

Title: Bootstrapping 3D CFTs

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Abstract: I will review recent results from applying the conformal bootstrap to 3D CFTs, including precise determinations of critical exponents and in the 3D Ising and $O(N)$ vector models, new constraints on 3D Gross-Neveu models, and general bounds on correlation function coefficients of currents and stress tensors.

Bootstrapping 3D CFTs

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August 26, 2016

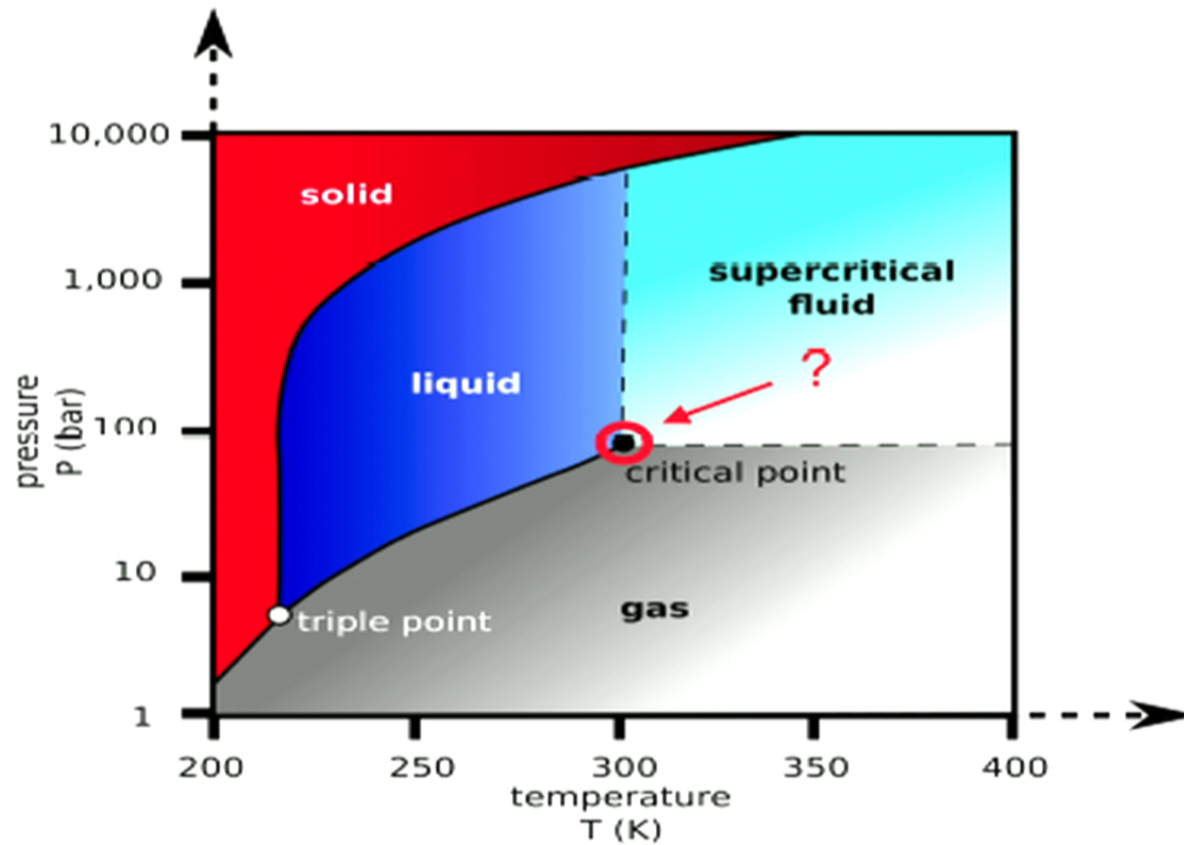
Low Energy Challenges for High Energy Physicists II

Why Study CFTs?

There are many interesting applications of conformal field theories:

- ▶ **2D**: String Theory
- ▶ **2D/3D**: Statistical and Condensed Matter Systems
- ▶ **4D**: Scenarios for Physics Beyond the Standard Model
- ▶ **6D**: Mysterious $(2, 0)$ Theory
- ▶ **Holography and AdS/CFT**: Study Quantum Gravity with CFTs

Phase Diagram for Fluids

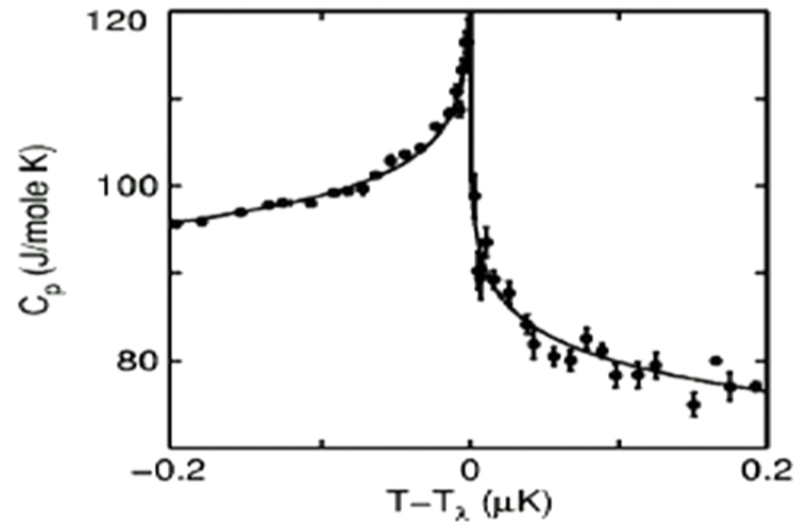


$$3\text{D Ising CFT} \leftrightarrow \mathcal{L}_{\text{Ising}} \sim \lambda_c \phi^4$$

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Bootstrapping 3D CFTs

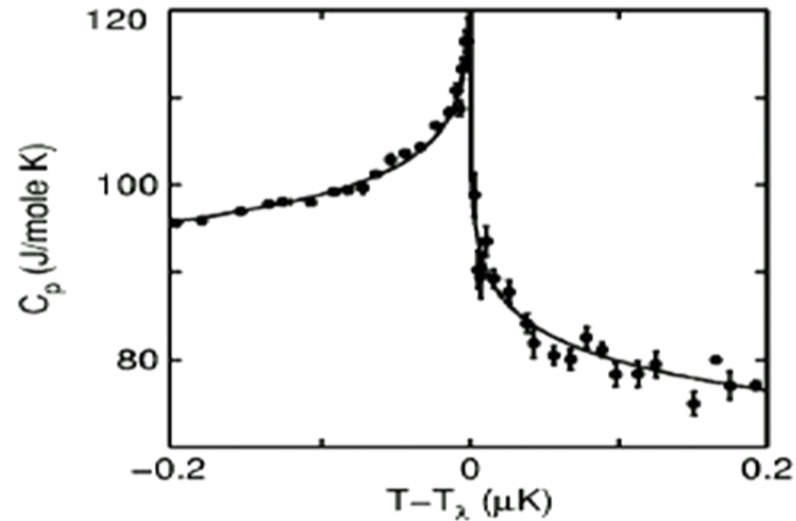
3D O(N) Models



Going to N scalars is also interesting! $\rightarrow \mathcal{L}_{O(N)} \sim \lambda_c (\phi_i \phi^i)^2$

- ▶ $N = 2$: Superfluid (λ) transition in ^4He [Lipa et al, '96; '03]
Quantum critical point in (2+1)D superconductors

3D $O(N)$ Models



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Quantum critical point in (2+1)D superconductors
- ▶ $N = 3$: Isotropic ferromagnets (Fe, Co, Ni, ...)
- ▶ Large N : Solvable in $1/N$ expansion

Main Goal

We would like to map out the space of CFTs and predict their observables

Conformal Bootstrap

- ▶ The **conformal bootstrap** aims to use **mathematical consistency** conditions to map out and solve the space of CFTs
 - ▶ Conformal Symmetry
 - ▶ Crossing Symmetry
 - ▶ Unitarity / Reflection Positivity

Conformal Bootstrap

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 - ▶ Conformal Symmetry
 - ▶ Crossing Symmetry
 - ▶ Unitarity / Reflection Positivity
- ▶ Beautiful success story in 2D
[Ferrara, Gatto, Grillo '73; Polyakov '74; Belavin, Polyakov, Zamolodchikov '83]
- ▶ Exciting progress in $D > 2$ starting in 2008
[Rattazzi, Rychkov, Tonni, Vichi '08; ...]

Conformal Block Expansion

Can probe spectrum by expanding 4-point functions in **conformal blocks**:

$$\langle \overline{\sigma(x_1)\sigma(x_2)} \overline{\sigma(x_3)\sigma(x_4)} \rangle = \frac{1}{x_{12}^{2\Delta_\sigma} x_{34}^{2\Delta_\sigma}} \sum_{\Delta, \ell} \lambda_{\mathcal{O}}^2 g_{\Delta, \ell}(u, v)$$

- ▶ Blocks $g_{\Delta, \ell}(u, v)$ are known functions of $u = \frac{x_{12}^2 x_{34}^2}{x_{13}^2 x_{24}^2}$, $v = \frac{x_{14}^2 x_{23}^2}{x_{13}^2 x_{24}^2}$
- ▶ Similar to expansion in spherical harmonics Y_ℓ^m , but for CFTs

Scalar Conformal Blocks

Explicit formulas in even D [Dolan, Osborn '00; '03]:

$$g_{\Delta, \ell}^{2D}(u, v) = k_{\Delta+\ell}(z)k_{\Delta-\ell}(\bar{z}) + z \leftrightarrow \bar{z}$$

$$g_{\Delta, \ell}^{4D}(u, v) = \frac{z\bar{z}}{z-\bar{z}}[k_{\Delta+\ell}(z)k_{\Delta-\ell-2}(\bar{z}) - z \leftrightarrow \bar{z}]$$

$$k_{\beta}(x) = x^{\beta/2} {}_2F_1(\beta/2, \beta/2, \beta; x)$$

where $u = z\bar{z}$ and $v = (1-z)(1-\bar{z})$

- ▶ Conformal blocks are eigenfunctions of $SO(D+1, 1)$ Casimir
- ▶ Outside of even D , can be computed recursively to arbitrary precision [El-Showk, Paulos, DP, Rychkov, Simmons-Duffin, Vichi; Kos, DP, Simmons-Duffin '13; '14]

Crossing Symmetry

$\langle \sigma(x_1)\sigma(x_2)\sigma(x_3)\sigma(x_4) \rangle$ is symmetric under permutations of x_i :

- ▶ Switching $x_1 \leftrightarrow x_3$ gives the **crossing symmetry** condition:

$$\sum \begin{array}{c} 1 \quad 4 \\ \diagdown \quad / \\ \text{O} \\ / \quad \diagdown \\ 2 \quad 3 \end{array} = \sum \begin{array}{c} 1 \quad 4 \\ / \quad \diagdown \\ \text{O} \\ \diagdown \quad / \\ 2 \quad 3 \end{array}$$

$$u^{-\Delta_\sigma} \sum_{\Delta, \ell} \lambda_{\text{O}}^2 g_{\Delta, \ell}(u, v) = v^{-\Delta_\sigma} \sum_{\Delta, \ell} \lambda_{\text{O}}^2 g_{\Delta, \ell}(v, u)$$

- ▶ Only **unknowns** are set of scaling dimensions and coefficients: $\{\Delta, \lambda_{\text{O}}\}$

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Numerical Approach

- ▶ By applying clever linear functionals α one can prove that some assumptions on $\{\Delta, \lambda_{\mathcal{O}}\}$ are **incompatible** with **crossing + unitarity**:

$$0 = \sum_{\Delta, \ell} \lambda_{\mathcal{O}}^2 \alpha [u^{-\Delta_\sigma} g_{\Delta, \ell}(u, v) - v^{-\Delta_\sigma} g_{\Delta, \ell}(v, u)] > 0$$

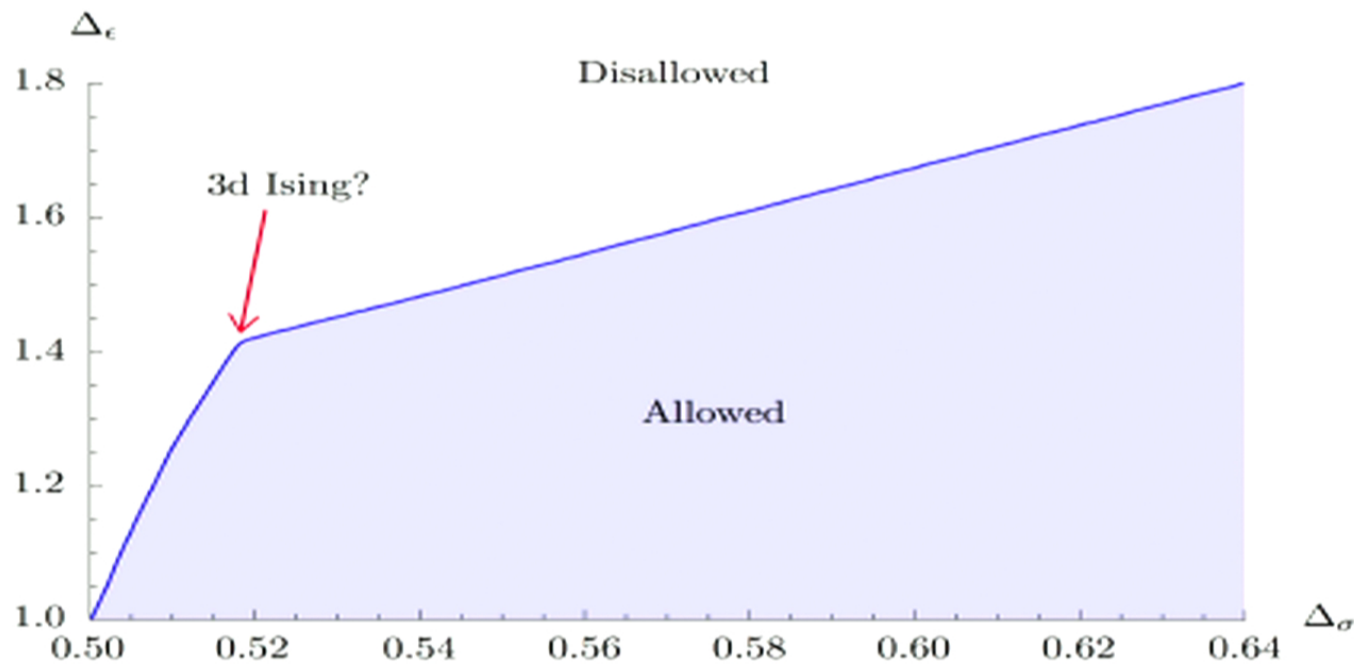
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- ▶ Find $\alpha \sim \sum_{mn} a_{mn} \partial_z^m \partial_{\bar{z}}^n |_{1/2, 1/2}$ using linear/semidefinite programming [Rattazzi, Rychkov, Tonni, Vichi '08; DP, Simmons-Duffin, Vichi '11]
 - ▶ Functional search space ranges from ~ 20 to ~ 1200 components
 - ▶ Each plot \leftrightarrow Solve $\mathcal{O}(1000)$ optimization problems on HPC clusters
 - ▶ State of the art: SDPB [Simmons-Duffin '15]

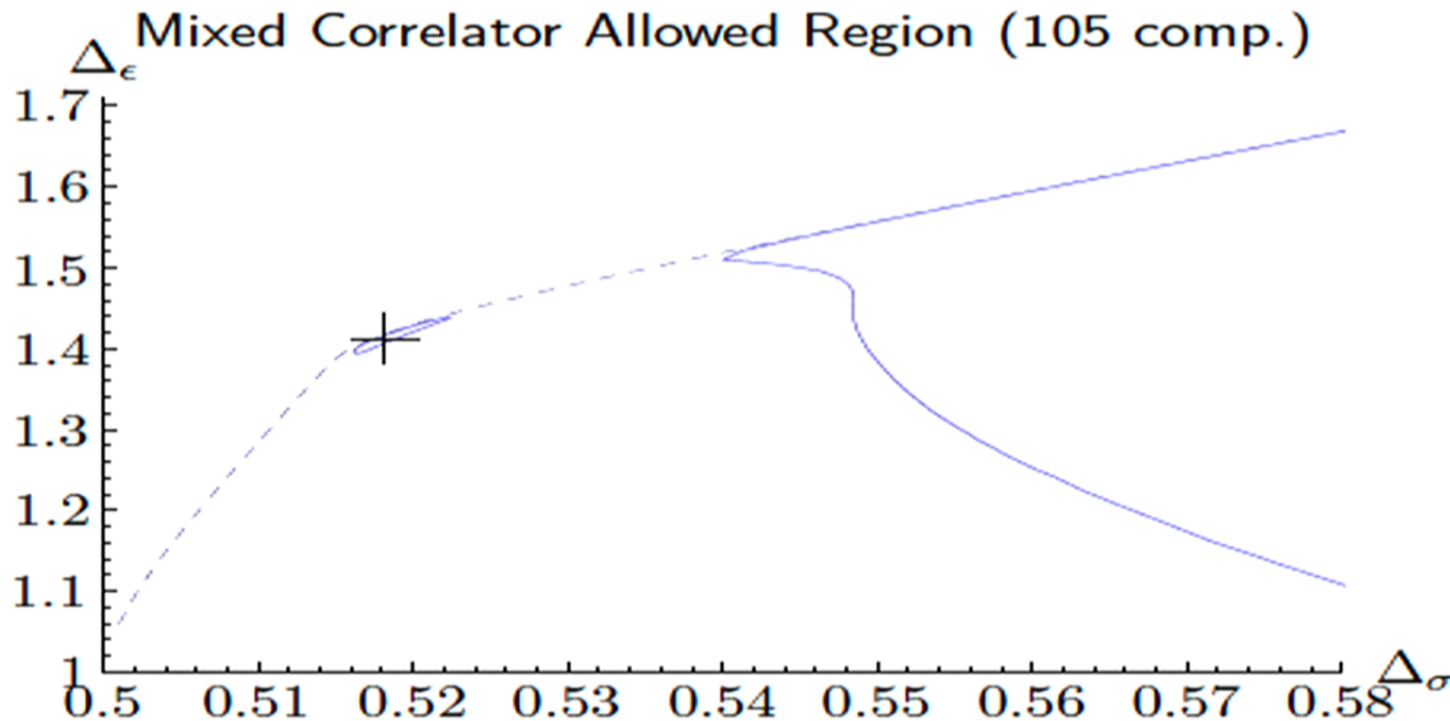
3D Dimension Bounds



[El-Showk, Paulos, DP, Rychkov, Simmons-Duffin, Vichi, '12; '14]

- ▶ Bound on leading scalar in $\sigma \times \sigma \sim \mathbb{1} + \epsilon + \dots$
- ▶ 3D Ising (Lattice): $\Delta_\sigma \simeq 0.51813(5)$, $\Delta_\epsilon \simeq 1.41275(25)$ [Hasenbusch '10]
(Leading critical exponents: $\eta = 2\Delta_\sigma - 1$, $\nu = \frac{1}{3-\Delta_\epsilon}$)

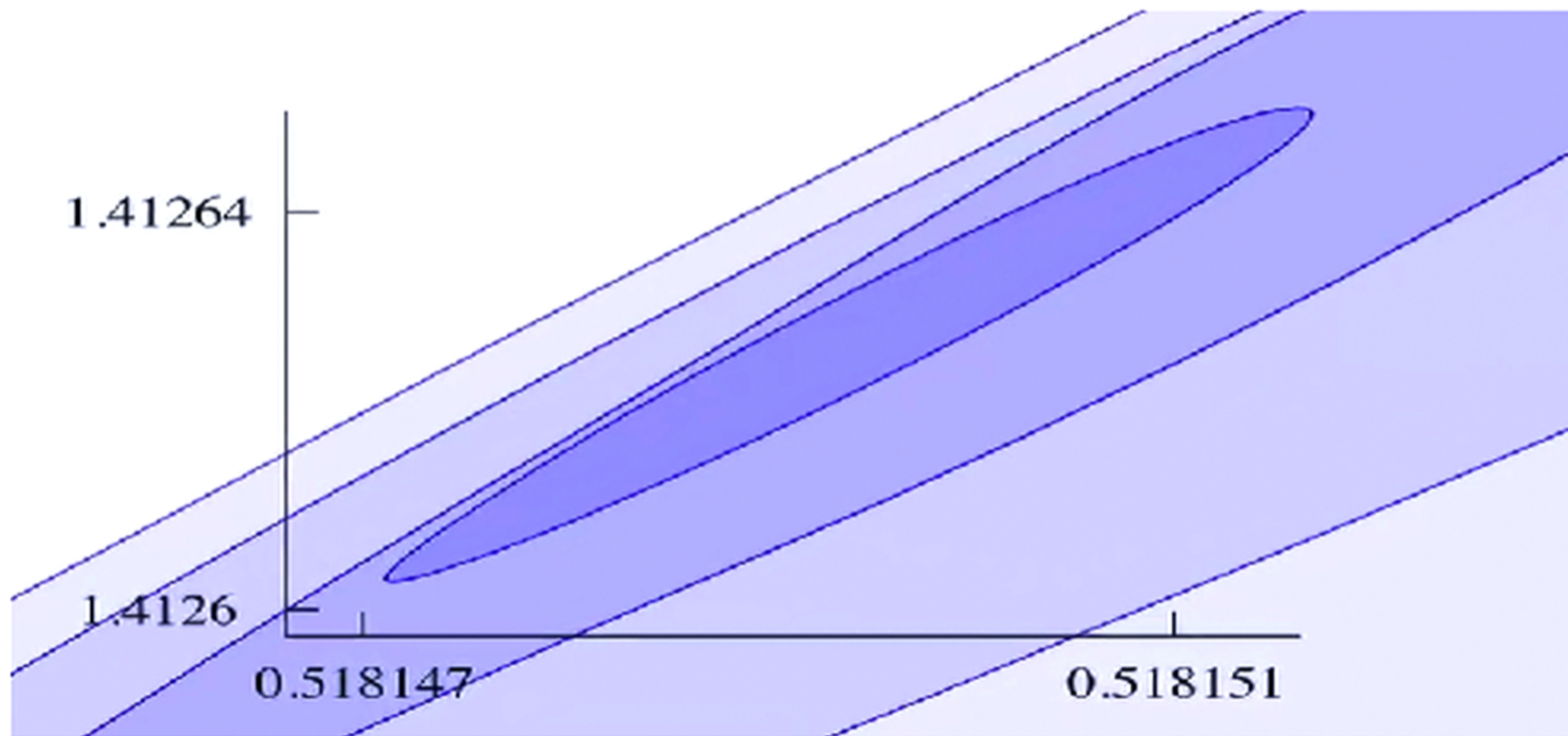
Mixed Correlator Islands



[Kos, DP, Simmons-Duffin '14]

- ▶ Combining constraints from $\langle \sigma\sigma\sigma\sigma \rangle$, $\langle \sigma\sigma\epsilon\epsilon \rangle$, $\langle \epsilon\epsilon\epsilon\epsilon \rangle$, can impose that σ and ϵ are only **relevant** ($\Delta < 3$) operators, yielding a **rigorous** island!

Mixed Correlator Islands

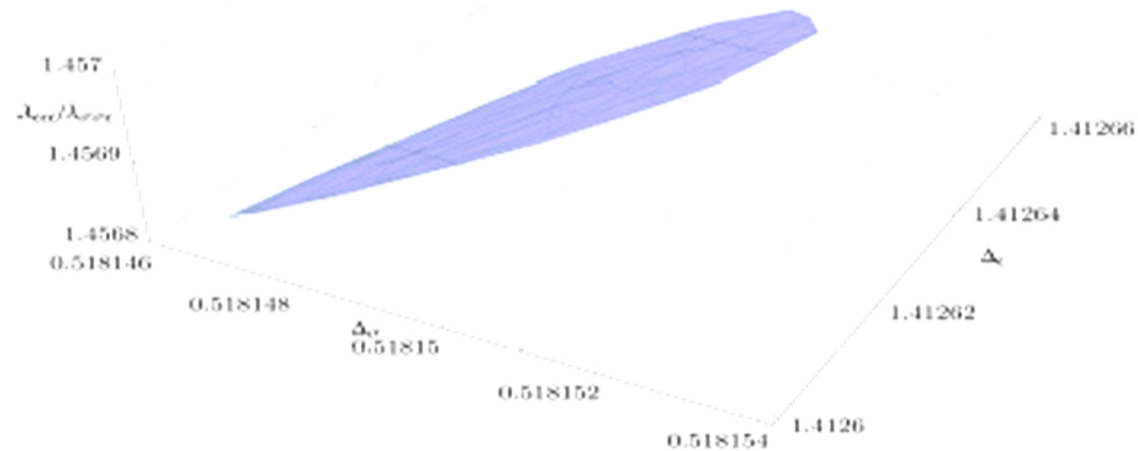


[Kos, DP, Simmons-Duffin '14; Simmons-Duffin '15; Kos, DP, Simmons-Duffin, Vichi '16]

- Increasing to 1265 components using SDPB, region keeps shrinking!

Mixed Correlator Islands

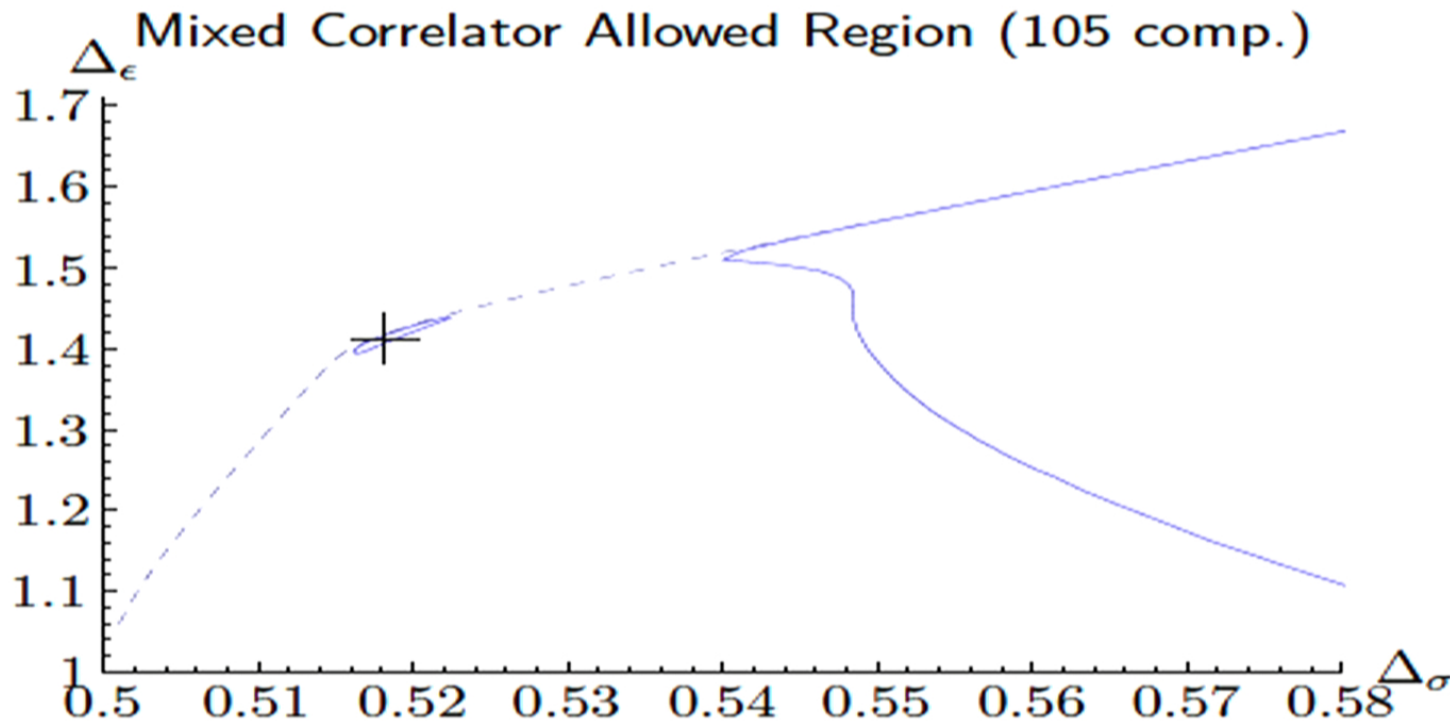
Ising: 3d Allowed Region



[Kos, DP, Simmons-Duffin, Vichi '16]

- ▶ Best bounds: first map out a 3d Island in $\{\Delta_\sigma, \Delta_\epsilon, \lambda_{\epsilon\epsilon\epsilon}/\lambda_{\sigma\sigma\epsilon}\}$
- ▶ Since the functional can be different for each choice of $\lambda_{\epsilon\epsilon\epsilon}/\lambda_{\sigma\sigma\epsilon}$, the $(\Delta_\sigma, \Delta_\epsilon)$ projection is smaller than having no assumption on $\lambda_{\epsilon\epsilon\epsilon}/\lambda_{\sigma\sigma\epsilon}$

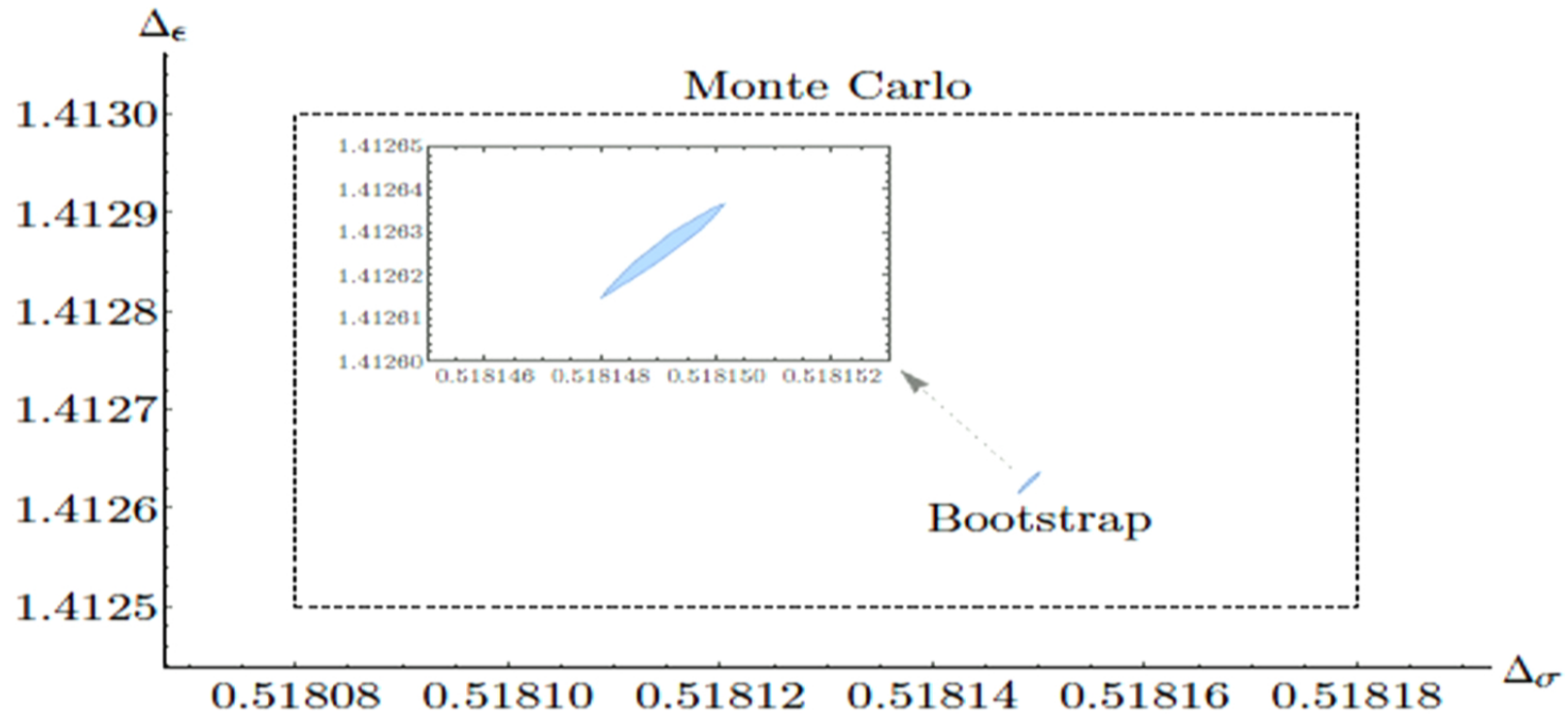
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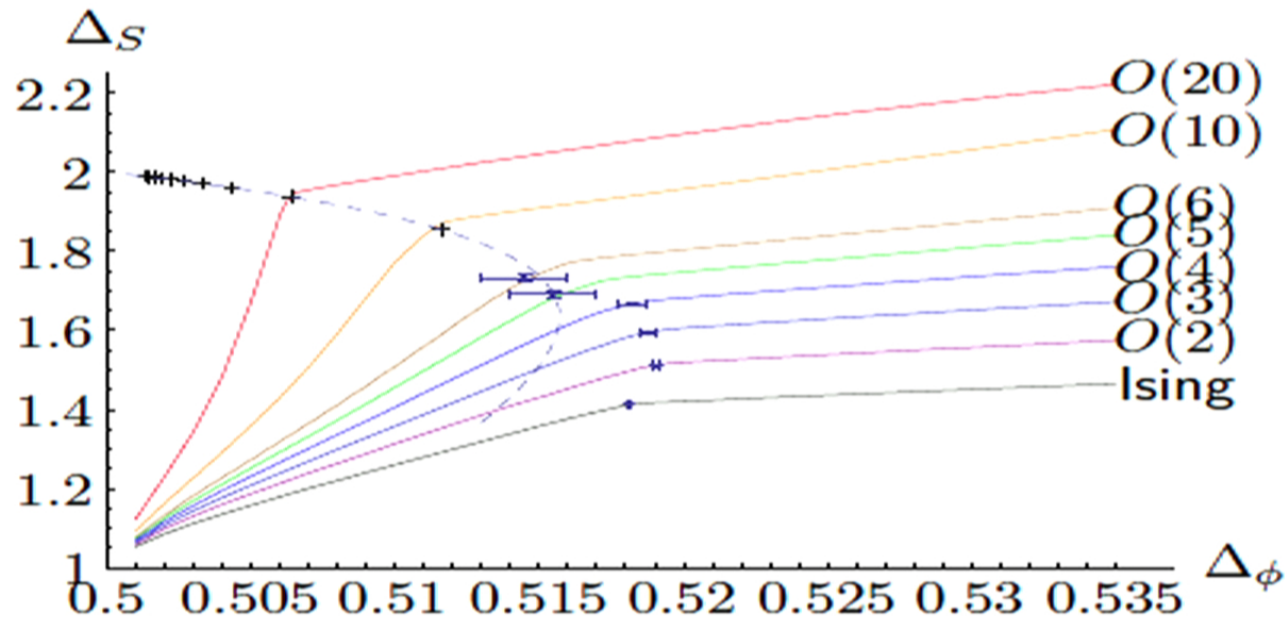
Mixed Correlator Islands



[Kos, DP, Simmons-Duffin, Vichi '16]

$$\begin{aligned}\{\Delta_\sigma, \Delta_\epsilon\} &= \{0.518149(1), 1.412625(10)\} \\ \{\lambda_{\sigma\sigma\epsilon}, \lambda_{\epsilon\epsilon\epsilon}\} &= \{1.0518537(41), 1.532435(19)\}\end{aligned}$$

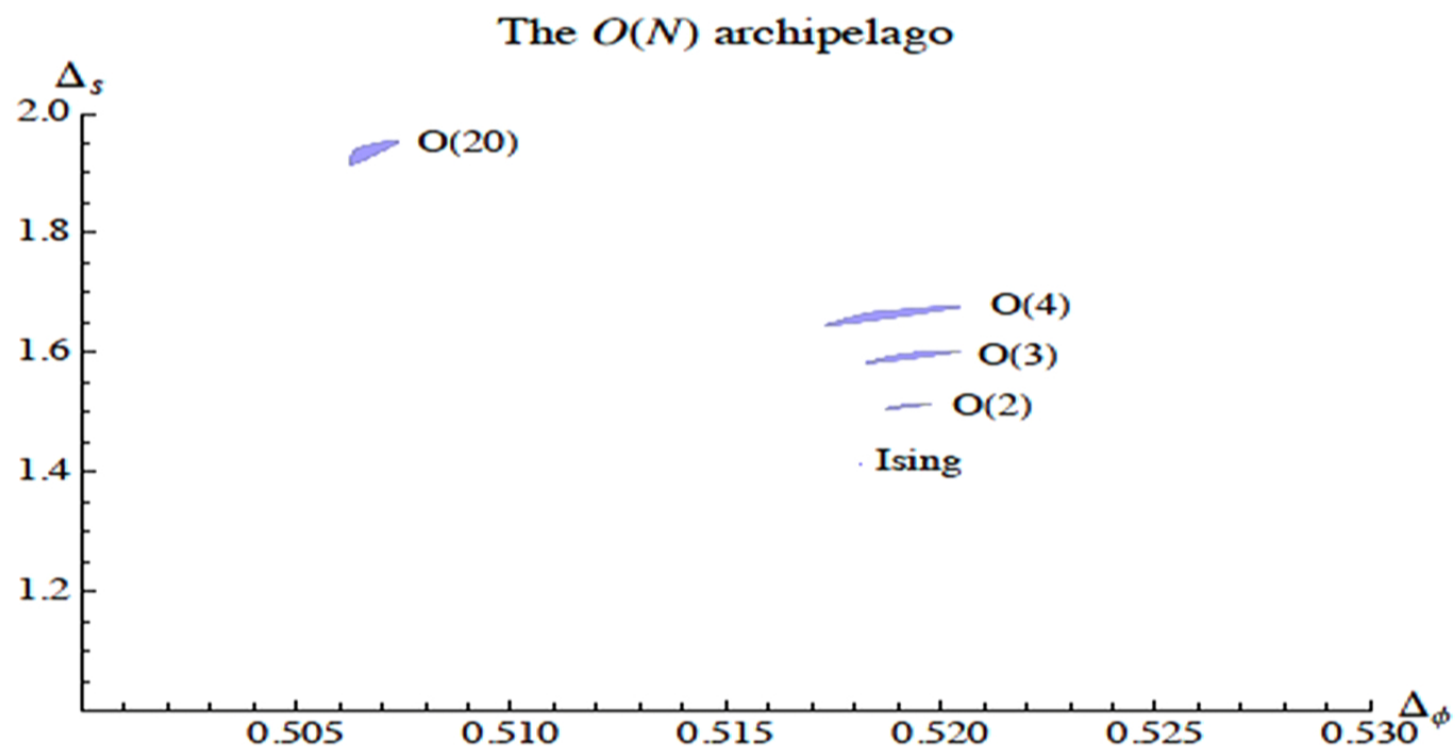
3D $O(N)$ Bounds



[Kos, DP, Simmons-Duffin '13]

- ▶ Extension to $\langle \phi_i \phi_j \phi_k \phi_l \rangle$, where ϕ_i is $O(N)$ vector
- ▶ Large N : matches $1/N$ expansion, Small N : matches experiment!

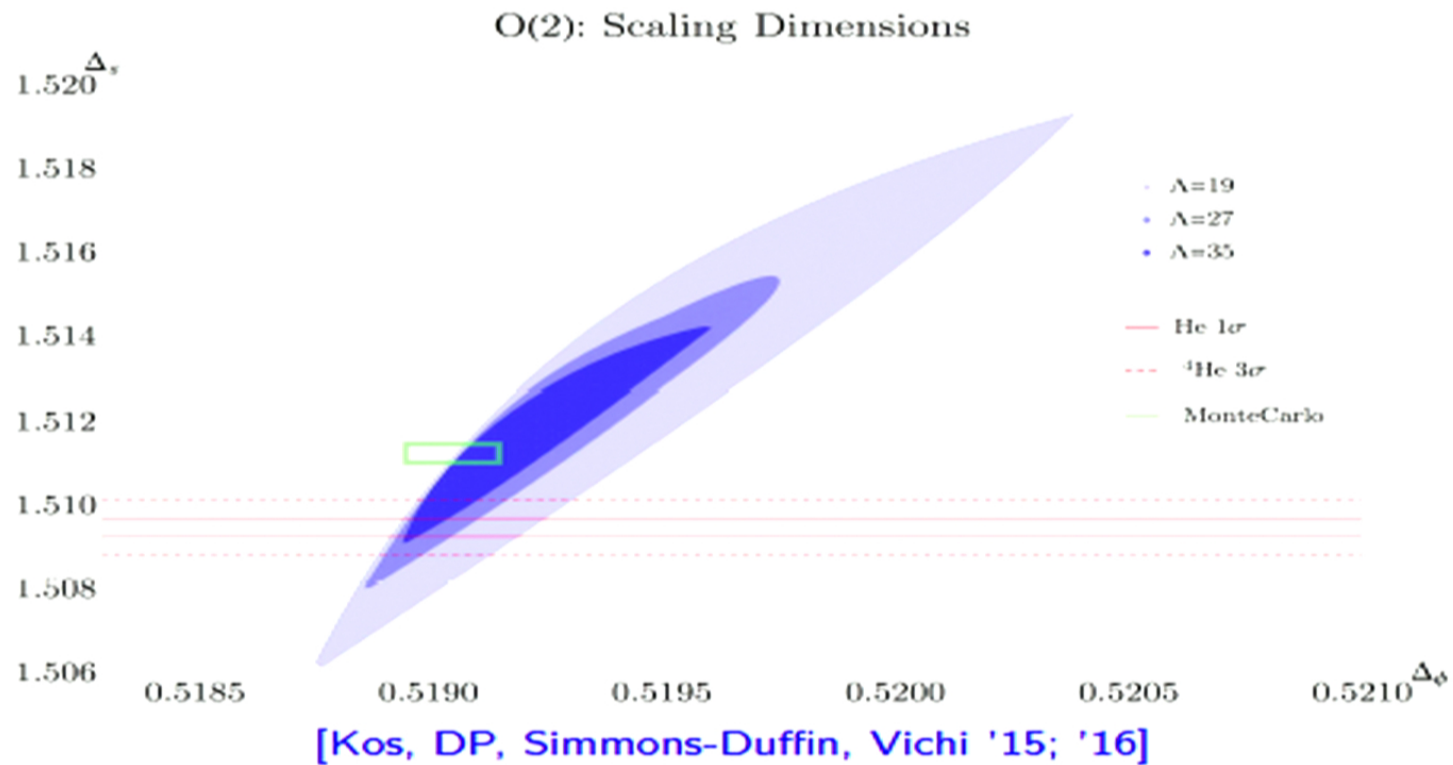
$O(N)$ Archipelago from Mixed Correlators



[Kos, DP, Simmons-Duffin, Vichi '15; '16]

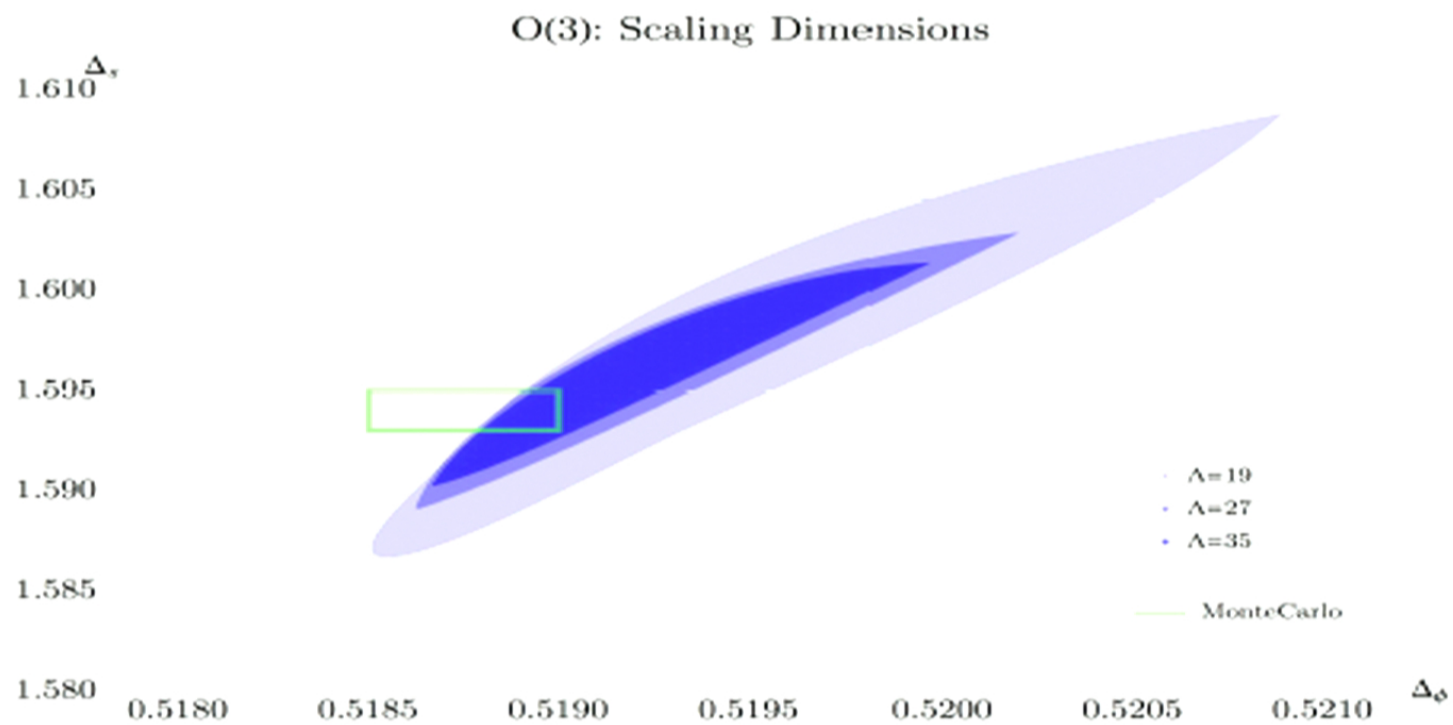
- ▶ Mixed system assuming only one relevant $O(N)$ vector ϕ_i and singlet s

$O(2)$ Zoom



- ▶ $\{\Delta_\phi, \Delta_s, \lambda_{\phi\phi s}, \lambda_{sss}\} = \{.51926(32), 1.5117(25), .68726(65), .8286(60)\}$
- ▶ Close to resolving 8σ discrepancy between lattice and expt

$O(3)$ Zoom

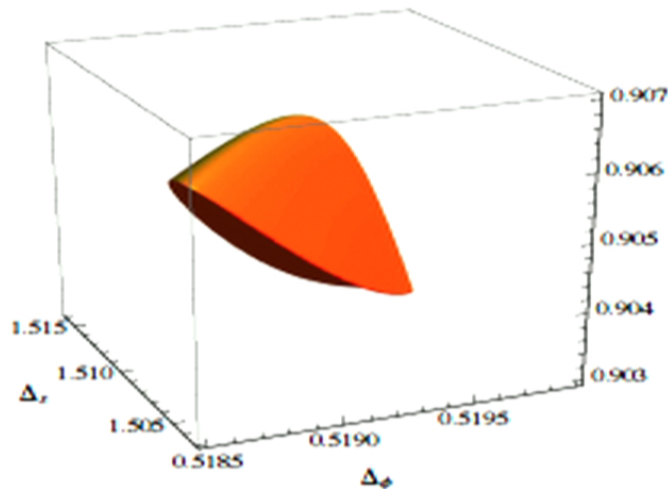


[Kos, DP, Simmons-Duffin, Vichi '15; '16]

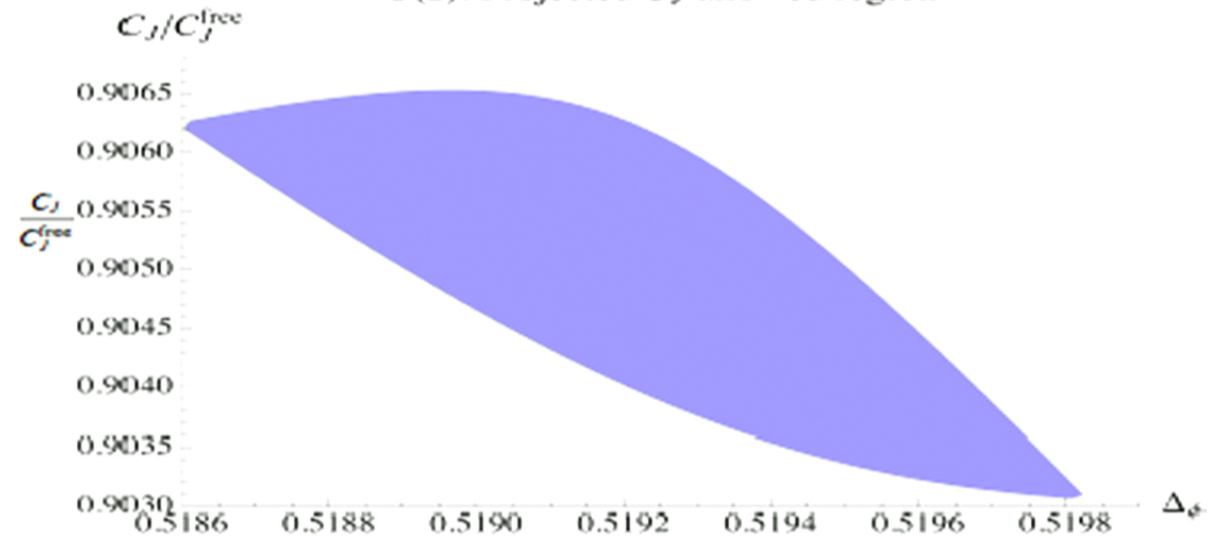
► $\{\Delta_\phi, \Delta_s, \lambda_{\phi\phi s}, \lambda_{sss}\} = \{.51928(62), 1.5957(55), .5244(11), .499(12)\}$

$O(2)$ Conductivity

$O(2)$: C_J allowed region



$O(2)$: Projected C_J allowed region

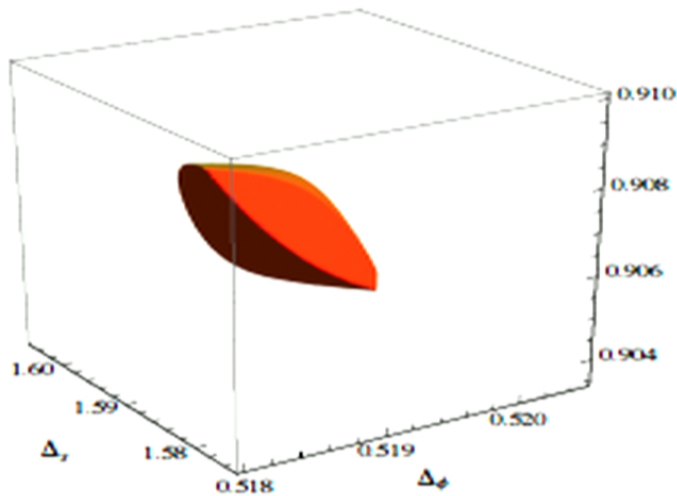


[Kos, DP, Simmons-Duffin, Vichi '15]

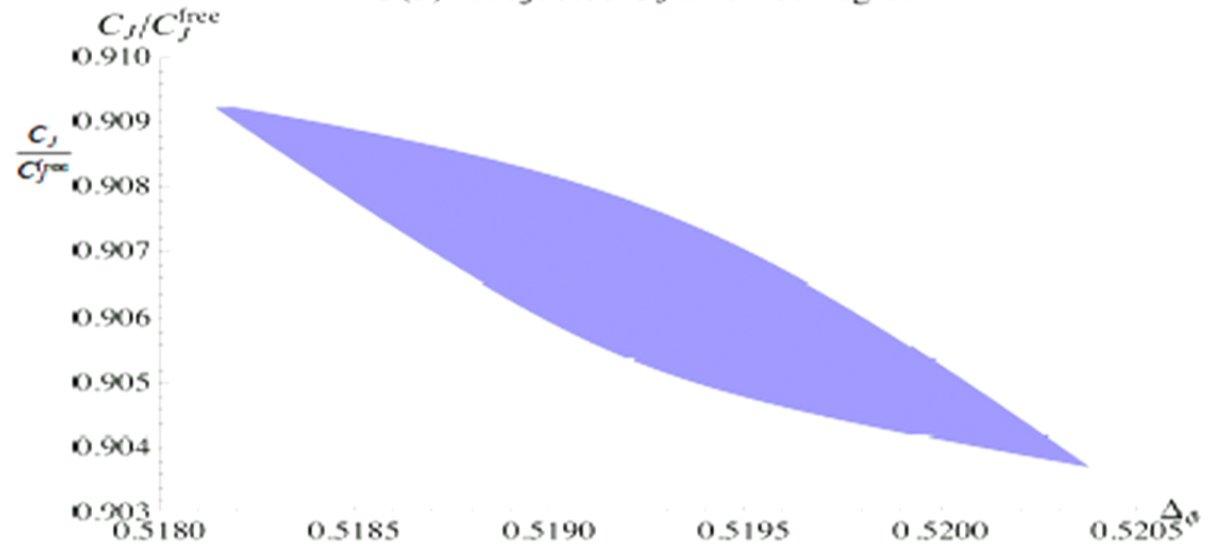
- ▶ Rigorous determination of $\langle JJ \rangle \propto C_J \propto \sigma_\infty$, giving high-frequency conductivity in $(2+1)$ D superconductors: $2\pi\sigma_\infty = 0.3554(6)$

$O(3)$ Conductivity

$O(3)$: C_J allowed region



$O(3)$: Projected C_J allowed region



[Kos, DP, Simmons-Duffin, Vichi '15]

- ▶ Similar determinations at large N , e.g. at $N = 20$:

$$C_J/C_J^{free} \Big|_{\text{bootstrap}} = 0.9674(8),$$

$$C_J/C_J^{free} \Big|_{\text{Large } N} \approx 1 - \frac{32}{9\pi^2} \frac{1}{N} = 0.964$$

Fermion Bootstrap

- ▶ Generalize to 4-point functions $\langle \psi\psi\psi\psi \rangle$ of a Majorana fermion in 3D ($SO(2, 1) \simeq Sp(2, \mathbb{R}) \rightarrow$ real two-component spinors)
- ▶ We will also assume a parity symmetry: $(x, y) \rightarrow (-x, y)$
- ▶ To classify 3-point and 4-point structures, we can work in an embedding space, where $SO(3, 2) \simeq Sp(4, \mathbb{R})$ is linearly realized

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Results:

- ▶ $\langle \psi\psi\mathcal{O}^{(\ell \text{ even})} \rangle$ has two structures of even parity and one of odd parity
- ▶ $\langle \psi\psi\mathcal{O}^{(\ell \text{ odd})} \rangle$ has one structure of odd parity
- ▶ $\langle \psi\psi\psi\psi \rangle$ has 5 independent tensor structures

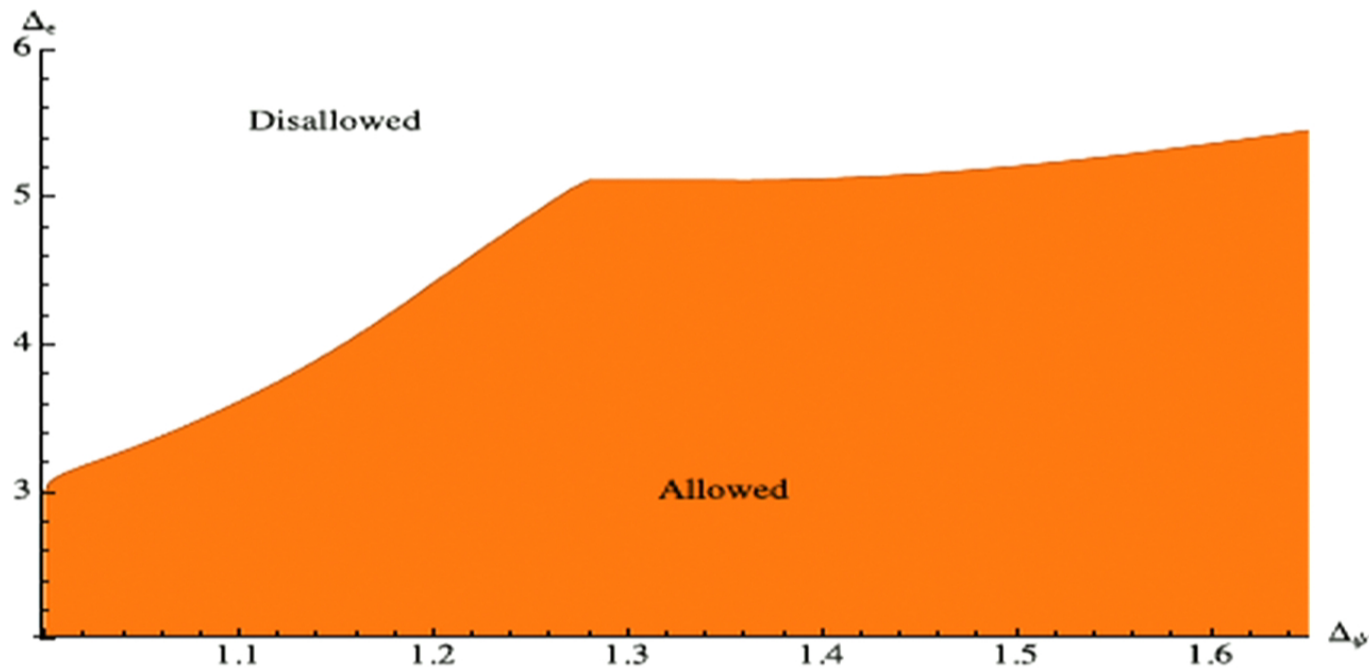
Fermion Bootstrap

- ▶ Crossing symmetry leads to a 5-vector of sum rules:

$$0 = \sum_{\mathcal{O}_+, \ell_+} \begin{pmatrix} \lambda_{\mathcal{O}_+}^1 & \lambda_{\mathcal{O}_+}^2 \end{pmatrix} \vec{F}_{++,\Delta,\ell}(u,v) \begin{pmatrix} \lambda_{\mathcal{O}_+}^1 \\ \lambda_{\mathcal{O}_+}^2 \end{pmatrix} \\ + \sum_{\mathcal{O}_-, \ell_+} (\lambda_{\mathcal{O}_-}^3)^2 \vec{F}_{-+,\Delta,\ell}(u,v) + \sum_{\mathcal{O}_-, \ell_-} (\lambda_{\mathcal{O}_-}^4)^2 \vec{F}_{--,\Delta,\ell}(u,v),$$

- ▶ To calculate the conformal blocks, we express $\langle \psi\psi\mathcal{O} \rangle_a = D_a \langle \phi\phi\mathcal{O} \rangle$, which lets us relate $\int \langle \psi\psi\mathcal{O} \rangle_a \langle \tilde{\mathcal{O}}\psi\psi \rangle_b$ to $\int \langle \phi\phi\mathcal{O} \rangle \langle \tilde{\mathcal{O}}\phi\phi \rangle$
- ▶ Bounds follow from applying functionals $\vec{\alpha}$ (SDP is mandatory)

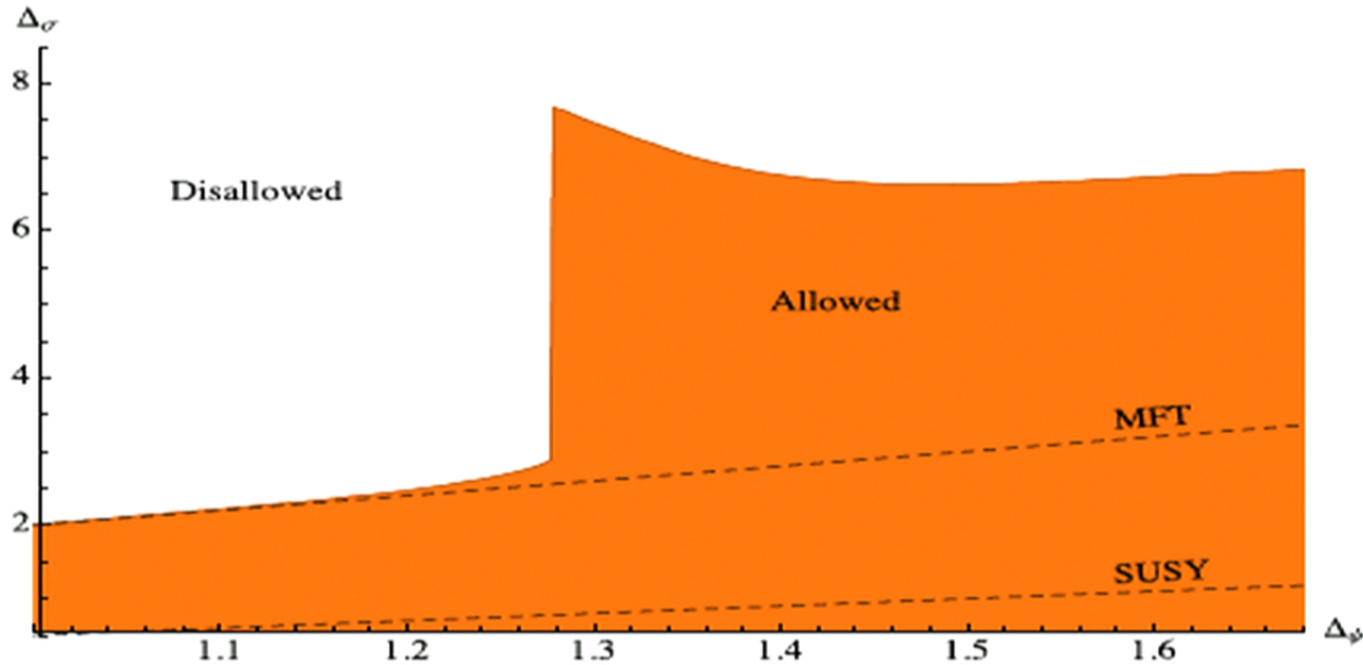
3D Fermion Bootstrap



[Iliesiu, Kos, DP, Pufu, Simmons-Duffin, Yacoby '15]

- ▶ Bound on leading parity-even scalar in $\psi \times \psi \sim \epsilon + \dots$

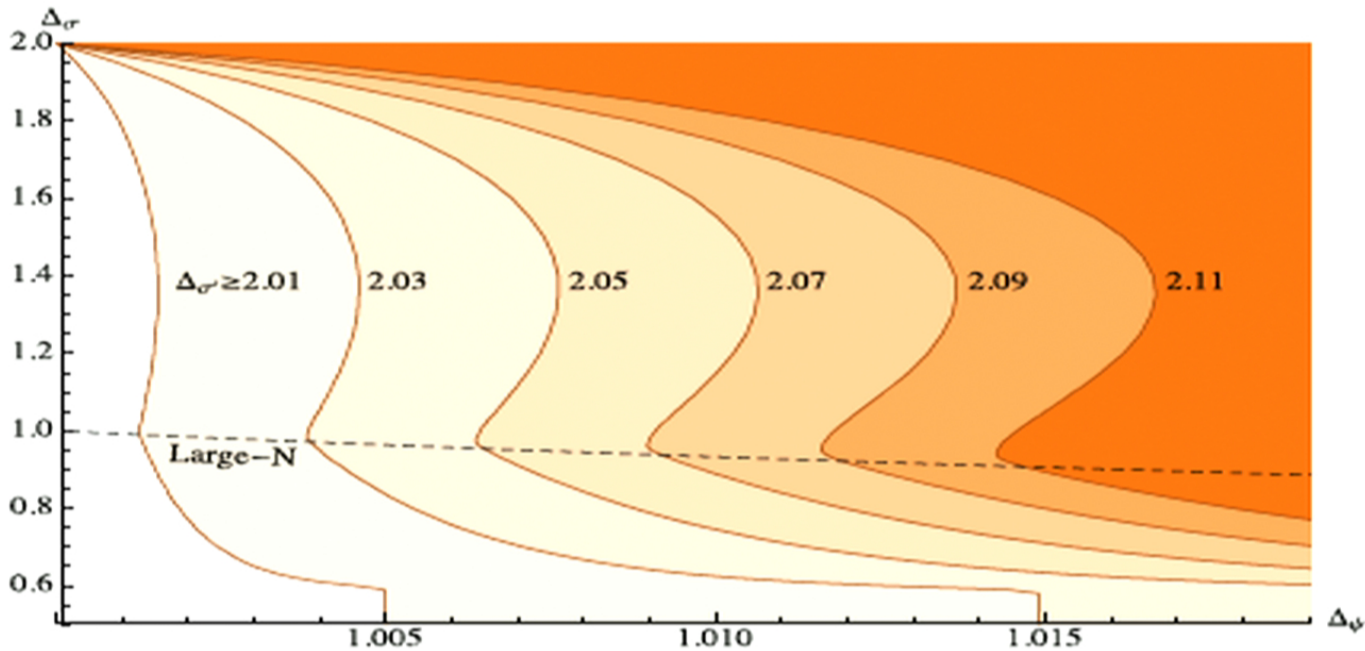
3D Fermion Bootstrap



[Iliesiu, Kos, DP, Pufu, Simmons-Duffin, Yacoby '15]

- ▶ Bound on leading parity-odd scalar in $\psi \times \psi \sim \sigma + \dots$
- ▶ Haven't yet found known CFT that coincides with jump

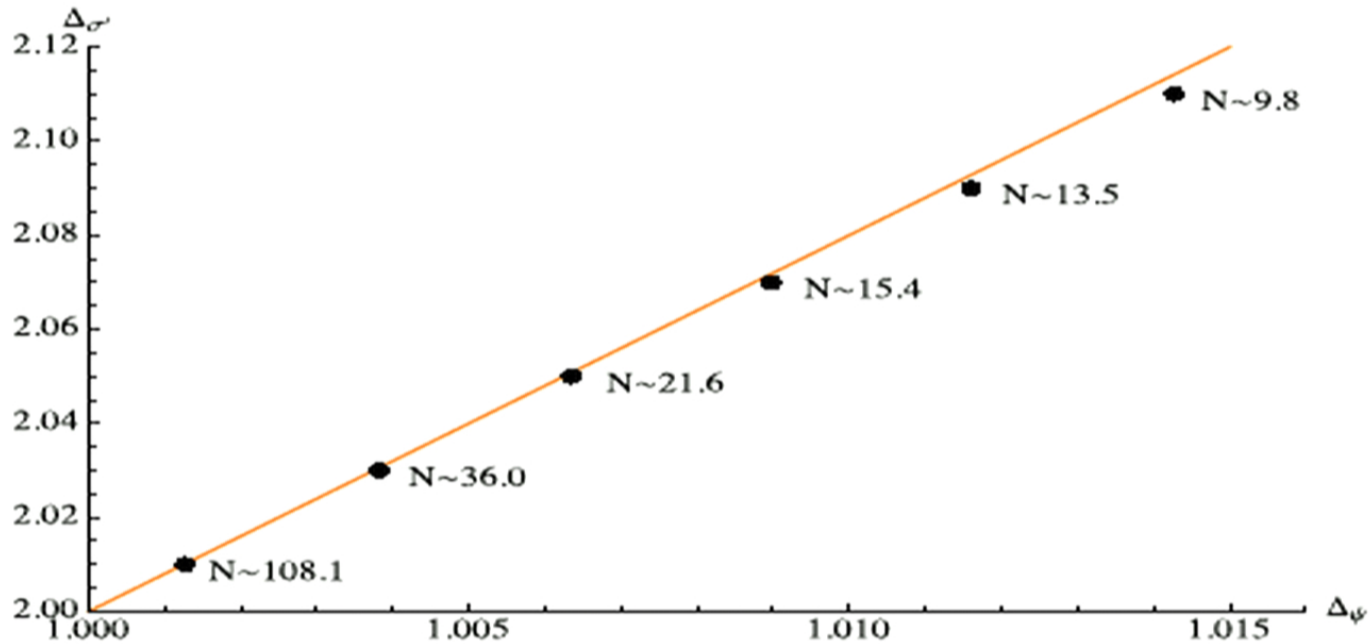
3D Fermion Bootstrap



[Iliesiu, Kos, DP, Pufu, Simmons-Duffin, Yacoby '15]

- ▶ However, adding a gap until 2nd parity-odd scalar σ' reveals features that coincide with Large N Gross-Neveu-Yukawa models ($\mathcal{L} \sim \sigma \bar{\psi}_i \psi^i$):

3D Fermion Bootstrap

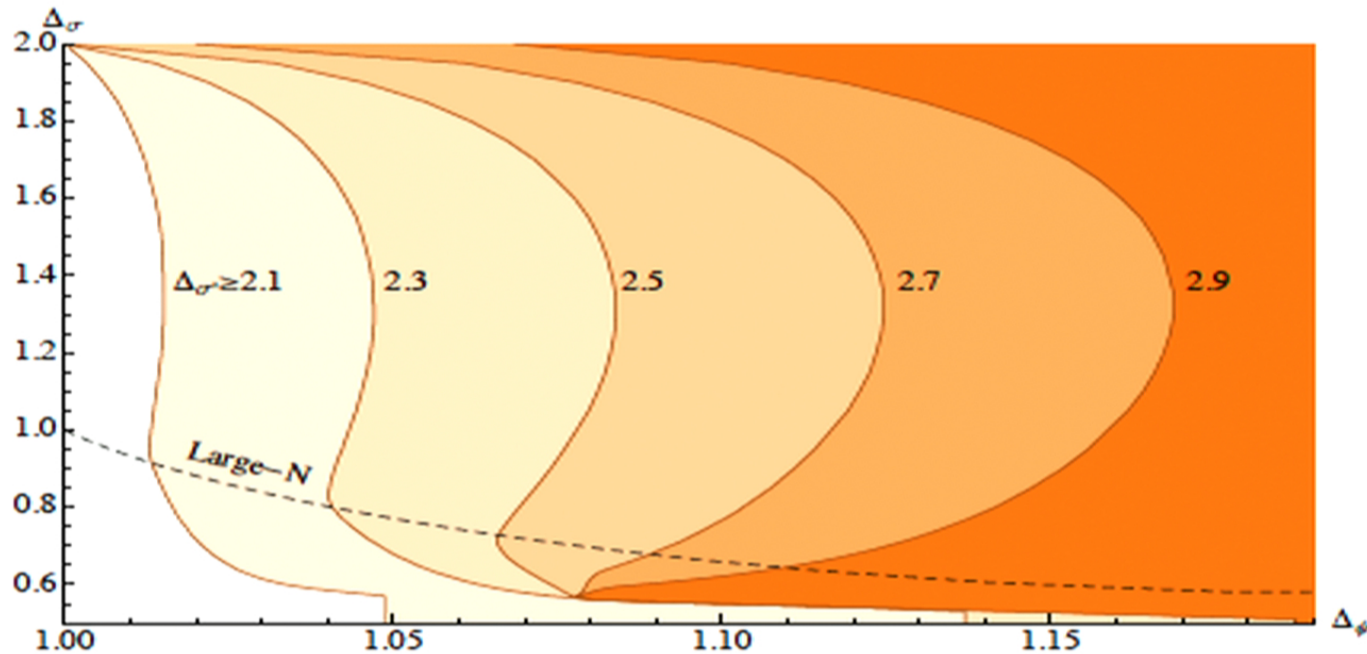


[Iliesiu, Kos, DP, Pufu, Simmons-Duffin, Yacoby '15]

Large N Gross-Neveu-Yukawa:

$$\blacktriangleright \Delta_\psi = 1 + \frac{4}{3\pi^2 N}, \quad \Delta_\sigma = 1 - \frac{32}{3\pi^2 N}, \quad \Delta_{\bar{\psi}\psi} = 2 + \frac{32}{3\pi^2 N}$$

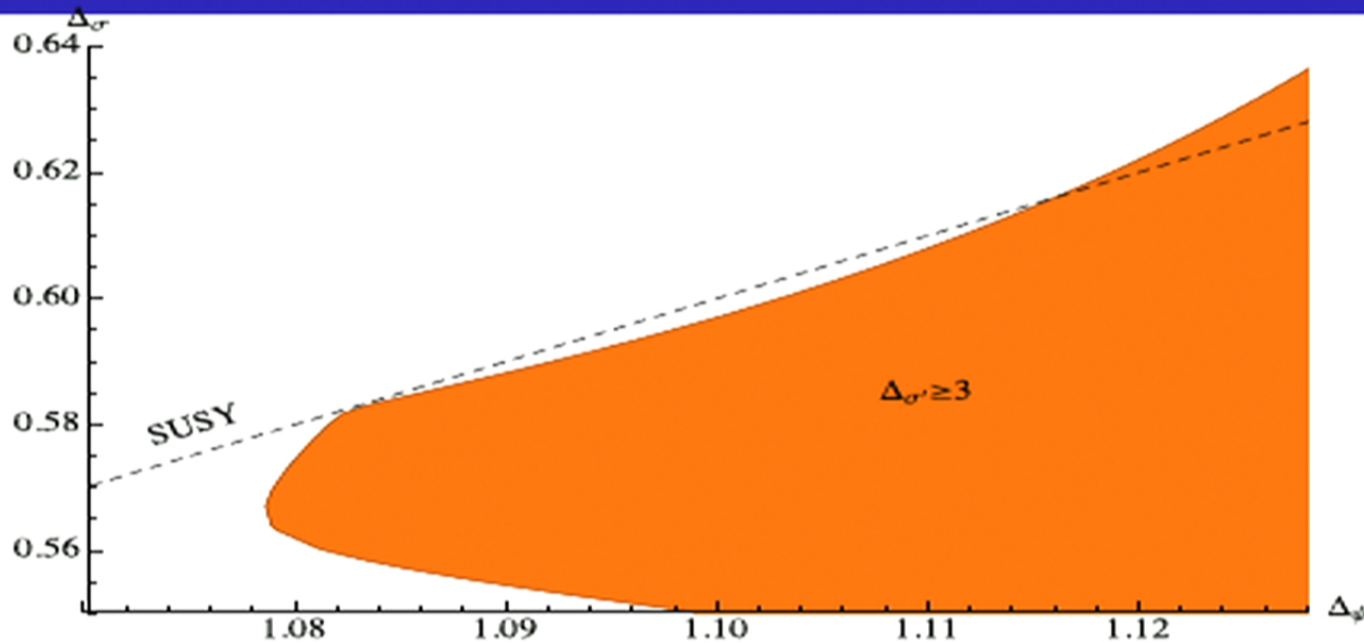
3D Fermion Bootstrap



[Iliesiu, Kos, DP, Pufu, Simmons-Duffin, Yacoby '15]

- ▶ Larger gaps in $\Delta_{\sigma'}$ probe the theories with smaller values of N
- ▶ Taking $N \rightarrow 1$ should reveal 3D $\mathcal{N} = 1$ SUSY Ising model ($W = \Phi^3$)

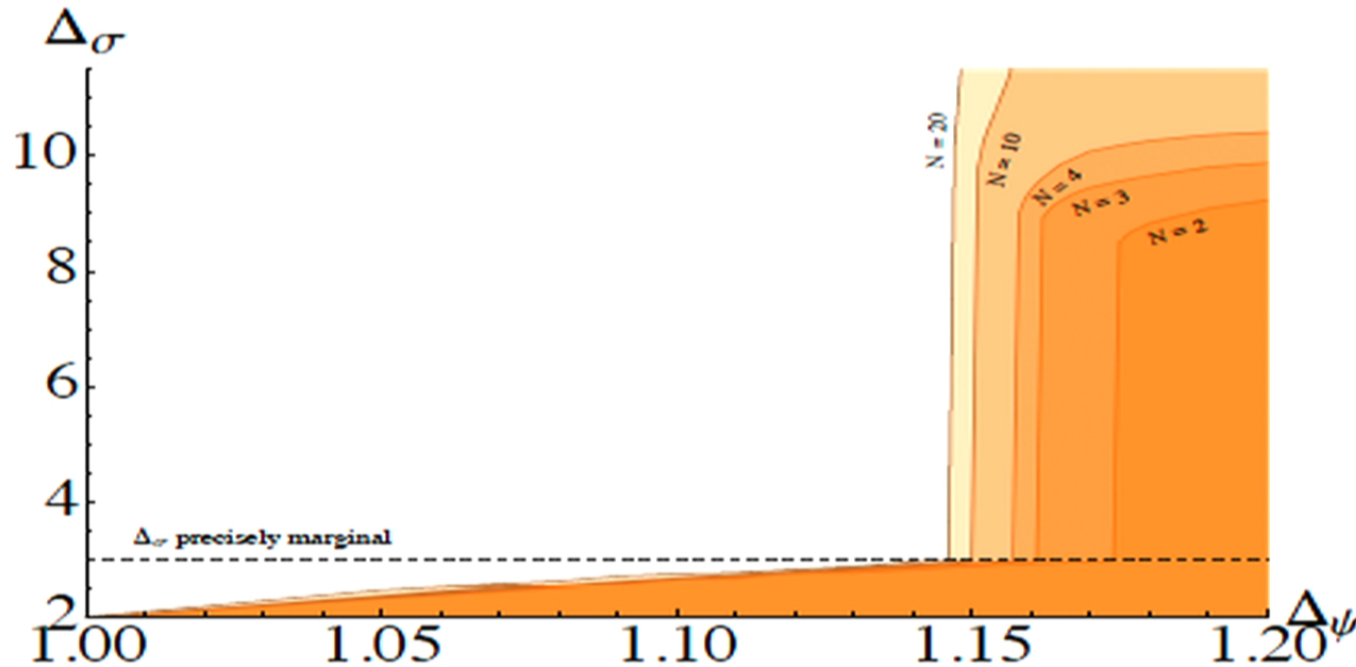
3D Fermion Bootstrap



[Iliesiu, Kos, DP, Pufu, Simmons-Duffin, Yacoby '15]

- ▶ Conjecture is that this theory sits at point where kink passes through SUSY line $\Delta_\psi = \Delta_\sigma + 1/2$ (occurs around $\Delta_\psi \sim 1.082$, $\Delta_{\sigma'} \sim 2.95$)
- ▶ Near ϵ -expansion estimate ($\Delta_\psi \approx 1.09$) [Fei, Giombi, Klebanov, Tarnopolsky '16]
- ▶ Boundary of topological superconductors? [Grover, Sheng, Vishwanath '13]

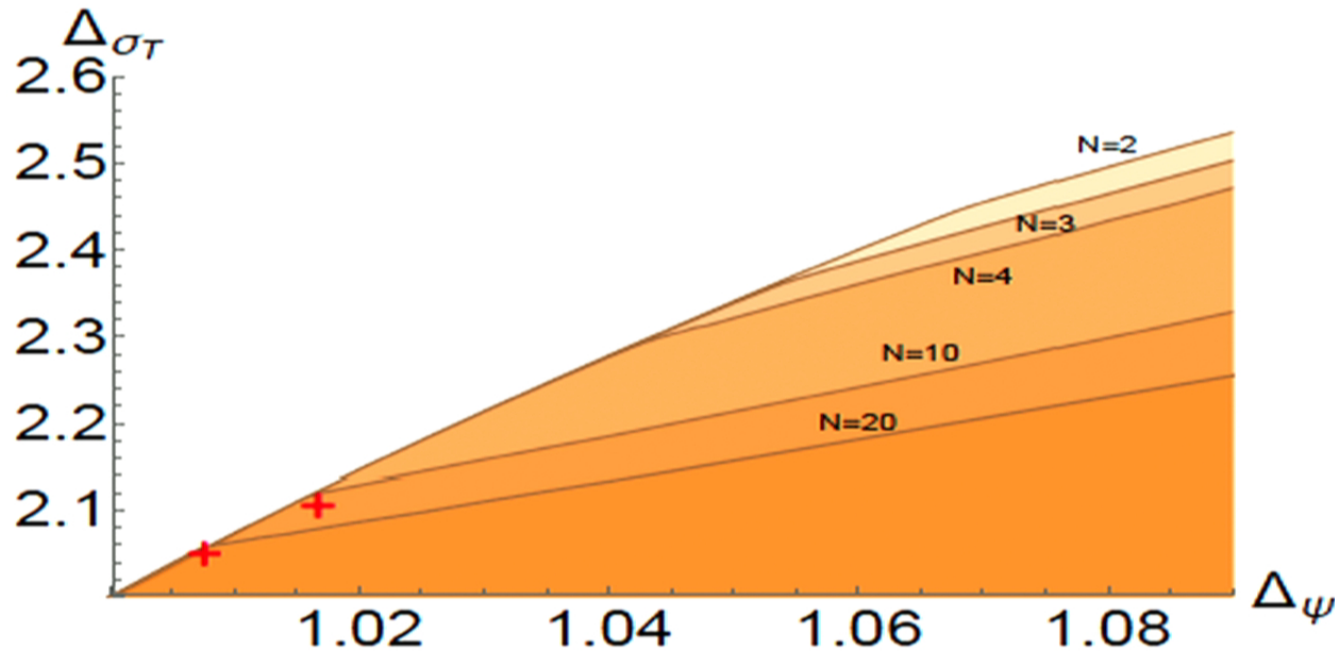
3D Fermions with $O(N)$ Symmetry



[Iliesiu, Kos, DP, Pufu, Simmons-Duffin, Yacoby, to appear]

- ▶ Generalization to $\langle \psi_i \psi_j \psi_k \psi_l \rangle$ where ψ_i is fundamental of $O(N)$
- ▶ Bound on leading parity-odd singlet σ still has mysterious jumps

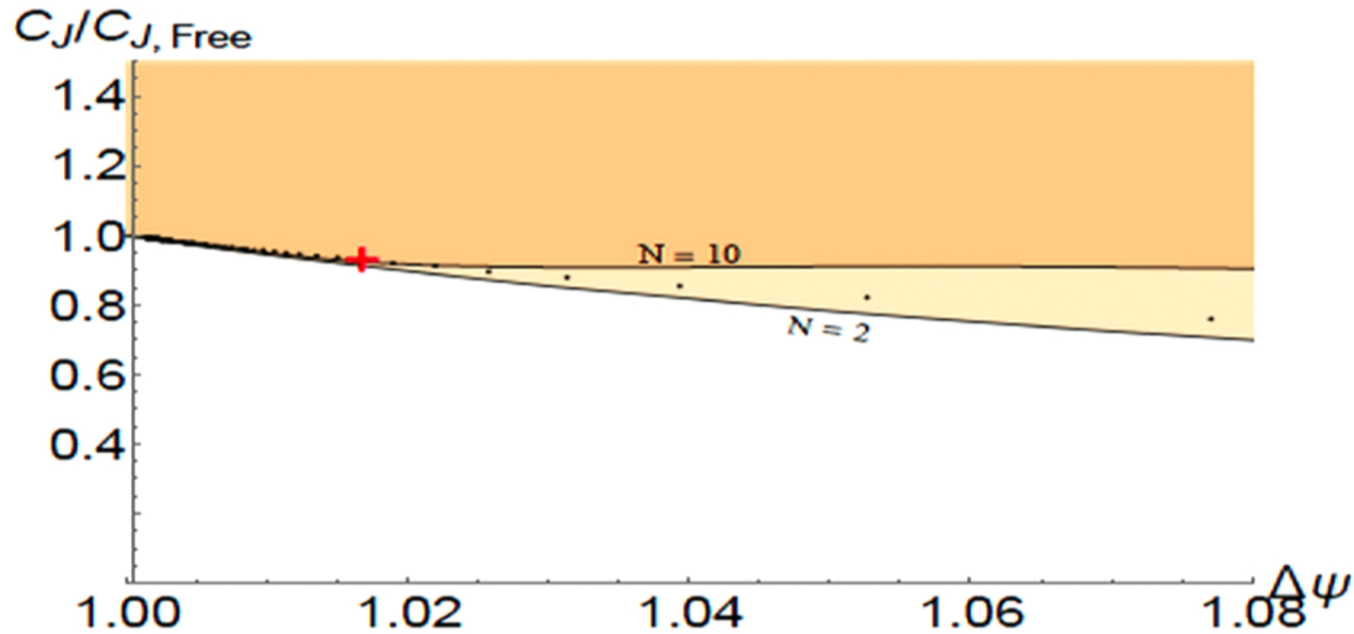
3D Fermions with $O(N)$ Symmetry



[Iliesiu, Kos, DP, Pufu, Simmons-Duffin, Yacoby, to appear]

- ▶ Bound on leading parity-odd symmetric tensor σ_T seems promising!
- ▶ Looks plausible that all models saturate $N = 2$ bound

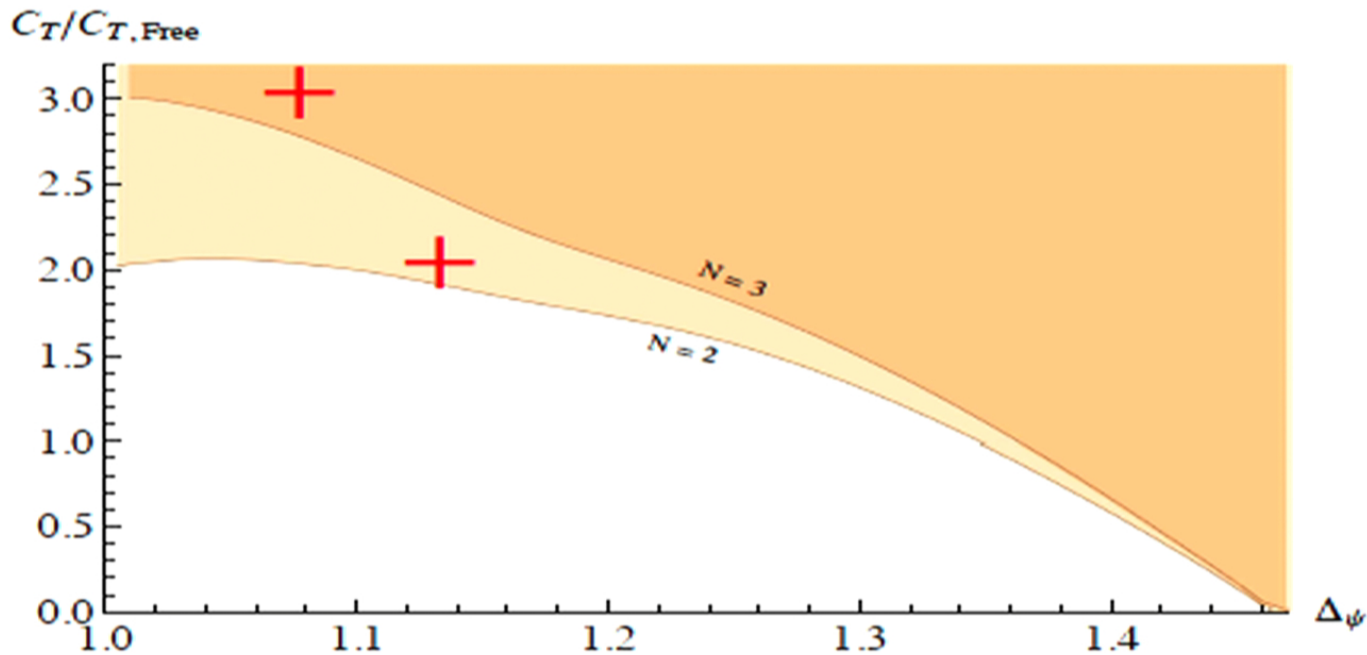
3D Fermions with $O(N)$ Symmetry



[Iliesiu, Kos, DP, Pufu, Simmons-Duffin, Yacoby, to appear]

- ▶ Preliminary bounds on $\langle JJ \rangle \propto C_J$ also follow large N curve
- ▶ Could again be possible that all models saturate $N = 2$ bound

3D Fermions with $O(N)$ Symmetry



[Iliesiu, Kos, DP, Pufu, Simmons-Duffin, Yacoby, to appear]

- ▶ Preliminary bounds on $\langle TT \rangle \propto C_T$ grow linearly with N , as expected

3D Bootstrap Future

Where do we go from here?

- ▶ Make $O(N)$ model predictions more **precise** (resolve 8σ discrepancy!)
 - ▶ Extend mixed correlator bootstrap to include external t^{ij} ?
 - ▶ Higher spectrum ($\{\phi', s', t'\}$, higher $O(N)$ reps, leading twist trajectory)

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 - ▶ 3D fermions: Study mixed $\{\psi, \sigma\}$ system
 - ▶ Gross-Neveu models, Hubbard model, $\mathcal{N} = 1, 2$ Ising model, ...
 - ▶ 3D QED, 3D QCD, Chern-Simons + matter, ...
 - ▶ Can we classify all the bootstrap solutions???

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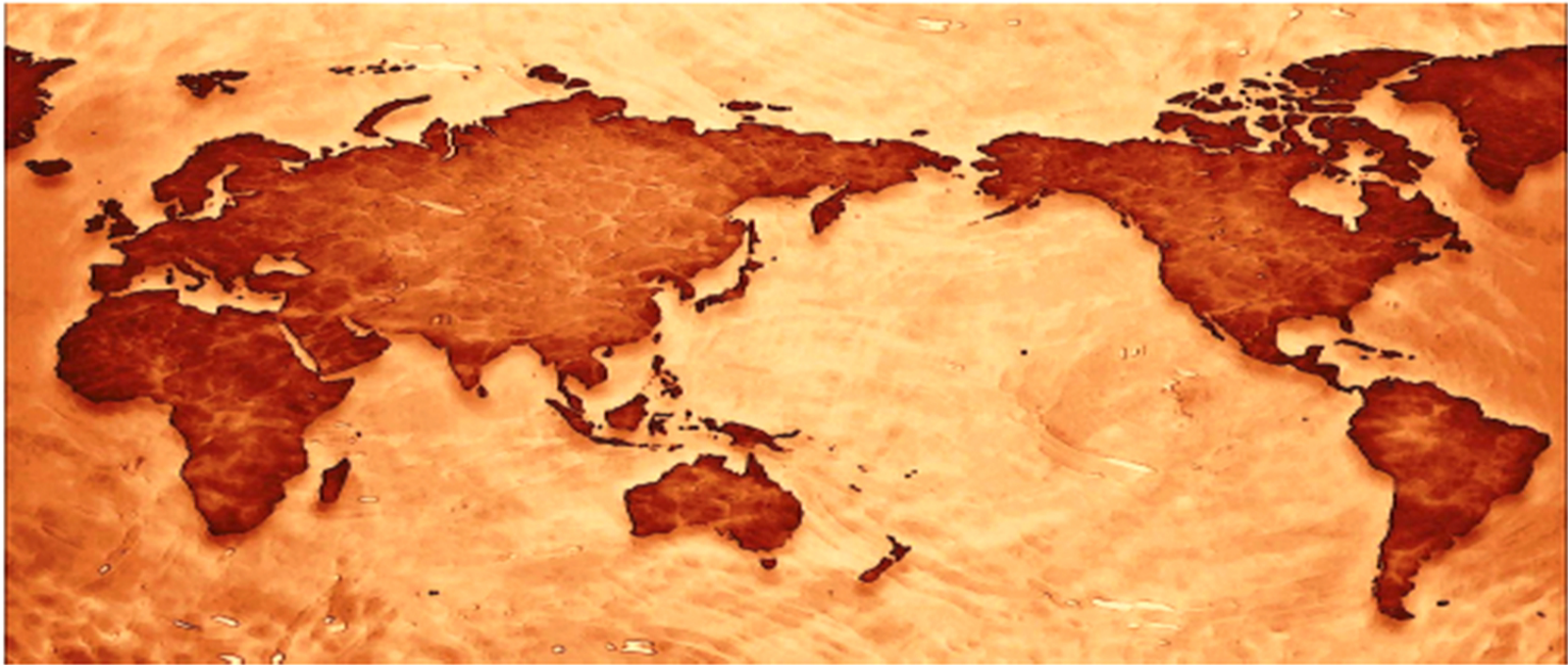
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- ▶ Bootstrap currents and stress tensor

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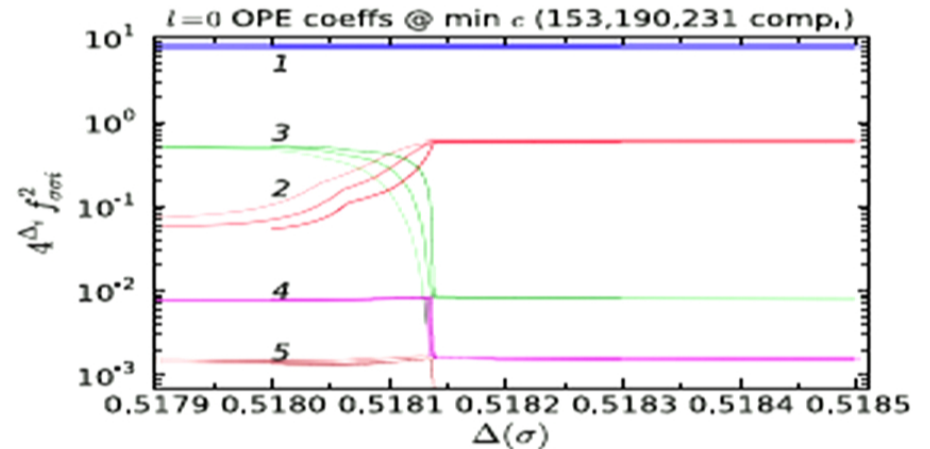
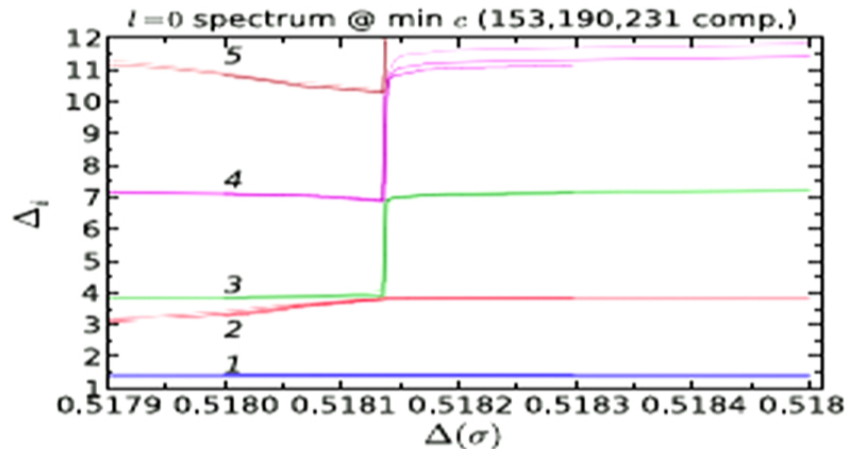
- ▶ Make $O(N)$ model predictions more **precise** (resolve 8σ discrepancy!)
 - ▶ Extend mixed correlator bootstrap to include external t^{ij} ?
 - ▶ Higher spectrum ($\{\phi', s', t'\}$, higher $O(N)$ reps, leading twist trajectory)
- ▶ Find **rigorous islands** for fermionic CFTs and gauge theories
 - ▶ 3D fermions: Study mixed $\{\psi, \sigma\}$ system
 - ▶ Gross-Neveu models, Hubbard model, $\mathcal{N} = 1, 2$ Ising model, ...
 - ▶ 3D QED, 3D QCD, Chern-Simons + matter, ...
 - ▶ Can we classify all the bootstrap solutions???
- ▶ Bootstrap currents and stress tensor
 - ▶ Analytic Bootstrap for $\langle JJ\phi\phi\rangle$ and $\langle TT\phi\phi\rangle$
 - Sum rules for coefficients in $\langle JJJ\rangle \propto \hat{n}_{s,f}$ and $\langle TTT\rangle \propto n_{s,f,t}$
 - Proof of Hofman-Maldacena bounds $\hat{n}_{s,f} \geq 0$ and $n_{s,f,t} \geq 0$
 - [Hartman, Jain, Kundu '15; '16; Hofman, Li, Meltzer, DP, Rejon-Barrera '16]
 - ▶ Can they be strengthened? Interplay with numerical studies?

Bootstrap Future



- ▶ With more work I believe we can create a **detailed map** of the space of conformal field theories...we may even discover a **new world!**

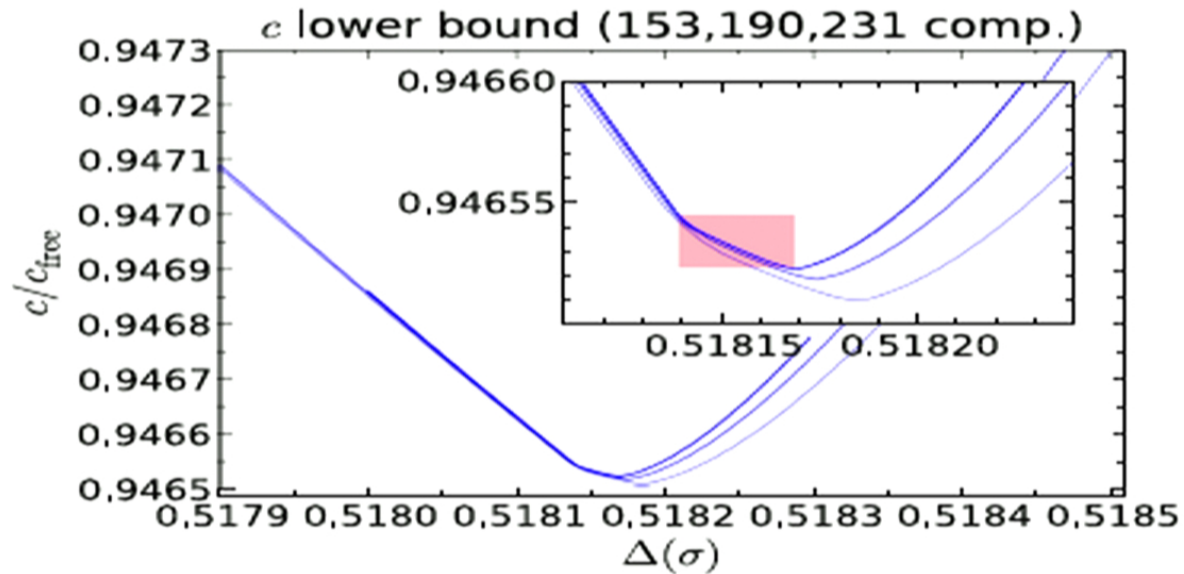
c -minimization and Non-rigorous Spectrum



[El-Showk, Paulos, DP, Rychkov, Simmons-Duffin, Vichi, '14]

- ▶ “Kink” \leftrightarrow operators merge and disappear from spectrum!
- ▶ Reminiscent of null states in 2D or equations of motion in $(4 - \epsilon)D$
 \rightarrow Non-perturbative equation of motion?
- ▶ E.g., in ϕ^4 theory, expect $\partial^2 \phi \sim \phi^3 \rightarrow$ gap in \mathbb{Z}_2 -odd spectrum...

c -minimization and Non-rigorous Spectrum



[El-Showk, Paulos, DP, Rychkov, Simmons-Duffin, Vichi, '14]

- ▶ Under the conjecture that the central charge $\langle TT \rangle \propto c$ is minimized, a precise spectrum in $\sigma \times \sigma \sim \mathbb{1} + \epsilon + \epsilon' + \dots$ can be extracted:

$$\Delta_\sigma \simeq 0.518154(15), \quad \Delta_\epsilon \simeq 1.41267(13), \quad \Delta_{\epsilon'} = 3.8303(18), \quad \dots$$