

Title: PSI 2016/2017 Standard Model (Review) - Lecture 12

Date: Jan 18, 2017 09:00 AM

URL: <http://pirsa.org/17010016>

Abstract:

representations

fundamental" of $SU(3)$.

"3"

$$q_i \rightarrow U_{ij} q_j$$

$$\bar{q}_i \rightarrow U_{ij}^* \bar{q}_j$$

fundamental."

" $\bar{3}$ "

Young tableaux & Irreducible representations

Single quark $q = \begin{pmatrix} u \\ d \\ s \end{pmatrix}$ is a "fundamental" of $SU(3)$

antiquark $\bar{q} = \begin{pmatrix} \bar{u} \\ \bar{d} \\ \bar{s} \end{pmatrix}$ is an "antifundamental".

Diquark (distinguishable): $q_i^{(1)} q_j^{(2)}$

Young tableaux & Irreducible representations

Single quark $q = \begin{pmatrix} u \\ d \\ s \end{pmatrix}$ is a "fundamental" of $SU(3)$.

antiquark $\bar{q} = \begin{pmatrix} \bar{u} \\ \bar{d} \\ \bar{s} \end{pmatrix}$ is an "antifundamental" $\bar{\mathbf{3}}$.

Diquark (distinguishable): $q_i^{(1)} q_j^{(2)} \rightarrow U_{ik} U_{jl} q_k^{(1)} q_l^{(2)}$

Product rep of two fundamentals: $3 \times 3 = 9$ dim rep \rightarrow

Symmetric rep: $q_i^{(1)} q_j^{(2)} + q_j^{(1)} q_i^{(2)}$

Antisym. rep: $q_i^{(1)} q_j^{(2)} - q_j^{(1)} q_i^{(2)}$

oo fundamentals: $3 \times 3 = 9$ dim. rep \rightarrow reducible.

$$g_i^{(1)} g_j^{(2)} + g_j^{(1)} g_i^{(2)} \rightarrow U_{ik} U_{jl} (g_k^{(1)} g_l^{(2)} + g_l^{(1)} g_k^{(2)})$$

$$g_i^{(1)} g_j^{(2)} - g_j^{(1)} g_i^{(2)} \rightarrow U_{ik} U_{jl} (g_k^{(1)} g_l^{(2)} - g_l^{(1)} g_k^{(2)})$$

available.

$$+ \begin{pmatrix} (1) & (2) \\ g & l & g & k \end{pmatrix}$$

6-dimensional rep (6)

$$\begin{pmatrix} (1) & (2) \\ g & l & g & k \end{pmatrix}$$

3-dimensional rep.

Compact notation for antisym. rep:

$$\epsilon_{ijk} g_j^{(1)} g_k^{(2)} \rightarrow \epsilon_{ijk} U_{jm} U_{kn} g_m^{(1)} g_n^{(2)} =$$

Use: $\epsilon_{ijk} U_{il} U_{jm} U_{kn} = \epsilon_{lmn}$

$$\Rightarrow \epsilon_{ijk} U_{jm} U_{kn} = U_{il}^* \epsilon_{lmn}$$

rep:

$$\epsilon_{ijk} U_{jm} U_{kn} g_m^{(1)} g_n^{(2)} = U_{il}^* \epsilon_{lmn} g_m^{(1)} g_n^{(2)}$$

$$)_{kn} = \epsilon_{lmn}$$

$$)_{kn} = U_{il}^* \epsilon_{lmn}$$

Antisym. rep transforms as $\bar{3}$

Compact notation for antisym. rep:

$$\epsilon_{ijk} q_j^{(1)} q_k^{(2)} \rightarrow \epsilon_{ijk} U_{jm} U_{kn} q_m^{(1)} q_n^{(2)} = U_{il}^*$$

Use: $\epsilon_{ijk} U_{il} U_{jm} U_{kn} = \epsilon_{lmn}$

Antisym

$$\Rightarrow \epsilon_{ijk} U_{jm} U_{kn} = U_{il}^* \epsilon_{lmn}$$

Diquark: $3 \times 3 = 6 + \bar{3}$

Young Tableaux

\boxed{i} = fundamental rep with index i

\boxed{j} = — " ————— " ————— j

$\begin{array}{|c|} \hline i \\ \hline j \\ \hline \end{array}$ = i, j antisymmetrized

$\boxed{i | j}$ = i, j Symmetrized

Put boxes together following rules:

1. Length of rows of boxes does not increase from top to bottom
2. length of columns does not increase from left to right.
3. No more than 3 boxes in a column for $SU(3)$ (N boxes)

Diquark: $\square \times \square = \begin{array}{|c|} \hline \square \\ \hline \square \\ \hline \end{array} + \begin{array}{|c|c|} \hline \square & \square \\ \hline \end{array}$

rules:

boxes does not increase from top to bottom.

does not increase from left to right.

3 boxes in a column for $SU(3)$ (N boxes for $SU(N)$)



Count dimension of rep. Number of ways to put indices $i_1, i_2, \dots = 1, 2, \dots$

1. indices in a row cannot decrease
2. indices in a column must increase.

$$\square = \boxed{1}, \boxed{2}, \boxed{3} \quad 3 \text{ dim.}$$

$$\begin{array}{|c|} \hline \square \\ \hline \square \\ \hline \end{array} =$$

Count dimension of rep. Number of ways to put indices $i_1, \dots, i_n = 1, 2, 3$

1. indices in a row cannot decrease
2. indices in a column must increase.

$$\square = \boxed{1}, \boxed{2}, \boxed{3} \quad 3 \text{ dim.}$$

$$\begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline \end{array}, \begin{array}{|c|} \hline 1 \\ \hline 3 \\ \hline \end{array}, \begin{array}{|c|} \hline 2 \\ \hline 3 \\ \hline \end{array}$$

s to put indices $i, j, \dots = 1, 2, 3$ into each box.

case

case.

$$\square = \begin{array}{|c|} \hline \square \\ \hline \end{array} = \begin{array}{|c|c|} \hline 1 & 1 \\ \hline \end{array}, \begin{array}{|c|c|} \hline 1 & 2 \\ \hline \end{array}, \begin{array}{|c|c|} \hline 1 & 3 \\ \hline \end{array}, \begin{array}{|c|c|} \hline 2 & 2 \\ \hline \end{array}, \begin{array}{|c|c|} \hline 2 & 3 \\ \hline \end{array}, \begin{array}{|c|c|} \hline 3 & 3 \\ \hline \end{array} = 6 \text{ dim.}$$

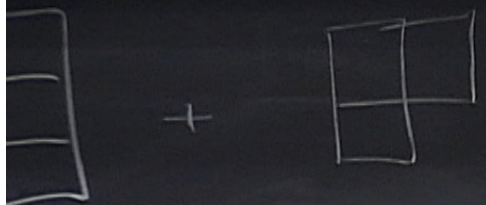
Meson: $q\bar{q}$

$$\square \times \begin{array}{|c|} \hline \square \\ \hline \square \\ \hline \end{array} = \begin{array}{|c|} \hline \square \\ \hline \square \\ \hline \square \\ \hline \end{array} + \begin{array}{|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array}$$

$3 \times \bar{3}$ $1 + 8$

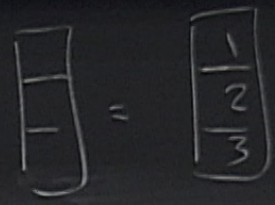
Singlet Octet

The diagram illustrates the decomposition of a meson, represented as a quark-antiquark pair ($q\bar{q}$), into its constituent parts. On the left, the quark is shown in a square, and the antiquark is in a vertical rectangle divided into three sections. An equals sign follows. To the right of the equals sign, a vertical rectangle divided into three sections represents the singlet state (labeled '1' and 'Singlet'), and a 2x2 grid represents the octet state (labeled '8' and 'Octet'). The numbers '3 x 3-bar' and '1 + 8' are positioned below the respective diagrams, indicating the mathematical representation of these states.

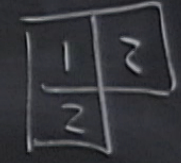
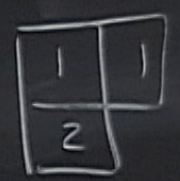


+

1 + 8
singlet + octet.



=



, etc.

Baryon: 888

$$\square \times \square \times \square = \square \times \left(\begin{array}{|c|} \hline \square \\ \hline \square \\ \hline \end{array} + \begin{array}{|c|c|} \hline \square & \square \\ \hline \end{array} \right) =$$
$$3 \times 3 \times 3 = 3 \times \left(\bar{3} + 6 \right) =$$

$$\begin{aligned} &+ \left(\begin{array}{|c|c|} \hline & \\ \hline \end{array} \right) \\ &+ 6 \end{aligned} = \begin{array}{|c|} \hline \\ \hline \\ \hline \end{array} + \begin{array}{|c|c|} \hline & \\ \hline & \\ \hline \end{array} + \begin{array}{|c|c|} \hline & \\ \hline & \\ \hline \end{array} + \begin{array}{|c|c|c|} \hline & & \\ \hline & & \\ \hline \end{array} \\ &= 1 + \underbrace{8 + 8}_{\text{octets}} + 10 \\ &\quad \text{singlet} \qquad \qquad \text{decuplet} \end{aligned}$$

Mesons: $q\bar{q}$ are distinguishable

Non-relativistic quark model: naive approx. that $q\bar{q}$ is hydro

Lowest lying states are $L=0$ spatial wavefunction. Spin

Apply Young tableaux to $su(2)$:

$$\begin{array}{c} \square \times \square \\ 2 \times 2 \end{array} = \begin{array}{c} \square \\ 1 \\ J=0 \end{array} + \begin{array}{c} \square \\ 3 \\ J=1 \end{array}$$

approx. that $q\bar{q}$ is hydrogen-like bound state.

radial wavefunction. Spin of meson due to spin of quarks.

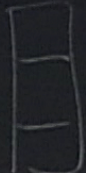
$$\square = \begin{array}{|c|} \hline \square \\ \hline \square \\ \hline \end{array} + \begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array}$$

$2 \qquad 1 \qquad + \qquad 3$
 $J=0 \qquad \qquad J=1$



always singlet under $SU(3)_{\text{color}}$.

$$(-1) \cdot P = (-1)^{L+1} = -1 \quad \text{for } L=0.$$


Color wavefunction (SU(3) color) \rightarrow  always singlet

$q\bar{q}$ has intrinsic parity (-1) . $P = (-1)^{L+1} = -$

Four possibilities:

pseudoscalar ($J^P = 0^-$) octet or singlet

vector ($J^P = 1^-$) octet or singlet

\rightarrow  always singlet under $SU(3)_{\text{color}}$.

(-1) . $P = (-1)^{L+1} = -1$ for $L=0$.

(5^-) octet or singlet (π, K, η, η')

(1^-) octet or singlet $(\rho, \omega, \phi, K^*)$

vector ($J = 1$) octet or singlet

Octet rep. (pseudoscalar)

$$M_{ij} \rightarrow U_{ik} U_{jl}^* M_{kl}$$

Write

$$M = M^A T^A$$

$$T^A = \frac{\lambda^A}{2} = \text{generators of } SU(3) \text{ flavor.}$$

$$= \frac{1}{2} \begin{pmatrix} M^3 + \frac{M^8}{\sqrt{3}} & M^1 - iM^2 & M^4 - iM^5 \\ M^1 + iM^2 & -M^3 + \frac{M^8}{\sqrt{3}} & M^6 - iM^7 \\ M^4 + iM^5 & M^6 + iM^7 & -\frac{2}{\sqrt{3}}M^8 \end{pmatrix} =$$

M_{kl}

$$M \rightarrow U M U^\dagger$$

generators of $SU(3)$ flavor.

$$\begin{pmatrix} M^5 \\ M^7 \\ M^8 \end{pmatrix} = \begin{pmatrix} \frac{\pi^0}{\sqrt{2}} + \frac{\eta}{\sqrt{6}} & \pi^+ & K^+ \\ \pi^- & -\frac{\pi^0}{\sqrt{2}} + \frac{\eta}{\sqrt{6}} & K^0 \\ K^- & K^0 & -\frac{2}{\sqrt{6}} \eta \end{pmatrix}$$

K^+

K^0

$-\frac{2}{\sqrt{6}}\eta$

$$\pi^{\pm} = \frac{M^1 \mp iM^2}{2}$$

etc.

Recall: $SU(2)$ has one diagonal generator T^3 . Label components
 For $SU(3)$, two diagonal generators T^3, T^8 . Label components

$$K^0: \left[T^3, \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & K^0 \\ 0 & 0 & 0 \end{pmatrix} \right] = -\frac{1}{2} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & K^0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\uparrow$$

$$\frac{1}{2} \begin{pmatrix} 1 & & \\ & -1 & \\ & & 0 \end{pmatrix}$$

Label components of rep. by eigenvalues under that generator.

Label components by eigenvalues.

$$\begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & K^0 \\ 0 & 0 & 0 \end{pmatrix},$$

$$[T^8, \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & K^0 \\ 0 & 0 & 0 \end{pmatrix}] = \frac{1}{\sqrt{3}} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & K^0 \\ 0 & 0 & 0 \end{pmatrix}$$

Where do isospin & strangeness fit in?

Isospin is $SU(2)$ subgroup of $SU(3)$ flavor

$$U = \begin{pmatrix} V & 0 \\ 0 & 1 \end{pmatrix}$$

$V = SU(2)$ isospin rotation.

3×3 $SU(3)$ flavor

Baryon: qqq

$$\square \times \square \times \square = \square \times (\square \oplus \square) = \square \oplus \square$$

$$MU^+ = \begin{pmatrix} V & 0 \\ 0 & 1 \end{pmatrix} M \begin{pmatrix} V^+ & 0 \\ 0 & 1 \end{pmatrix}$$

$$= \begin{pmatrix} V \begin{pmatrix} \frac{\pi^0}{\sqrt{2}} & \pi^+ \\ \pi^- & -\frac{\pi^0}{\sqrt{2}} \end{pmatrix} V^+ + \frac{\eta}{\sqrt{6}} & V \begin{pmatrix} K^+ \\ K^0 \end{pmatrix} \\ \hline (K^-, \bar{K}^0) V^+ & -\frac{2}{\sqrt{6}} \eta \end{pmatrix}$$

$$M: \quad M \rightarrow U M U^\dagger = \begin{pmatrix} V & 0 \\ 0 & 1 \end{pmatrix} M \begin{pmatrix} V^\dagger & 0 \\ 0 & 1 \end{pmatrix}$$

$$= \begin{pmatrix} V \begin{pmatrix} \frac{\pi_0}{\sqrt{2}} & \pi^+ \\ \pi^- & -\frac{\pi_0}{\sqrt{2}} \end{pmatrix} V^\dagger + \frac{3}{\sqrt{6}} \\ \dots \\ (K^-, K^0) V^\dagger \end{pmatrix}$$

$$V \begin{pmatrix} K^+ \\ K^0 \\ -\frac{2}{\sqrt{6}} \eta \end{pmatrix}$$

$\pi^\pm, \pi^0 =$ isospin triplet
(adjoint rep of $SU(2)$)

$\begin{pmatrix} K^+ \\ K^0 \end{pmatrix}, \begin{pmatrix} \bar{K}^0 \\ K^- \end{pmatrix} =$ isospin doublet

$\eta =$ isospin singlet.

$I_3 = T^3$ isospin 3-component.

Baryon number is additional conserved quantum number.

In $SU(3)$ flavor, $B = \frac{1}{3} \begin{pmatrix} 1 & & \\ & 1 & \\ & & 1 \end{pmatrix}$

Strangeness: $S = \frac{1}{\sqrt{3}} T^8 - B$

Baryon number is additional conserved quantum number.

In $SU(3)$ flavor, $B = \frac{1}{3} \begin{pmatrix} 1 & & \\ & 1 & \\ & & 1 \end{pmatrix}$

Strangeness: $S = \frac{1}{\sqrt{3}} T^8 - B$ ($B=0$ for mesons)

<u>States</u>	<u>I_3</u>	<u>I</u>	<u>T^8</u>	<u>S</u>
π^+	+1	1	0	0
π^0	0	1	0	0
π^-	-1	1	0	0
η, η'	0	0	0	0
K^+	$1/2$	$1/2$	$\sqrt{3}$	1
K^0	$-1/2$	$1/2$	$\sqrt{3}$	1
\bar{K}^0	$1/2$	$1/2$	$-\sqrt{3}$	-1
K^-	$-1/2$	$1/2$	$-\sqrt{3}$	-1

mesons)

Lagrangian for pseudoscalars:

$$\mathcal{L} = \frac{1}{2} \text{Tr} [\partial_\mu M \partial^\mu M] - \frac{1}{2} M_8^2 \text{Tr}[M^2] + \frac{1}{2} \partial_\mu \eta^i$$

Kinetic terms are in canonical form:

$$m_\pi^2 = m_K^2 = m_\eta^2 = M_8^2$$

$$m_{\eta'}^2 = M_1^2$$

vars:

$$M] - \frac{1}{2} M_8^2 \text{Tr}[M^2] + \frac{1}{2} \partial_\mu \eta^I \partial^\mu \eta^I - \frac{1}{2} M_1^2 \eta^{I2} + \mathcal{L}_{int.}$$

form:

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