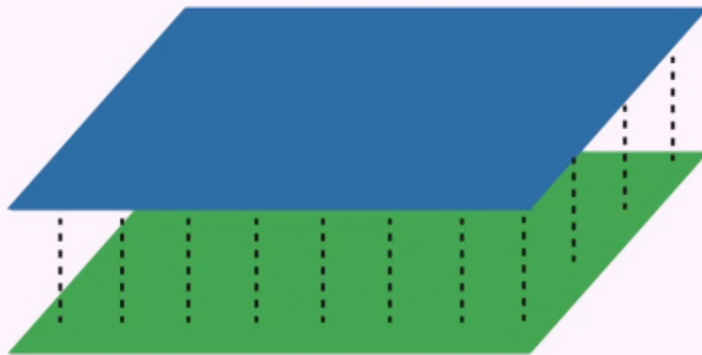


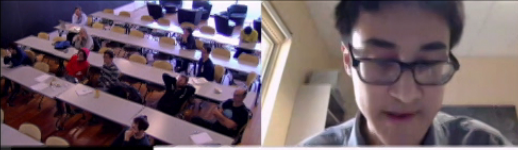
# Part III: Topological Bootstrap



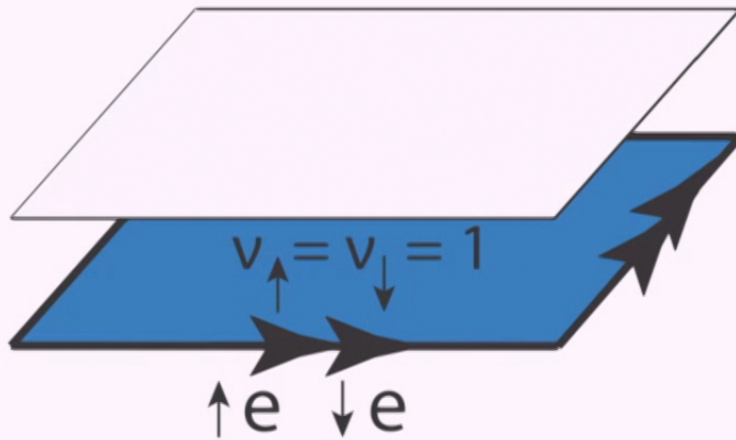


# Part III: Topological Bootstrap





# General Setup II



B: Free spin-1/2s

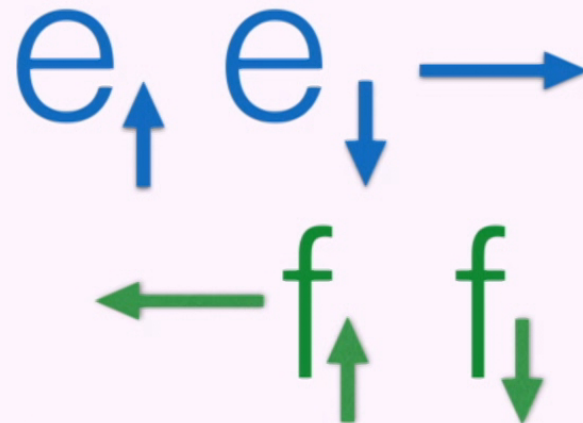
A: Free fermions with spin  
in Chern band

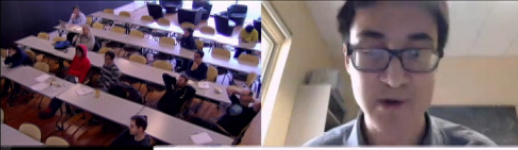
$$\begin{aligned} H &= H_0 + H_K = \\ &= H_{\nu=1}(c_{\uparrow}) + H_{\nu=1}(c_{\downarrow}) + g \sum_i (c_i^{\dagger} \vec{\sigma} c_i) \cdot \vec{S}_i \end{aligned}$$



# Not the Usual Kondo Lattice

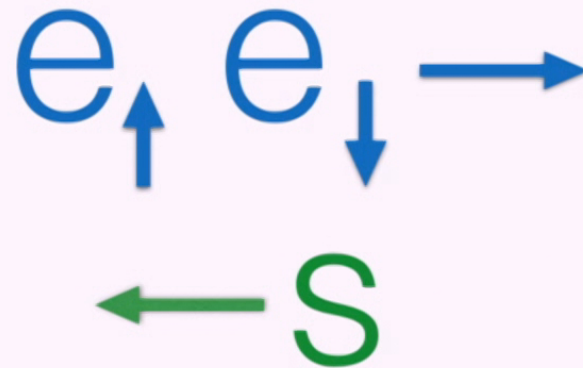
- Gapped electron system  $\rightarrow$  singlet formation evaded for small coupling ( $g \ll \Delta_A$ )
- Intuition from topological proximity effect: spins want to screen out Chern insulator





# Not the Usual Kondo Lattice

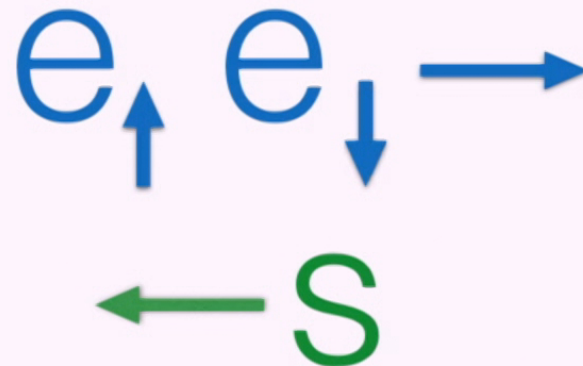
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# Not the Usual Kondo Lattice

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spins form  
analogue  
of Chern  
insulator

$$\vec{S} = \frac{1}{2} f^\dagger \vec{\sigma} f$$

slave fermions  
inherit  
inverse Chern #



# Chiral Spin Liquid

- No local magnetic order
- Breaks time-reversal symmetry; chiral edge states
- Fractionalized phase
- Excitations are “semions”



shaken and stirred

V. Kalmeyer and R. Laughlin, PRL 59, 2095 (1987)  
X. Wen, F. Wilczek, and A. Zee, PRB 39, 11413 (1989)





# An Exact Limit

AFM Kondo coupling = local projectors onto singlets

$$|\psi_{AB}\rangle = \begin{array}{cccc} \bullet & \bullet & \bullet & \bullet \\ | & | & | & | \\ \bullet & \bullet & \bullet & \bullet \end{array}$$

Replace Kondo coupling with global projector:

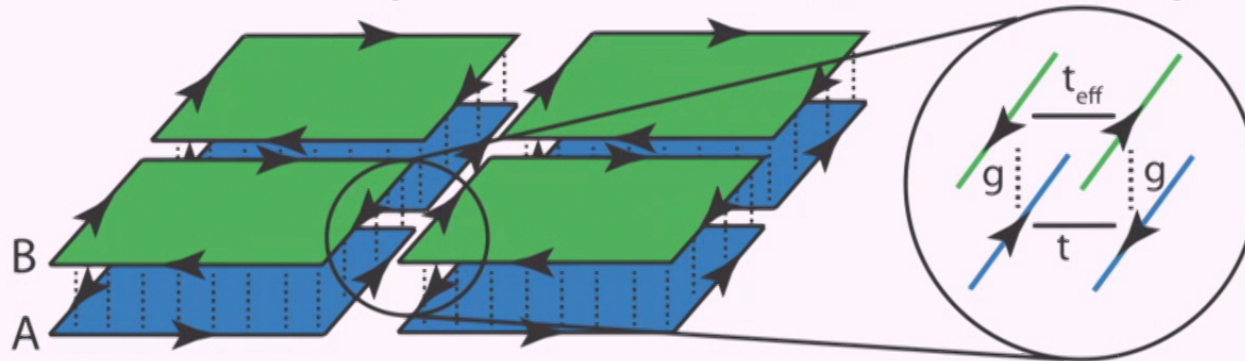
$$H = H_{\nu=1}(c_{\uparrow}) + H_{\nu=1}(c_{\downarrow}) - g|\psi_{AB}\rangle\langle\psi_{AB}|$$
$$H_{eff}^B = -g|\overline{\psi}_0^s\rangle\langle\overline{\psi}_0^s|$$

$|\overline{\psi}_0^s\rangle =$  Time-reversed Gutzwiller projection of Chern bands



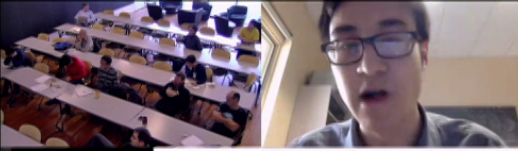
# Interpolate from Global to Local Kondo

Partition A,B into regions  $r$  of size  $l > \xi$  (correlation length of A)

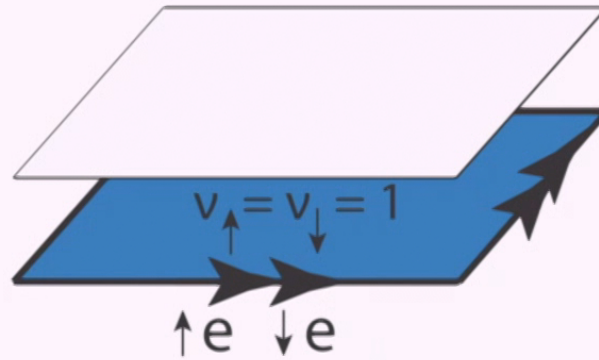


Quasi-local projector  $P_r = -g|\psi_{AB,r}\rangle\langle\psi_{AB,r}|$

To interpolate all the way, require  $\xi$  to be minimal ( $\rightarrow 0$ )  
Expect CSL to be induced for range of Kondo coupling  $g$



# Phenomenology

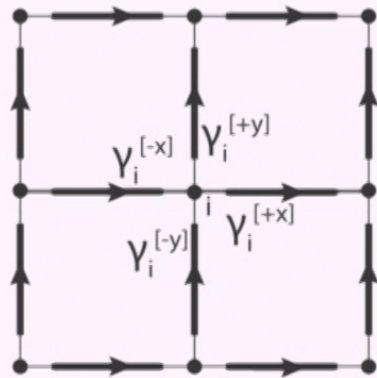




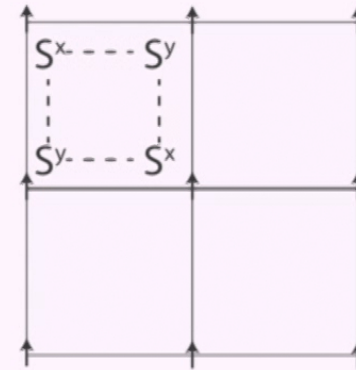
# $Z_2$ Spin Liquid from Superconductor



# $\mathbb{Z}_2$ Spin Liquid from Superconductor



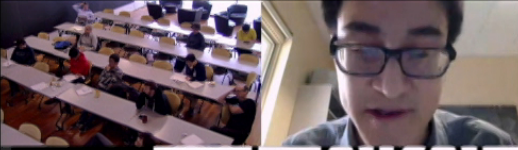
$$H_{AB} = g \sum_{i,\alpha,\beta} \vec{S}_i \cdot (c_{i,\alpha}^\dagger \vec{\sigma}_{\alpha\beta} c_{i,\beta})$$



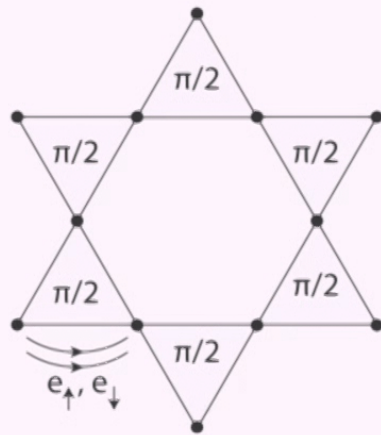
$$H_A = \sum_i \sum_{\mu=+x,+y} i \gamma_i^{[+\mu]} \gamma_{i+\mu}^{[-\mu]}$$

$$H_B^{eff} \propto g^4 \sum_i S_i^y S_{i+x}^x S_{i+x+y}^y S_{i+y}^x$$

Wen plaquette model  
(toric code)

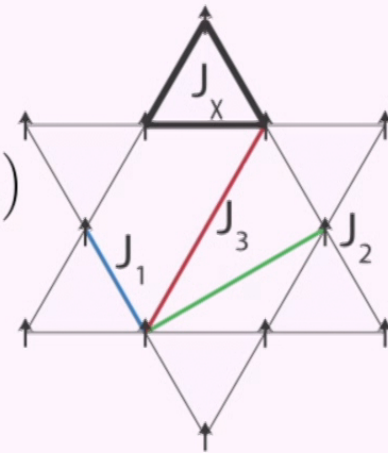


# Chiral Spin Liquid from Chern Insulator



$$|\nu_{\uparrow} = 1\rangle \otimes |\nu_{\downarrow} = 1\rangle$$

$$H_{AB} = g \sum_{i,\alpha,\beta} \vec{S}_i \cdot (c_{i,\alpha}^{\dagger} \vec{\sigma}_{\alpha\beta} c_{i,\beta})$$



$$J_1 = 1, J_2 = 0.51, J_3 = 0.56$$

S-S. Gong, W. Zhu, and D. Sheng, Sci.Reports 4, 6317 (2014)

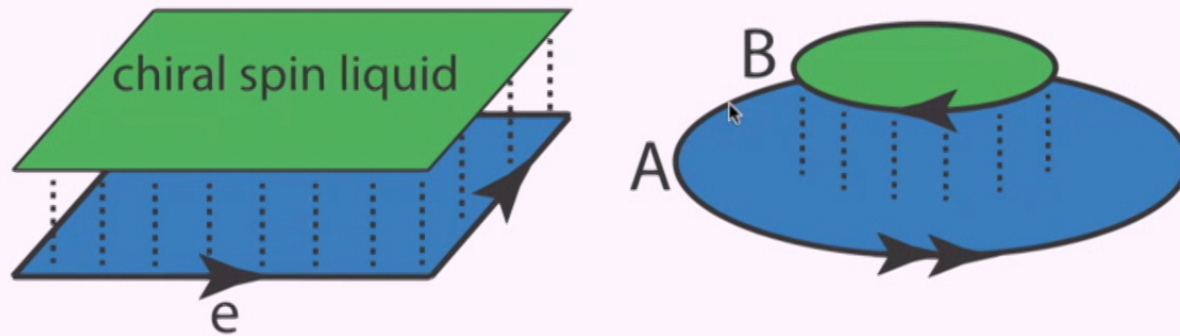
B. Bauer et.al. Nat. Comm. 5, 5137 (2014)

Chiral Spin Liquid



# Topological Bootstrap

More is different and less: Fractionalization via Kondo coupling

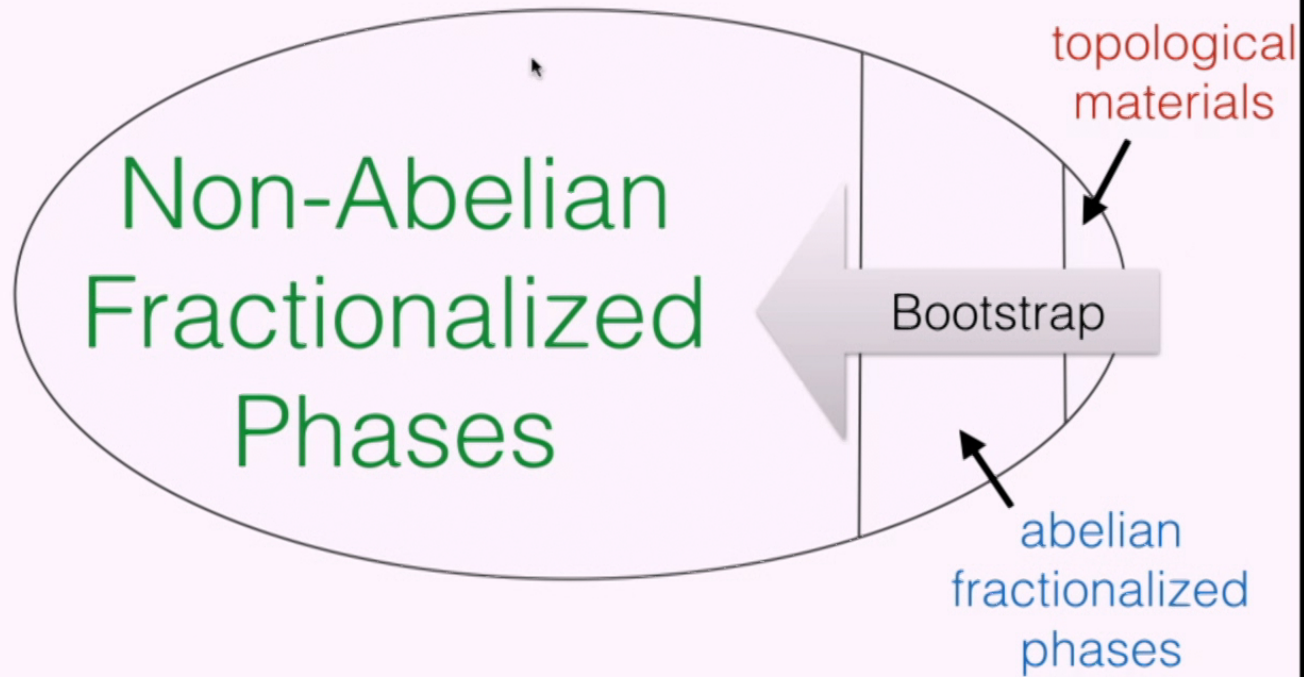


Potential realization in synthetic quantum matter

Apply techniques to Kondo lattices in heavy fermion systems,  
e.g.  $\text{SmB}_6$



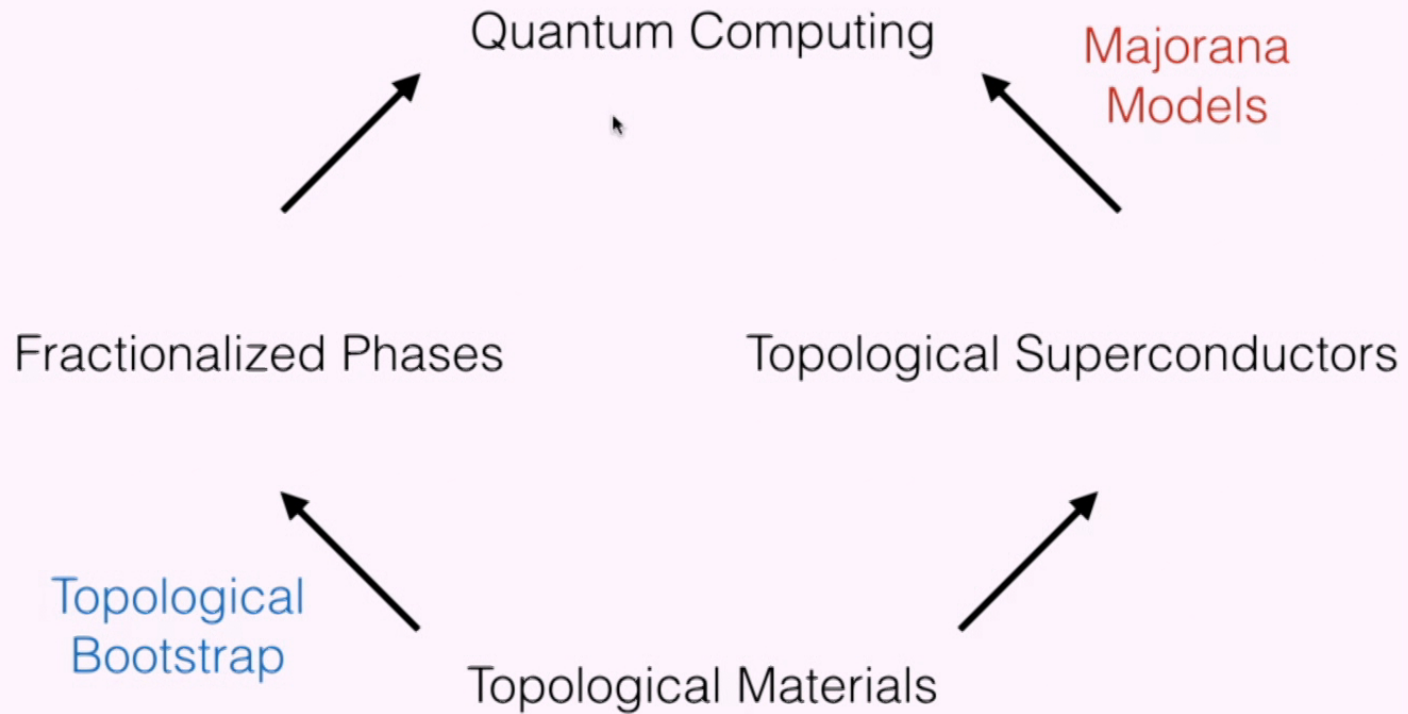
# Pushing the Topological Frontier via Bootstrap







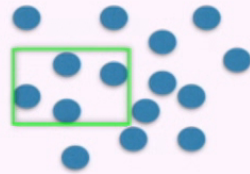
# The Bigger Picture





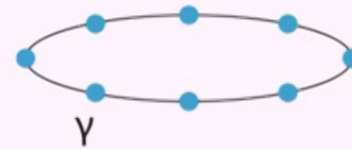
# Majorana Models

0d: quantum chaos, black holes



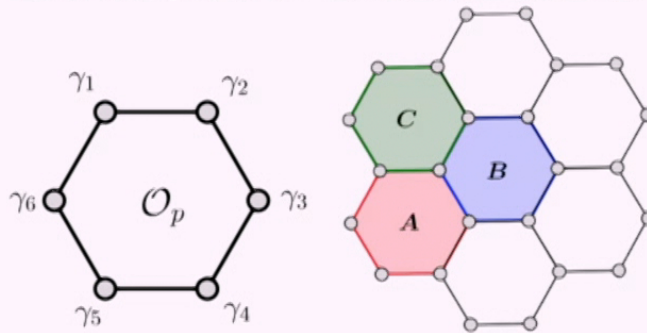
Sachdev, Ye (1993), Kitaev (2015)

1d: supersymmetry



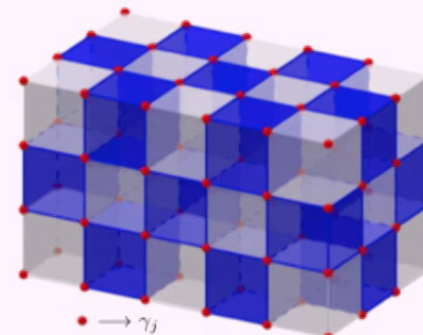
**TH**, Halasz, Grover (2016)

2d: Majorana surface codes



Vijay, **TH**, Fu (2015)

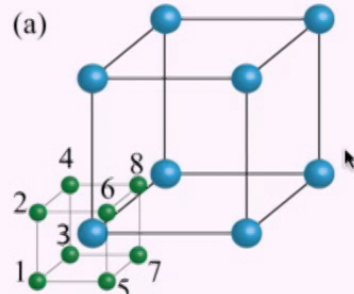
3d: Fracton topological order



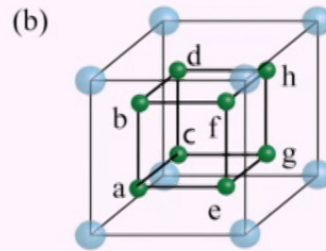
Vijay, Haah, Fu (2016)



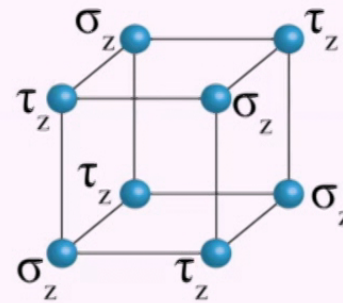
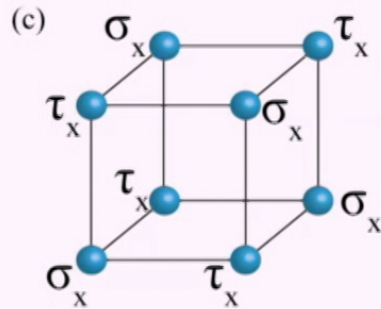
# Fractons from Partons



$$\gamma_1 \gamma_4 \gamma_6 \gamma_7 = \gamma_2 \gamma_3 \gamma_5 \gamma_8 = 1$$



$$\begin{aligned} \gamma_a \gamma_b \gamma_c \gamma_d &= \gamma_a \gamma_c \gamma_e \gamma_g = 1 \\ \gamma_a \gamma_b \gamma_e \gamma_f &= \gamma_c \gamma_f \gamma_g \gamma_h = 1 \end{aligned}$$

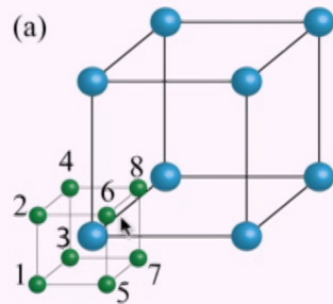


TH and G. Halasz,  
arXiv:1703.02973

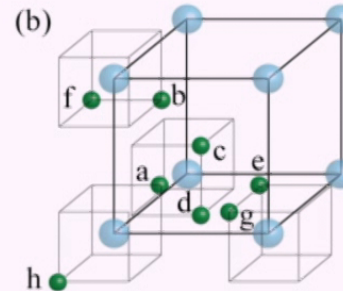
New Type-I fracton w/ string and membrane logical operators



# Fractons from Partons

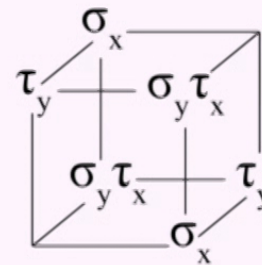
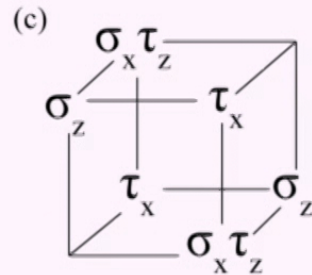


$$\gamma_1 \gamma_4 \gamma_6 \gamma_7 = \gamma_2 \gamma_3 \gamma_5 \gamma_8 = 1$$



$$\gamma_a \gamma_b \gamma_c \gamma_d = \gamma_a \gamma_c \gamma_e \gamma_g = 1$$

$$\gamma_a \gamma_b \gamma_e \gamma_f = \gamma_e \gamma_f \gamma_g \gamma_h = 1$$

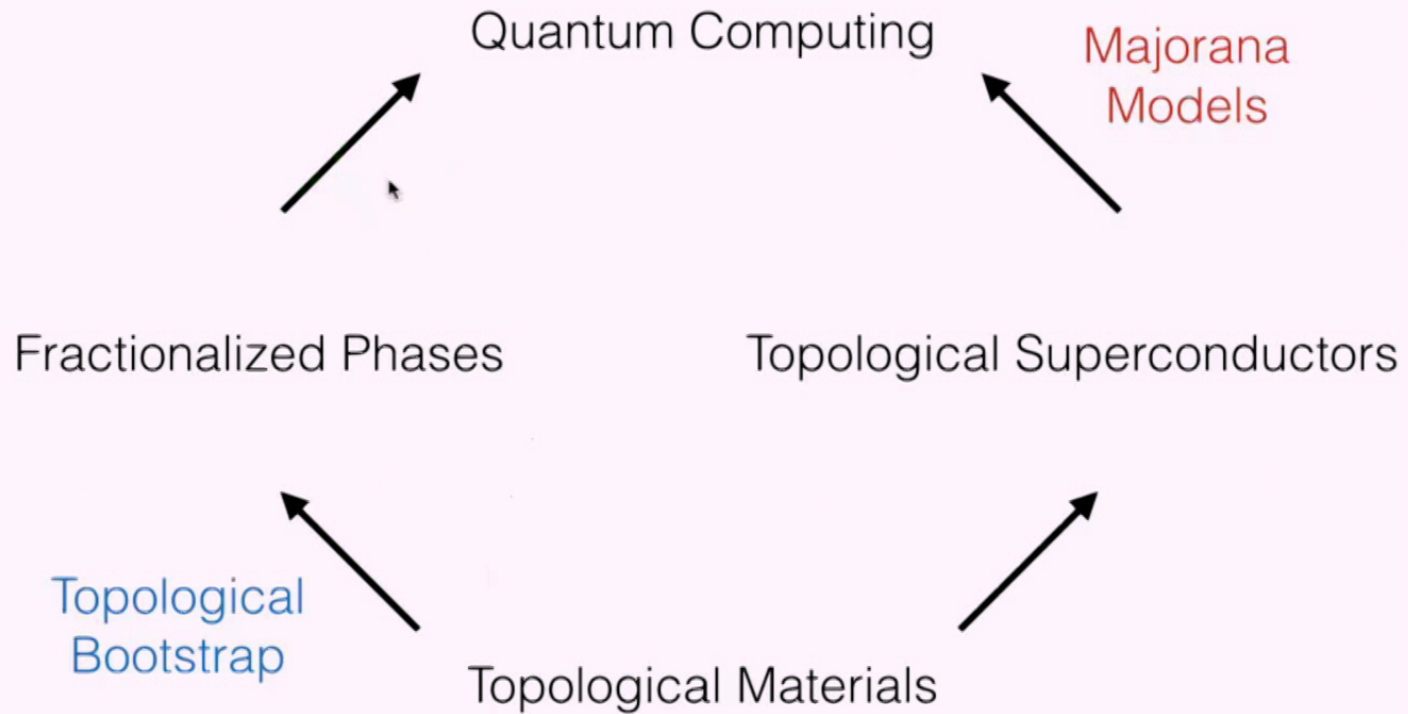


TH and G. Halasz,  
arXiv:1703.02973

New Type-II non-CSS fracton w/ fractal logical operators

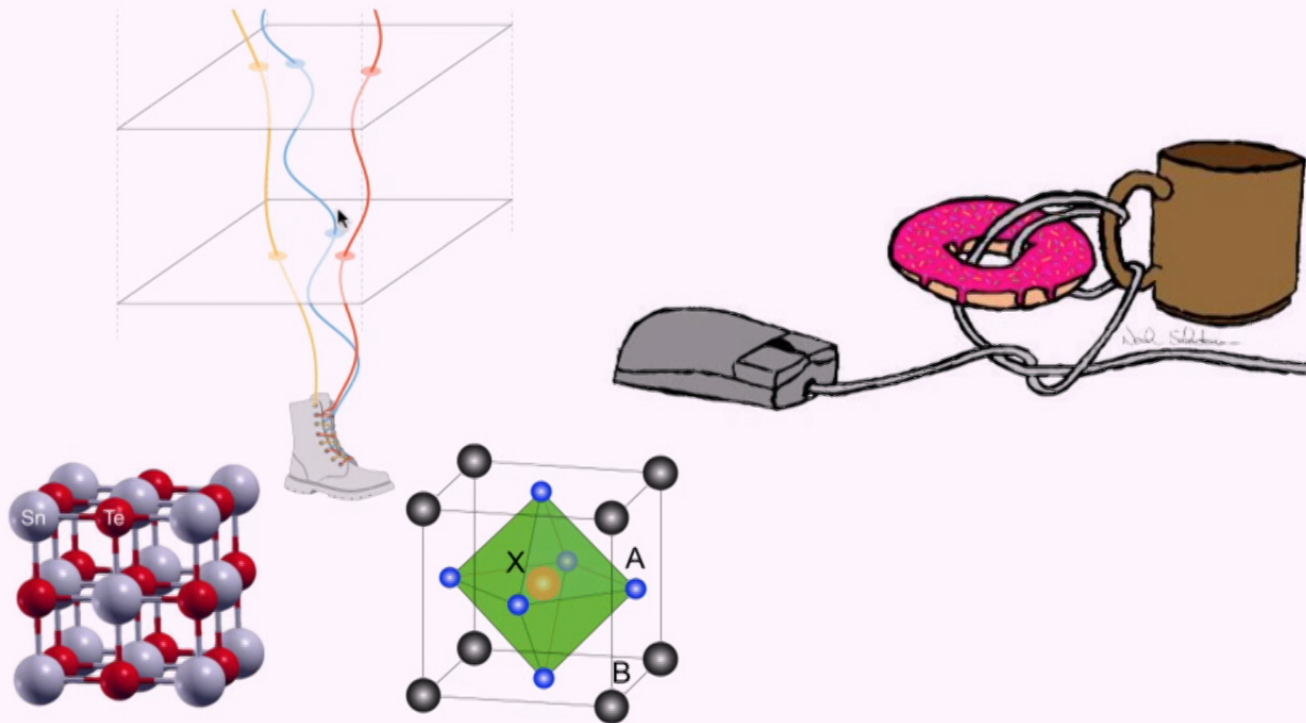


# The Bigger Picture





# Transforming Topology



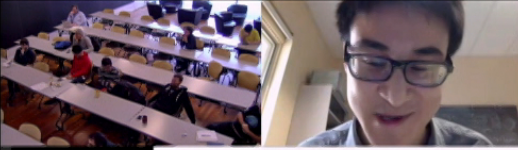


# Collaborations

Theory: L. Fu (MIT), L. Balents (KITP), A. Ludwig (UCSB),  
T. Hughes (UIUC), Y. Lu (Ohio State), X. Qi (Stanford)

Experiment: S. Nakatsuji (ISSP Tokyo), P. Ong (Princeton)

Project supervision: W. Ho (Geneva -> Harvard),  
T. Iadecola (BU -> JQI Maryland)



# Thanks!

SnTe TCIs	Nature Communications 3:982 (2012)
SnTe Thin Film TCIs	Nature Materials 13:178 (2014)
Bulk Topological Proximity Effect	PRL 116, 086802 (2016)
Entangled Cloning	PRB 94, 161112 (R) (2016)
Topological Bootstrap	arXiv:1610.04614 (2016)
Majorana Surface Code	PRX 5, 041038 (2015)
Majorana Supersymmetry	PRL 117, 166802 (2016)
Fractons from Partons	arXiv:1703.02973 (2017)

