

Title: PSI 2016/2017 Condensed Matter - Lecture 8

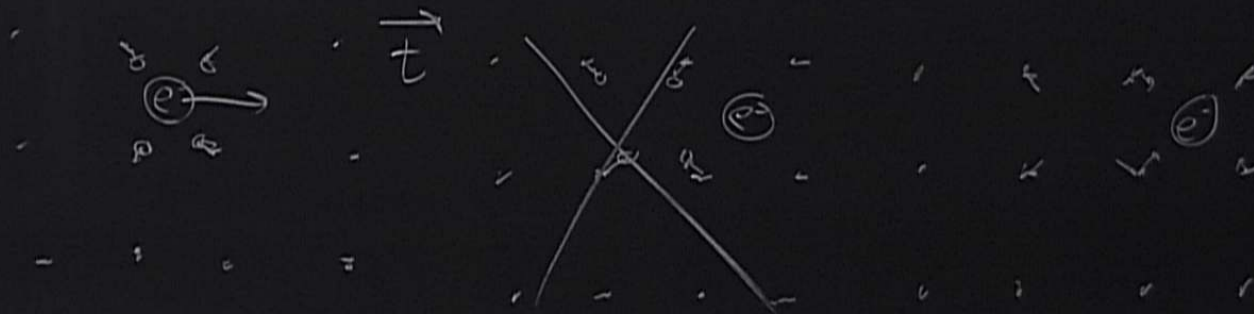
Date: Nov 16, 2016 10:45 AM

URL: <http://pirsa.org/16110059>

Abstract:

polaron : electron interacting with phonons

when an electron travels through a lattice,
it distorts the lattice as it travels



tigh
H_{ee}
H_{pp}

tight-binding

$$H_{ee} = -\sum_{\langle ij \rangle} t_{ij} (c_{i\sigma}^\dagger c_{j\sigma} + h.c.)$$

$$H_{ph} = \sum_i \frac{p_i^2}{2M} + \frac{1}{2} \sum_{\langle ij \rangle} K_{ij} (u_i - u_j)^2$$

$$H_{ee}^{(0)} = -\sum_{\langle ij \rangle} t_{ij}^{(0)} (c_{i\sigma}^\dagger c_{j\sigma} + h.c.)$$

$$H_{ed-ph} = \left[-\sum_{\langle ij \rangle} \frac{\partial t}{\partial r} (\bar{u}_i - \bar{u}_j) (c_{i\sigma}^\dagger c_{j\sigma} + h.c.) \right]$$

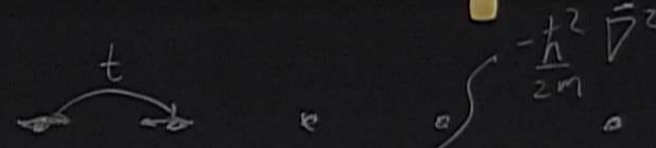


Diagram illustrating hopping t between sites i and j , and a potential well with depth $-\frac{\hbar^2}{2m} |\nabla^2|$.

$$t_{ij} = \int d^3r \Phi_i^*(r) h(r) \Phi_j(r)$$

$$= t(\bar{r}_i - \bar{r}_j) = t(\bar{r}_i + \bar{u}_i - \bar{r}_j - \bar{u}_j)$$

$$= \underbrace{t(\bar{r}_i - \bar{r}_j)}_{t_{ij}^{(0)}} + \nabla t(\bar{u}_i - \bar{u}_j)$$

$\text{with } \bar{r} = \bar{r}_i - \bar{r}_j$

$$H_{el-ph} = - \sum_{\langle i,j \rangle} \frac{\partial t}{\partial r} (\bar{u}_i - \bar{u}_j) (c_{i\sigma}^\dagger c_{j\sigma} + c_{i\sigma} c_{j\sigma}^\dagger)$$

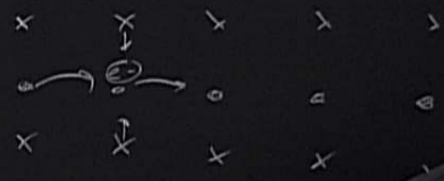
SSH model.

↖ Su, Schrieffer + Heeger
PRL 1979
polyacetylene

BLF-SSH model

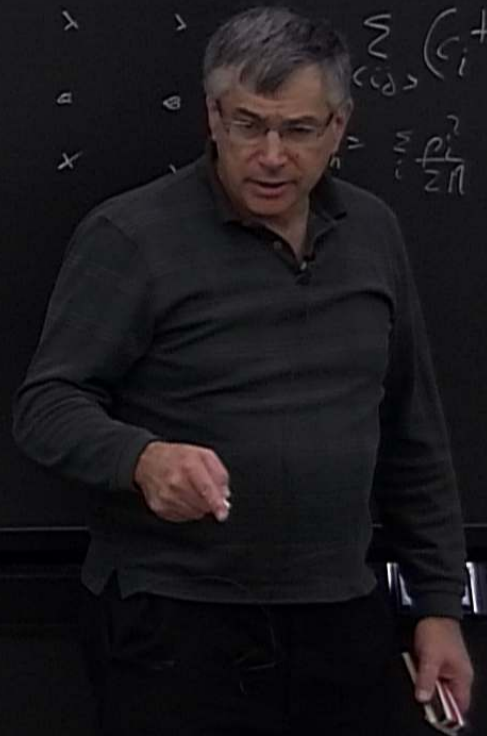
↖ Borjesson-Labbe-Friedel
PRL 1970

1959 Molecular Crystal Model (Holstein)

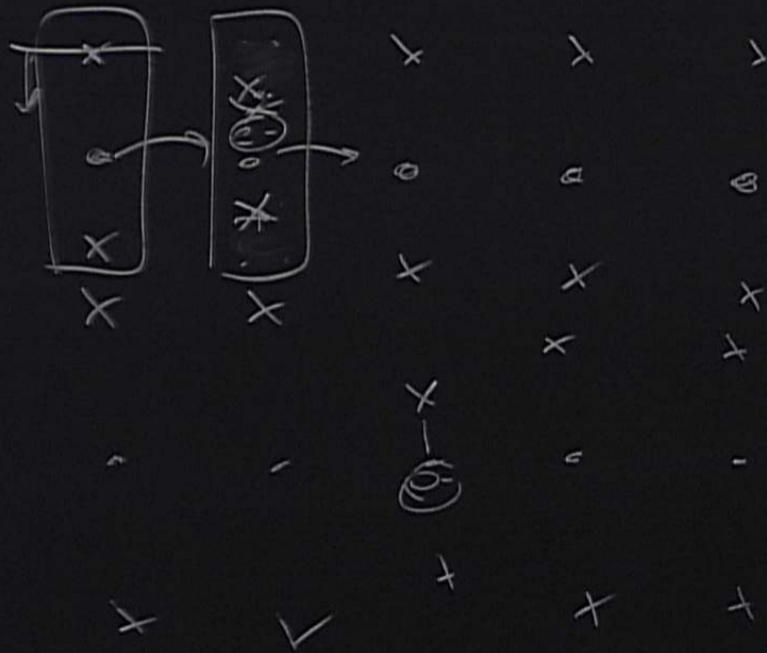


$$\sum_{\langle i,j \rangle} (c_i^\dagger c_j + c_j^\dagger c_i)$$

$$= \sum_i \frac{p_i^2}{2m} + \frac{1}{2} K \sum_i x_i^2$$



1959 Molecular Crystal Model (Holstein)



$$-t \sum_{\langle ij \rangle} (c_i^\dagger c_j + c_j^\dagger c_i)$$

$$H_{ph} = \sum_i \frac{p_i^2}{2M} + \frac{1}{2} K \sum_i u_i^2$$

$$H_{elph} = \alpha \sum_i u_i n_i$$

\uparrow
 $c_i^\dagger c_i$

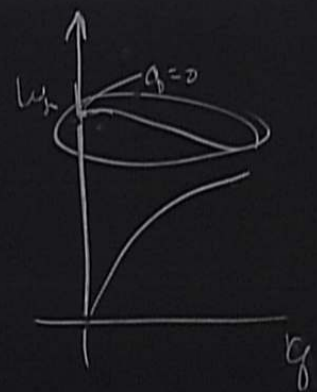
(stein)

$$\sum_{\langle ij \rangle} (c_i^\dagger c_j + c_j^\dagger c_i)$$

$$= \sum_i \frac{p_i^2}{2m} + \frac{1}{2} K \sum_i u_i^2$$

$$= \propto \sum_i u_i n_i$$

\uparrow
 $c_i^\dagger c_i$



Fröhlich model (1950)

Weak-coupling perturbation.

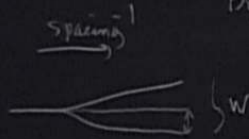
$$U_i = \sqrt{\frac{\hbar}{2M\omega_E}} (a_i + a_i^\dagger)$$

$$g \equiv \propto \sqrt{\frac{\hbar}{2M\omega_E}}$$

$$P_i = -i \sqrt{\frac{\hbar M \omega_E}{2}} (a_i - a_i^\dagger)$$

$$H = -t \sum_{\langle ij \rangle} (c_i^\dagger c_j + \text{h.c.}) + \omega_E \sum_i a_i^\dagger a_i + g \omega_E \sum_i c_i^\dagger c_i (a_i + a_i^\dagger)$$

$\lambda \equiv \frac{g^2 W t}{\frac{W}{2}}$ ← energy gained if the electron localizes
 where $W = 2zt \equiv$ electron bandwidth
 $\frac{W}{2}$ ← energy gained by free electrons
 z ← coordination #
 cubic systems
 $z = 2$ in 1D
 $z = 4$ in 2D
 $z = 6$ in 3D (sc)



"site" Hamiltonian

$$\begin{aligned}
 & \frac{1}{2} K u_i^2 + \alpha u_i n_i \\
 & = \frac{1}{2} K \left(u_i + \frac{\alpha n_i}{K} \right)^2 - \left(\frac{\alpha^2}{2K} \right) n_i^2 \\
 & \tilde{u}_i = u_i + \frac{\alpha n_i}{K} \quad n_i = n_{i\uparrow} + n_{i\downarrow} \\
 & (n_{i\uparrow} + n_{i\downarrow})^2 = n_{i\uparrow} + n_{i\downarrow} + 2n_{i\uparrow}n_{i\downarrow}
 \end{aligned}$$

$$\frac{\alpha^2}{2K} = \frac{\hbar^2 g^2 \partial^2 \mu_{i\uparrow} \mu_{i\downarrow}}{2M W E^2} = g^2 W t$$

$$H = \underbrace{\sum_{\mathbf{k}} \epsilon_{\mathbf{k}} c_{\mathbf{k}}^{\dagger} c_{\mathbf{k}}}_{\hat{H}_0} + W_E \sum_{\mathbf{q}} a_{\mathbf{q}}^{\dagger} a_{\mathbf{q}} + \underbrace{\frac{g W_E}{\sqrt{N}} \sum_{\mathbf{k}, \mathbf{q}} c_{\mathbf{k}+\mathbf{q}}^{\dagger} c_{\mathbf{k}} (a_{\mathbf{q}} + a_{-\mathbf{q}}^{\dagger})}_{H_{\text{int.}}} \quad \text{2nd order}$$

$$\Delta E_p^{(2)} = \sum_{\eta} \frac{|\langle \eta | H_{\text{int.}} | \phi_p^{(0)} \rangle|^2}{E_p^{(0)} - E_{\eta}}$$

$$|\phi_p^{(0)}\rangle = c_p^{\dagger} |0\rangle \otimes |0\rangle$$

$$\Delta E^{(1)} = 0 = \langle \phi_p^{(0)} | H_{\text{int.}} | \phi_p^{(0)} \rangle$$

$\eta \equiv$ excited states

$$c_{\mathbf{k}}^{\dagger} |0\rangle \otimes |n_{\mathbf{q}_1}\rangle |n_{\mathbf{q}_2}\rangle |n_{\mathbf{q}_3}\rangle \dots |n_{\mathbf{q}_n}\rangle$$

only need $c_p^{\dagger} a_{\mathbf{q}}^{\dagger} |0\rangle$

$$\langle 0 | c_k a_q - \frac{q \omega_E}{\sqrt{N}} \sum_{k'} c_{k+q}^\dagger c_k (a_{k'}^\dagger + a_{-q}^\dagger) c_{p'} | 0 \rangle \int_{-\frac{\pi}{2a}}^{\frac{\pi}{2a}} \frac{dk'}{2\pi}$$

$$q = -q'$$

$$R = p$$

$$k' = k + q$$

$$k' = p - q'$$

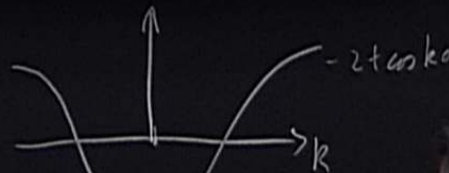
$$q' = p - k'$$

$$\Delta E_p^{(2)} = \frac{q^2 \omega_E^2}{N} \sum_{k'} \frac{1}{\epsilon_p - (\epsilon_{k'} + \omega_{p-k'})}$$

$$= - \frac{q^2 \omega_E}{2t} 2t \omega_E \frac{\partial(\omega + \epsilon_0 - \epsilon_p)}{\sqrt{(\epsilon_p - \omega_E)^2 - \epsilon_0^2}}$$

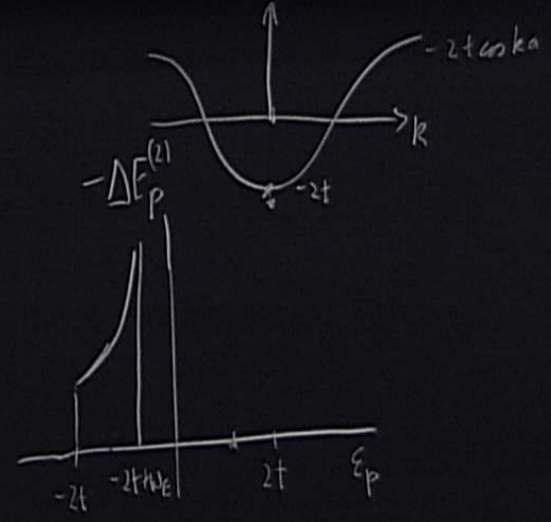
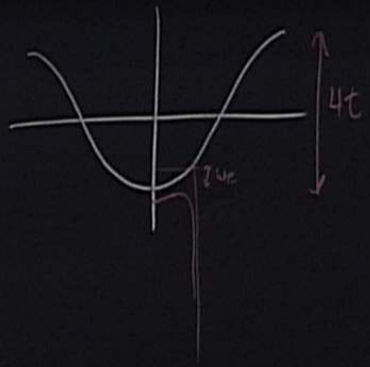
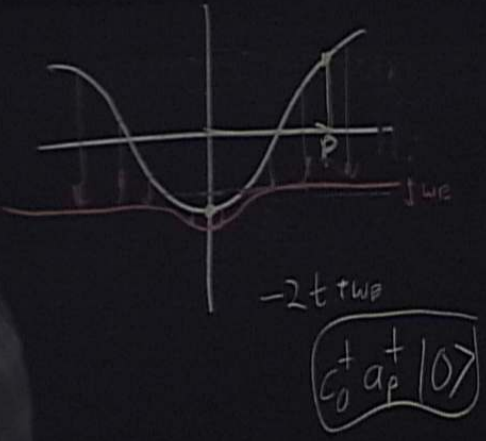
$$= \lambda$$

$$= -2\lambda t \omega_E \frac{\partial(\omega + \epsilon_0 - \epsilon_p)}{\sqrt{(\epsilon_p - \omega_E)^2 - \epsilon_0^2}}$$



$$q_j^1 = p \cdot k = \lambda \left(\frac{c_p \omega_E}{\epsilon_p} \right)$$

$$= -2\lambda t \omega_E \frac{\Theta(\omega + \epsilon_0 - \epsilon_p)}{\sqrt{(\epsilon_p - \omega_E)^2 - \epsilon_0^2}} \quad \epsilon_p = \epsilon_0 + \omega_E$$



$$\Delta E^{(1)} = 0 \quad \langle \phi_p^{(m)} | H_{\text{pert}} | \phi_p^{(0)} \rangle$$

only need $\langle \phi_p^{(0)} | a_p^\dagger | 0 \rangle$

$$m_{\text{band}}: -2t \cos ka \sim -2t \left(1 - \frac{1}{2} (ka)^2 \right) = -2t + (ka)^2 \quad p=0$$

$$\frac{\hbar^2 k^2}{2m} = a^2 \hbar^2$$

$$m_{\text{band}} = \frac{1}{2ta^2}$$

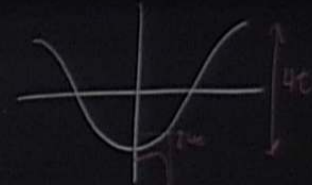
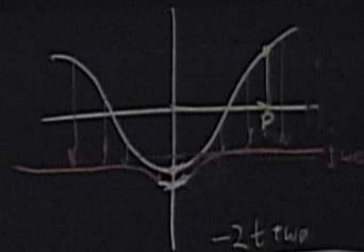
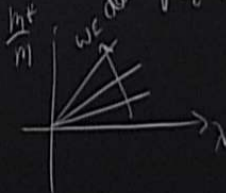
$$\Delta E_{p=0}^{(2)} = -2t \lambda \sqrt{\frac{W_E}{W_E + 4t}}$$

$$\frac{\hbar^2 E_p^{(2)}}{2p^2} = \frac{1}{m^*}$$

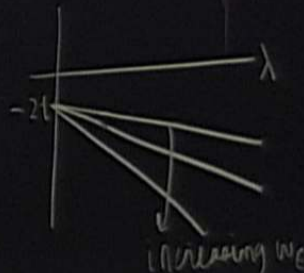
$$E_p^{(2)} \approx \frac{\hbar^2 k^2}{2m^*}$$

$$E^{(2)} \approx -2t - 2t \lambda \sqrt{\frac{W_E}{W_E + 4t}}$$

$$\frac{m^*}{m_{\text{band}}} = 1 + \frac{1}{2} \sqrt{\frac{t}{W_E}} \lambda \frac{1 + \frac{W_E}{2t}}{\left(1 + \frac{W_E}{4t} \right)^{3/2}}$$



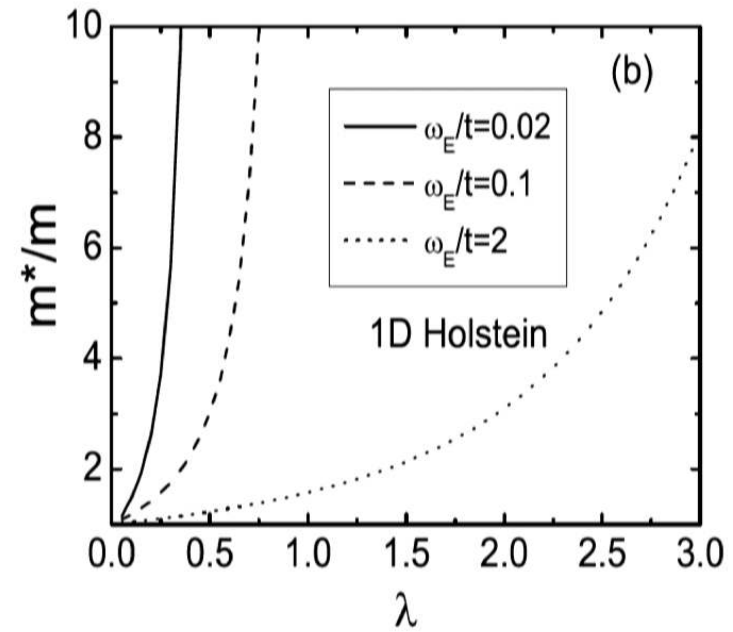
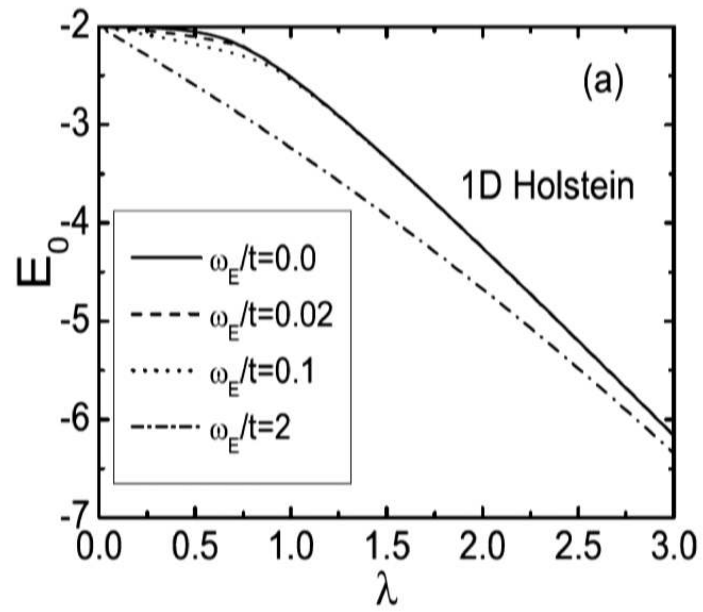
$$-2t \frac{W_E}{W_E + 4t}$$



$$|\Phi_p^{(1)}\rangle = \sqrt{z_0} \left\{ c_p^\dagger |0\rangle + \sqrt{\frac{2\lambda w_c t}{N}} \sum_k \frac{1}{\epsilon_p - (\epsilon_k + w_c)} c_k^\dagger a_{p-k}^\dagger |0\rangle \right\}$$

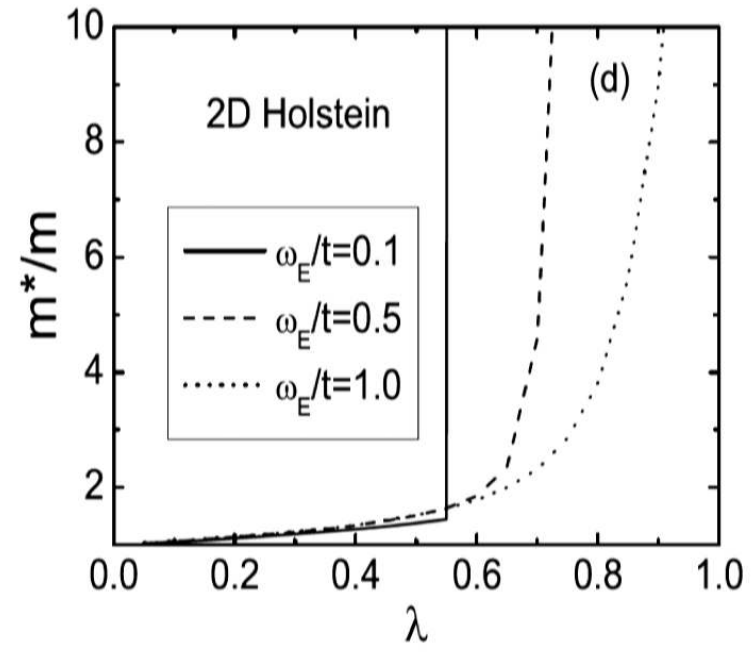
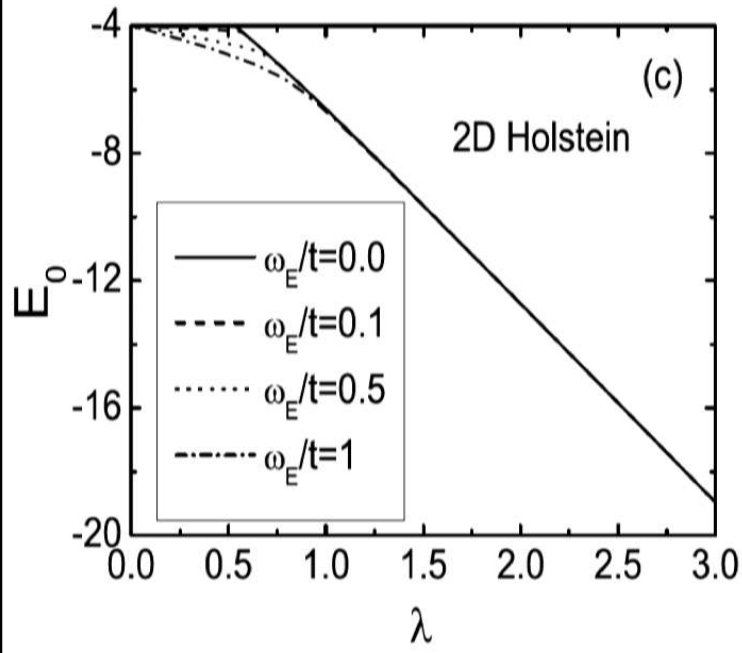
P.W. Anderson, Phys. Today Oct. 1997 p 42
"When the electron falls apart"

Polarons: 1D



J Supercond Nov Magn (2012) 25:1313–1317
DOI 10.1007/s10948-012-1601-6

Polarons: 2D



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DOI 10.1007/s10948-012-1601-6

