

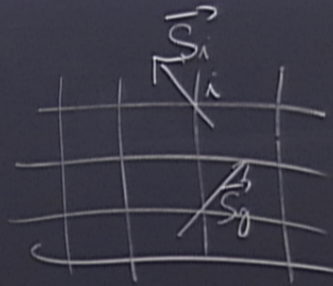
Title: PSI 2016/2017 Statistical Mechanics - Lecture 11

Date: Oct 25, 2016 10:45 AM

URL: <http://pirsa.org/16100025>

Abstract:

YESTERDAY
m-VECTOR MODEL
(d-DIM)

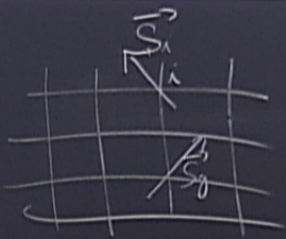


\vec{S}_i ... m-VECTOR

$$H = -\frac{1}{2} \sum_{ij} J_{ij} \vec{S}_i \cdot \vec{S}_j$$

O(m) CONTINUOUS SYMMETRY

Y
R MODEL
(M)



\vec{S}_i ... m-VECTOR

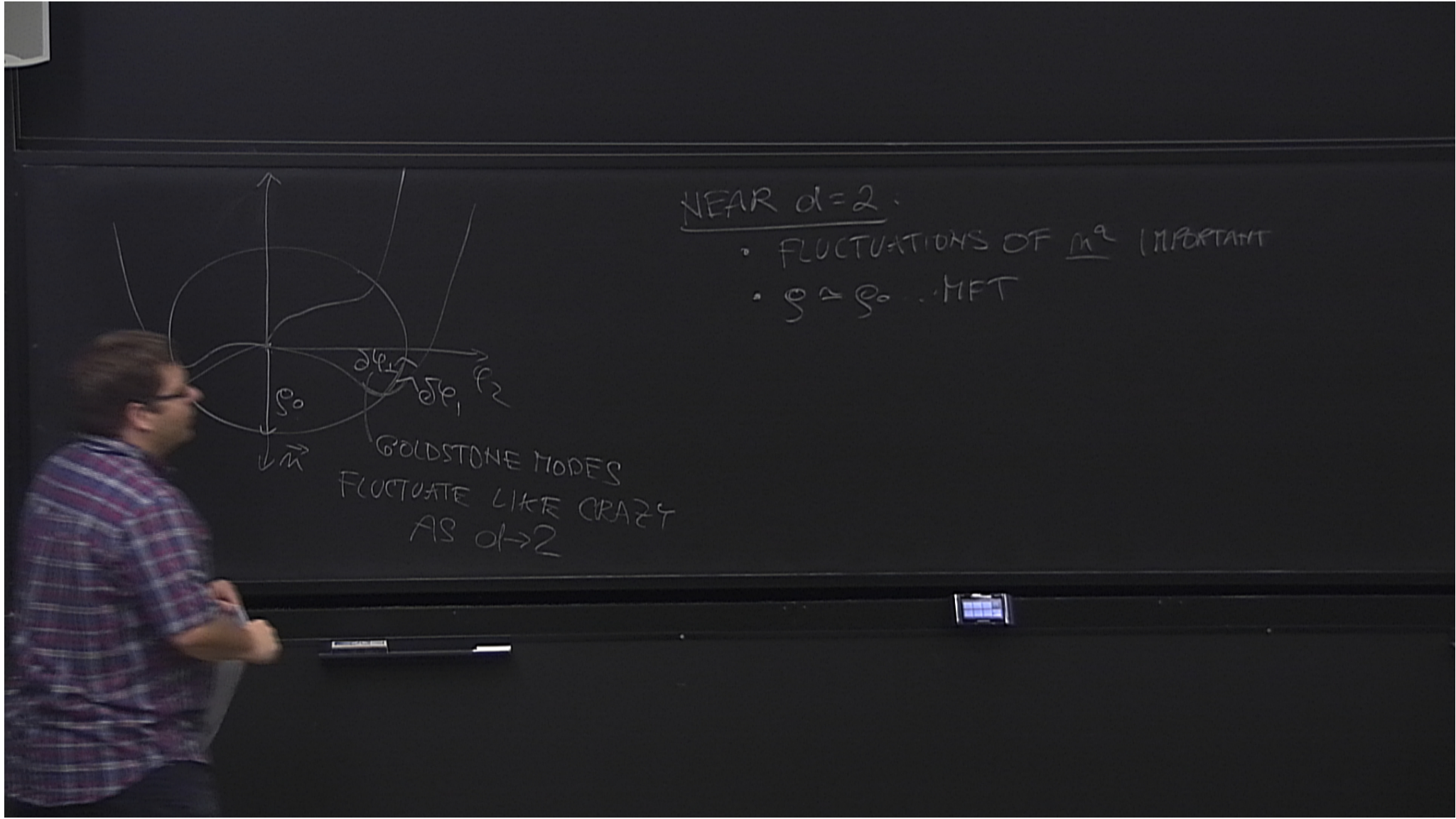
L-G

$$\varphi^a = \rho n^a \quad a=1, \dots, m$$
$$|\vec{n}| = 1$$

$$H = -\frac{1}{2} \sum_{ij} J_{ij} \vec{S}_i \cdot \vec{S}_j$$

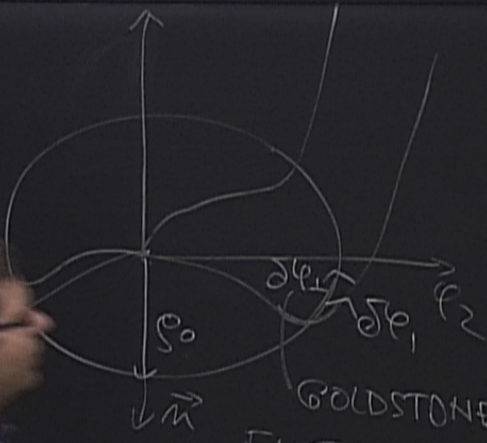
$$S[\rho, \vec{n}] = \int d^d x \left(\frac{1}{2} |\nabla \rho|^2 + \frac{1}{2} \rho^2 (\nabla n^a)^2 + \frac{\Lambda}{2} \rho^2 + \frac{\mu}{4} \rho^4 \right)$$

O(m) CONTINUOUS SYMMETRY

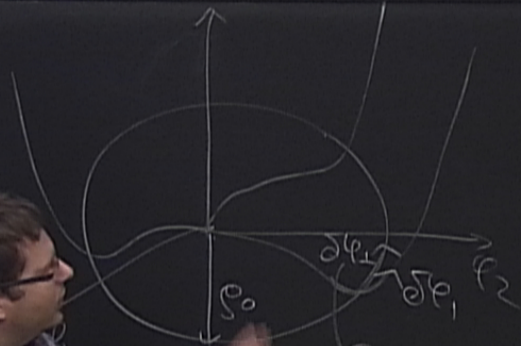


NEAR $d=2$.

- FLUCTUATIONS OF m^2 IMPORTANT
- $S \approx S_0 \dots$ MFT



GOLDSTONE MODES
FLUCTUATE LIKE CRAZY
AS $d \rightarrow 2$



GOLDSTONE MODES
FLUCTUATE LIKE CRAZY
AS $d \rightarrow 2$

NEAR $d=2$.

- FLUCTUATIONS OF m^2 IMPORTANT
- $\phi \approx \phi_0$ MFT

$$\Rightarrow S[\vec{m}] = \frac{\phi_0^2}{2} \int d^d x (\nabla m^a)^2$$

• SET UP $d=2+\epsilon$ EXPANSION
↑ LOWER CRITICAL DIM.

$$\frac{d\tilde{T}}{d\epsilon} = -\epsilon\tilde{T} + (m-2)K_d\tilde{T}^2$$

$$\tilde{T} = \Lambda^\epsilon T, \quad K_d = \frac{S_d}{(2-\epsilon)^d}$$

SIGMA

$$\frac{d}{d\epsilon} (\Lambda^\epsilon T)$$

• SET UP $d=2+\epsilon$ EXPANSION
 \uparrow LOWER CRITICAL DIM.

$$\frac{d\tilde{T}}{d\epsilon} = -\epsilon\tilde{T} + (m-2)K_d\tilde{T}^2$$

2 FIXED POINTS

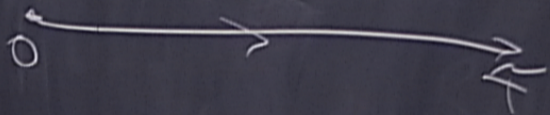
$$\begin{aligned} \tilde{T}^* &= 0 \\ \tilde{T}^* &= \frac{\epsilon}{(m-2)K_d} \end{aligned}$$

$\tilde{T} = \Lambda^\epsilon T$, $K_d = \frac{S_d}{(2\pi)^d}$

d=2.

$$\frac{dT}{dx} = (n-2)k_d T^2$$

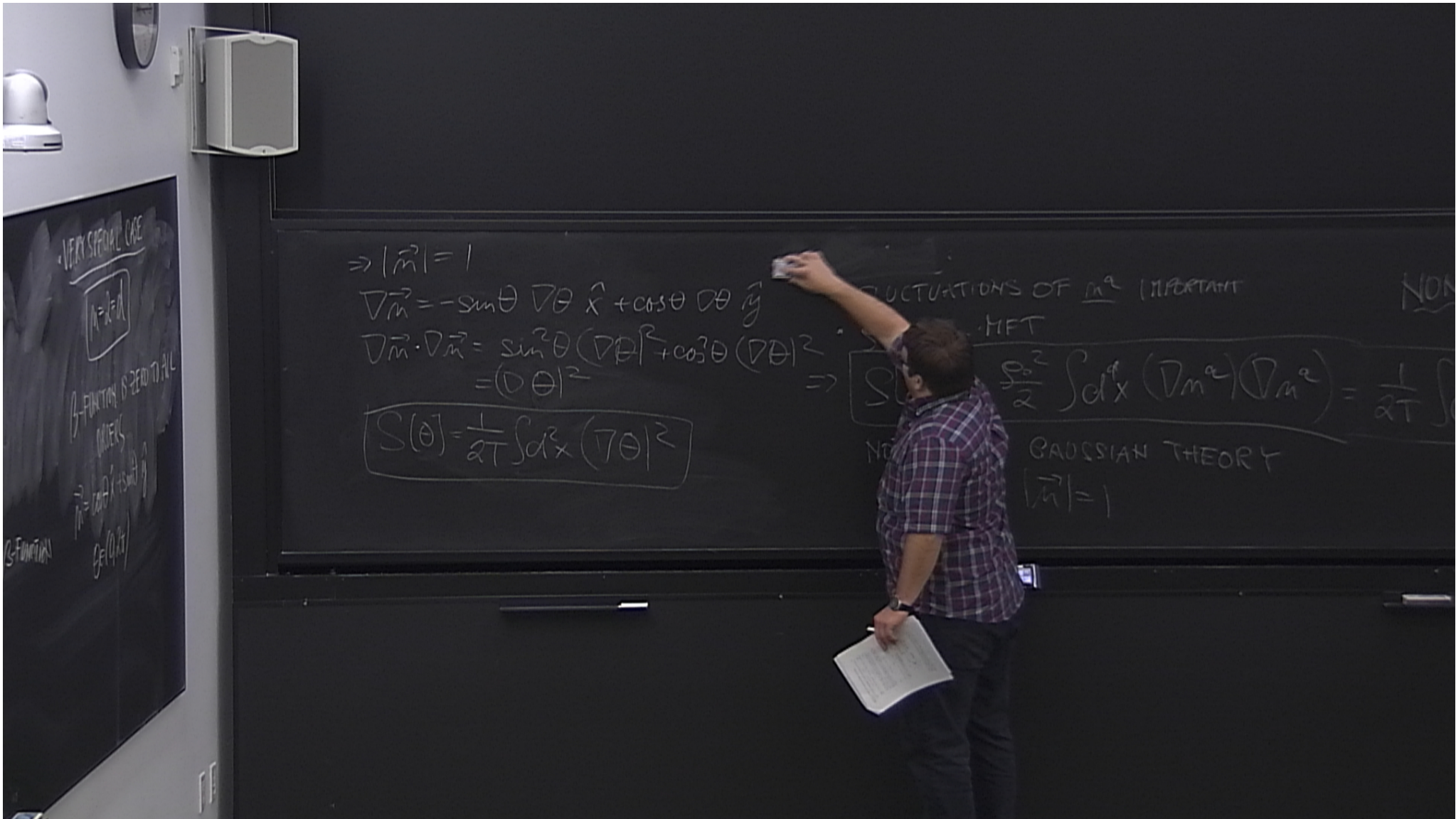
LACK OF ORDER



L-G

φ_a

$\ln^d \left(\frac{1}{2} \right)$
+



$$\Rightarrow |\vec{m}| = 1$$

$$\nabla_{\vec{m}} = -\sin\theta \nabla\theta \hat{x} + \cos\theta \nabla\theta \hat{y}$$

$$\nabla_{\vec{m}} \cdot \nabla_{\vec{m}} = \sin^2\theta (\nabla\theta)^2 + \cos^2\theta (\nabla\theta)^2 = (\nabla\theta)^2$$

$$S(\theta) = \frac{1}{2T} \int dx (\nabla\theta)^2$$

FLUCTUATIONS OF m^a (IMPORTANT)
MFT

$$\int \mathcal{D}m^a \frac{\rho_0^2}{2} \int dx (\nabla m^a)^2 = \frac{1}{2T} \int$$

GAUSSIAN THEORY
 $|\vec{m}| = 1$

VERY SPECIAL CASE
 $m = \hat{z}$
B-FUNCTION IS ZERO TO ALL
ORDER
 $\vec{m} = \cos\theta \hat{z} + \sin\theta \hat{y}$
 $\hat{z} = (0, 0, 1)$
 $\hat{y} = (0, 1, 0)$
B-FUNCTION

• LET'S PRETEND THAT Θ IS NOT COMPACTIFIED.

$$G(\vec{x}) = \left\langle e^{i\Theta(x)} e^{-i\Theta(0)} \right\rangle = = = \left(\frac{a}{x} \right)^{\frac{T}{2\pi}} \propto \frac{e^{-x/3}}{|x|^{d-2+\gamma}}$$

$\frac{dT}{dx} = 0$ $\beta \rightarrow \infty$ $\gamma = \frac{T}{2\pi}$

SEEMS LIKE A LINE OF FIXED POINTS
(CRITICAL POINTS)
BUT THERE IS NO LONG-RANGE ORDER

$$\langle e^{i\theta} \rangle = 0$$

$\frac{l^{-x/3}}{|x|^{d-2+\eta}}$
 POINTS
 $\langle \theta \rangle = M=0$

PHYSICALLY FOR LARGE TEMPERATURES
 THE CORRELATION LENGTH MUST BE FINITE.

$\Rightarrow \exists$ CRITICAL TEMPERATURE T_c

$T < T_c$
 $G(x) \sim \frac{1}{x^\eta}$
 $\xi = \infty$

$T > T_c$
 $G(x) \sim \frac{l^{-x/3}}{x^\eta}$
 $\xi \dots$ FINITE



PHYSICALLY FOR LARGE TEMPERATURES
THE CORRELATION LENGTH MUST BE FINITE.

$\Rightarrow \exists$ CRITICAL TEMPERATURE T_c

$T < T_c$

$$G(x) \sim \frac{1}{x^\eta}$$

$$\xi = \infty$$

$T > T_c$

$$G(x) \sim \frac{e^{-x/\xi}}{x^\eta}$$

$\xi \dots$ FINITE

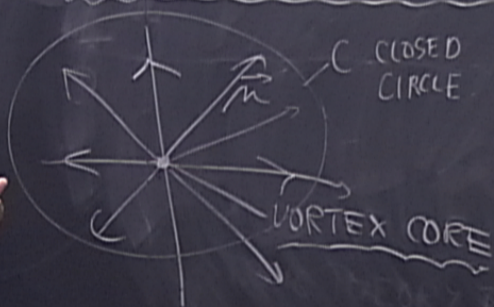
HAS NOTHING TO DO WITH LONG-RANGE ORDER

$$\frac{e^{-x/\xi}}{|x|^{d-2+\eta}}$$

D POINTS

$$\langle \theta \rangle = M = 0$$

• BY ASSUMING THAT Θ IS NON-COMPACT
 WE HAVE NEGLECTED A POSSIBILITY
 OF VORTEX-LIKE CONFIGURATIONS



$$\oint_C \nabla \Theta \cdot d\mathbf{l} = 2\pi m$$

WINDING NUMBER

• VERY SPECIAL CASE

$$m = 2 = d$$

β -FUNCTION IS ZERO TO ALL
 ORDERS.

$$\vec{m} = \cos \Theta \hat{x} + \sin \Theta \hat{y}$$

$$\Theta \in (0, 2\pi)$$

IS NON-COMPACT
AND A POSSIBILITY
OF CONFIGURATIONS

WINDING NUMBER

$$\oint_C \nabla \theta \cdot d\mathbf{l} = 2\pi m$$

C CLOSED
CIRCLE

VORTEX CORE

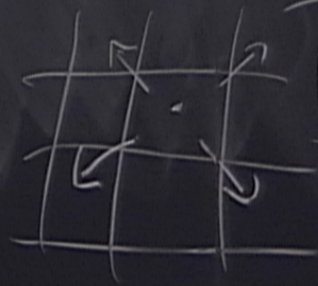
VORTEX CAN BE
DETECTED FROM FAR
AWAY

COSTS INFINITE ENERGY

EXAMPLE OF A TOPOLOGICAL
DEFECT

$$\nabla \theta \sim \frac{1}{r} \xrightarrow{r \rightarrow 0} \infty$$

CONTINUUM LIMIT BREAKS
DOWN NEAR VORTEX CORE



$$|\theta_i - \theta_j| = \frac{\pi}{2}$$

• WHY CAN WE HAVE VORTICES?

• E SINGLE VORTEX.

$$H = \frac{1}{2} \int d^2x |\nabla\theta|^2$$

WANT TO MINIMIZE
SUBJECT TO

$$\oint_C \nabla\theta \cdot d\ell = 2\pi n$$

HAVE VORTICES?
VORTEX.

$$\int d^2x |\nabla\theta|^2$$

MINIMIZE
CT TO

$$\oint \theta \cdot d\ell = 2\pi n$$

$$\delta H = \frac{1}{2} \int d^2x 2 \nabla\theta \cdot \nabla\delta\theta = - \int d^2x (\nabla^2\theta) \delta\theta = 0$$

$$\nabla^2\theta = 0$$

$$\Rightarrow \theta = m\varphi$$

POLAR ANGLE.

$$\Rightarrow \nabla\theta = \frac{m\hat{\varphi}}{r} \rightarrow$$

$$E = \frac{1}{2} \int d^2x |\nabla\theta|^2 = \frac{1}{2} \int r dr d\varphi \frac{m^2}{r^2}$$

$$= \pi m^2 \int_a^L \frac{dr}{r} = \pi m^2 \log \frac{L}{a}$$

↑ EXCLUDE THE CORE

• ACTUALLY WE WANT TO MINIMIZE
FREE ENERGY

$$F = E - TS$$

$$S \sim \log \# \text{POSSIBILITIES} \approx \log \left(\frac{L}{a} \right)^2 = 2 \log \frac{L}{a}$$

$$F = \left(\pi m^2 - 2T \right) \log \frac{L}{a}$$

• ACTUALLY WE WANT TO MINIMIZE
FREE ENERGY

$$F = E - TS$$

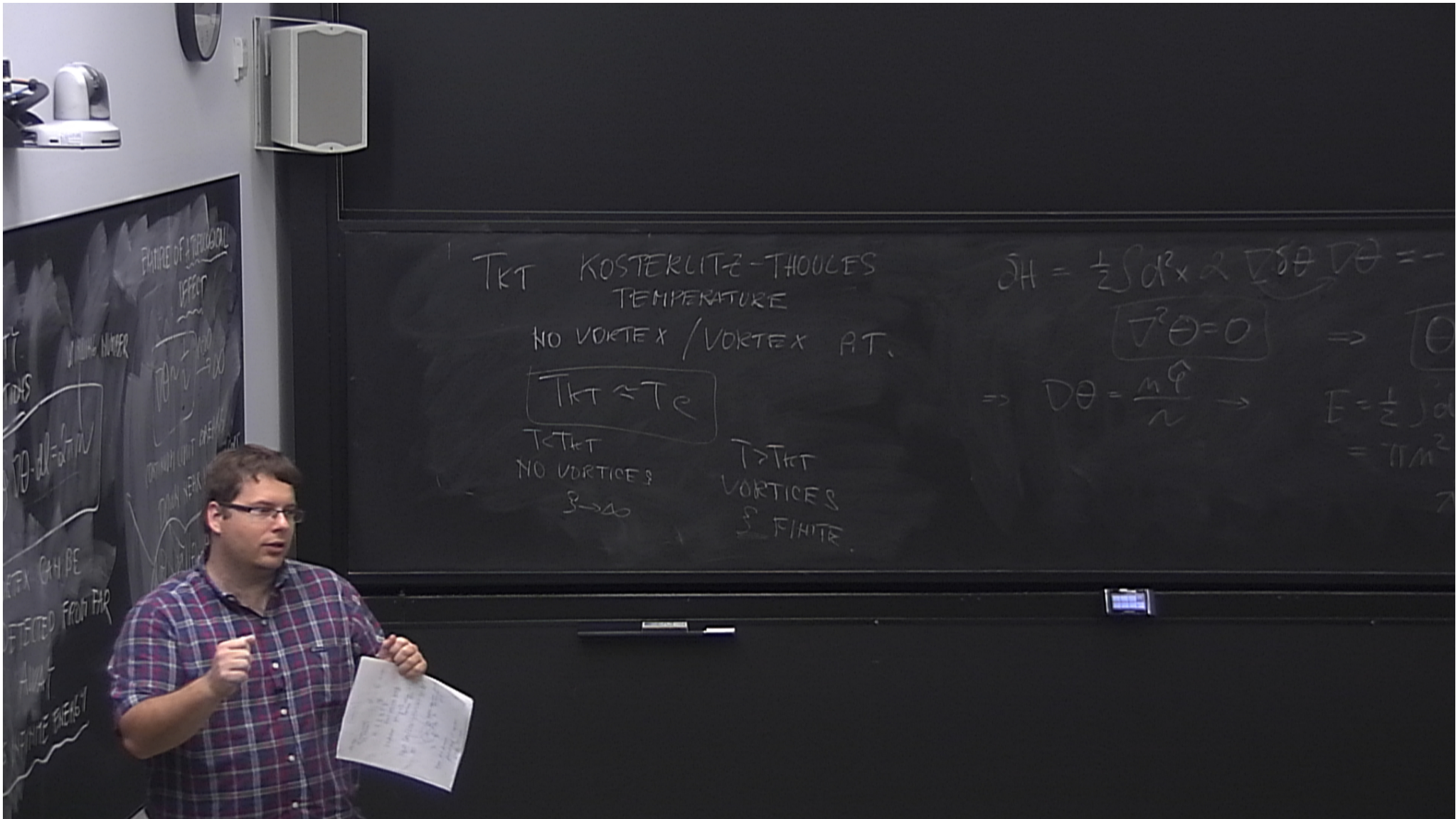
$$S \sim \log \# \text{POSSIBILITIES} \approx \log \left(\frac{L}{a} \right)^2 = 2 \log \frac{L}{a}$$

$$F = (\pi m^2 - 2T) \log \frac{L}{a}$$

$F < 0$.. VORTEX IS PREFERABLE

$$m=1$$

$$T > T_K = \frac{T}{2}$$



T_{KT} KOSTERLITZ-THOULES
TEMPERATURE

NO VORTEX / VORTEX P.T.

$$T_{KT} \approx T_c$$

$T < T_{KT}$
NO VORTICES
 $\xi \rightarrow \infty$

$T > T_{KT}$
VORTICES
 ξ FINITE

$$\delta H = \frac{1}{2} \int d^2x \left[\epsilon \nabla^2 \theta \right]^2 = -$$

$$\nabla^2 \theta = 0 \Rightarrow \theta$$

$$\Rightarrow \nabla \theta = \frac{n \hat{\phi}}{2} \rightarrow$$

$$E = \frac{1}{2} \int d^2x \left(\frac{1}{2} \right) = \pi n^2$$

T_{KT} KOSTERLITZ-THOULES
TEMPERATURE

TOPOLOGICAL

NO VORTEX / VORTEX P.T.

$$T_{KT} \leq T_c$$

$T < T_{KT}$
NO VORTICES
 $\xi \rightarrow \infty$

$T > T_{KT}$
VORTICES
 ξ FINITE DISTANCE BETWEEN V.

TOPOLOGICAL PHASE TRANSITION

2016 NOBEL PRIZE

FINAL REMARK $T=0$ NO VORTICES AT ALL
EVEN AT $T < T_{KT}$... CAN HAVE SOME.

DISTANCE BETWEEN V.

• ACTUALLY WE
FREE ENERGY

$$F = E - TS$$

$$S \sim \log \#$$

$$F = (\pi m^2 - 2T)$$

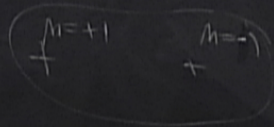
TOPOLOGICAL PHASE TRANSITION

2016 NOBEL PRIZE

FINAL REMARK $T=0$ NO VORTICES AT ALL

EVEN AT $T < T_{KT}$... CAN HAVE SOME

IN THE FORM OF VORTEX-ANTI VORTEX PAIRS



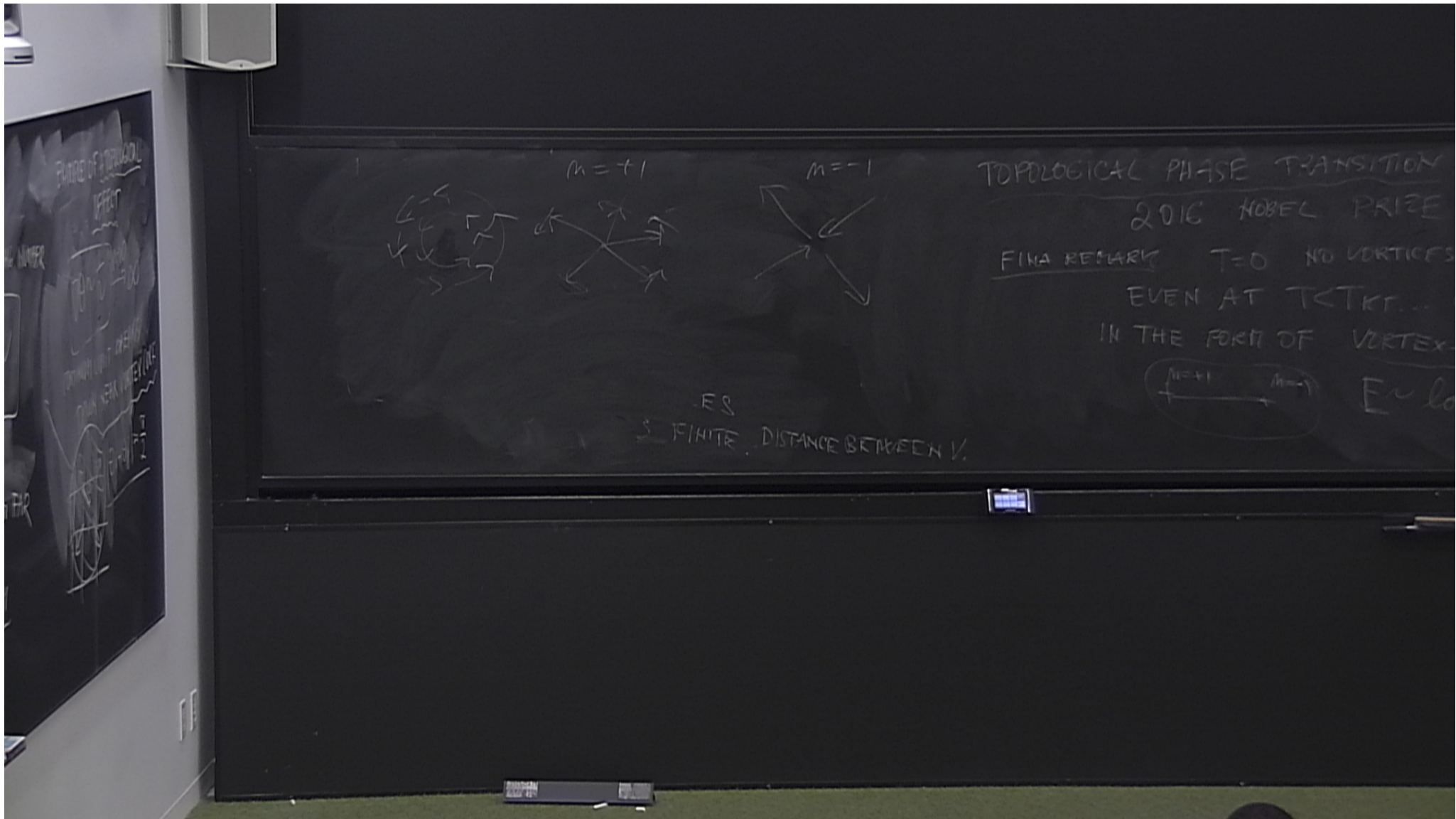
$$F \sim \log\left(\frac{x_i - x_j}{a}\right)$$

• ACTUALLY WE
FREE ENERGY

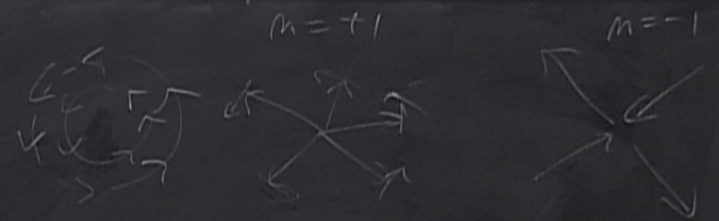
$$F = E - TS$$

$$S \sim \log \#$$

$$F = (\pi m^2 - 2T)$$



PHASE OF TOPOLOGY
EFFECT
TOPOLOGICAL ANOMALY
QUANTUM HALL EFFECT
1/2



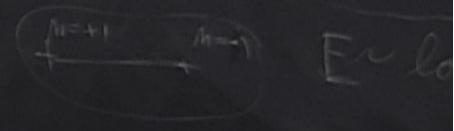
TOPOLOGICAL PHASE TRANSITION

2016 NOBEL PRIZE

FINA REMARK $T=0$ NO VORTICES

EVEN AT $T < T_c$...

IN THE FORM OF VORTEX



ES
FINITE DISTANCE BETWEEN V.