

Title: Gravitational waves via Kerr/CFT

Date: Nov 17, 2016 01:00 PM

URL: <http://pirsa.org/16110029>

Abstract: <p>Massive objects orbiting a near-extreme Kerr black hole plunge into the horizon after passing the innermost stable circular orbit, producing a potentially observable signal of gravitational radiation. The near horizon dynamics of such rapidly rotating black holes is governed by a conformal symmetry. In the talk I will show how this symmetry can be exploited to analytically compute the gravitational waves produced by a variety of orbits. I will also discuss an application to gravitational self-force and comment on the holographic interpretation of the process.</p>

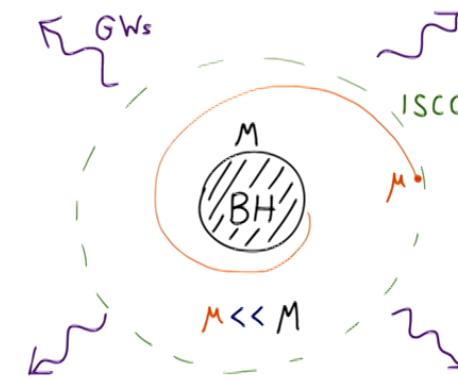
# Motivation

- Enhanced symmetry:  
Kerr/CFT

Guica, Hartman, Song, Strominger '08



- GWs in EMRI & plunge



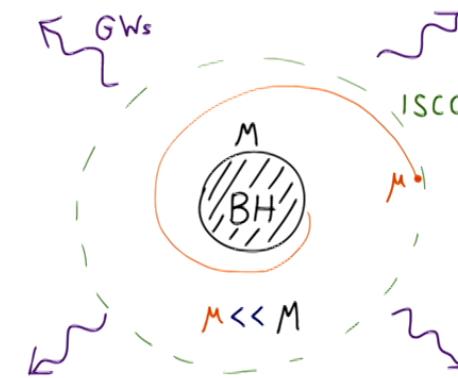
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- **Applied holography for astrophysics**

# Kerr/CFT

- Conjecture ("strong Kerr/CFT"): Kerr BH is holographically dual to a (1+1)D CFT.
- Tests: BH entropy, scattering off Kerr.
- Problems: GHSS boundary conditions, CFT details unknown.
- Fact ("weak Kerr/CFT"): All physics in (near-extreme) Kerr constrained by conformal symmetry.

# Gravity Waves

- From idea to observed signal in 100 years
- New era of experimental BH physics: "**spectroscopy**"
- Our analysis: relevant for **eLISA**



LIGO: stellar mass

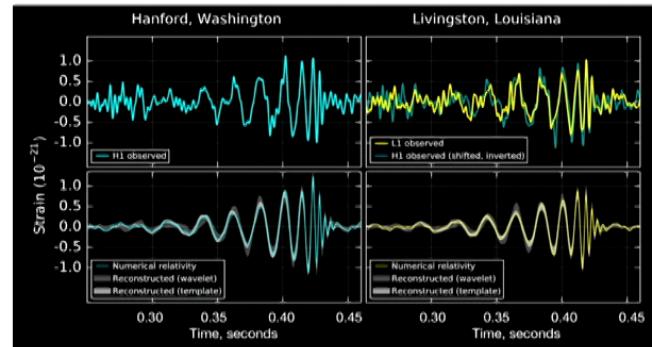
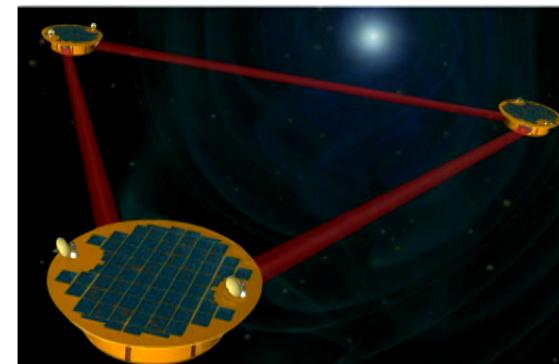


IMAGE CURTESY OF LIGO, NSF



eLISA: supermassive

Q: observational  
consequences of the  
symmetry?

# Applications:

Compere, Gralla, SH, Hughes, Lupsasca,  
Oliveri, Porfyriadis, Rodriguez, Shi,  
Strominger, Warburton

- Gravitational waves
- Magnetospheres: solve Force-free electrodynamics.

$$F_{\mu\nu} \nabla_\rho F^{\nu\rho} = 0$$

- EM emissions: Event Horizon Telescope, iron lines.
- ?



Today:

Focus on

**GW emission in final stages of  
EMRI, via conformal symmetry**

# Rapidly rotating BHs

- (near-) Extremal Kerr  $\sqrt{1 - \frac{\alpha^2}{M^2}} \ll 1$
- Close to horizon  $\epsilon \equiv r_{\text{BL}} - r_+ \ll 1$

NHEK = near horizon extreme Kerr

$$ds^2 = 2M^2\Gamma(\theta) \left[ -R^2dT^2 + \frac{dR^2}{R^2} + d\theta^2 + \Lambda(\theta)^2(d\Phi + RdT)^2 \right]$$

Bardeen & Horowitz '99



- Enhanced isometry

- Kerr:

- NHEK:

$$\begin{array}{c} \mathbb{R} \times U(1) \\ \downarrow \\ SL(2, \mathbb{R}) \times U(1) \end{array}$$

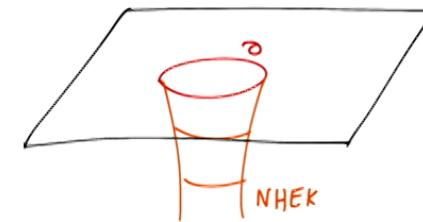
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## Near-horizon, near-extremal limit

Coordinate  
transformation,  
"zoom-in":

$$R = \frac{\hat{r} - r_+}{\epsilon r_+} ; \quad T = \frac{\epsilon \hat{t}}{2M} ; \quad \Phi = \hat{\phi} - \frac{\hat{t}}{2M}$$

Near extremal BH:

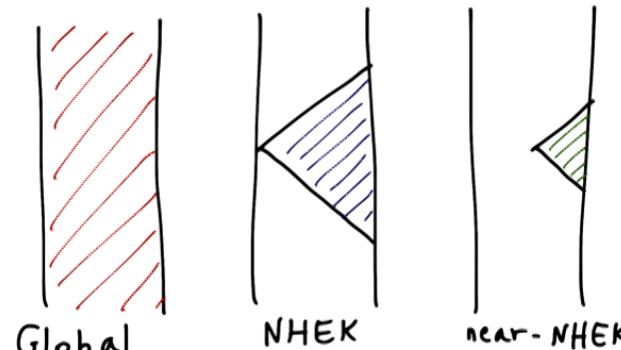
$$\gamma \equiv \sqrt{1 - \frac{a^2}{M^2}} \ll 1$$



# NHEK & near-NHEK

**NHEK:**  $\delta = 0, \epsilon \ll 1$

$$ds^2 = 2M^2\Gamma(\theta) \left[ -R^2dT^2 + \frac{dR^2}{R^2} + d\theta^2 + \Lambda(\theta)^2(d\Phi + RdT)^2 \right]$$



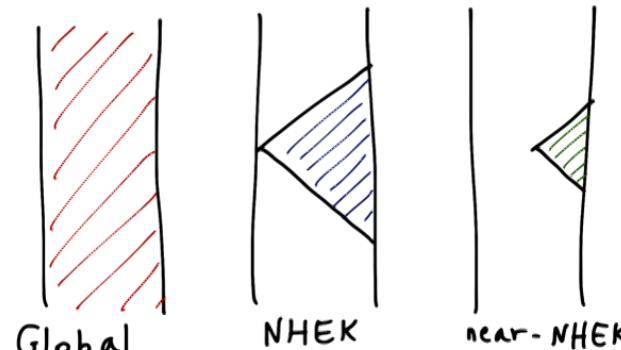
**Near-NHEK:**  $\delta = \kappa \epsilon \ll 1$

$$ds^2 = 2M^2\Gamma(\theta) \left[ -r(r+2\kappa)dt^2 + \frac{dr^2}{r(r+2\kappa)} + d\theta^2 + \Lambda(\theta)^2(d\phi + (r+\kappa)dt)^2 \right]$$

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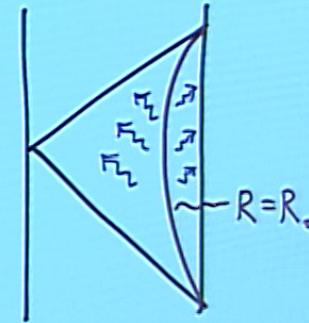
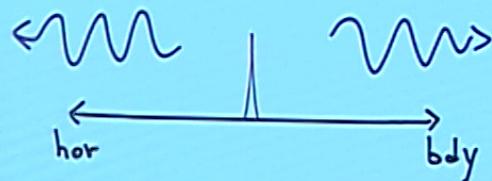


**Near-NHEK:**  $\delta = K \epsilon \ll 1$

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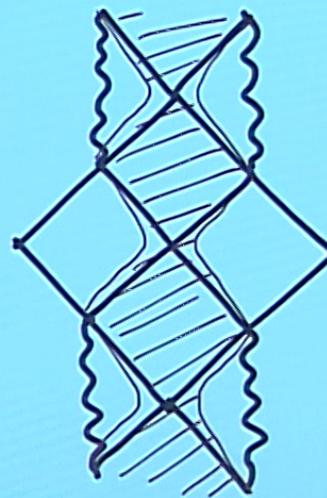
# Circular orbit

Reduction to 1D



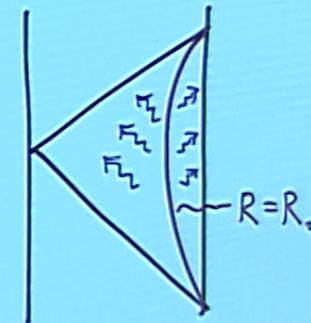
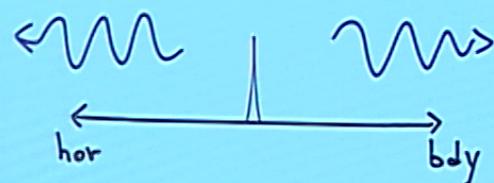
ISCO orbital parameters:  $e = 0$ ,  $l = 2M/\sqrt{3}$

# Embedding in Kerr



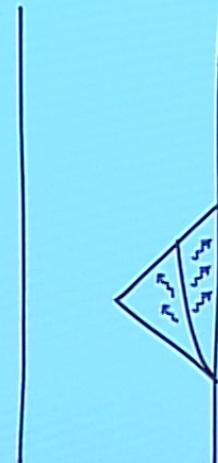
# Circular orbit

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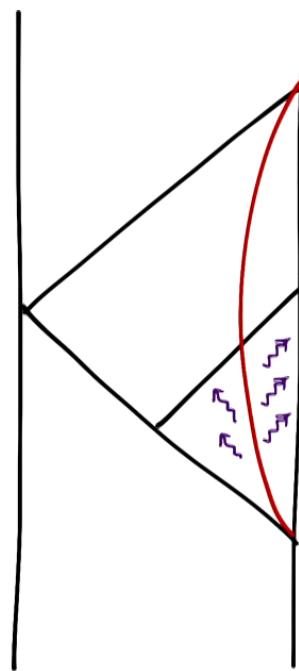


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# Plunge



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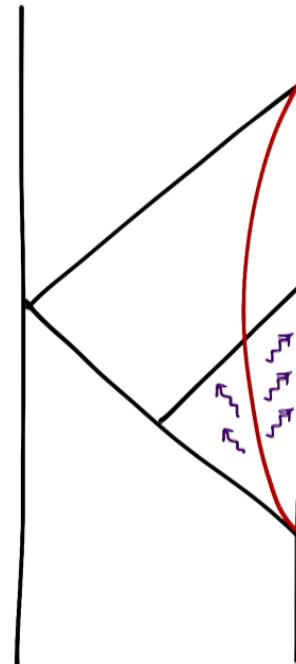
# Plunge

Transformation

$$T = -e^{-\kappa t} \frac{r + \kappa}{\sqrt{r(r + 2\kappa)}},$$

$$R = \frac{1}{\kappa} e^{\kappa t} \sqrt{r(r + 2\kappa)},$$

$$\Phi = \phi - \frac{1}{2} \ln \frac{r}{r + 2\kappa}.$$



Circular in NHEK

$$R = R_0,$$

$$\Phi(T) = -\frac{3}{4}R_0 T + \Phi_0,$$

Plunge in near-NHEK

$$t(r) = \frac{1}{2\kappa} \ln \frac{1}{r(r + 2\kappa)} + t_0,$$

$$\phi(r) = \frac{3r}{4\kappa} + \frac{1}{2} \ln \frac{r}{r + 2\kappa} + \phi_0,$$

"Post-ISCO plunge" parameters:  $e = 0$ ,  $l = 2M/\sqrt{3}$

# CFT interpretation

- Orbiting object dual to external driving:

$$S = S_{\text{CFT}} + \int \theta J$$

- Induces nonzero transition rate out of the vacuum state of CFT, dual to flux into horizon.

Flux:  $\mathcal{F} = \int \sqrt{-g} J^r d\theta d\phi, \quad J^\mu = \frac{i}{8\pi} (\Psi^* \nabla^\mu \Psi - \Psi \nabla^\mu \Psi^*)$

Transition rate:  $\mathcal{R} = \int d\omega \sum_{\ell,m} |J_{\ell m \omega}|^2 \int d\phi dt e^{-im\phi+i\omega t} \langle \mathcal{O}^\dagger(\phi, t) \mathcal{O}(0, 0) \rangle_T$

$$\mathcal{R} = \mathcal{F}$$

# CFT interpretation

- Plunge CFT: via boundary conformal transformation!
- Mapping on boundary:

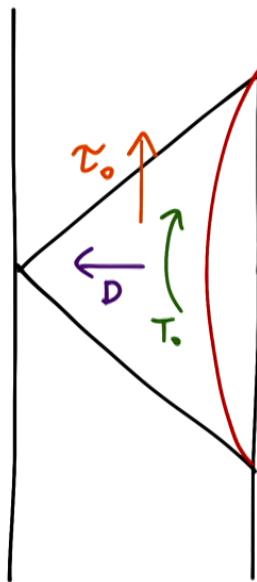
$$\begin{aligned} T &= -e^{-\kappa t}, \\ \Phi &= \phi, \end{aligned} \Rightarrow J \rightarrow (\kappa e^{-\kappa t})^{1-h} J$$

h - conformal weight of  $\phi$

$$R = \mathcal{T}$$

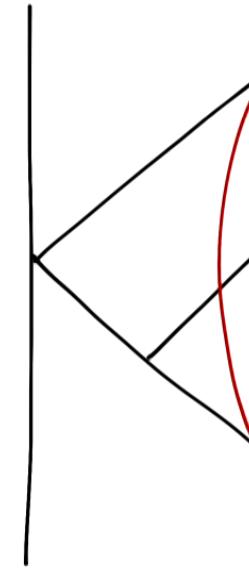
# More mapping fun

SL(2,R) operation



- Global time translations
- Poincare time translations
- Dilatations

"Rindler" cutting



# Fast plunges

Fast NHEK plunge



Fast near-NHEK plunge

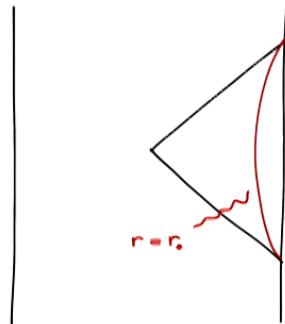


1-parameter family:  
 $e = \frac{4M}{\sqrt{3}R_0}$  ,     $l = \frac{2M}{\sqrt{3}}$

# Generic orbits

- Generalize to any equatorial plunge?
- L invariant under  $SL(2, \mathbb{R})^*U(1)$

Unstable circular orbits  
in near-NHEK

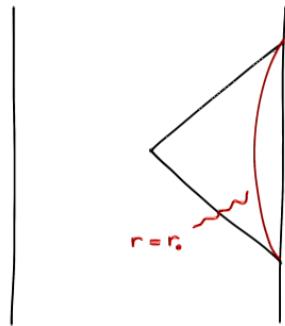


1-parameter family

$$l = \frac{2M(R_0 + \kappa)}{\sqrt{3R_0^2 + 6R_0\kappa - \kappa^2}}$$

# Generic orbits

Unstable circular orbits

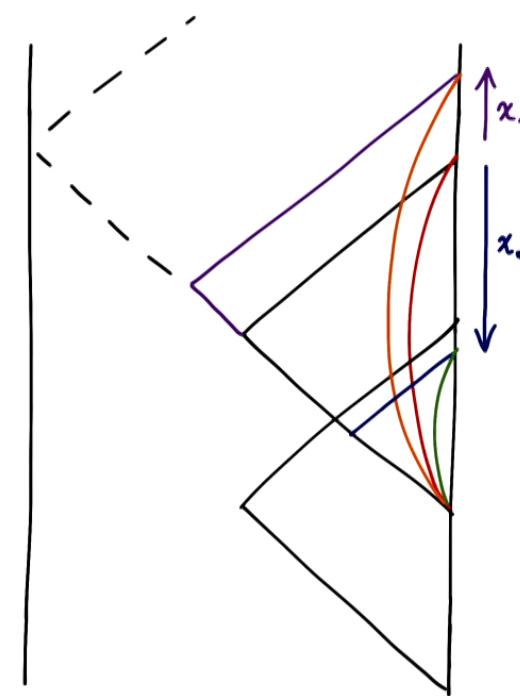


2-parameter family:

$$l = \frac{2M(R_0 + \kappa)}{\sqrt{3R_0^2 + 6R_0\kappa - \kappa^2}}$$

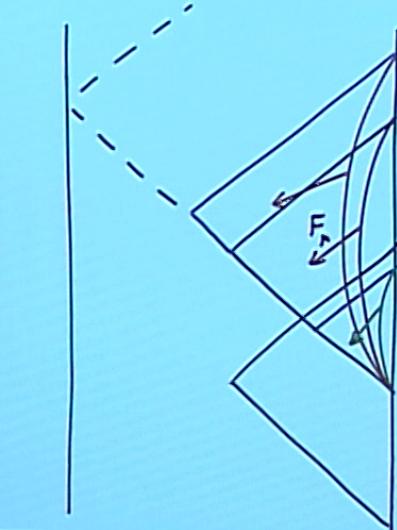
$$e = \frac{2M\kappa^2 \chi_0}{\sqrt{3R_0^2 + 6R_0\kappa - \kappa^2}}$$

Generic near-NHEK plunge



# Self-force

Generic equatorial orbits:



$$F_\mu = \nabla_\mu \Psi^R$$

$\Psi^R$   
Regular piece of field  
(Detweiler-Whiting)

- Any orbit from circular
- Relevant for: trajectories,  
"overspinning",...

# GW signal

- Frequency domain GW given by matching.
- Time domain (ringdown regime): quasinormal modes.

Frequencies (weakly damped):  $\hat{\omega}_{N\ell m} = \frac{1}{2M} [m - i\delta(N + h)]$

- QNM amplitudes summable analytically:

Time domain waveform:

$$\Psi^{far}(r \rightarrow \infty) \sim \exp \left[ -i \frac{m - i\delta h}{2M} (\hat{t} - \hat{t}_0) - \frac{3im}{4\delta} e^{-\frac{\delta}{2M}(\hat{t} - \hat{t}_0)} \right]$$



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# Outlook

- More astrophysical applications!
- ASG w/ astrophysical BC.
- Backreaction: various BH physics applications.
- Relating to horizon instability.

