Title: Higher-order interference doesn't help in searching for a needle in a haystack

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Abstract:

Pirsa: 16080030 Page 1/37

# Higher-order interference doesn't help in finding a needle in a haystack

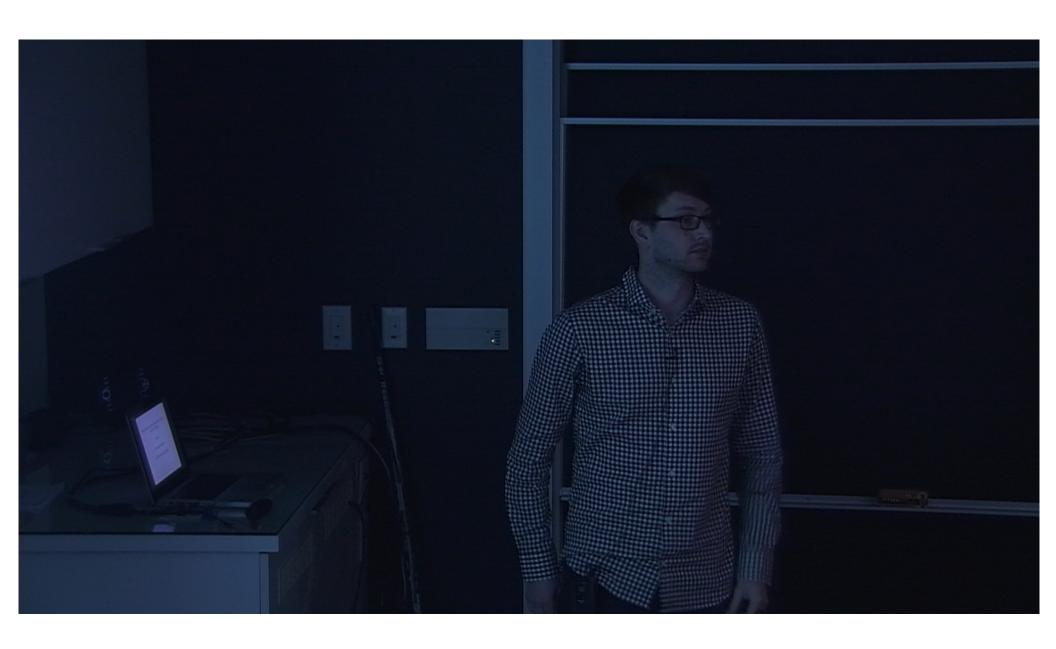
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Joint work with John Selby

arXiv:1604.03118 & arXiv:1510.04699



Pirsa: 16080030 Page 2/37

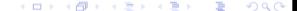


Pirsa: 16080030

## Consequences of higher-order interference?

► Absence of third-order interference, in conjunction with other physical principles, uniquely specifies quantum theory.

► Does post-quantum interference imply post-quantum features? Information-theoretic advantages?



Pirsa: 16080030 Page 4/37

#### The search problem

- ▶ Items indexed 1, ..., x, ..., N, with x the 'marked' item.
- ▶ One has access to an **oracle**, which when asked if item i = x returns 'yes' or 'no'.
- ▶  $f: \{1, ..., N\} \rightarrow \{0, 1\}$  satisfies f(i) = 1 if and only if i = x.
- ▶ What is minimal number of queries to this oracle to find *x* in the worst case?



#### The search problem

▶ Classical algorithms require O(N) queries to find marked item in worst case.

▶ There exists a quantum algorithm which finds item in  $O(\sqrt{N})$  queries.

▶  $O(\sqrt{N})$  queries is **optimal** for quantum theory.



#### Quantum oracles

In quantum theory, an oracle corresponds to a controlled unitary transformation

$$U|i\rangle|q\rangle = |i\rangle|q \oplus f(i)\rangle$$

where  $f: \{1, ..., N\} \rightarrow \{0, 1\}$  satisfies f(i) = 1 if and only if i = x.



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#### Quantum oracles

▶ Inputting  $|-\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$  results in a "kicked-back" phase:

$$U|i\rangle|-\rangle = (-1)^{f(i)}|i\rangle|-\rangle$$

ightharpoonup Discarding  $|-\rangle$  reduces the action of the oracle to

$$O_{\times}|i\rangle=(-1)^{f(i)}|i\rangle.$$



#### Operational theories and Physical principles

► We study the connection between higher-order interference and the search problem in the setting of operational theories.

► An operational theory specifies a set of laboratory devices which can be connected together to form experiments and assigns probabilities to experimental outcomes.



Pirsa: 16080030 Page 10/37

## Principle 1: Purification

- ▶ Process  $\{E_j\}_{j\in Y}$  refines process  $\{F_k\}_{k\in X}$  if  $F_k = \sum_{j\in X_k} E_j$ .
- ► A process is pure if it has trivial refinements.
- ► A pure process is one about which we have "maximal information".



# Principle 1: Purification

All states can be 'purified' by including an appropriate environment:

$$\psi$$
 =  $\sigma$ 



Pirsa: 16080030 Page 12/37

## Principle 1: Purification

if, 
$$\psi$$
 =  $\tilde{\psi}$ 

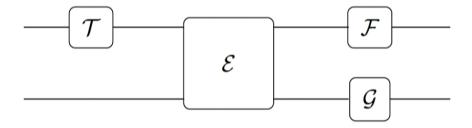
$$\Rightarrow$$
  $\psi$   $=$   $\tilde{\psi}$   $T$ 



Pirsa: 16080030 Page 13/37

## Principle 2: Purity preservation

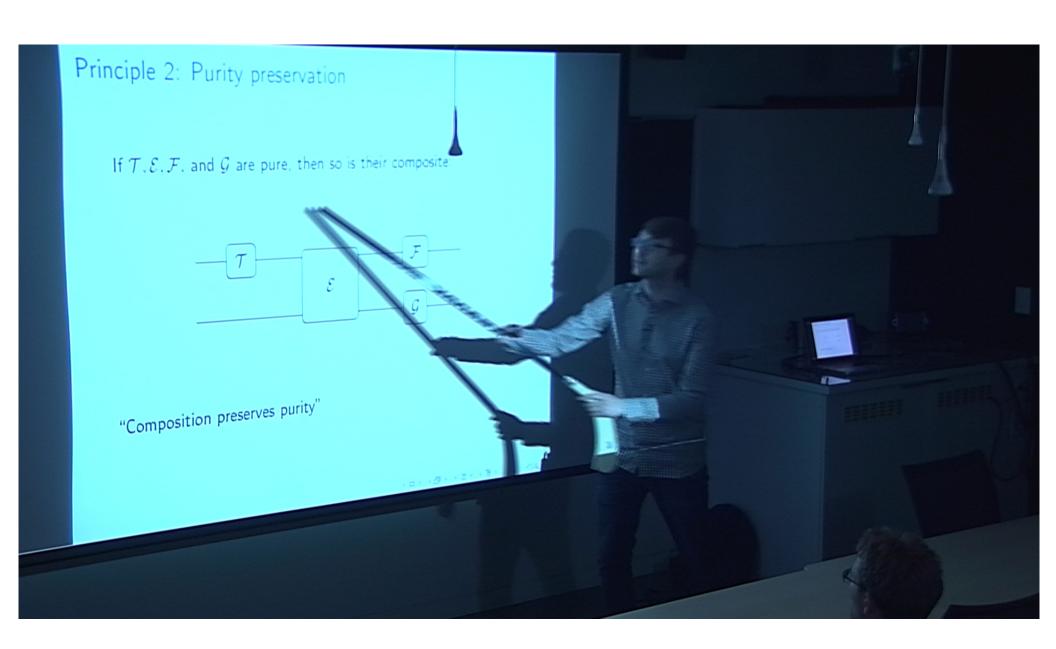
If  $\mathcal{T}, \mathcal{E}, \mathcal{F}$ , and  $\mathcal{G}$  are pure, then so is their composite:



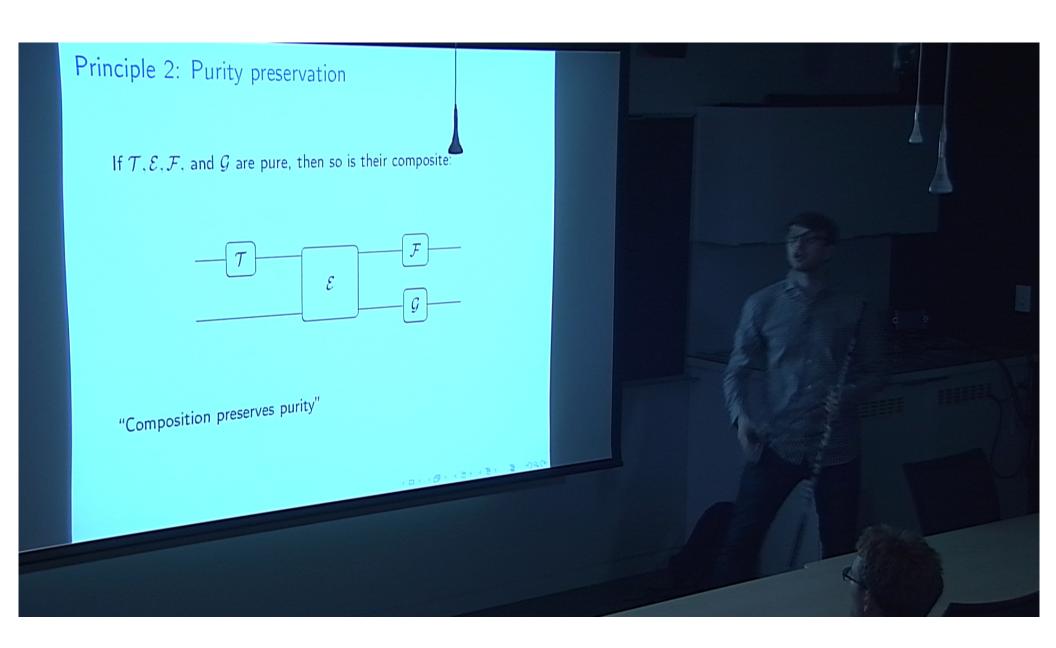
"Composition preserves purity"



Pirsa: 16080030 Page 14/37



Pirsa: 16080030



Pirsa: 16080030 Page 16/37

## Principle 3: Strong symmetry

Given two sets of perfectly distinguishable pure states

$$\{\sigma_i\}$$
 and  $\{\rho_i\}$ 

there exists reversible  $\mathcal{T}$  such that:

$$\sigma_i$$
  $=$   $\rho_i$   $\forall i$ 



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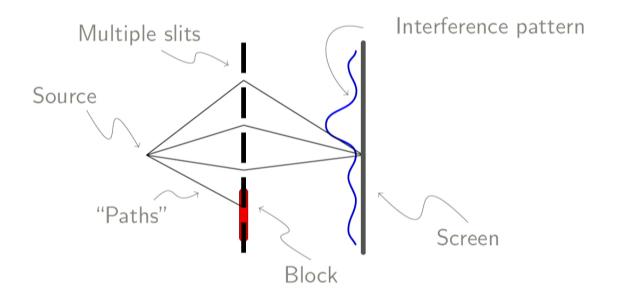
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there exists reversible  $\mathcal{T}$  such that:

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## Higher-order interference in the presence of these principles



Blocking some slits, but leaving subset  $I \subseteq \{1, ..., N\}$  open, corresponds to applying the projector  $P_I$ , satisfying  $P_I P_J = P_{I \cap J}$ .

Pirsa: 16080030 Page 19/37

#### Higher-order interference in the presence of these principles

Maximal order of interference *h* corresponds to:

$$\mathbb{1}_{N} = \sum_{\substack{I \subseteq \mathbf{N} \\ |I| \le h}} (-1)^{h-|I|} \left( \begin{array}{c} N - |I| - 1 \\ h - |I| \end{array} \right) P_{I}$$

The case of N = h + 1 corresponds to  $(-1)^{h-|I|}$ . In quantum theory, this is:

$$\mathbb{1}_{N} = \sum_{i < j} P_{\{ij\}} - (N-2) \sum_{i} P_{\{i\}},$$



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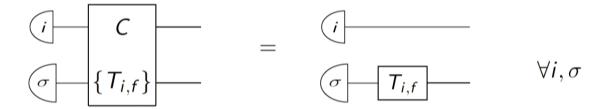
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## Oracles in operational theories



Making the set of transformations  $\{T_{i,f(i)}\}$  depend on the function  $f:\{i\} \to \{0,1\}$  encoding the search problem allows us to define a computational oracle.



Pirsa: 16080030 Page 22/37

## Oracles in operational theories

There exists a state *s* such that:

Moreover,  $\mathcal{O}_{s,f}$  is phase transformation:

$$-\mathcal{O}_{s,f}$$
  $=$   $i$   $\forall i$ 



Pirsa: 16080030

▶ In quantum theory, to query the oracle about i one applies the oracle to state  $|i\rangle\langle i|$ .

 $|i\rangle\langle i|$  can be prepared by passing a uniform superposition through an N-slit experiment with all but the ith slit blocked.

► The oracle can be implemented by placing a phase shifter behind slit *x*.



Pirsa: 16080030 Page 24/37

$$i = 0$$

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- 1.  $\mathcal{O}_{x}P_{I}=P_{I}$ , if  $x\notin I$  or |I|=1
- 2.  $\mathcal{O}_X$  can act non-trivially on  $P_I$  with  $X \in I$  and |I| > 1, but must satisfy  $\mathcal{O}_X P_I = P_I \mathcal{O}_X$ , for all  $P_I$



The requirement  $\mathcal{O}_{\times}P_I = P_I\mathcal{O}_{\times}$  ensures one cannot gain any information about item i when querying with a state that doesn't pass through slit i, i.e. a state s such that  $P_I s = s$  where  $i \notin I$ .



Pirsa: 16080030 Page 26/37

A reversible transformation is a search oracle, denoted  $\mathcal{O}_{\times}$ , if and only if:

i) 
$$\mathcal{O}_{x}P_{I}=P_{I}$$
 for all  $x\notin I$  or  $|I|=1$  and,

ii) 
$$\mathcal{O}_X P_I = P_I \mathcal{O}_X$$
, for all  $P_I$ .



Given a search oracle  $\mathcal{O}_{\times}$  and an arbitrary collection of reversible transformations  $\{G_i\}$ , what is the minimal  $k \in \mathbb{N}$  such that

$$G_k \mathcal{O}_X G_{k-1} \dots G_1 \mathcal{O}_X s$$

can be found, with high probability, to be in the state x, for arbitrary input state s, averaged over all possible marked items?

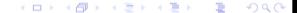


Pirsa: 16080030 Page 28/37

#### Main result

In theories satisfying our principles, with (finite) maximal order of interference h, the number of queries needed to solve the search problem is

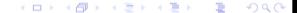
$$\Omega\left(\sqrt{\frac{N}{h}}\right)$$
.



#### Main result

In theories satisfying our principles, with (finite) maximal order of interference h, the number of queries needed to solve the search problem is

$$\Omega\left(\sqrt{\frac{N}{h}}\right)$$
.



The projector  $P_{\{0,1\}}$  acts as:

$$P_{\{0,1\}} :: \begin{pmatrix} \rho_{00} & \rho_{01} & \rho_{02} \\ \rho_{10} & \rho_{11} & \rho_{12} \\ \rho_{20} & \rho_{21} & \rho_{22} \end{pmatrix} \mapsto \begin{pmatrix} \rho_{00} & \rho_{01} & 0 \\ \rho_{10} & \rho_{11} & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

whilst the 'coherence-projector'  $\omega_{\{0,1\}}$  acts as:

$$\omega_{\{0,1\}} :: \begin{pmatrix} \rho_{00} & \rho_{01} & \rho_{02} \\ \rho_{10} & \rho_{11} & \rho_{12} \\ \rho_{20} & \rho_{21} & \rho_{22} \end{pmatrix} \mapsto \begin{pmatrix} 0 & \rho_{01} & 0 \\ \rho_{10} & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}.$$

That is,  $\omega_{\{0,1\}}$  corresponds to the linear combination of projectors:

$$P_{\{0,1\}} - P_{\{0\}} - P_{\{1\}}.$$



Can define coherence projectors for any 1:

$$\omega_I := \sum_{\widetilde{I} \subseteq I} (-1)^{|I|+|\widetilde{I}|} P_{\widetilde{I}}.$$

► Alternate definition of maximal order *h*:

$$\mathbb{1}_N = \sum_{I,|I|=1}^h \omega_I, \text{ for all } N \geq h$$



Apply  $\mathbb{1}_N$  to a state s

$$\Rightarrow s = \sum_{I,|I|=1}^{h} s_I$$
, with  $s_I := \omega_I s$ .

Think of  $s_I$  as the "coherences" between the slits in I.



▶ Oracle defined to act on projectors  $P_I$ , hence acts on  $\omega_I$ .

▶ Oracle only acts non-trivially on parts of the decomposition "coherently linked" to x, i.e. those  $s_I$  with  $x \in I$ .



Pirsa: 16080030 Page 34/37

In quantum theory

$$\frac{\text{\# terms coherently linked to x}}{\text{\# total terms}} = \frac{N-1}{N^2} \sim \frac{1}{N}$$

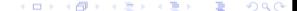
Intuitively speaking, a quantum oracle can only "move" a given state a small amount in a single query.



#### Conclusion and further work

► As far as the search problem goes, all non-trivial (finite) orders of interference are asymptotically equivalent.

▶ Derivation of quadratic lower bound to search from simple physical principles similar to derivations of Tsirelson's bound from information causality, local orthogonality, etc.



Pirsa: 16080030 Page 36/37

#### Conclusion and further work

► Would existence of post-quantum interference compromise security of quantum protocols?

► Verification of delegated quantum computation needs to be proven secure against adversaries with post-quantum quantum dynamics, for example higher-order phase transformations.



Pirsa: 16080030 Page 37/37