

Title: PSI 2016/2017 Standard Model (Review) - Lecture 9

Date: Jan 13, 2017 09:00 AM

URL: <http://pirsa.org/17010013>

Abstract:

Higgs decays & production

$$m_h \approx 125 \text{ GeV}$$

Largest coupling (largest mass): t, W, Z

$h \rightarrow t\bar{t}, W^+W^-, ZZ$ are kinematically forbidden.

Higgs decays & production

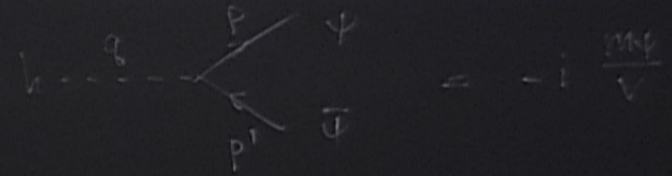
$$M_h \approx 125 \text{ GeV}$$

Largest coupling (largest mass): t, W^\pm, τ

$h \rightarrow t\bar{t}, W^+W^-, \tau\tau$ are kinematically forbidden.

Tree-level decays:

$h \rightarrow \text{fermions}$ $\mathcal{L}_{\text{int}} = - \frac{m_\psi}{v} h \bar{\Psi} \Psi$



$i\mathcal{M} = -i \frac{m_\psi}{v} \bar{u}(p) v(p')$ $2m_\psi \ll m_h$

$\sum_{\text{spins}} |\mathcal{M}|^2 = \frac{m_\psi^2}{v^2} 4 p \cdot p' = \frac{2m_\psi^2 m_h^2}{v^2}$ $q^2 = m_h^2 = 2p \cdot p'$

$\Gamma(h \rightarrow \psi \bar{\psi}) = \frac{1}{16\pi m_h} \sum |\mathcal{M}|^2 = \frac{m_\psi^2 m_h}{8\pi v^2}$ (after doing 2-body)

$\Gamma(h \rightarrow \tau \dots)$

$$\Gamma(h \rightarrow \psi \bar{\psi}) = \frac{1}{16\pi m_h} \sum |g_{h\psi\bar{\psi}}|^2 = \frac{m_\psi^2 m_h}{8\pi v^2}$$

$$\Gamma(h \rightarrow \tau \bar{\tau}) = \frac{m_\tau^2 m_h}{8\pi v^2} \approx 0.26 \text{ MeV}$$

$$\Gamma(h \rightarrow b \bar{b}) = \frac{m_b^2 m_h}{8\pi v^2}$$

(after doing 2-body integral w/ $m_4=0$)

$$\sum_{\text{spins}} |M|^2 = \frac{1}{v^2} 4 p \cdot p' = \frac{2 m_\psi m_h}{v^2}$$

$$g^2 = m_i^2 = v p \cdot p'$$

$$\Gamma(h \rightarrow \psi \bar{\psi}) = \frac{1}{16\pi m_h} \sum |M|^2 = \frac{m_\psi^2 m_h}{8\pi v^2} \quad (\text{after doing 2-body integral of})$$

$$\Gamma(h \rightarrow \tau \bar{\tau}) = \frac{m_\tau^2 m_h}{8\pi v^2} \approx 0.26 \text{ MeV} \quad (6.3\%)$$

$$\Gamma(h \rightarrow b \bar{b}) = \frac{3 m_b^2 m_h}{8\pi v^2} \approx 4.3 \text{ MeV} \times (58\%) \quad \times \text{RGE of } m_b, m_c \text{ (or } y_b, y_c)$$

reduces m_b, m_c by factor of ~ 2 (or ~ 1.5)

$$\Gamma(h \rightarrow c \bar{c}) = \frac{3 m_c^2 m_h}{8\pi v^2} \approx 0.4 \text{ MeV} \times (2.9\%)$$

$$\Gamma(h \rightarrow \psi\bar{\psi}) = \frac{1}{16\pi m_h} \sum |g_{h\psi\bar{\psi}}|^2 = \frac{m_\psi^2 m_h}{8\pi v^2}$$

$$g_{h\psi\bar{\psi}} = 0.1 \psi$$

$$\Gamma(h \rightarrow \psi\bar{\psi}) = \frac{1}{16\pi m_h} \sum |g_{h\psi\bar{\psi}}|^2 = \frac{m_\psi^2 m_h}{8\pi v^2} \quad (\text{after doing 2-body integral w/ } m_\psi \ll m_h)$$

$$\Gamma(h \rightarrow \tau\bar{\tau}) = \frac{m_\tau^2 m_h}{8\pi v^2} \approx 0.26 \text{ MeV} \quad (6.3\%)$$

$$\Gamma(h \rightarrow b\bar{b}) = \frac{3m_b^2 m_h}{8\pi v^2} \approx 4.3 \text{ MeV} \quad (55\%) \quad \times \text{RGE of } m_b, m_c \text{ (or } g_b, g_c)$$

reduces m_b, m_c by factor of ~ 2 (~ 4 GeV)

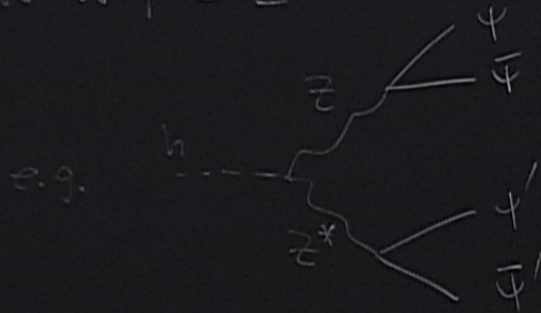
$$\Gamma(h \rightarrow c\bar{c}) = \frac{3m_c^2 m_h}{8\pi v^2} \approx 0.4 \text{ MeV} \quad (2.7\%)$$

• $h \rightarrow W^+W^-, Z\bar{Z}$

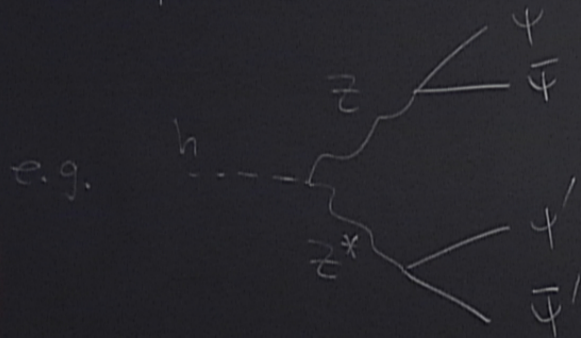
(allowed if one/both vector bosons off-shell)

"golden mode"

$h \rightarrow Z\bar{Z}^* \rightarrow 4 \text{ leptons}$
 leptons = e^+e^- or $\mu^+\mu^-$



• $h \rightarrow W^+W^-, Z Z$ (allowed if one/both vector bosons off-shell)



"golden mode" $h \rightarrow Z Z^* \rightarrow 4 \text{ leptons}$
 leptons = e^+e^- or $\mu^+\mu^-$

$$\Gamma(h \rightarrow Z Z^*) \approx 0.11 \text{ MeV} \quad (2.6\%)$$

$$\Gamma(h \rightarrow W W^*) \approx 0.88 \text{ MeV} \quad (22\%)$$

Loop decays: very important

$$\Gamma(h \rightarrow gg) \approx 0.35 \text{ MeV} \quad (8.8\%)$$

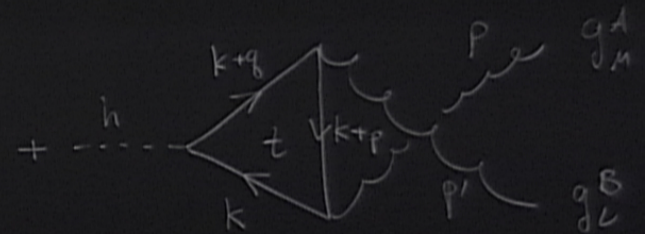
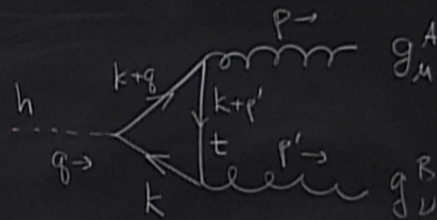
$$\Gamma(h \rightarrow \gamma\gamma) \approx 6 \times 10^{-3} \text{ MeV} \quad (0.15\%)$$

$h \rightarrow \gamma\gamma$ is rare, but experimentally clean (other "golden mode")
 $h \rightarrow gg$ easier to compute (difficult to measure)

$$\sum_{\text{polarizations}} |M|^2 = \frac{m_4}{v^2} 4 p \cdot p' = \frac{2 m_4 m_h^2}{v^2} \quad g^2 = m_h^2 = 2 p \cdot p'$$

Compute $h \rightarrow gg$:

$$iM = iM_1 + iM_2 =$$



$$i\mathcal{M}_1 = \int \frac{d^d k}{(2\pi)^d} (-1) \text{Tr} \left[\frac{i(k+m_t)}{k^2 - m_t^2} i g_s T^B \gamma^\nu \frac{i(k+p'+m_t)}{(k+p')^2 - m_t^2} i g_s T^A \gamma^\mu \frac{i(k+q+m_t)}{(k+q)^2 - m_t^2} \left(-\frac{i m_t}{v}\right) \right] \epsilon_\mu \epsilon_\nu$$

$$= -g_s^2 \text{Tr} [T^A T^B] \left(\frac{m_t}{v}\right) \epsilon_\mu(p) \epsilon_\nu(p') \int \frac{d^d k}{(2\pi)^d} \frac{\text{Tr} \left[(k+m_t) \gamma^\nu (k+p'+m_t) \gamma^\mu (k+q+m_t) \right]}{(k^2 - m_t^2)((k+p')^2 - m_t^2)((k+q)^2 - m_t^2)}$$

$$T^A = \frac{\lambda^A}{2}$$

$$\sum |\mathcal{M}|^2 = \frac{m_4^2 m_h}{8\pi v^2} \quad (\text{after doing 2-body integral w/ } m_4 \rightarrow 0)$$

$$\approx 0.26 \text{ MeV} \quad (6.3\%)$$

$$\frac{g_h}{2} \approx 4.3 \text{ MeV}^* \quad (5.8\%) \quad \times \text{RGE of } m_b, m_c \text{ (or } y_b, y_c)$$

reduces m_b, m_c by factor of ~ 2 (~ 4)

$$\frac{2 m_h}{3} \approx 0.4 \text{ MeV}^* \quad (2.9\%)$$

Use Feynman parameters:

$$i\mathcal{M}_1 = -g_s^2 \text{Tr}[T^A T^B] \frac{m_t}{v} \epsilon_\mu \epsilon_\nu$$

$$\times \int dx dy dz \, 2 \delta(1-x-y-z) \frac{\text{Tr}[\dots]}{\left[(k^2 - m_t^2)x + ((k+p')^2 - m_t^2)y + ((k+q)^2 - m_t^2)z \right]^3}$$

Use Feynman parameters:

$$i\mathcal{M}_i = -g_s^2 \text{Tr}[T^A T^B] \frac{m_t}{V} \epsilon_\mu \epsilon_\nu \int \frac{d^d k}{(2\pi)^d} \int dx dy dz \, 2 \delta(1-x-y-z) \frac{\text{Tr}[\dots]}{[(k^2 - m_t^2)x + (k^2 - m_t^2)y + (k^2 - m_t^2)z]}$$

$$= \underbrace{(k + y p' + z q)^2}_{l^2} - (y p' + z q)^2 + m_h^2 z - m_t^2$$

$$= l^2 - \underbrace{\left(z^2 m_h^2 - m_h^2 z + m_t^2 + 2 y z p' \cdot q \right)}_{M^2}$$

$$p' \cdot q = p' \cdot (p + p') = p \cdot p' = \frac{1}{2} q_0^2 = \frac{1}{2} m_h^2$$

$$M^2 = m_t^2 + (z^2 - z) m_h^2 + yz m_h^2$$

$$= m_t^2 - m_h^2 \underbrace{(1 - y - z)}_x z = m_t^2 - zx m_h^2$$

$$l = k + y p' + z q = k + (1 - x) p' + z p$$

$$= -g_s^2 \text{Tr}[T^A T^B] \left(\frac{m_t}{v}\right) \bar{E}_\mu(p) \bar{E}_\nu(p') \int \frac{d^d k}{(2\pi)^d} \frac{\text{Tr}[(\not{k} + m_t) \gamma^\mu (\not{k} + \not{p}' + m_t) \gamma^\nu (\not{k} + \not{q} + m_t)]}{(k^2 - m_t^2)((k+p')^2 - m_t^2)((k+q)^2 - m_t^2)}$$

Numerator

$$\text{Tr}[\dots] = 16 m_t l^\mu l^\nu - 4 l^2 m_t \gamma^{\mu\nu} + 4 m_t^3 \gamma^{\mu\nu} - 2 m_t m_h^2 (1-2xz) \gamma^{\mu\nu} + 4(1-4xz) p_1^\nu p_2^\mu m_t$$

discarded terms $\sim p^\mu$ or p'^ν
and terms linear in l .

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In d -dimensions: $l^4 l^\mu \rightarrow \frac{1}{d} \eta^{\mu\nu} l^2$ under the integral.

$$\eta^{\mu\nu} l_\mu l_\nu = l^2 = \frac{1}{d} l^\mu l^\mu$$

$$i\mathcal{M}_1 = -g^2 \text{Tr}[T^A T^B] \frac{m_t^2}{v} E_\mu E_\nu$$

$$\times \int \frac{d^d l}{(2\pi)^d} \int dx dy dz 2\delta(1-x-y-z) \frac{1}{(l^2 - M^2)^3}$$

$$\times \left(\left(\frac{16}{d} - 4 \right) l^2 \eta^{\mu\nu} + 2 \left(2m_t^2 - m_b^2 (1-2xz) \right) \eta^{\mu\nu} + 4(1-4xz) p_1^\nu p_2^\mu \right)$$

Momentum integrals:

$$(1) = \int \frac{d^d l}{(2\pi)^d} \frac{l^2}{(l^2 - M^2)^3} = \frac{-i}{(4\pi)^{d/2}} \frac{d}{2} \frac{\Gamma(3 - \frac{d}{2} - 1)}{\Gamma(3)} \left(\frac{1}{M^2}\right)^{3 - \frac{d}{2} - 1}$$

$$(2) = \int \frac{d^d l}{(2\pi)^d} \frac{1}{(l^2 - M^2)^3} = \frac{i}{(4\pi)^{d/2}} \frac{\Gamma(3 - \frac{d}{2})}{\Gamma(3)} \left(\frac{1}{M^2}\right)^{3 - d/2}$$

$$\begin{aligned}
i\mathcal{M}_1 = & -\frac{i}{16\pi^2} g_s^2 \text{Tr}[T^A T^B] \frac{m_t^2}{v} \epsilon_\mu \epsilon_\nu \int_0^1 dx \int_0^{1-x} dz \\
& \times \left\{ -4 \left(\frac{4-d}{d} \right) \frac{2}{4-d} 2 (m_t^2 - xz m_h^2) \eta^{\mu\nu} \right. \\
& \quad \left. + (4m_t^2 - 2m_h^2 + 4xz m_h^2) \eta^{\mu\nu} \right. \\
& \quad \left. + 4(1-4xz) p_1^\nu p_2^\mu \right\} \frac{1}{M^2}
\end{aligned}$$

$$+ 4(1-4xz) p_1^{\nu} p_2^{\mu} \int \frac{1}{M^2}$$

$$= - \frac{i}{16\pi^2} g_s^2 \text{Tr}[T^A T^B] \frac{m_t^2}{V} \sum_u(p) E_b(p') \int_0^1 dx \int_0^{1-x} dz (4 p_1^{\nu} p_2^{\mu} - 2 m_h^2 \eta^{\mu\nu}) \frac{1-4xz}{m_t^2 - xz m_h^2}$$