

Title: PSI 2016/2017 Quantum Theory - Lecture 8

Date: Sep 15, 2016 10:45 AM

URL: <http://pirsa.org/16090022>

Abstract:

Quantum Mechanics

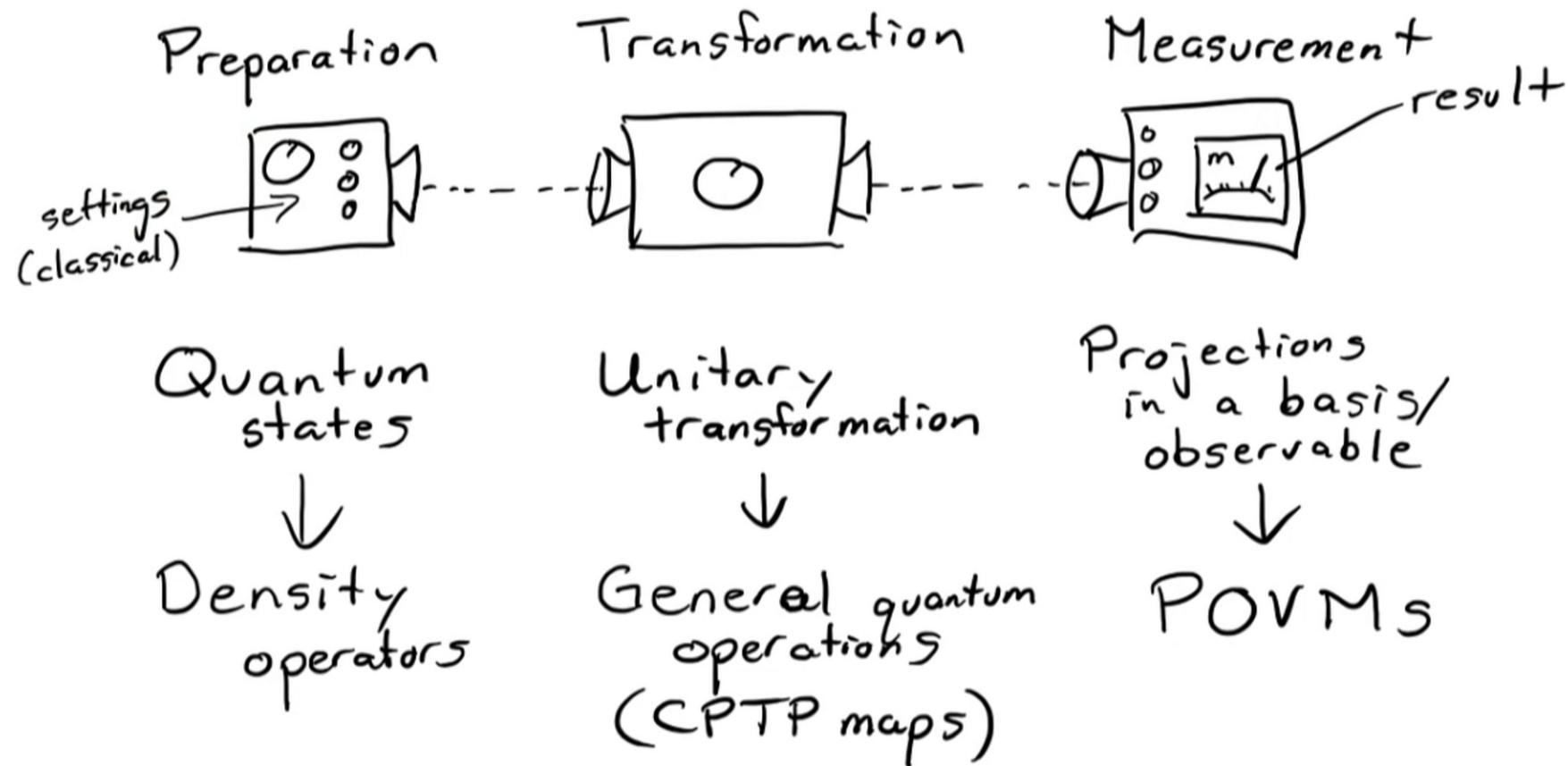
Quantum Operations and

Quantum Noise

Lecture 8



Operational quantum mechanics



Example: Projective measurement and more

$$P \xrightarrow{\quad} \boxed{P} \xrightarrow{\quad} \begin{matrix} \text{E} \\ \text{Im}^2 \end{matrix} \xrightarrow{\quad} P_m = \langle m | \rho | m \rangle$$

Many possible update rules

$$\text{Lüder's rule: } \rho \rightarrow P_m = |m\rangle\langle m|$$

Example: Projective measurement and more

$$\boxed{\rho} \xrightarrow{\quad} \boxed{\begin{matrix} \rho \\ \downarrow m \\ \rho' \end{matrix}} \xrightarrow{\quad}$$

Many possible update rules

$$\text{Lüder's rule: } \rho \rightarrow \rho_m = |m\rangle\langle m| \text{ with prob } p_m$$

Example: Projective measurement and more



ε_{lmz^3}

Many possible update rules

$$P_m = \langle m | \rho | m \rangle$$



not unitary

Lüder's rule: $\rho \rightarrow P_m = |m\rangle\langle m|$ with prob P_m

$$= \frac{|m\rangle\langle m|\rho|m\rangle\langle m|}{\langle m|\rho|m\rangle}$$

Example: Projective measurement and more

$$\boxed{P} \rightarrow \boxed{\text{update}} \quad P_m = \langle m | \rho | m \rangle$$

ε_{lmz^3}

Many possible update rules

✓ no⁺ unitary

Lüder's rule: $\rho \rightarrow P_m = |m\rangle\langle m|$ with prob P_m

Define measurement operators
 $\{M_m = |m\rangle\langle m|\}$ satisfying $\sum_m M_m^\dagger M_m = I$

$$\text{then } \rho \rightarrow P_m = \frac{M_m \rho M_m^\dagger}{\text{Tr}(M_m \rho M_m^\dagger)}$$

Example: Projective measurement and more

$$\boxed{P} \rightarrow \boxed{\text{update}} \quad P_m = \langle m | \rho | m \rangle$$

$\sum m_i^2 = 1$

Many possible update rules

✓ no⁺ unitary

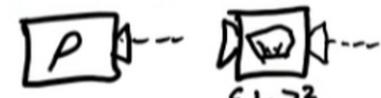
Lüders rule: $\rho \rightarrow P_m = |m\rangle\langle m|$ with prob P_m

Define measurement operators
 $\{M_m = |m\rangle\langle m|\}$ satisfying $\sum_m M_m^\dagger M_m = I$

then $\rho \rightarrow P_m = \frac{M_m \rho M_m^\dagger}{\text{Tr}(M_m \rho M_m^\dagger)}$

Note $\{M_m^\dagger M_m\}$ set of positive op.

Example: Projective measurement and more


 $\rho \rightarrow \rho_m = \langle m | \rho | m \rangle$

$\sum m \in \mathbb{Z}^3$

Many possible update rules

Lüders rule: $\rho \rightarrow \rho_m = |m\rangle\langle m|$ with prob p_m

Define measurement operators
 $\{M_m = |m\rangle\langle m|\}$ satisfying $\sum_m M_m^\dagger M_m = I$

then $\rho \rightarrow \rho_m = \frac{M_m \rho M_m^\dagger}{\text{Tr}(M_m \rho M_m^\dagger)}$

Note $\{M_m^\dagger M_m\}$ set of positive op. satisfy $\sum_m M_m^\dagger M_m = I \Rightarrow p_m$

Generalize: Given a POVM $\{E_m\}$ find a set $\{M_m\}$
 s.t. $E_m = M_m^\dagger M_m$

then $\rho \rightarrow \rho_m = \frac{M_m \rho M_m^\dagger}{\text{Tr}(M_m \rho M_m^\dagger)}$ is a possible update rule

Example: Projective measurement and more



ϵ_{lmz^3}
Many possible update rules

$$\rho_m = \langle m | \rho | m \rangle$$

✓ no + unitary

Lüder's rule: $\rho \rightarrow \rho_m = |m\rangle\langle m|$ with prob p_m

Define measurement operators

$$\{M_m = |m\rangle\langle m|\}$$
 satisfying $\sum_m M_m^\dagger M_m = I$

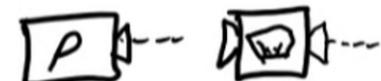
then $\rho \rightarrow \rho_m = \frac{M_m \rho M_m^\dagger}{\text{Tr}(M_m \rho M_m^\dagger)}$

Note $\{M_m^\dagger M_m\}$ set of positive op. satisfy $\sum_m M_m^\dagger M_m = I \Rightarrow p_m$

Generalize: Given a POVM $\{E_m\}$ find a set $\{M_m\}$
s.t. $E_m = M_m^\dagger M_m$ $p(m) = \text{Tr}(M_m \rho M_m^\dagger) = \text{Tr}(E_m \rho)$

then $\rho \rightarrow \rho_m = \frac{M_m \rho M_m^\dagger}{\text{Tr}(M_m \rho M_m^\dagger)}$ ✓ is a possible update rule

Example: Projective measurement and more

 $\rho \rightarrow \rho_m = \langle m | \rho | m \rangle$

$\sum m_i^2 = 1$

Many possible update rules

Lüders' rule: $\rho \rightarrow \rho_m = |m\rangle\langle m|$ with prob p_m

Define measurement operators
 $\{M_m = |m\rangle\langle m|\}$ satisfying $\sum_m M_m^\dagger M_m = I$

then $\rho \rightarrow \rho_m = \frac{M_m \rho M_m^\dagger}{\text{Tr}(M_m \rho M_m^\dagger)}$

Note $\{M_m^\dagger M_m\}$ set of positive op. satisfy $\sum_m M_m^\dagger M_m = I \Rightarrow p_m$

Generalize: Given a POVM $\{E_m\}$ find a set $\{M_m\}$
s.t. $E_m = M_m^\dagger M_m$ $p(m) = \text{Tr}(M_m \rho M_m^\dagger) = \text{Tr}(E_m \rho)$

then $\rho \rightarrow \rho_m = \frac{M_m \rho M_m^\dagger}{\text{Tr}(M_m \rho M_m^\dagger)}$ ✓ is a possible update rule

The University of Sydney "Forget" meas. result: $\rho \rightarrow \rho' = \sum_m p(m) \rho_m = \sum_m M_m \rho M_m^\dagger$

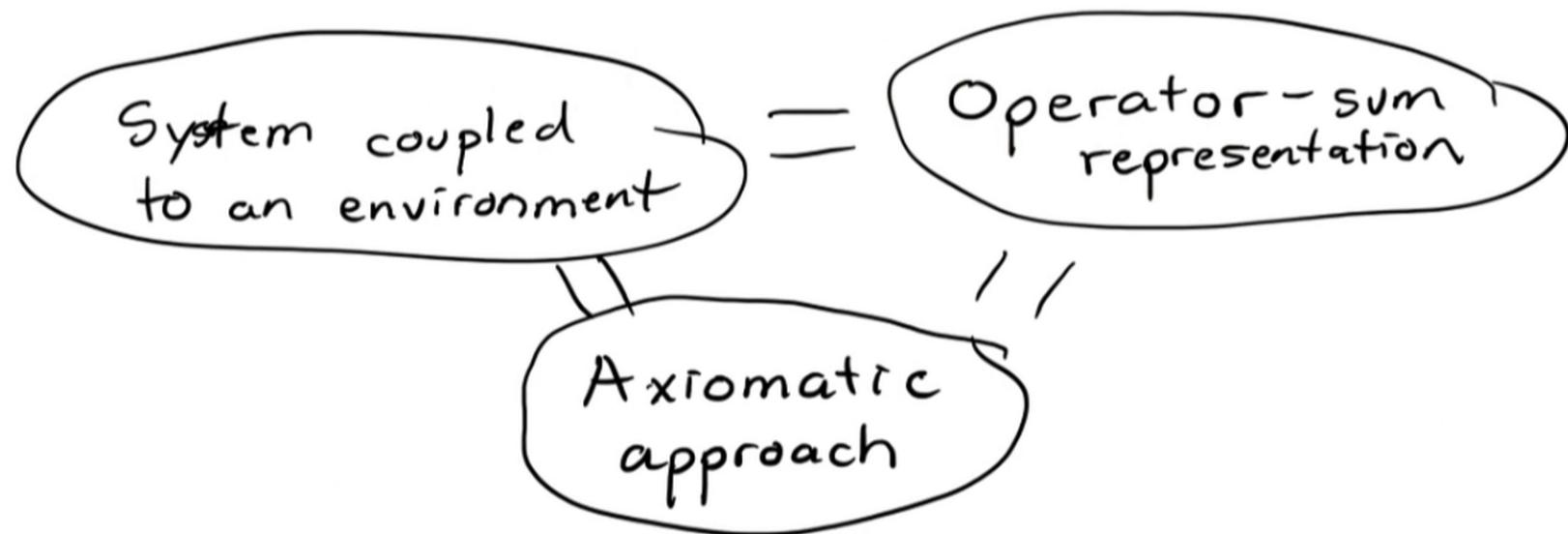
Three approaches to quantum operations

System coupled
to an environment

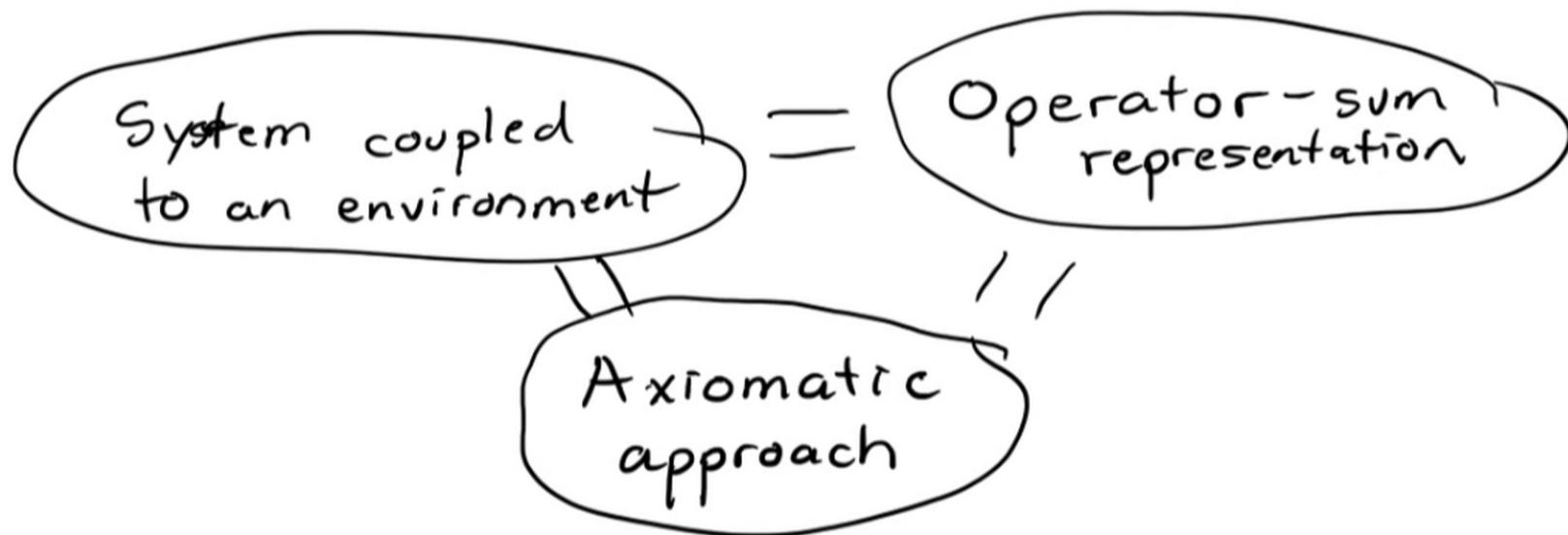
Operator-sum
representation

Axiomatic
approach

Three approaches to quantum operations



Three approaches to quantum operations



Quantum operation: $\rho' = \mathcal{E}(\rho)$

map from operators to
operators
satisfying some requirements

Page 5

Interaction with an environment

System : initial state ρ

Environment : initial state ρ_E

Interaction with an environment

System: initial state ρ

Environment: initial state ρ_E

How large?
Could be anything.
For gen. operation, needs
 z, d_{sys}^2

Interaction with an environment

System : initial state ρ

Environment : initial state ρ_E

System + Env. $\rho \otimes \rho_E$

How large?
Could be anything.
For gen. operation, needs
 z, d_{sys}^2

Interaction with an environment

System: initial state ρ

Environment: initial state ρ_E

How large?

Could be anything.

For gen. operation, needs

d_{sys}^2

System + Env. $\rho \otimes \rho_E$ uncorrelated

Unitary coupling U on sys + env

Interaction with an environment

System: initial state ρ

Environment: initial state ρ_E ← How large?
Could be anything.
For gen. operation, needs d_{sys}^2

System + Env. $\rho \otimes \rho_E$ uncorrelated

Unitary coupling U on sys + env

$$E(\rho) = \text{Tr}_E[U(\rho \otimes \rho_E) U^\dagger]$$

Operator-sum representation

Let's derive this from the coupling with the env.

Let E be fin. dim. Choose orthonormal basis $\{|e_k\rangle_E\}$

Let $\rho_{SE} = |e_0\rangle_E \langle e_0|$

$$E(\rho) = \sum_k \langle e_k | U(\rho_S \otimes |e_0\rangle_E \langle e_0|) U^\dagger |e_k\rangle_E$$

$$= \sum_k E_k \rho E_k^\dagger$$

where $E_k = \langle e_k | U_S | e_0 \rangle_E$ operators on system

Operator-sum representation

Let's derive this from the coupling with the env.

Let E be fin. dim. Choose orthonormal basis $\{|e_k\rangle_E\}$

Let $\rho_{SE} = |e_0\rangle_E \langle e_0|$

$$\begin{aligned}\mathcal{E}(\rho) &= \sum_k \langle e_k | U (\rho_{SE} | e_0 \rangle_E \langle e_0|) U^\dagger | e_k \rangle_E \\ &= \sum_k E_k \rho E_k^\dagger\end{aligned}$$

where $E_k = \langle e_k | U | e_0 \rangle_E$ operators on system

What properties does $\{E_k\}$ have?

$$\textcircled{1} \quad \mathcal{E}(\rho) \geq 0 \quad \textcircled{2} \quad \text{Tr}[\mathcal{E}(\rho)] = 1 = \text{Tr}(\sum_k E_k \rho E_k^\dagger) = \text{Tr}(\sum_k E_k^\dagger E_k \rho)$$

Operator-sum representation

Let's derive this from the coupling with the env.

Let E be fin. dim. Choose orthonormal basis $\{|e_k\rangle_E\}$

Let $\rho_{SE} = |e_0\rangle_E \langle e_0|$

$$\begin{aligned}\mathcal{E}(\rho) &= \sum_k \langle e_k | U (\rho_{SE} | e_0 \rangle_E \langle e_0|) U^\dagger | e_k \rangle_E \\ &= \sum_k E_k \rho E_k^\dagger\end{aligned}$$

where $E_k = \langle e_k | U_{SE} | e_0 \rangle_E$ operators on system

What properties does $\{E_k\}$ have?

$$\begin{aligned}① \mathcal{E}(\rho) &\geq 0 \quad ② \text{Tr}[\mathcal{E}(\rho)] = 1 = \text{Tr}(\sum_k E_k \rho E_k^\dagger) = \text{Tr}(\sum_k E_k^\dagger E_k \rho)\end{aligned}$$
$$\Rightarrow \sum_k E_k^\dagger E_k = I \quad \text{Trace preserving}$$

Operator-sum representation

Let's derive this from the coupling with the env.

Let E be fin. dim. Choose orthonormal basis $\{ |e_k\rangle_E \}$

Let $\rho_{\text{tot}} = |e_0\rangle_E \langle e_0|$

$$E(\rho) = \sum_k \langle e_k | U(\rho_{\text{tot}} |e_0\rangle_E \langle e_0|) U^\dagger |e_k\rangle_E$$

$$= \sum_k E_k \rho E_k^\dagger$$

where $E_k = \langle e_k | U_{SE} | e_0 \rangle_E$ operators on system

What properties does $\{ E_k \}$ have?

① $E(\rho) \geq 0$ ② $\text{Tr}[E(\rho)] = 1 = \text{Tr}(\sum_k E_k \rho E_k^\dagger) = \text{Tr}(\sum_k E_k^\dagger E_k \rho)$

$$\Rightarrow \sum_k E_k^\dagger E_k = I \quad \text{Trace preserving}$$

Definition: A map $E(\rho) = \sum_k E_k \rho E_k^\dagger$ is a quantum operation if $\sum_k E_k^\dagger E_k = I$

Physical interpretation of operator-sum representation

Consider coupling to an environment which is subsequently measured in the basis $|e_k\rangle$

$$\rho_k \propto \text{Tr}_E [|e_k\rangle\langle e_k| U (\rho \otimes |e_0\rangle\langle e_0|) U^\dagger |e_k\rangle\langle e_k|] = E_k \rho E_k^\dagger$$

Physical interpretation of operator-sum representation

Consider coupling to an environment which is subsequently measured in the basis $|e_k\rangle$

$$\rho_k \propto \text{Tr}_E [|e_k\rangle\langle e_k| U (\rho \otimes |e_0\rangle\langle e_0|) U^\dagger |e_k\rangle\langle e_k|] = E_k \rho E_k^\dagger$$

Normalize: $\rho_k = \frac{E_k \rho E_k^\dagger}{\text{Tr}[E_k \rho E_k^\dagger]}$

Prob. of outcome k occurring is

$$p(k) = \text{Tr} [|e_k\rangle\langle e_k| (U (\rho \otimes |e_0\rangle\langle e_0|) U^\dagger) |e_k\rangle\langle e_k|] \\ = \text{Tr} [E_k \rho E_k^\dagger]$$

The ensemble of outcomes is represented as a density operator

$$\mathcal{E}(\rho) = \sum_k p(k) \rho_k = \sum_k E_k \rho E_k^\dagger$$

Example: amplitude damping

Consider a 2-level system. Excited state $|1\rangle$
Ground state $|0\rangle$
If it's in the excited state $|1\rangle$, there is
some probability γ of transitioning to the g.s. $|0\rangle$
Qubit environment initially in $|0\rangle$ state.

$$(\alpha|0\rangle_s + \beta|1\rangle_s)|0\rangle_E \rightarrow \alpha|00\rangle_{SE} + \beta(\sqrt{1-\gamma}|10\rangle_{SE} + \sqrt{\gamma}|01\rangle_{SE})$$

Example: amplitude damping

Consider a 2-level system. Excited state $|1\rangle$
Ground state $|0\rangle$
If it's in the excited state $|1\rangle$, there is some probability γ of transitioning to the g.s. $|0\rangle$
Qubit environment initially in $|0\rangle$ state.

$$(\alpha|0\rangle_s + \beta|1\rangle_s)|0\rangle_E \rightarrow \alpha|00\rangle_{SE} + \beta(\sqrt{1-\gamma}|10\rangle_{SE} + \sqrt{\gamma}|01\rangle_{SE})$$

$$\mathcal{E}_{AD}(\rho) = E_0 \rho E_0^\dagger + E_1 \rho E_1^\dagger$$

$$E_0 = \begin{pmatrix} 1 & 0 \\ 0 & \sqrt{1-\gamma} \end{pmatrix} \quad E_1 = \begin{pmatrix} 0 & \sqrt{\gamma} \\ 0 & 0 \end{pmatrix}$$

Coupling to environment from the operator-sum representation

Start with $\{E_k\}$ s.t. $\sum_k E_k^\dagger E_k = I$

Choose env. with dimension equal to the cardinality of $\{E_k\}$,
pick a basis $\{|e_k\rangle\}$, inc. initial $|e_0\rangle$

Define U satisfying

$$U |f\rangle_s |e_0\rangle_E = \sum_k E_k |f\rangle_s |e_k\rangle_E \quad \forall |f\rangle_s$$

Coupling to environment from the operator-sum representation

Start with $\{E_k\}$ s.t. $\sum_k E_k^\dagger E_k = I$

Choose env. with dimension equal to the cardinality of $\{E_k\}$,
pick a basis $\{|e_k\rangle\}$, inc. initial $|e_0\rangle$

Define U satisfying

$$U|\psi\rangle_s |e_0\rangle_E = \sum_k E_k |\psi\rangle_s |e_k\rangle_E \quad \forall |\psi\rangle_s$$

Note $\forall |\psi\rangle_s, |\phi\rangle_s$

$$\langle \psi | \langle e_0 | U^\dagger U |\phi\rangle |e_0\rangle = \sum_k \langle \psi | E_k^\dagger E_k |\phi\rangle = \langle \psi | \phi \rangle$$

Coupling to environment from the operator-sum representation

Start with $\{E_k\}$ s.t. $\sum_k E_k^\dagger E_k = I$

Choose env. with dimension equal to the cardinality of $\{E_k\}$,
pick a basis $\{|e_k\rangle\}$, inc. initial $|e_0\rangle$

Define U satisfying

$$U|\psi\rangle_s |e_0\rangle_E = \sum_k E_k |\psi\rangle_s |e_k\rangle_E \quad \forall |\psi\rangle_s$$

Note $\forall |\psi\rangle_s, |\phi\rangle_s$

$$\langle \psi | \langle e_0 | U^\dagger U |\phi\rangle |e_0\rangle = \sum_k \langle \psi | E_k^\dagger E_k |\phi\rangle = \langle \psi | \phi \rangle$$

Then U can be extended to a unitary op. on $S+E$

Then

$$\text{Tr}_E(U(\rho \otimes |e_0\rangle\langle e_0|)U^\dagger) = \sum_k E_k \rho E_k^\dagger$$

Freedom in the operator-sum representation

Can $\{E_k\}$ and $\{F_k\}$ be such that $\mathcal{E} = \tilde{\mathcal{F}}$?
In "coupling with env" picture,

Freedom in the operator-sum representation

Can $\{E_k\}$ and $\{F_k\}$ be such that $E = F$?

In "coupling with env" picture, a final unitary on env. (or equiv. a change of basis $\{|e_k\rangle\}$) changes the op. sum rep but not the q. operation.

$$\begin{aligned} E_k &= \langle e_k | U | e_0 \rangle & F_k &= \langle e_k | (I \otimes U') U | e_0 \rangle \\ & & &= \sum_j [I \otimes \langle e_k | U' | e_j \rangle] \langle e_j | U | e_0 \rangle \\ & & &= \sum_j U'_{kj} E_k \end{aligned}$$

Freedom in the operator-sum representation

Can $\{E_k\}$ and $\{F_k\}$ be such that $E = F$?

In "coupling with env" picture, a final unitary on env. (or equiv. a change of basis $\{|e_k\rangle\}$) changes the op. sum rep but not the q. operation.

$$\begin{aligned} E_k &= \langle e_k | U | e_0 \rangle & F_k &= \langle e_k | (I \otimes U') U | e_0 \rangle \\ & & &= \sum_j [I \otimes \langle e_k | U' | e_j \rangle] \langle e_j | U | e_0 \rangle \\ & & &= \sum_j U'_{kj} E_{kj} \end{aligned}$$

Thm: $E = F$ iff \exists $m \times m$ unitary U' s.t. $\sum_k F_k = \sum_j U'_{kj} E_j$

Freedom in the operator-sum representation

Can $\{E_k\}$ and $\{F_k\}$ be such that $E = F$?

In "coupling with env" picture, a final unitary on env. (or equiv. a change of basis $\{|e_k\rangle\}$) changes the op. sum rep but not the q. operation.

$$\begin{aligned} E_k &= \langle e_k | U | e_0 \rangle & F_k &= \langle e_k | (I \otimes U') U | e_0 \rangle \\ & & &= \sum_j [I \otimes \langle e_k | U' | e_j \rangle] \langle e_j | U | e_0 \rangle \\ & & &= \sum_j U'_{kj} E_{kj} \end{aligned}$$

Thm: $E = F$ iff \exists $m \times m$ unitary U' s.t. $F_k = \sum_j U'_{kj} E_j$
(Note: pad out smaller sets with zero operators)