

Title: Supersymmetric Field Theories for Mathematicians

Date: Nov 16, 2016 03:30 PM

URL: <http://pirsa.org/16110023>

Abstract:

7d gauge theory \implies K-theoretic Donaldson Thomas invariants.

7d gauge theory \implies K-theoretic Donaldson Thomas invariants.
On flat spaces twist of 7d gauge theory is CS
on $\mathbb{C}P^3 \times \mathbb{R}$

7d gauge theory \Rightarrow K-theoretic Donaldson Thomas invariants.
On flat spaces twist of 7d gauge theory is CS

Fields

$$\varphi \in \Omega^{0,1}(\mathbb{C}^3) \otimes \Omega^1(\mathbb{R})[\varepsilon] \otimes g[\mathbb{1}]$$
$$\int CS(e) dz_1 dz_2 dz_3 d\varepsilon$$

$$\int_{\text{CY3}} dz_1 dz_2 dz_3 d\epsilon$$

X a CY3

What's Hilbert space on $X \times \mathbb{R}$?

$\int_{\text{CY3}} dz_1 dz_2 dz_3 d\epsilon$

X a CY3 What's Hilbert space on $X \times \mathbb{R}$?

1st Understand phase space

Phase space = solns to EOM on a small interval $X \times (-\epsilon, \epsilon)$

Always symplectic.

$\int_{\text{CY3}} dz_1 dz_2 dz_3 d\epsilon$

X a CY3 What's Hilbert space on $X \times \mathbb{R}$?

1st Understand phase space

Phase space = solns to EOM on a small interval $X \times (-\epsilon, \epsilon)$

Always symplectic.

Phase space = solⁿs to EOM which are independent of t and dt

Phase space:

$$\varphi \in \Omega^{\text{odd}}(X) \otimes_{\mathbb{R}} \mathbb{C}$$

satisfying an equation

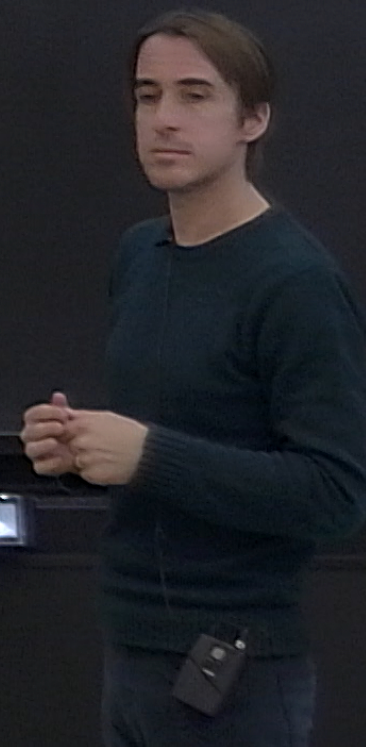
ε is of degree -1 , then degree 0 part is

$$A \in \Omega^{\text{odd}}(X) \otimes_{\mathbb{R}} \mathfrak{g} \quad (\text{coeff. of } \varepsilon)$$

$$B \in \Omega^{\text{even}}(X) \otimes_{\mathbb{R}} \mathfrak{g}$$

Satisfying $F^{\text{odd}}(A) = 0, \bar{\partial}_A B = 0$

Modulo gauge, $A \rightarrow XA X^{-1} + X^{-1} \bar{\partial} X$
 $B \rightarrow X^{-1} B X$
 $B \rightarrow B + \bar{\partial}_A \eta$
 $\eta \in \Omega^0(\mathfrak{g})$



Modulo gauge, $A \rightarrow X A X^{-1} + X^{-1} \bar{\partial} X$

$$B \rightarrow X^{-1} B X$$

$$B \rightarrow B + \bar{\partial}_A \eta$$

$\eta \in \Omega^{0,1}(X, \mathfrak{g})$

A: holomorphic bundle on X

B: element of $H^2(X, \mathfrak{g}_{ad})$
↳ adjoint bundle

Phase space:

$$\mathcal{P} \in \mathcal{S}^0(X) \otimes \mathfrak{g} \otimes \mathbb{C}^2$$

satisfying an equation
 \mathcal{E} is of degree -1

Satisfying $F^{0,2}(A) = 0$, $\bar{\partial}_A B = 0$

This is $T^* \text{Bun}_G(X)$

At a bundle P ,

$$T_P \text{Bun}_G(X) = H^1_{\bar{\partial}}(X, \mathfrak{g}_{\text{ad}})$$

dual by Serre duality to $H^2_{\bar{\partial}}(X, \mathfrak{g}_{\text{ad}})$

$$K_X = 0$$

7d gauge theory on $X \times \mathbb{R}$



Quantum mechanics on

$T^*(\text{Bun}_G(X))$

7d gauge theory on $X \times \mathbb{R}$

→ Quantum mechanics on

$T^* \text{Bun}_G(X)$

Hilbert space is

$$H^0(\text{Bun}_G(X), K^{\otimes 1/2})$$

7d gauge theory on $X \times \mathbb{R}$



Quantum mechanics on

$T^* \text{Bun}_G(X)$

Hilbert space is

$H_{\frac{1}{2}}^0(\text{Bun}_G(X), K^{\otimes 1/2})$

$Z(X \times S^1) = \text{Euler char. of this}$

By RP, $Z(X, S^1) = \int_{\text{Bun}_G(X)} \text{Td}(\text{Bun}_G(X)) e^{\frac{1}{2}c_1}$
 $= K\text{-theoretic DT invariants}$

8 supercharges

To give a theory we need
a Lie alg. \mathfrak{g} and a symplectic representation V .

16 supercharges is special case $V = T^*\mathfrak{g} = \mathfrak{g} \oplus \mathfrak{g}^*$

Given (\mathfrak{g}, V)

Given (g, V)

Let

$\mathfrak{g}_V = \text{graded Lie algebra}$

$x \in \mathfrak{g}$

$v \in V$

$x \in \mathfrak{g}^*$

$[X, -]$ action of \mathfrak{g} on everything

$[v_1, v_2] = \partial_{v_1} \partial_{v_2} \mu$ $\mu: V \rightarrow \mathfrak{g}^*$ is
moment map (quadratic fn)

An element $\alpha \in \mathfrak{g}_V$ of degree 1, α satisfies MC equation if $[\alpha, \alpha] = 0$

An element $\alpha \in \mathfrak{g}_V$ of degree 1, α satisfies MC equation if $[\alpha, \alpha] = 0$
In this case $\alpha \in V$ satisfies MC $\Leftrightarrow \alpha \in \mu^{-1}(0)$

An element $\alpha \in \mathfrak{g}_V$ of degree 1, α satisfies MC equation if $[\alpha, \alpha] = 0$
 In this case $\alpha \in V$ satisfies MC $\Leftrightarrow \alpha \in \mu^{-1}(0)$
 $\chi \in \mathfrak{g}_V^0$ act as symmetries of set of MC solⁿs
 $\alpha \rightarrow \alpha + \varepsilon[\chi, \alpha]$

An element $\alpha \in \mathfrak{g}_V$ of degree 1, α satisfies MC equation if $[\alpha, \alpha] = 0$
 In this case $\alpha \in V$ satisfies MC $\Leftrightarrow \alpha \in \mu^{-1}(0)$
 \mathfrak{g}_V act as symmetries of set of MC solⁿs
 $\alpha \rightarrow \alpha + \varepsilon [\chi, \alpha]$
 MC solⁿs/gauge = $\mu^{-1}(0)/G$ = symplectic reduction

7d gauge theory on $X \times \mathbb{R}$

Theorem $N=(1,0)$ theory in 6 dimensions
matter in V is the theory (after twisting)
whose fields are in
 $\Omega^{0,2}(\mathbb{P}^3) \otimes g_V [1]$

Theorem $N=(1,0)$ theory in 6 dimensions
 matter in V is the theory (after twisting)
 whose fields are in

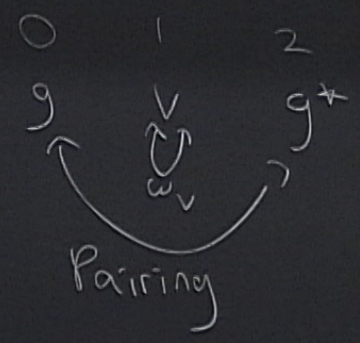
$$\varphi \in \Omega^{0,2}(\mathbb{P}^3) \otimes g_V [1]$$

$$S(\varphi) = \int dz_1 dz_2 dz_3 \left(\frac{1}{2} \langle \varphi, d\varphi \rangle + \frac{1}{6} \langle \varphi, [\varphi, \varphi] \rangle \right)$$

where $\langle \rangle$ is the invariant pairing on g_V of deg. -2

0
 9
 ↓

s
(isting)
(eT)
v of deg. -2



Theorem $N=(1,0)$ theory in 6 dimensions
 matter in V is the theory (after twisting)
 whose fields are in

$$\varphi \in \Omega^{0,2}(\mathbb{C}^3) \otimes g_V [1]$$

$$S(\varphi) = \int dz_1 dz_2 dz_3 \left(\frac{1}{2} \langle \varphi, d\varphi \rangle + \frac{1}{6} \langle \varphi, [\varphi, \varphi] \rangle \right)$$

where $\langle \cdot \rangle$ is the invariant pairing on g_V of deg. -2

This is holomorphic Rozansky-Witten theory (gauged)

0
 \curvearrowright
 g
 Pair

If X is a CY3
What are solⁿs to EOM?

If X is a CY3

What are solⁿs to EOM?

$$\Omega^{0,2}(X) \otimes \begin{pmatrix} g & V & g^* \\ & -1 & 0 \\ & & 1 \end{pmatrix}$$

g : we will find a hol. G -bundle on X

V part: \leadsto a section φ of assoc. G -bundle
 $\partial_A \varphi = 0$

Also $\mu(\varphi) = 0 \in \Omega^{0,2}(X, g_0)$

If $U \subseteq V$ preserved by G , and G acts
 freely. If φ lands in U , then
 Solⁿs to EOM \equiv hol. maps X to $U//G$.
 " $V//G = \text{Higgs branch}$ "



If we reduce $S^1 \times \mathbb{R}$ to 3d, we find same theory built from
Fields are \mathfrak{g}_V

$\Omega^{0,1}(\sigma) \otimes \Omega^*(\mathbb{R})[\mathfrak{g}_V]$
A twist of 3d $N=4$ gauge theory

the theory built from

Remaining SUSY

is given by the odd v. fields
 $\epsilon \partial_z$ and ∂_ϵ on \mathbb{P}^1

the theory built from

Remaining SUSY

is given by the odd v fields

$\epsilon \partial_z$ and ∂_ϵ on \mathbb{P}^1

$\epsilon \partial_z \rightsquigarrow$ Rozansky-Witten twist

$\partial_\epsilon \rightsquigarrow$ Twisted RW twist

the theory built from

Remaining SUSY

is given by the odd v. fields

$\epsilon \partial_z$ and ∂_ϵ on \mathbb{P}^1

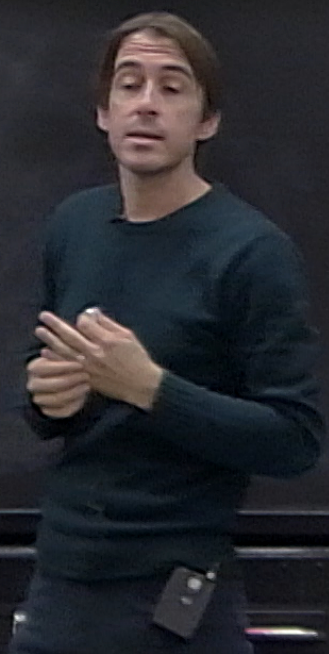
$\epsilon \partial_z \rightsquigarrow$ Rozansky-Witten twist (B)

$\partial_\epsilon \rightsquigarrow$ Twisted RW twist (A)

3d mirror symmetry

$$g = g \oplus g^*$$

Start with $g = 0$
 $\varphi \in \Omega^0(\mathbb{R}) \otimes \Omega^1(\mathbb{R}^n) \otimes V(\mathbb{R})$
 $(\bar{\partial}_z + d_{dR})\varphi = 0$



Start with $\eta = 0$.
 $\varphi \in \Omega^0(\mathbb{R}) \otimes \Omega^1(\mathbb{R}) \otimes V[\varepsilon]$
 $(\bar{\partial}_z + d_{dK})\varphi = 0$

$$\varphi = \varphi_0 + \varepsilon \varphi_1$$

Operators we can build are functions of

$$\partial_z^k \varphi_0(0), \quad \partial_z^k \varphi_1(0)$$

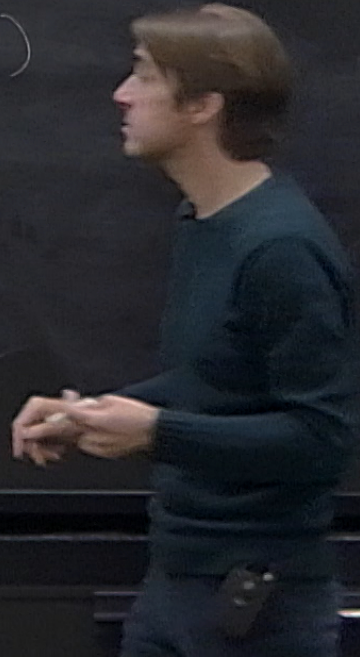
↑
bosonic

↑
fermionic

$\partial_z \varphi_0(0)$, $\partial_z^k \varphi_1(0)$
↑ bosonic ↑ fermionic

$\varepsilon \partial_z$ sends operator

$\partial_z^k \varphi_1(0) \rightsquigarrow \partial_z^{k+1} \varphi_0(0)$
 ∂_ε sends
 $\partial_z^k \varphi_0(0) \rightsquigarrow \partial_z^k \varphi_1(0)$



$\partial_z \varphi_0(0)$, $\partial_z^k \varphi_1(0)$
 ↗ bosonic ↗ fermionic

$\mathcal{E}\partial_z$ sends operator

$$\partial_z^k \varphi_1(0) \rightsquigarrow \partial_z^{k+1} \varphi_0(0)$$

∂_ε sends

$$\partial_z^k \varphi_0(0) \rightsquigarrow \partial_z^k \varphi_1(0)$$

∂_ε no cohomology

$\mathcal{E}\partial_z$: Only $\varphi_0(0)$ survives in cohomology

$\mathcal{E}\partial_z$ is RW twist
deduce

$\mathcal{E}D_2$ is RW twist
deduce local operators
= $S \cdot V^\circ$

When there's a gauge group,
operators in RW twist
= hol. functions on $V//G$

$V//G = \text{Higgs branch}$