

Title: Building a Gauge-Invariant Holographic Dictionary

Date: Oct 18, 2016 02:00 PM

URL: <http://pirsa.org/16100060>

Abstract: <p>What natural CFT quantities can be seen in the interior of the bulk AdS in a diffeomorphism invariant way? And how can we use them to learn about the emergence of local bulk physics? Inspired by the Ryu-Takayanagi relation, we construct a class of simple non-local operators on both sides of the duality and demonstrate their equivalence. Integrals of free bulk fields along geodesics/minimal surfaces are dual to what we will call "OPE blocks": Individual conformal family contributions to the OPE of local operators. Our findings can be utilized to reconstruct local bulk operators and unify a number of previously disconnected AdS/CFT results. We extend this kinematic correspondence to incorporate gravitational interactions in AdS₃ by relating geodesic operators to the Virasoro OPE blocks, hence providing a natural CFT prescription for "gravitational dressing". We conclude with discussion of preliminary results on an interesting CFT structure that generalizes our dictionary to include arbitrary local bulk interactions. </p>

Building a Gauge-Invariant Holographic Dictionary

Lampros Lamprou
Stanford University

Perimeter Institute, October 2016

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Based on:

[1604.03110], [1608.06282] B.Czech, L.L., S.McCandlish, B.Mosk, J.Sully
L.L., S.McCandlish *to appear*

See also: [1606.03307] J.de Boer, F.Haehl, M. Heller, R.Myers

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Philosophical Overture

Gravity **disfavors locality**:

- Diffeomorphism invariance.
- Black hole resonances in high energy scattering.
- Holographic bound.

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- What mathematics describe their **dynamics**?
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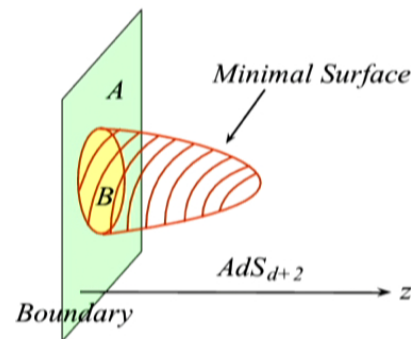
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Our goal: Make progress via a **gauge-invariant** approach to the **AdS/CFT dictionary**.

A Clue from AdS/CFT

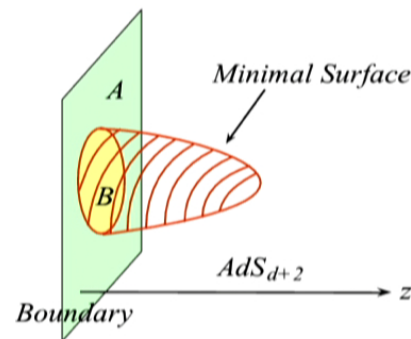
New approach to observables in gravity: **Ryu-Takayanagi** relation.



Entanglement entropy = Area of minimal surfaces

A Clue from AdS/CFT

New approach to observables in gravity: **Ryu-Takayanagi** relation.



Entanglement entropy = Area of minimal surfaces

Has all the right properties:

- Probes geometry** of bulk interior.
- Diff invariant** – makes reference only to the boundary.
- Natural non-local** object in the **CFT**.

A Gauge-Invariant Dictionary

Apply RT philosophy to the **operator dictionary**:

- Identify **natural CFT operators** dual to **extended bulk probes**.
- Reconstruct the familiar **local bulk fields** and describe their **dynamics**.

A Gauge-Invariant Dictionary

Apply RT philosophy to the **operator dictionary**:

- Identify **natural CFT operators** dual to **extended bulk probes**.
- Reconstruct the familiar **local bulk fields** and describe their **dynamics**.
- **Ultimate goal**: What is the **CFT principle** responsible for the approximate **bulk locality** at low energies?

Our Approach

- **Step 1:** Gauge-invariant dictionary for **free bulk fields**.
(“A Stereoscopic Look into the Bulk”)
- **Step 2:** Include **gravitational** interactions in **2+1** dimensions.
(See also Sam’s talk in a couple of weeks!)
- **Step 3:** Incorporate arbitrary **local interactions**.
(In progress...)

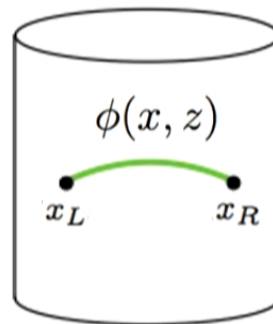
Geodesic Operators

- **Geodesics** and **minimal surfaces** are special bulk objects.
- Both defined with reference to the **asymptotic boundary**.
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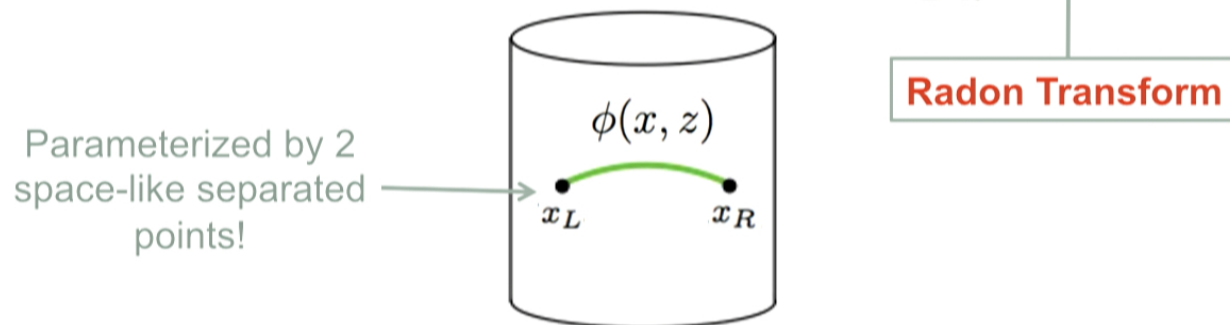
- Restrict to AdS_3 for simplicity:
$$\tilde{\phi}(\gamma_{x_L x_R}) = \int_{\gamma_{x_L x_R}} ds \phi(x, z) \Big|_{\gamma}$$



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Global OPE Blocks

- We are looking for a **natural CFT bi-local**.
- The **OPE** provides with an **organizational principle**.
- Product of any two operators is expanded as:

$$\mathcal{O}_i(x) \mathcal{O}_j(0) = \sum_k C_{ijk} |x|^{\Delta_k - \Delta_i - \Delta_j} (1 + b_1 x^\mu \partial_\mu + b_2 x^\mu x^\nu \partial_\mu \partial_\nu + \dots) \mathcal{O}_k(0)$$

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Global OPE Blocks

- **OPE blocks** admit an **integral representation**.

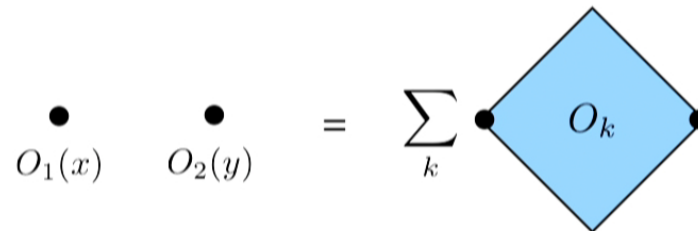
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- Smear the primary operator $\mathcal{O}_k(w, \bar{w})$ over the **causal diamond** defined by the two points.

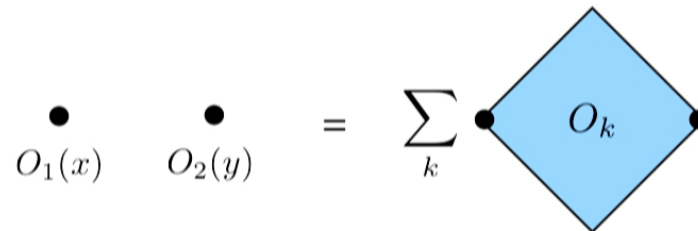


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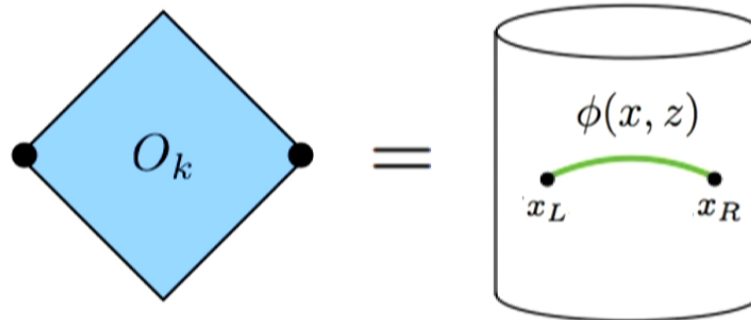


- Smearing function: $G_k(w, \bar{w}; x, y) = \frac{\langle \mathcal{O}_H(x) \mathcal{O}_H(y) \tilde{\mathcal{O}}_k(w, \bar{w}) \rangle}{\langle \mathcal{O}_H(x) \mathcal{O}_H(y) \rangle}$

Gauge-Invariant Dictionary 1.0

OPE Blocks are dual to **geodesic operators**:

$$\mathcal{B}_{h,\bar{h}}(x_L, x_R) = \tilde{\phi}(\gamma_{x_L x_R}) = \int_{\gamma_{x_L x_R}} ds \phi(x, z) \Big|_{\gamma}$$

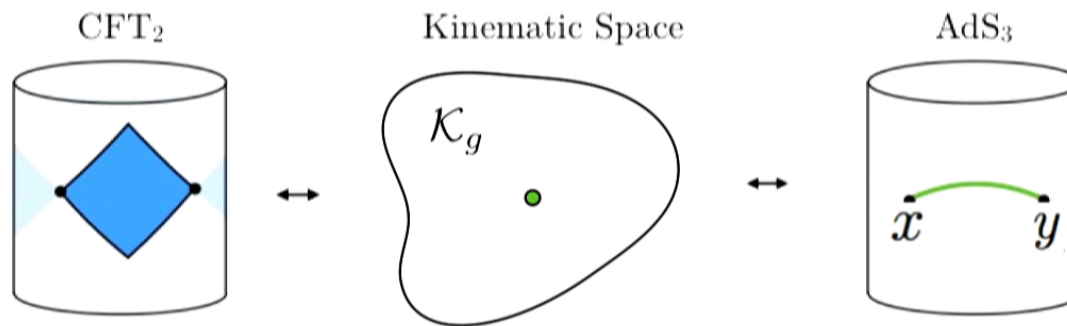


Kinematic Space

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- They live on the space of **pairs of points** or, equivalently, the space of **causal diamonds**.

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- This is **kinematic space** and it has **geometric** structure.

Kinematic Space

- Utilize **conformal symmetry** to define a **metric** on kinematic space.

$$ds^2 = \frac{I_{\mu\nu}(x-y)}{|x-y|^2} dx^\mu dy^\nu \quad \text{with:} \quad I_{\mu\nu} \equiv \eta_{\mu\nu} - 2 \frac{x_\mu x_\nu}{x^2}$$

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- Signature (d,d)**: If 2 pairs of points **share** a point, their kinematic distance is zero (no invariant cross-ratio for 3 points)
- For a **CFT₂** kinematic space decomposes to a product of **two de Sitter** spaces.

$$ds^2 = \underbrace{\frac{du_1 du_2}{|u_2 - u_1|^2}}_{dS_2} + \underbrace{\frac{dv_1 dv_2}{|v_2 - v_1|^2}}_{dS_2}$$



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- Blocks are **eigenvectors** of the **conformal Casimirs:**

$$[L^2, \mathcal{B}_k(x, y)] = C_k \mathcal{B}_k(x, y)$$

$$\text{where: } C_k = (\Delta_k + \ell_k)(\Delta_k + l_k - 2)$$

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$$[\square_{dS_2} + \square_{\bar{dS}_2} - m_{\Delta_k}^2] \tilde{\phi}(\gamma) = 0$$

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Redundancy: Not all functions on kinematic space are consistent Radon transforms!

We need **constraints:** John's equations.

$$(\square_{dS_2} - \bar{\square}_{dS_2}) \tilde{\phi} = 0$$

Proof of the Dictionary

- **One last step:** Boundary conditions.
- Determined by the coincident limit.

OPE Block

$$\mathcal{B}_k(x, y) \rightarrow |x - y|^{\Delta_k} \mathcal{O}_k(x)$$

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$$\tilde{\phi}(x \rightarrow y) \sim |x - y|^{\Delta} \mathcal{O}_{\Delta}(x)$$

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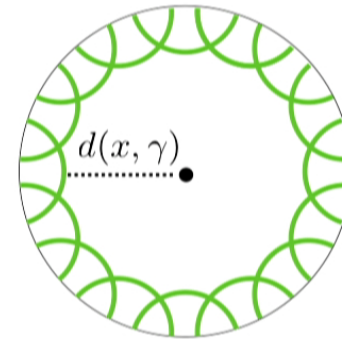
- Both objects obey the **same equations** with the **same boundary conditions**.

$$\mathcal{B}_k(x, y) = \tilde{\phi}(\gamma) = \int ds \phi(x, z) \Big|_\gamma$$

Local Bulk Operators

- Use geodesic operators to **reconstruct** the **local** operators:
Inverse Radon transform.

$$f(x) = -\frac{1}{\pi} \int_0^{\infty} \frac{dp}{\sinh p} \frac{d}{dp} \left(\text{average}_{d(x,\gamma)=p} \tilde{f}_g(\gamma) \right)$$



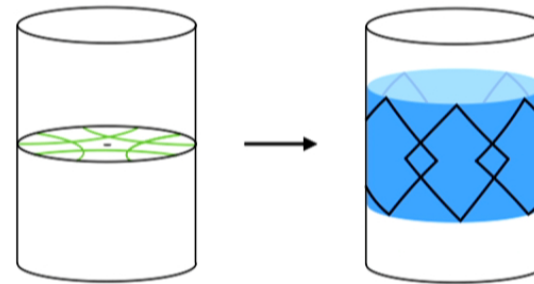
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- Manifestly **diffeomorphism invariant**.
- Reconstruction **point** defined by **set of geodesics** that cross it.

Local Bulk Operators

- **Result:** Integral over the **spacelike** separated **boundary region**:

$$\phi(\rho = 0) = \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} dt \int_0^{2\pi} d\theta K_{\Delta}(t) \mathcal{O}_{\Delta}(t, \theta)$$

$$K_{\Delta}(t) = -\frac{k}{\pi} (\cos t)^{\Delta-2} \log \cos t$$

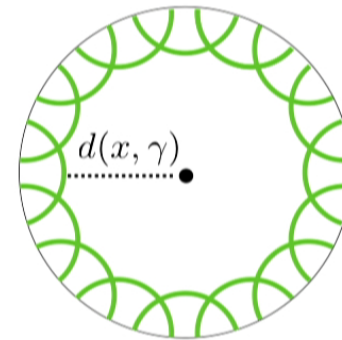


- This matches the **HKLL global smearing** function.

Local Bulk Operators

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- Principle behind the correspondence: **conformal symmetry**.
- **The idea**: Repackaging our data by utilizing symmetries can manifest connections previously invisible.
- Indeed: It helps us “tidy up” seemingly disconnected holographic results:
 - a) **Conformal blocks** as **geodesic Witten diagrams**.
 - b) **First law of entanglement entropy** and **linearized Einstein’s equations**.

What is it good for?

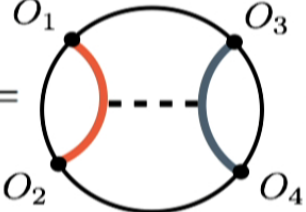
- **Conformal blocks** as **geodesic Witten diagrams**:

$$\text{Recall: } \mathcal{O}_i(x) \mathcal{O}_j(0) = |x|^{-\Delta_i - \Delta_j} \sum_k C_{ijk} \mathcal{B}_k^{ij}(x, 0)$$

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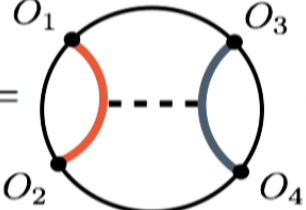
$$\langle 0 | \mathcal{B}_k(x_1, x_2) \mathcal{B}_k(x_3, x_4) | 0 \rangle = g_{k|1234}(u, v) =$$


[Hijano, Kraus, Perlmutter, Snively]

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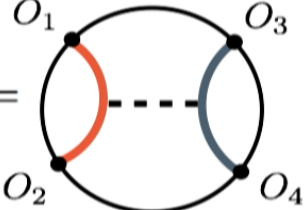
- **Entanglement and Einstein's equations:**

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- Entanglement and **Einstein's equations**:

$$\delta S = \delta \langle H_{\text{mod}} \rangle$$

$$H_{\text{mod}} = -\frac{1}{6} \mathcal{B}_{T_{00}}$$

$$(\square_{\mathcal{K}} + 2d)\delta H = - \int_{\gamma} (\square_{AdS} - 2d)\delta g_{\mu\nu} \dot{x}^{\mu} \dot{x}^{\nu} = 0$$

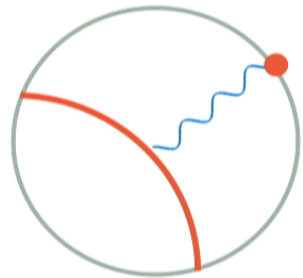
[Faulkner, Guica, Hartman, Lashkari, McDermont, Myers, Swingle, Van Raamsdonk]

STEP 2: GRAVITATIONAL DRESSING

Geodesic Operators and Virasoro OPE Blocks

A Glimpse of Interactions

- Without interactions: $\tilde{\phi}_\Delta(\gamma) = \mathcal{B}_\Delta(\gamma)$

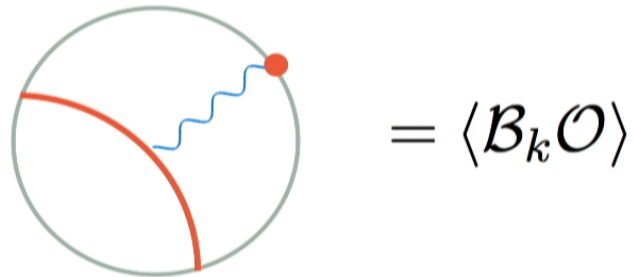


A Feynman diagram enclosed in a circle. It features a red curved line on the left side, representing a fermion propagator. A blue wavy line, representing a boson propagator, forms a loop that starts from the red line, goes up and right, then down and right, and finally up and right to end at a red dot on the red line. This diagram represents a self-energy correction to the fermion propagator.

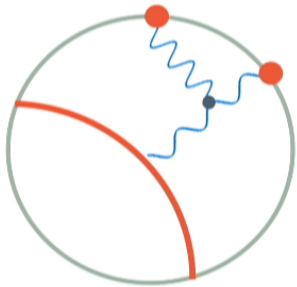
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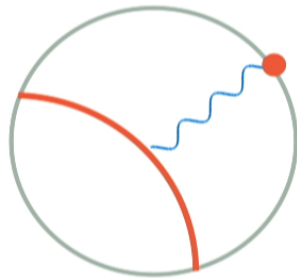


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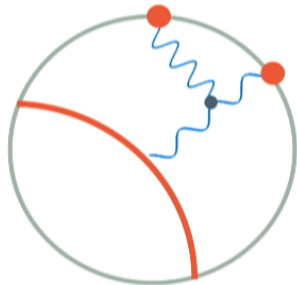
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Multi-trace Blocks

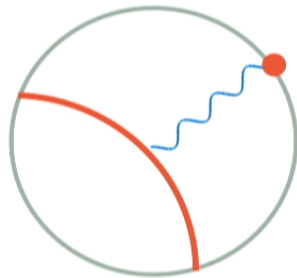
- With interactions: $\tilde{\phi}_\Delta(\gamma) = \mathcal{B}_\Delta(\gamma) + \frac{1}{N} \sum_n a_n^{CFT} \mathcal{B}_n(\gamma)$



$$\neq \langle \mathcal{B}_k \mathcal{O} \mathcal{O} \rangle$$

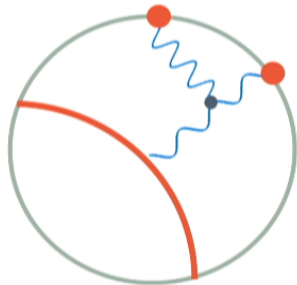
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$$\neq \langle \mathcal{B}_k \mathcal{O} \mathcal{O} \rangle$$

Multi-trace Blocks

- What **CFT principle** determines these “dressed blocks”?

Virasoro OPE Blocks

- **OPE** of 2D CFTs is organized in **larger block structures**.
- These are fixed by **local** conformal symmetry.

$$O_{\Delta}(x_1)O_{\Delta}(x_2) = \mathcal{V}_{\mathbb{I}}^{\Delta\Delta} + \sum_k c_k^{\Delta\Delta} \mathcal{V}_k^{\Delta\Delta}$$

- We will call them **Virasoro OPE blocks**.
- They contain the contributions from all local descendants.

Virasoro OPE Blocks

Identity Block: $\mathcal{V}_{\mathbb{I}}^{\Delta\Delta}(x_1, x_2) = \frac{\mathbb{I}}{x_{12}^{2\Delta}} \left(1 + \frac{2\Delta}{c} \mathcal{B}_T + \# \mathcal{B}_{T\partial^n T} + \dots \right)$

$$= \frac{1}{x_{12}^{2\Delta}} \left(1 + \text{[Diagram: diamond with } T \text{]} + \text{[Diagram: diamond with } TT \text{]} + \dots \right)$$

The diagram shows two blue diamond shapes representing operator blocks. The first diamond contains the symbol T and has two black dots on its left and right vertices. The second diamond contains the symbol TT and also has two black dots on its left and right vertices. These diamonds are connected by plus signs and an ellipsis, all enclosed within large parentheses.

Virasoro OPE Blocks

Identity Block: $\mathcal{V}_I^{\Delta\Delta}(x_1, x_2) = \frac{\mathbb{I}}{x_{12}^{2\Delta}} \left(1 + \frac{2\Delta}{c} \mathcal{B}_T + \# \mathcal{B}_{T\partial^n T} + \dots \right)$

$$= \frac{1}{x_{12}^{2\Delta}} \left(1 + \text{diamond}(T) + \text{diamond}(TT) + \dots \right)$$

Arbitrary family: $\mathcal{V}_k^{\Delta\Delta}(x_1, x_2) = \frac{1}{x_{12}^{2\Delta}} (\mathcal{B}_k + \# \mathcal{B}_{T\partial^n \mathcal{O}_k} + \dots)$

$$= \frac{1}{x_{12}^{2\Delta}} \left(\text{diamond}(\mathcal{O}_k) + \text{diamond}(T\mathcal{O}_k) + \dots \right)$$

Gauge Invariant Dictionary 2.0

- What I would like to claim:

$$\mathcal{V}_{\mathbb{I}}^{\Delta\Delta}(x_1, x_2) = e^{-\frac{6\Delta}{c} \ell(x_1, x_2)}$$

$$\mathcal{V}_k^{\Delta\Delta}(x_1, x_2) = e^{-\frac{6\Delta}{c} \ell(x_1, x_2)} \tilde{\phi}_k(x_1, x_2)$$

Semi-classical States

- **Semi-classical states:** CFT states with well-defined **energy**.

$$\langle T \rangle \sim \mathcal{O}(c) \quad \frac{\langle T^n \rangle - \langle T \rangle^n}{\langle T \rangle^n} \sim \mathcal{O}\left(\frac{1}{c^k}\right)$$

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- Necessary (not sufficient) requirement for smooth bulk geometry.
- It is convenient to **decompose** the **stress tensor** to a **background** piece and a **fluctuation** piece:

$$T = \langle T \rangle + \delta T$$

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Sketch of the proof

- **Semi-classical** states are characterized **only** by their energy density:

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- Then we can **find a frame** in which the state looks like the **vacuum**.
- If **geodesic operators** and **Virasoro blocks** transform the same way then the **proof follows** from previous discussion.

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- Equation for length:
$$\frac{\partial^2 \omega}{\partial u_1 \partial u_2} = \frac{c}{6} e^{-\frac{12}{c} \omega} + J(u_1, u_2)$$

[J.de Boer, F.Haehl, M. Heller, R.Myers]

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- Equation for geodesic operators:

$$e^{-\frac{12}{c} \omega} \partial_{u_1} \partial_{u_2} \tilde{\phi}_k + m_k^2 \tilde{\phi}_k = 0$$

- This is a Liouville theory!

Taking stock...

- **Geodesic integrals** of free fields are dual to **global OPE blocks**.
- **Gravitational dressing** for geodesic operators becomes the “**Virasoro dressing**” for the OPE.
- These operators have a life of their own: Can be described as **fields on kinematic space** in terms of a **Liouville theory**.

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- Can we use this dictionary to “**see**” inside **black holes**?
- Can this teach us how to think of **operators** in gravity on **flat** and/or **de Sitter** space?



Bartek Czech



Sam McCandlish

Thank you!



Benjamin Mosk



James Sully