

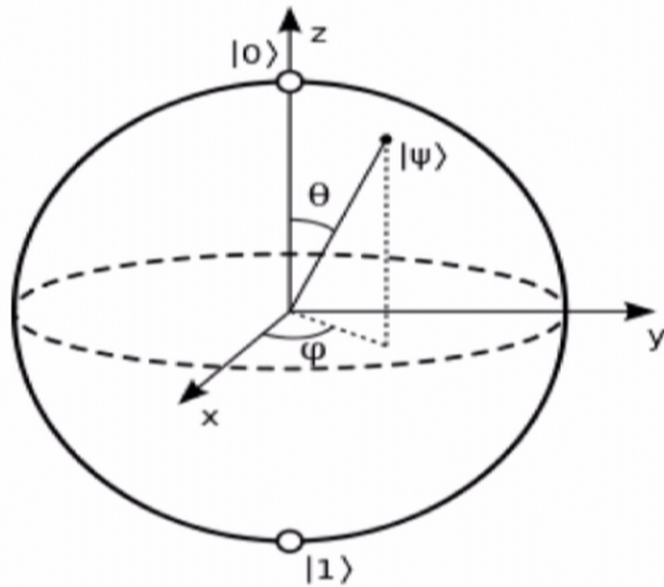
Title: PSI 2016/2017 Quantum Theory - Lecture 4

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Abstract:

The Bloch sphere for mixed states of a qubit



Note $\{\sigma_0, \sigma_1, \sigma_2, \sigma_3\}$ form an orthonormal basis of operators

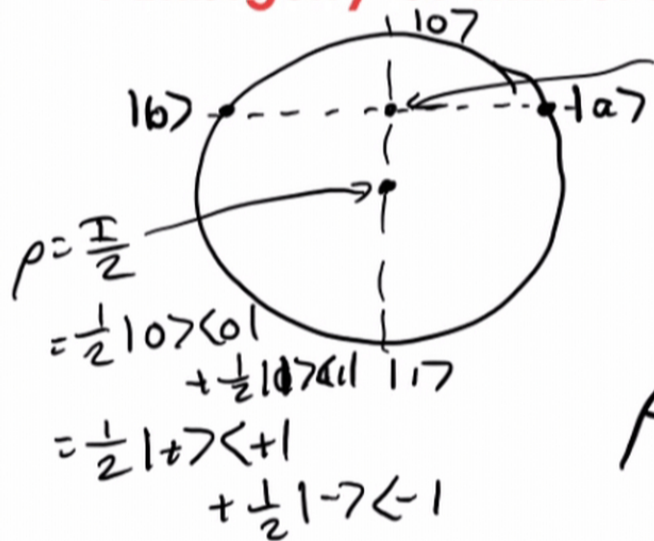
$$\text{Tr}[\sigma_i \sigma_j] = 2\delta_{ij} \quad \vec{\sigma} = (X, Y, Z)$$

$$\rho = \frac{1}{2}(\mathbf{I} + \vec{r} \cdot \vec{\sigma}) \quad |\vec{r}| \leq 1$$

$$\text{So } \text{Tr}[\rho] = 1$$

$$\text{Tr}[\rho \vec{\sigma}] = \vec{r}$$

Ambiguity of mixtures



$$\rho = \frac{3}{4} |0\rangle\langle 0| + \frac{1}{4} |1\rangle\langle 1|$$

Define $|a\rangle = \sqrt{\frac{3}{4}} |0\rangle + \sqrt{\frac{1}{4}} |1\rangle$

$$|b\rangle = \sqrt{\frac{3}{4}} |0\rangle - \sqrt{\frac{1}{4}} |1\rangle$$

$$\begin{aligned} \rho' &= \frac{1}{2} |a\rangle\langle a| + \frac{1}{2} |b\rangle\langle b| \\ &= \frac{3}{4} |0\rangle\langle 0| + \frac{1}{4} |1\rangle\langle 1| \end{aligned}$$

* Ensembles with the same ρ are indistinguishable by any measurement
 No privileged ensemble

Many such equivalent decompositions

The reduced density operator

Two quantum systems A & B , joint system AB

Density operator ρ_{AB} of joint system. How do I describe just A ?

$$\rho_A = \text{Tr}_B[\rho_{AB}] \quad \text{Tr}_B \text{ partial trace}$$

Defined by

$$\text{Tr}_B[|a_1\rangle_A \langle a_2| \otimes |b_1\rangle_B \langle b_2|] = |a_1\rangle_A \langle a_2| \text{Tr}(|b_1\rangle_B \langle b_2|) = |a_1\rangle_A \langle a_2| \langle b_2 | b_1 \rangle$$

and extend by linearity to full operator space

$$\text{So } \text{Tr}_B[\rho_A \otimes \rho_B] = \rho_A (\text{Tr}(\rho_B)) = \rho_A$$

Why the partial trace?

For any POVM $\{E_m^A\}$ on just system A

$$\begin{aligned} p_m &= \text{Tr}[(E_m^A \otimes I_B) \rho_{AB}] \\ &= \text{Tr}(E_m^A \rho_A) \end{aligned}$$

$$\text{with } \rho_A = \text{Tr}_B(\rho_{AB})$$

Reduced density operators of Bell states

What about $\rho_{AB} = |\Psi_{00}\rangle\langle\Psi_{00}|$ $|\Psi_{00}\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$

$$\begin{aligned} \rho_A &= \text{Tr}_B \left(\frac{1}{2} |00\rangle\langle 00| + \frac{1}{2} |11\rangle\langle 00| + \frac{1}{2} |00\rangle\langle 11| + \frac{1}{2} |11\rangle\langle 11| \right) \\ &= \frac{1}{2} |0\rangle_A\langle 0| \langle 0|0\rangle + \frac{1}{2} |1\rangle_A\langle 0| \langle 0|1\rangle + \frac{1}{2} |0\rangle_A\langle 1| \langle 1|0\rangle + \frac{1}{2} |1\rangle_A\langle 1| \langle 1|1\rangle \\ &= \frac{1}{2} (|0\rangle_A\langle 0| + |1\rangle_A\langle 1|) = \frac{1}{2} \mathbf{I}_A \quad \text{completely mixed state} \end{aligned}$$

Bob measure his system B in Z basis, gets outcome $|0\rangle$ (+1)
 If Alice measures her system A in Z basis, will get $|0\rangle$ (+1) with certainty
 Which description should Alice use?

Depends on her knowledge of Bob's meas.
 $\rho_A = \frac{1}{2} \mathbf{I}_A$ or $\rho_A = |0\rangle_A\langle 0|$

If Bob measured in X basis, gets $|+\rangle_B$ (+1)
 $\rho_A = \frac{1}{2} \mathbf{I}_A$ or $\rho_A = |-\rangle_A\langle -|$

The Schmidt decomposition

Suppose $|\psi\rangle_{AB}$ is a pure state of a composite system AB .
Then there exist orthonormal states $|i\rangle_A$ for A , $|i\rangle_B$ for B ,
s.t. Schmidt coefficients

$$|\psi\rangle_{AB} = \sum_i \sqrt{\lambda_i} |i\rangle_A |i\rangle_B$$

where $\lambda_i \geq 0$ real nonneg. numbers satisfying $\sum_i \lambda_i = 1$.

Proof: Pick any bases $|j\rangle_A$; $|k\rangle_B$

$$|\psi\rangle_{AB} = \sum_{j,k} a_{j,k} |j\rangle_A |k\rangle_B \quad \text{SVD on } A = UDV$$

Define $|i\rangle_A = \sum_j u_{ji} |j\rangle_A$ $|i\rangle_B = \sum_k v_{ik} |k\rangle_B$ and $\sqrt{\lambda_i} = d_i$
↑ unitaries
diagonal, not neg.
 orthonormal due to unitarity of u, v

$$\rho_A = \text{Tr}_B(\rho_{AB}) = \sum_i \lambda_i |i\rangle_A \langle i| \quad \rho_B = \text{Tr}_A(\rho_{AB}) = \sum_i \lambda_i |i\rangle_B \langle i|$$

Purifications

Given ρ_A , find orthogonal decomp. $\rho_A = \sum_i p_i |i\rangle_A \langle i|$

Introduce new system R , choose any $r = \text{rank}(\rho_A)$ orthogonal states $|i\rangle_R$

$$|\Psi\rangle_{AR} = \sum_i \sqrt{p_i} |i\rangle_A |i\rangle_R \quad \text{satisfies}$$
$$\text{Tr}_R [|\Psi\rangle_{AR} \langle \Psi|] = \rho_A$$

Different choices of basis on $R \Leftrightarrow$ unitary U_R

$$|k\rangle_R = \sum_i U_{ki} |i\rangle_R$$

$|\Psi\rangle_{AR}$ and $|\Psi'\rangle_{AR} = (I_A \otimes U_R) |\Psi\rangle_{AR}$
different purifications both satisfying

$$\rho_A = \text{Tr}_R [|\Psi\rangle_{AR} \langle \Psi|] = \text{Tr}_R [|\Psi'\rangle_{AR} \langle \Psi'|]$$