

Title: Disorder in spin liquids

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Abstract: Beyond their deceptively featureless ground states, spin liquids are particularly remarkable in the exotic nature of their (fractionalised and gauge charged) excitations. Quenched disorder can be instrumental in nucleating or localising defects with unusual properties, revealing otherwise hidden features of these topological many-body states. This talk discusses how to turn the nuisance of disorder into a powerful probe and origin of new collective behaviour.

Disorder in frustrated magnets

Roderich Moessner

MPI-PKS

John Chalker and Adam Willans

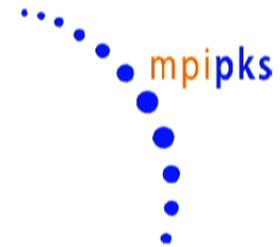
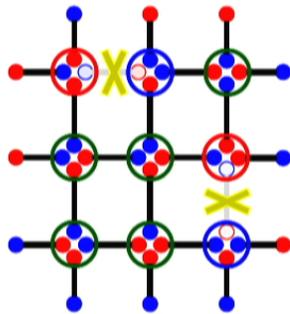
Oxford

Kedar Damle and Arnab Sen

TIFR/IACS Kolkata

Jorge Rehn

Dresden



Outline

Disorder in magnetism

- ▶ distraction, probe, source of new physics

Disorder and frustration

- ▶ history, disorder-free glassiness?

Diluted spin ice

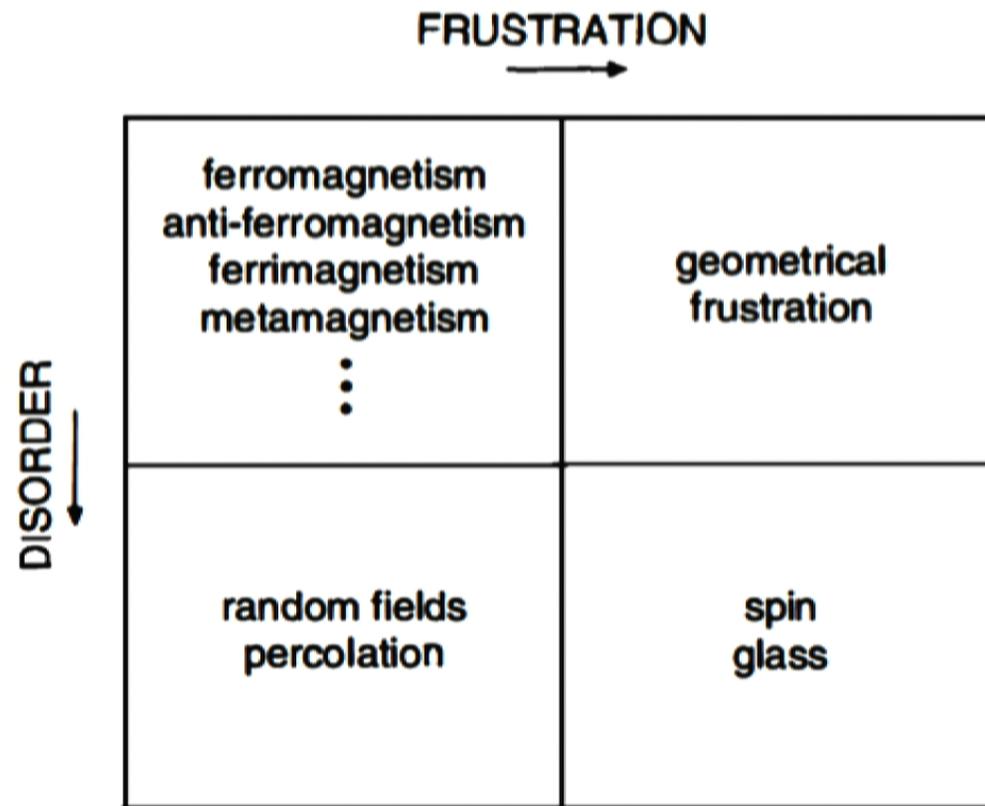
- ▶ phenomenology
- ▶ effective impurity Hamiltonian
 - ▶ topological spin glass

Disorder and topology

- ▶ fractionalisation made visible
- ▶ example: Kitaev spin liquids



Disorder and frustration



Ramirez, 1994



Frustration and glassiness

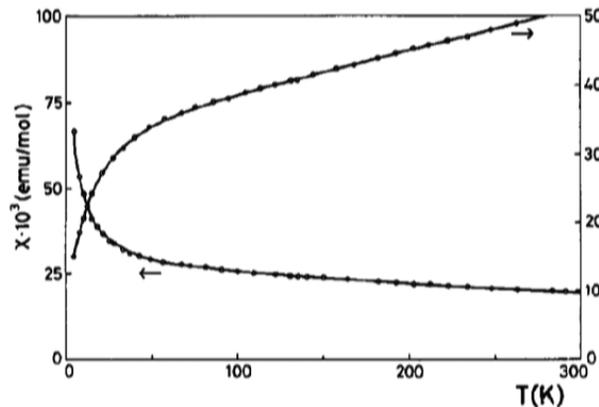
Combination is very common overview: Cepas, Canals PRB 86, 024434 (2012)

- ▶ Ising/Heisenberg
- ▶ $d = 2, d = 3$
- ▶ distortion, dilution

Origin?

- ▶ intrinsic: disorder-free glassiness
- ▶ extrinsic: a little disorder → a little glassiness
 - ▶ freezing often incomplete
- ▶ many possible mechanisms for slow dynamics

SCGO, the archetypal frustrated magnet Obradors 88⁺



started 'modern' frustrated magnetism

- ▶ frustration ratio $f > 117$
- order by disorder?

- ▶ kagome: yes

Chalker et al., Chandra et al., Berlinsky et al.,
Huse et al. (early 90s)

- ▶ pyrochlore+SCGO: no

freezes at low temperature?

e.g. Lee et al. (1996), Limot et al. (2000)

temperature range. The high temperature susceptibility follow a Curie-Weiss law

$$\chi = \frac{C}{T - \theta}$$

C

T-θ

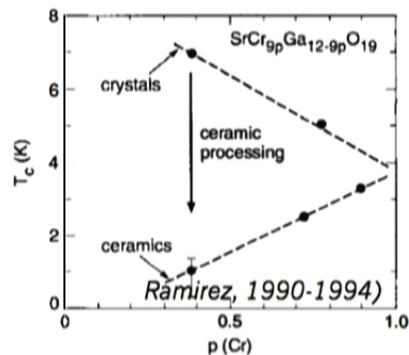
with an effective magnetic moment corresponding to spin-only ($S=3/2$) Cr^{3+} magnetic moment and a Curie temperature $\theta = -492\text{K}$ which shows that strong antiferromagnetic interactions do exist among the Cr ions. It is noteworthy that an antiferromagnetic or spin glass transition below 4.2K would give a very high $|\theta|/T_g$ ratio ($|\theta|/T_g > 117$). For the



SCGO: frustration and freezing

Freezing without disorder?

▶ $\lim_{x \rightarrow 0} T_F(x) > 0$



SCGO \implies kagome: dilution x

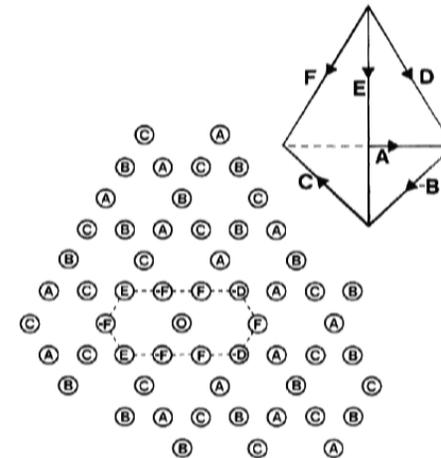
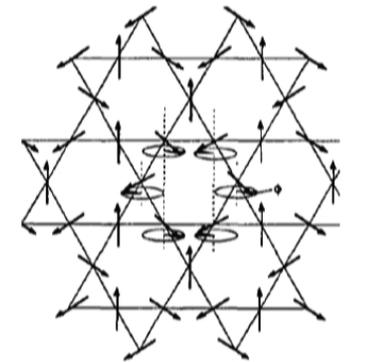
▶ weathervane modes / spin origami

Chandra et al., Shender et al., 1993

▶ classical kagome orders Chern+RM

Disorder level in SCGO underestimated

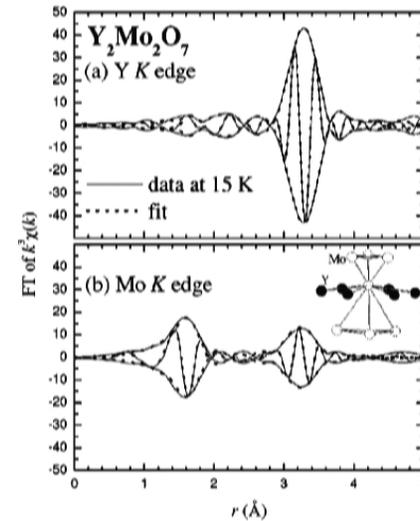
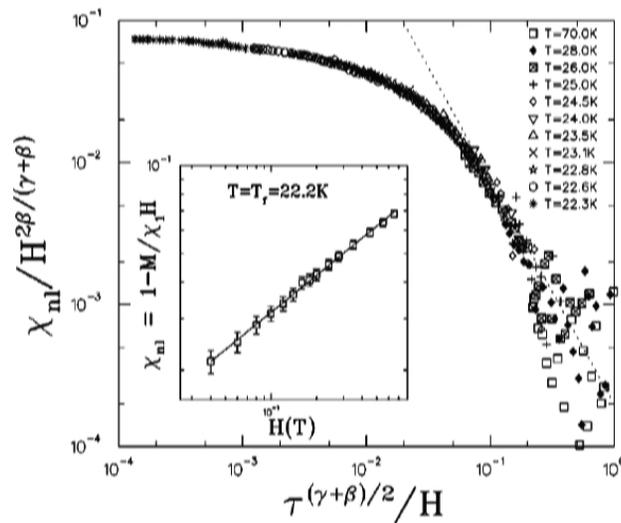
▶ diluted SCGO does not freeze Rehn et al.



Y₂Mo₂O₇ pyrochlore: bona fide spin glass

Conventional freezing for immeasurably small disorder Gingras et al., 1997

- ▶ disorder level in Y₂Mo₂O₇ underestimated? Booth et al., 2000
- ▶ bond disorder leads to freezing Bellier-Castella, 2001; ..., Andreev et al., 2010



Open questions

How well do we estimate disorder?

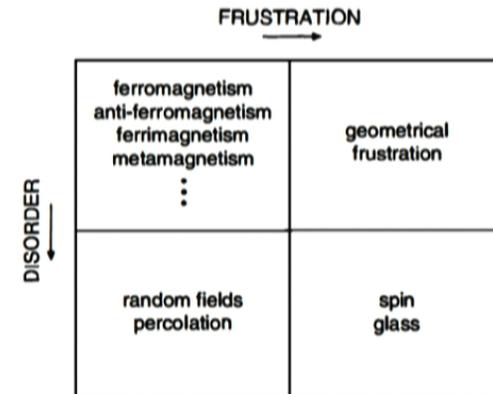
- ▶ important input parameter!

Coexistence of liquidity and glassiness

- ▶ “who” freezes?

Features of freezing specific to spin liquids?

- ▶ topology + disorder



Spin freezing in spin ice

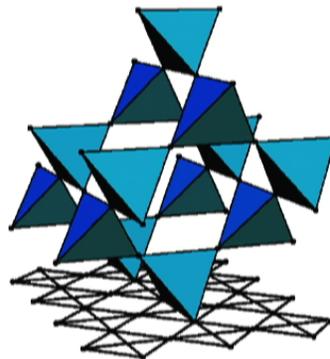
Many experimental studies

- ▶ “how spin ice freezes” Matsuhira, Schiffer, Bramwell, ..., Petrenko, Gaulin, Pomaranski
- ▶ slow dynamics not properly understood

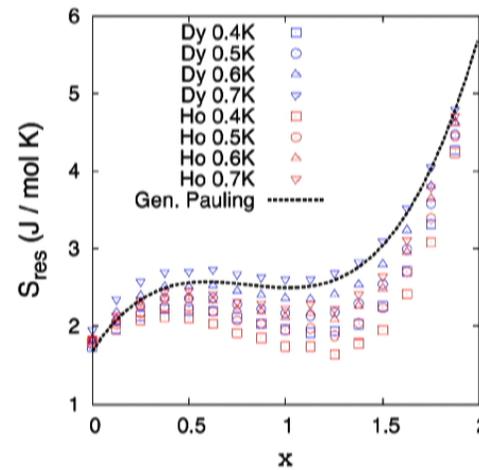
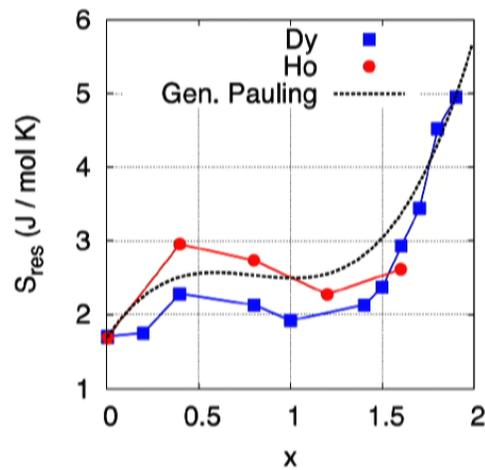
Pristine dipolar spin ice

- ▶ (elusive) thermodynamic ordering transition Melko et al., 2004
- ▶ Coulomb phase can persist to $T = 0$ upon fine-tuning

Isakov et al., 2005



Diluted spin ice



Non-monotonic “zero-point entropy” with dilution x Ke et al., 2007

- ▶ need to consider interactions between defects Lin et al., 2013
 - ▶ effective impurity Hamiltonian?

Analytical approach: dumbbell model

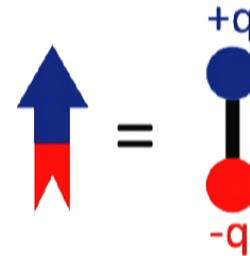
Spin = dipole = pair of charges

- ▶ $|\mu| = |q|a$ ($a \dots$ premedial lattice constant)

Transparently accounts for Coulomb phase

Castelnovo et al., 2008

- ▶ easily visualised
- ▶ stable to $T = 0$
- ▶ monopole excitations



Genesis of ghost spins

Visualize in $d = 2$

- ▶ tetrahedron = vertex

Satisfied tetrahedron

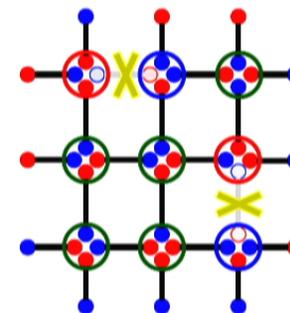
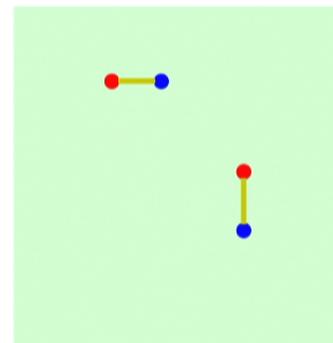
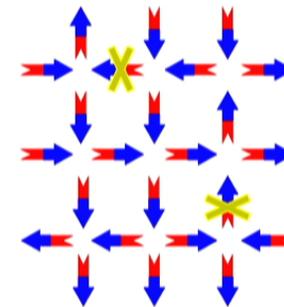
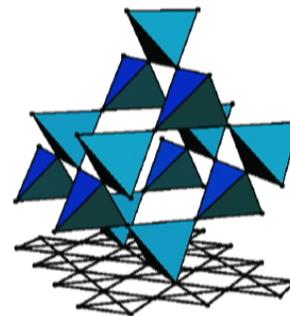
- ▶ charge neutral

Missing spin

- ▶ pair of charged tetrahedra
 - ▶ with dipole moment

Delete neutral sites

- ▶ ghost spins left over
 - ▶ moment opposite to missing spin
- ▶ embedded in medium



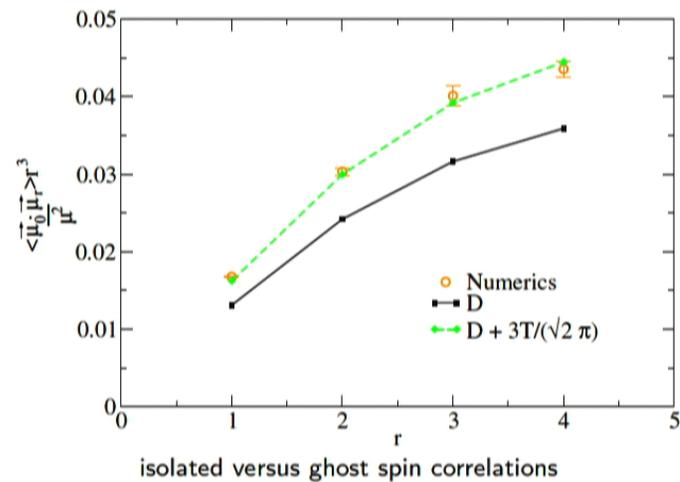
Interactions between ghost spins

Compare free and **ghost** spin pair

- ▶ enhanced (entropic) dipolar interactions

Entropic interaction

- ▶ as in pristine n.n. spin ice
- ▶ missing dipoles interact as if
 - ▶ not missing
 - ▶ not dipoles



Freezing of ghost spins

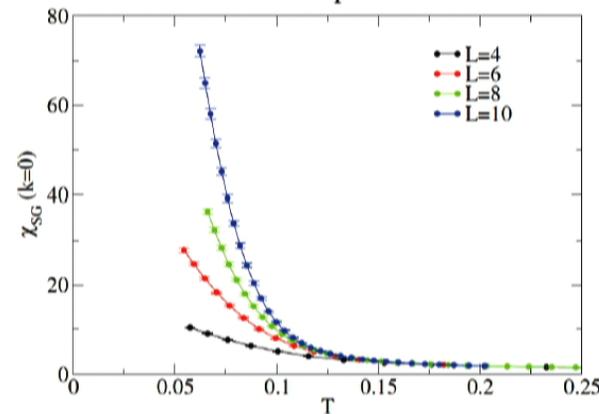
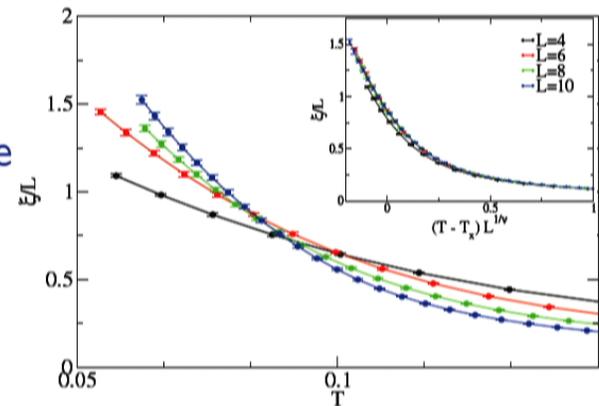
Low dilution = sparse ghost spins

- ▶ simulation with effective \mathcal{H} possible

Spin freezing of ghost spins

- ▶ dipolar spin glass on pyrochlore lattice
- ▶ scaling collapse of ξ
- ▶ non-linear susceptibility

- ▶ $T_x = 0.95Dx$
- ▶ $T_c(x) = T_x / (1 - 3T_x / \sqrt{2\pi D})$
- ▶ $D \approx 1.41K$ for Dy/Ho spin ice



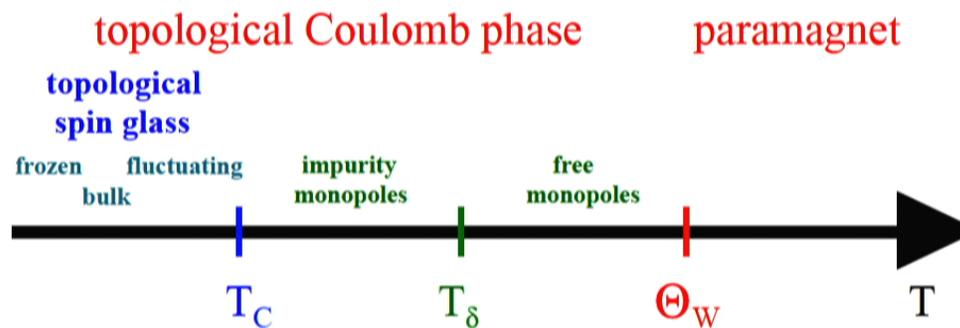
Phase diagram

Different regimes and crossovers

- ▶ paramagnet \rightarrow regular spin ice
- ▶ bulk \rightarrow impurity monopoles

Glass phase

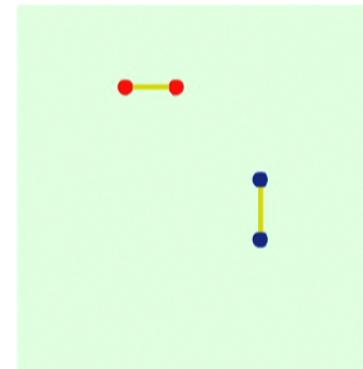
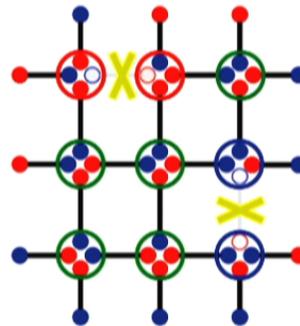
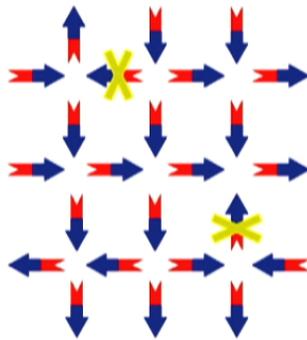
- ▶ continuous onset of frozen moment
- ▶ competition: entropy vs. random frozen fields



Impurity monopoles

Flip “Dirac string” between ghost spins

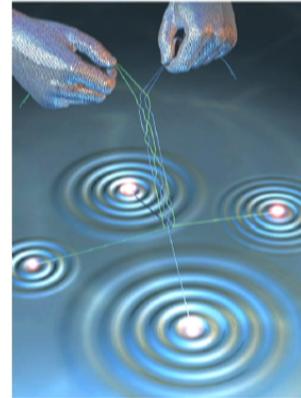
- ▶ monopoles with usual magnetic charge
- ▶ cheaper excitation energy δ than bulk Δ
 - ▶ dominate for $xe^{-\delta/T} > e^{-\Delta/T}$



Disorder in topological physics

Impurities/defects can be (multiply) gauge charged

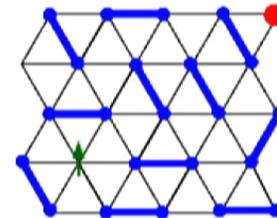
- ▶ interactions (poss. quantum statistical) via bulk
- ▶ underpins topological quantum computation



Spin liquids are such topological phases

- ▶ impurity physics generally interesting
- ▶ diagnosing spin liquids
- ▶ access fractional particles

Scientific American



Topological excitations in Kitaev spin liquid

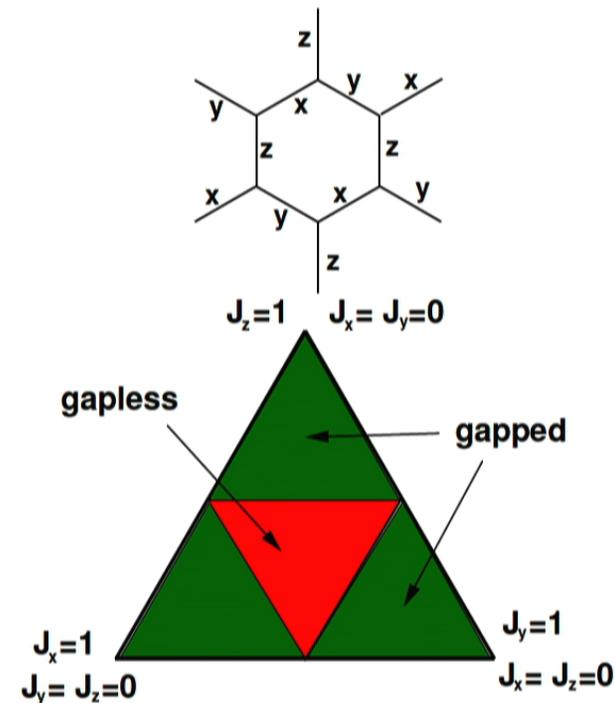
$$H = -J_x \sum_{x\text{-links}} \sigma_j^x \sigma_k^x - J_y \sum_{y\text{-links}} \sigma_j^y \sigma_k^y - J_z \sum_{z\text{-links}} \sigma_j^z \sigma_k^z$$

Spins-1/2 on honeycomb lattice

- ▶ includes strong “spin-orbit coupling”

Two **spin liquid phases**

- ▶ one gapped, one gapless
- ▶ short-range spin correlations throughout Baskaran et al.
- ▶ **model is exactly soluble**

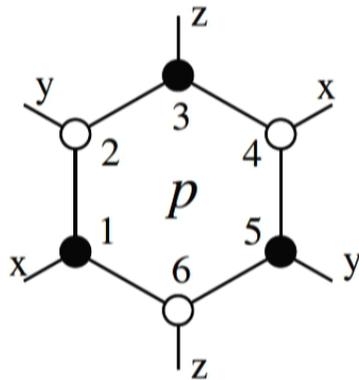


Emergent degrees of freedom: fluxes and fermions

Local conserved quantities:

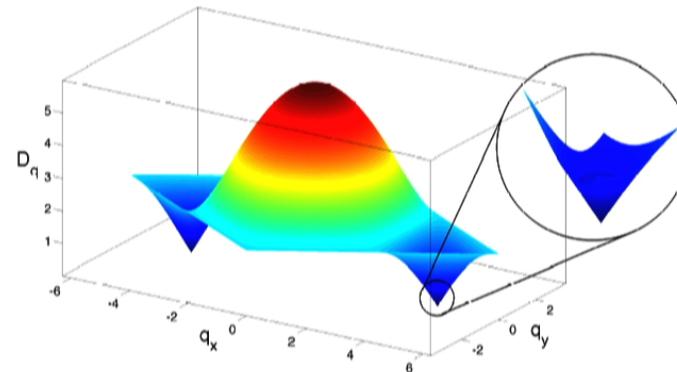
- ▶ plaquette fluxes $[W_b, H] = 0$
- ▶ emergent Z_2 gauge field

$$W_p = \sigma_1^x \sigma_2^y \sigma_3^z \sigma_4^x \sigma_5^y \sigma_6^z$$



In each flux sector

- ▶ quadratic Majorana Fermion Hamiltonian
- ⇒ soluble honeycomb tight-binding model



Signatures???

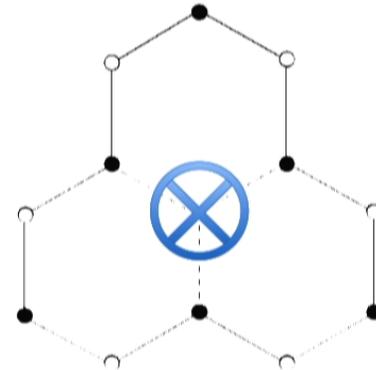
T/S matrices — dynamics — disorder



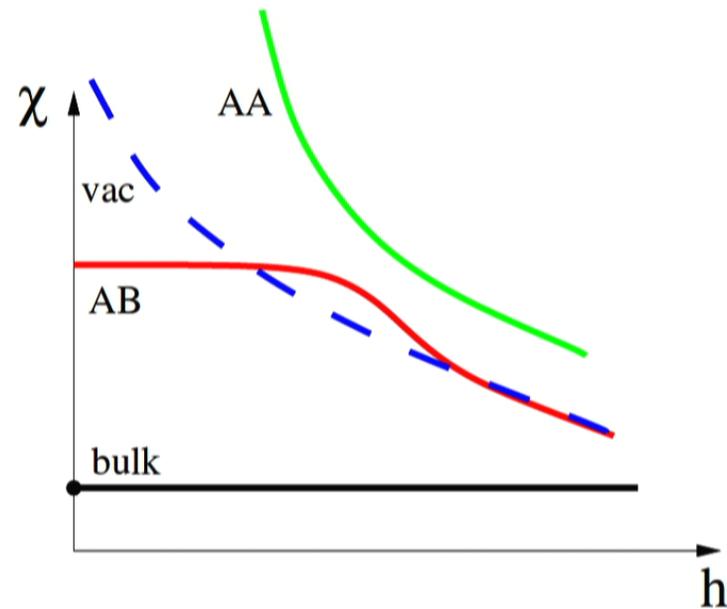
Spontaneous flux attachment

Ground state flux sector

- ▶ No vacancies: flux free Lieb
 - ▶ Each vacancy binds one Z_2 flux
- ⇒ elegant way to see emergent fluxes in ground state!

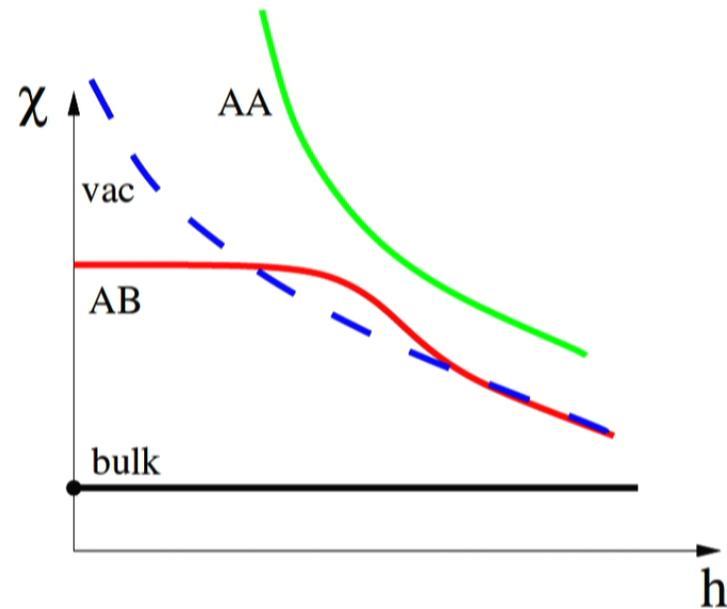


Susceptibility of vacancy configurations



<i>bulk</i>	<i>single vacancy</i>	<i>AB pair</i>	<i>AA pair</i>
$\chi \sim \text{const}$	$\chi \sim \log h $	$\chi \sim (\text{large}) \text{ constant}$	$\chi \sim 1/[h \log h ^{3/2}]$

Susceptibility of vacancy configurations



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$\chi \sim \text{const}$	$\chi \sim \log h $	$\chi \sim (\text{large}) \text{ constant}$	$\chi \sim 1/[h \log h ^{3/2}]$

Local moments and their interactions: gapped phase

Vacancy zero mode: uncompensated bipartite lattice

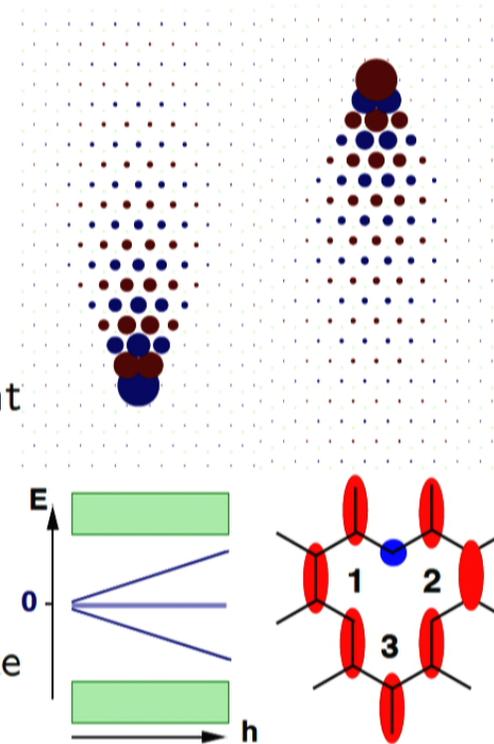
- ▶ confined to a wedge
- ▶ amplitude exponentially decaying along wedge
- ▶ wavefunction given by Pascal's triangle

Vacancy moment given by adjacent zero-mode amplitude

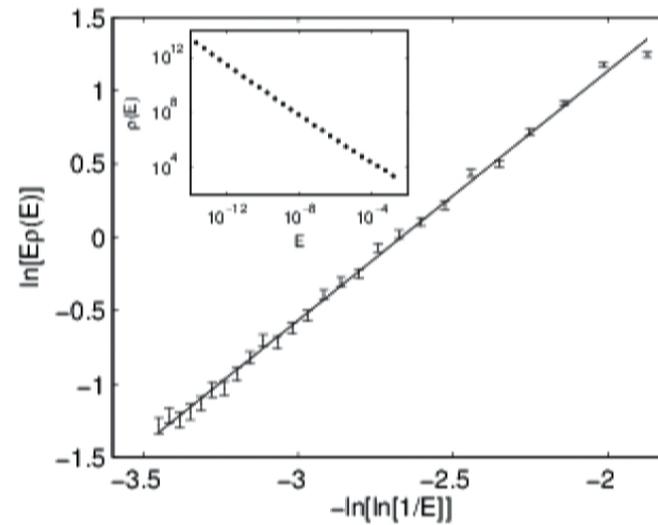
- ▶ $g_{eff} \rightarrow 1$ for $J_{x,y}/J_z \rightarrow 0$
- ▶ $g_{eff} \rightarrow 0$ for $J_x + J_y \rightarrow J_z$

Many vacancies: random bipartite hopping problem

- ▶ $\rho(E) \sim 1/[E(\log E)^{1.7}]$



Numerical results for d.o.s. of the Kitaev model



- ▶ divergent density of states near zero energy:
 $\rho(E) \approx 1/[E(\log E)^{1.7}]$
- ▶ divergent correlation length near zero energy

Bond randomness in gapless phase, $\sigma = \langle (\delta J)^2 \rangle$

Fermionic excitations: massless Dirac particles

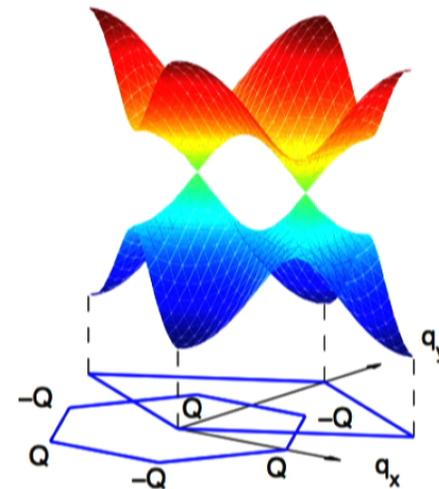
- ▶ exchange disorder translates into random vector potential

Power law exponents vary with disorder strength σ Ludwig et al, 1994

- ▶ density of states: $\rho(E) \sim E^{\frac{1-\sigma}{1+\sigma}}$
- ▶ specific heat: $C \sim T^{\frac{2}{1+\sigma}}$

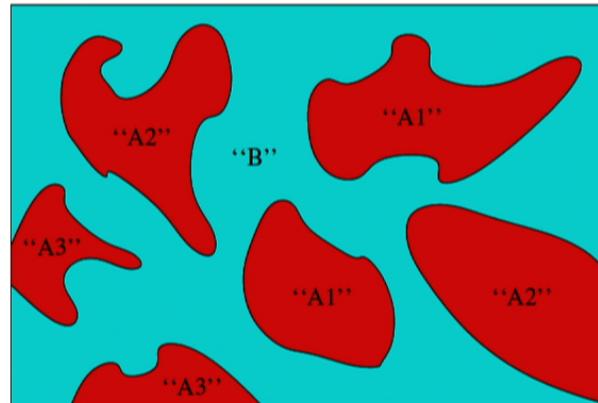
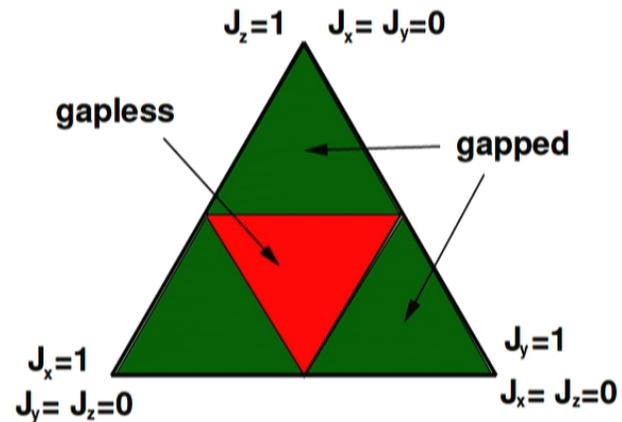
No qualitative change to susceptibility

⇒ implications for other spin liquids?



More disorder

- ▶ $\pm J$ disorder \Rightarrow can be 'gauged away'
- ▶ Slowly varying, strong disorder - percolating gapless background



\Rightarrow implications for transport?

Topological spin glass in diluted spin ice

Dilution leaves Coulomb phase intact

- ▶ but generates ghost spins

Ghost spins interact like normal spins

- ▶ dipolar interactions
- ▶ entropic interactions

Ghost spins freeze

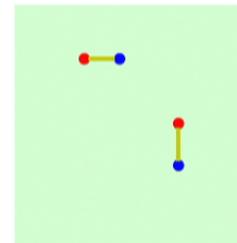
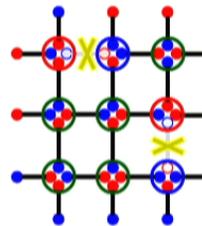
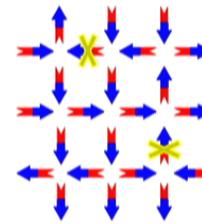
- ▶ glass + topological phase

Coulomb phase acts as substrate for

- ▶ emergent ghosts
- ▶ entropic interactions

General feature of topological phases

- ▶ cf. topological quantum computing



Disorder in frustrated magnets

Many roles and types

- ▶ useful, interesting, nuisance

Topological glass in spin ice

- ▶ dipolar glass of ghost spins

Vacancies in Kitaev QSL

- ▶ ground state: flux nucleation
- ▶ excitations: local moment formation
- ▶ new collective physics

Many other instances

- ▶ orphan physics in classical SL
 - ▶ fractionalised $1/2$ or $1/3$ moment
 - ▶ random Coulomb antiferromagnet

