

Title: A (Nearly) Weaker-Than-Gravity Bound on Dark Matter Electromagnetism

Date: May 10, 2016 01:00 PM

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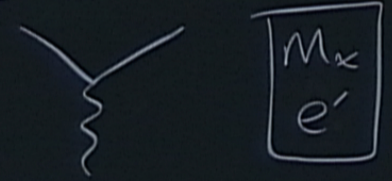
Abstract: 

An unbroken  $U(1)$  is a minimal possibility for a dark matter self interaction, and may even be associated with dark matter stability. However, such an interaction faces incredibly strong constraints due to collective plasma effects, which dominate over 2-to-2 scattering by an order-of-magnitude of orders-of-magnitude. I will discuss the physics of these collective effects, and show preliminary results of simulation. The constraint of such a self interaction is estimated to be nearly as weak as gravity.

# A (nearly) weaker-than-gravity bound on DM EM

W. B. Feldstein & Luis O. Silva |xxxxxxx

- DM with  $U(1)'$  (unbroken)



- DM stability?
- Rich DM astrophysics

BC:  $\frac{\sigma}{m} \lesssim \frac{1 \text{ cm}^2}{g}$  (hard)

→ Coulomb scattering  
(Kamionkowski 08)

$\frac{\langle \sigma \Delta E \rangle}{m \Delta E} \lesssim \frac{1 \text{ cm}^2}{g}$

BC:  $t_{\text{scatter}} \gtrsim t_{\text{merge}} \sim 9 \text{ Gyr}$

Subhalo:  $t_{\text{scatter}} \gtrsim t_{\text{turn}} \sim 10 \text{ Gyr}$

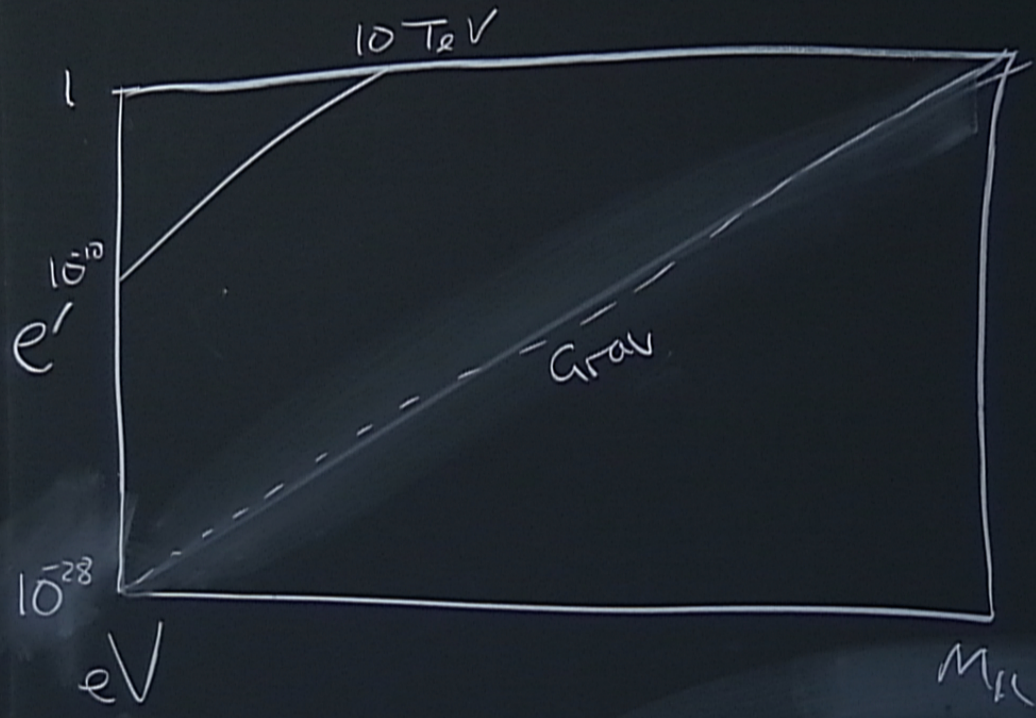
$$\sigma_{\text{Coulomb}} \approx \frac{e^4}{4\pi m^2 v^4} \ln(\dots)$$

$$\frac{t_{\text{univ}}}{t_{\text{scatter}}} \approx n \sigma v t_{\text{univ}} \approx \frac{e^4 \rho M R}{m^3 v^3 \sqrt{\rho}} \lesssim 1$$

$$\frac{\rho}{m} \quad \frac{1}{H_0} \approx \frac{M R}{\sqrt{\rho}}$$

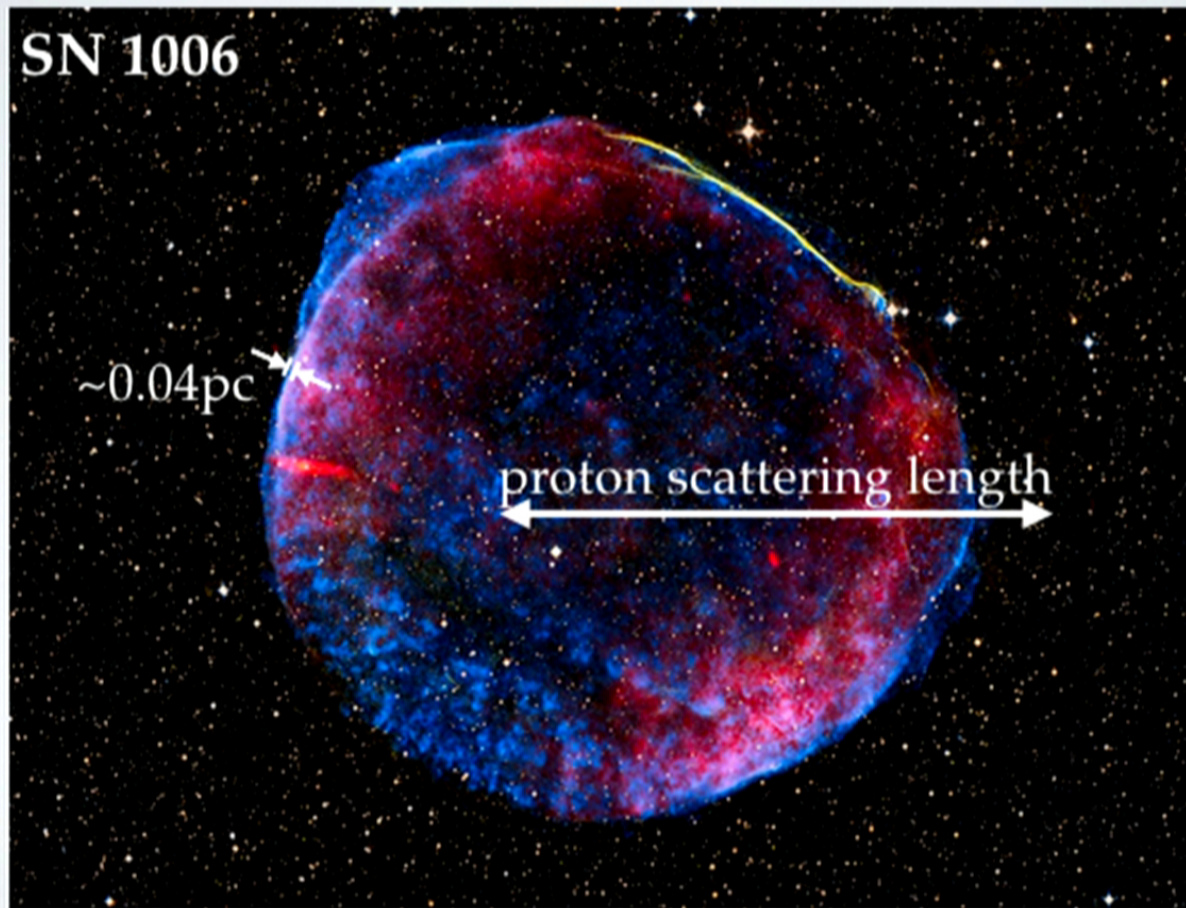
$$v \sim v_{\text{vir}} \sim \sqrt{\frac{GM}{R}}, \quad R^3 = \frac{M}{\rho}$$

$$\frac{e}{m} \lesssim \frac{1}{M_{\text{pl}}} \left( \frac{\rho}{m} \right)^{\frac{1}{3}} \left( \frac{M}{m} \right)^{\frac{1}{4}} \left( \frac{10^{67} \text{ GeV}}{m} \right)^{\frac{1}{4}}$$



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 w. B  
 • D  
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# SNR SHOCK FRONTS



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# SNR SHOCK FRONTS

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- Plasma of free electrically charged particles
- No particle collisions
- **Collisionless plasma regime**
  - Text-book physics + active research
  - Complex dynamics

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# PLASMA INSTABILITIES

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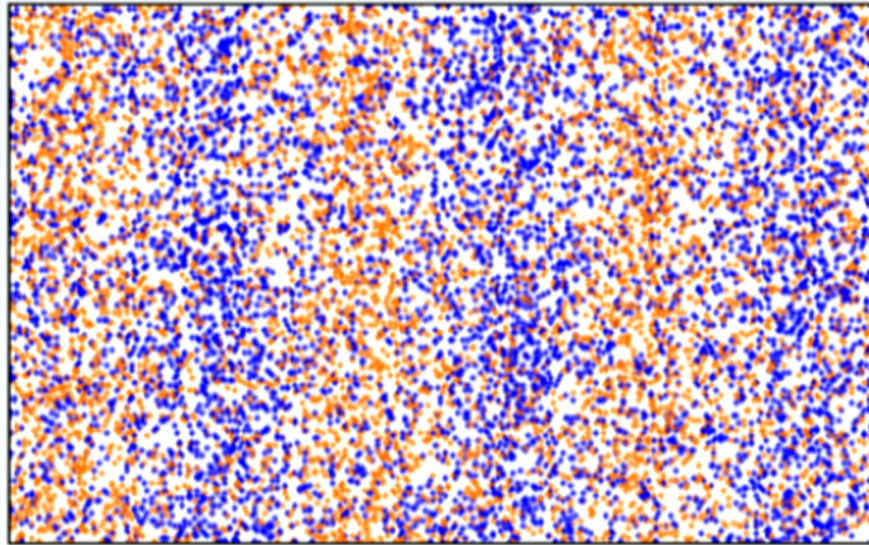
# LONGITUDINAL INSTABILITY

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Single cold plasma

$\chi^+$   $\chi^-$

moving plasma oscillation



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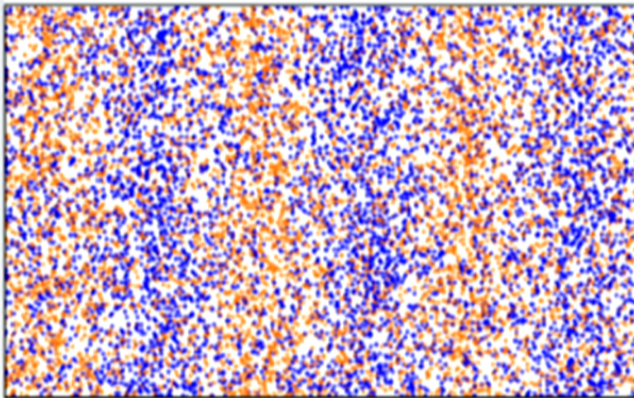


# LONGITUDINAL INSTABILITY

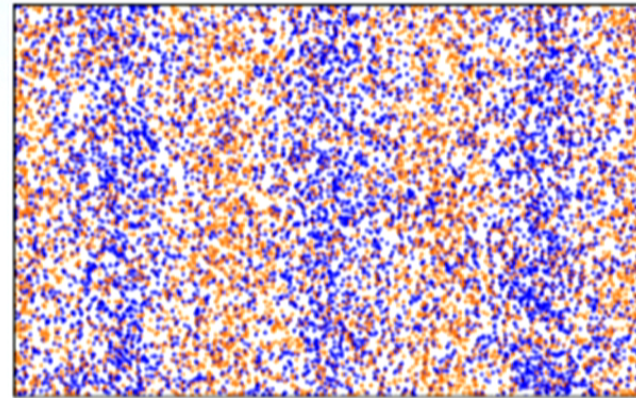
Single cold plasma

$\chi^+$   $\chi^-$

plasma rest frame



wave frame



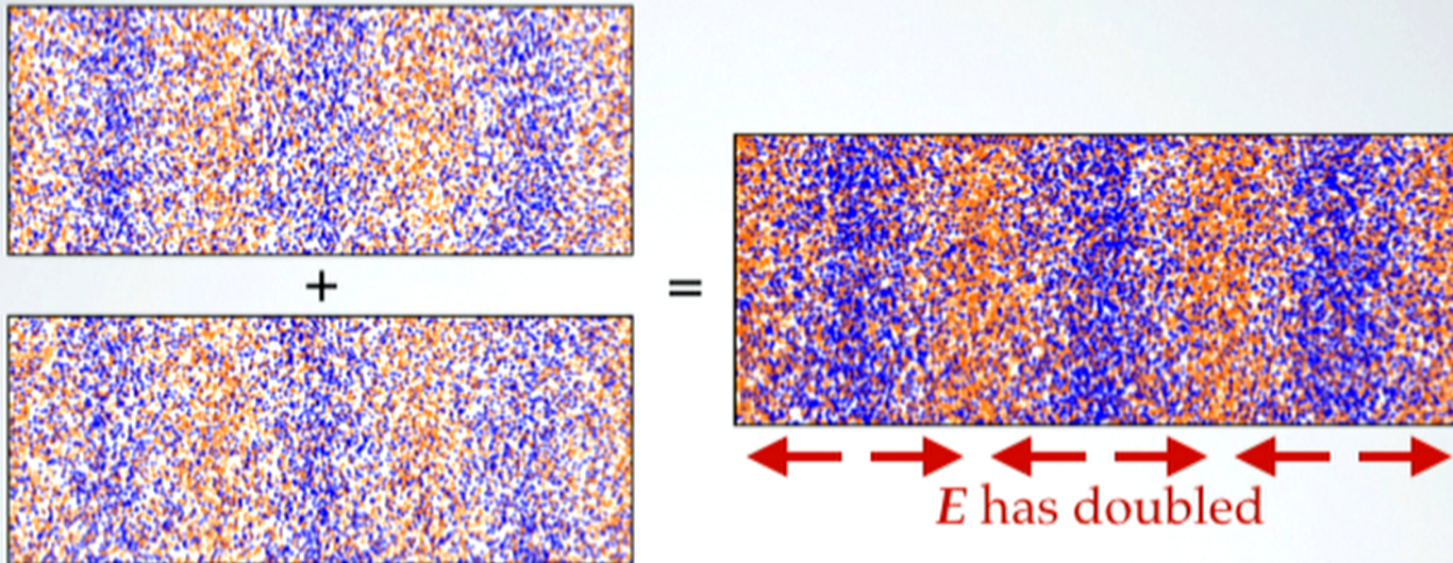
$E$   $\longleftrightarrow$   $\longleftrightarrow$   $\longleftrightarrow$   $\longleftrightarrow$   $\longleftrightarrow$   $\longleftrightarrow$



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# LONGITUDINAL INSTABILITY

Two oppositely moving plasmas



Bunching generates *twice* required E field  
E field causes *twice* as much bunching

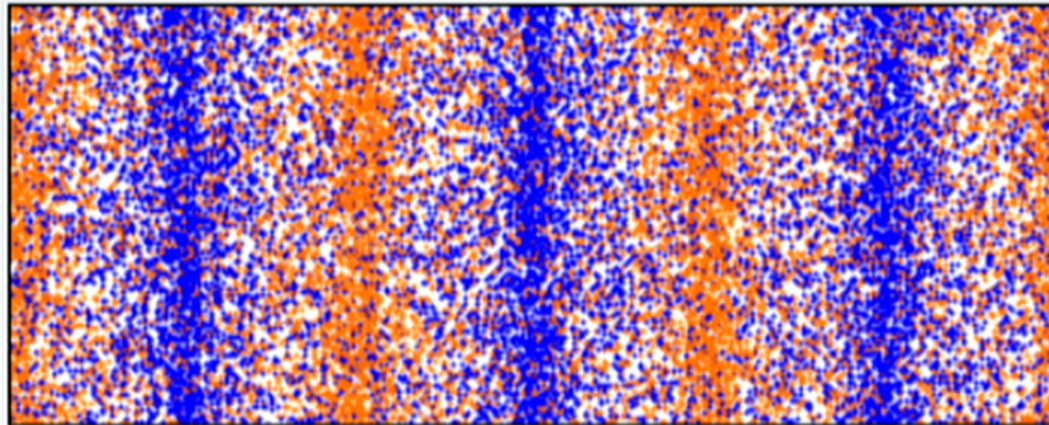
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# LONGITUDINAL INSTABILITY

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Two oppositely moving plasmas:

**exponential runaway**



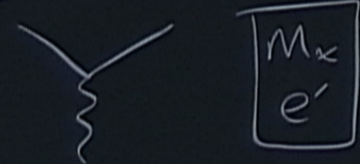
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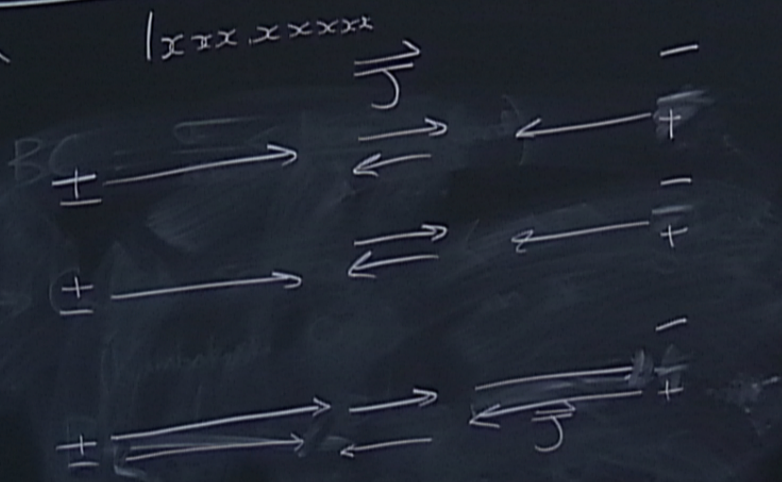
# A (nearly) weaker-than-gravity bond on DM EM

W. B. Feldstein & Luis O'Silva

DM with  $U(1)'$  (unbroken)



- DM stability?
- Rich DM astrophysics



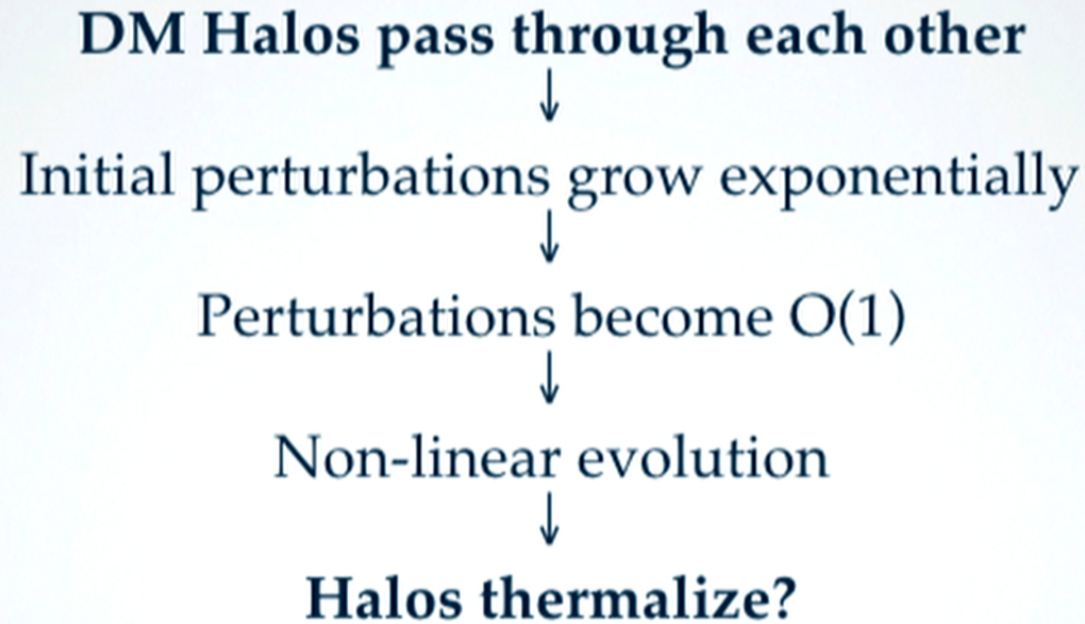
$$\sigma_{\text{Coulomb}} \approx \frac{e^2}{4\pi m^2}$$

$$\frac{t_{\text{univ}}}{t_{\text{scatter}}} \approx n \sigma \frac{v}{m}$$

$$V \sim V_{\text{vir}} \sim \sqrt{\dots}$$

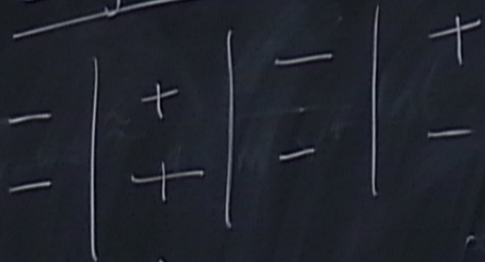
# WHAT HAPPENS TO DM HALOS

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Long Inst



$$\Sigma = e' n \quad s_x$$

↑  
 $x_+ - x_-$

$$E = \Sigma$$

$$\ddot{s}_x = \frac{e'}{m} E = \frac{e'^2 n}{m} s_x$$

↑  
 $\omega_{pe}^2$

$\Gamma_{\text{long}} \sim \omega_{pe}$



Trans

$$J = e' n \frac{s_y}{\lambda} v$$

$$\nabla \times B = J \Rightarrow B \approx \lambda J \approx e' n s_y v$$

$$s_y \approx \frac{e'}{m} v \times B = \omega_{pe}^2 v^2 s_y$$

$\Gamma_{\text{trans}} \sim \omega_{pe} v$

**Q:** What happens in the non-lin. regime?

**Expectation:** Halos merge / disrupted / thermalize

**Approach:** **SIMULATION**

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# WHY SIMULATE

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## True answer 1:

Buried section in Kamionkowski et al '08

Discusses plasma instabilities

Makes similar estimate to us

Almost no one knows about it...

## True answer 2:

3 physics reasons

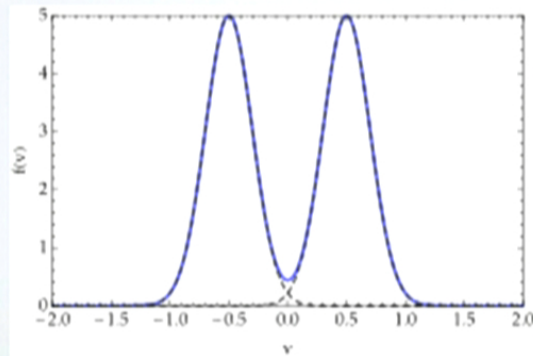


# 3 REASONS TO SIMULATE

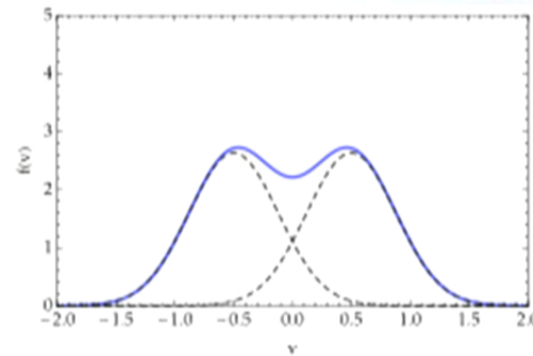
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## 1: Are we sure that the DM gets disrupted?

- can't tell analytically what happens in non-lin. regime
- instabilities can shut off when velocity dispersion large enough
- velocity dispersion  $\sigma$  in galaxy cluster  $\sim 30\%$  of relative velocity
- longitudinal instability shuts off at  $\sigma \sim 38\%$



unstable



stable

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# 3 REASONS TO SIMULATE

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## 2. Exactly where is the bound?

Which astrophysical situation does it come from?

- Longitudinal or transverse instability?
- How far into non-lin regime before halos disrupted?
- Cluster mergers / sub-cluster survival / galaxy survival?

# 3 REASONS TO SIMULATE

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## 3. Are there signatures below the bound?

- Understanding of the astrophysics is “imperfect”
- Transverse instability sets conservative bound
- Longitudinal instability turns on 100 times faster (then saturates)
- Longitudinal instability has observable effect? (e.g. heating)

# 3 REASONS TO SIMULATE

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## What existing simulations do differently:

Different types of plasma (with different dynamics)

- relativistic plasma
- electron-ion
- large background B-field

Aren't looking at overall slowdown

Don't include gravity (nor do we — yet)

# SIMULATIONS

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by N. Shukla in Luis O. Silva group

2D and 3D simulations of collisionless  $e^+ e^-$  plasmas

Use “particle-in-cell” (N-body) method

Simulate 2 bodies of cold plasma colliding

preliminary results

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# VLASOV-MAXWELL EQNS

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$$[\partial_t + v \cdot \nabla_x \pm (E + v \times B) \cdot \nabla_v] f(x, v, t) = 0$$

$$\nabla \cdot E = (n_+ - n_-) / n_0$$

$$\nabla \cdot B = 0$$

$$\partial_t B = -\nabla \times E$$

$$\partial_t E = \nabla \times B - (\langle n_+ v_+ \rangle - \langle n_- v_- \rangle) / n_0$$

times & lengths measured in  $\omega_{pl}^{-1}$

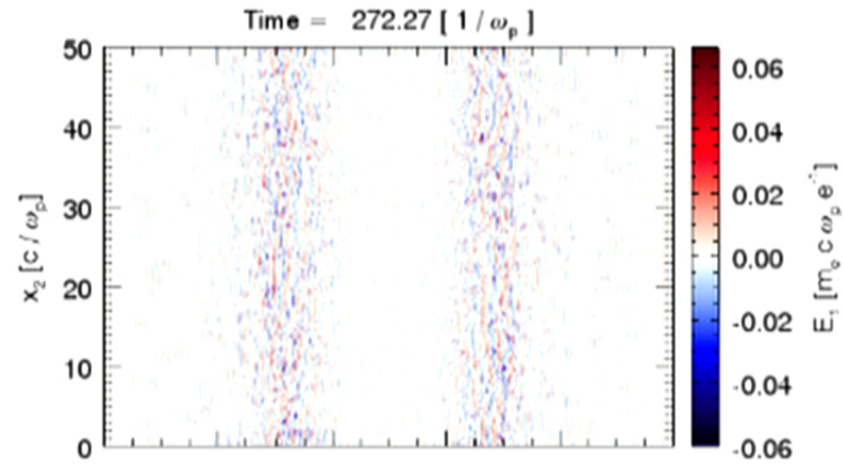
fields in  $(m \omega_{pl} / e)$

**DM parameters enter only in measuring stick**

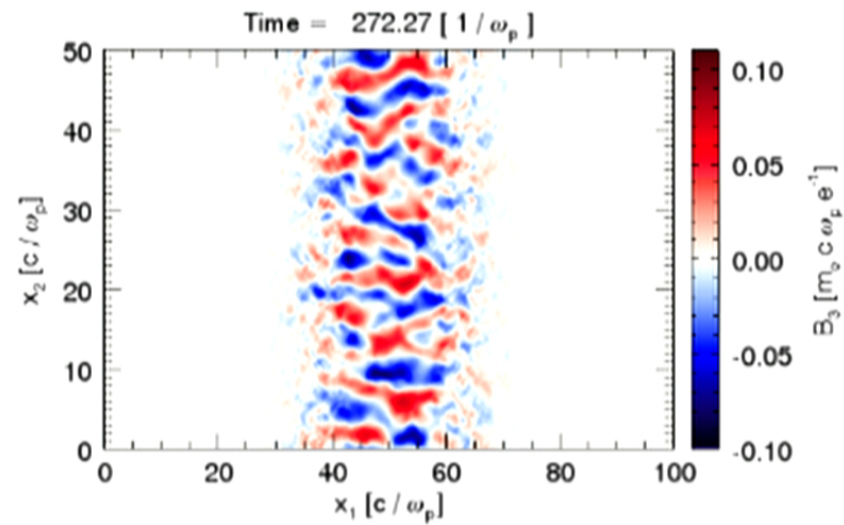
Makes simulations very universal

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longitudinal E field

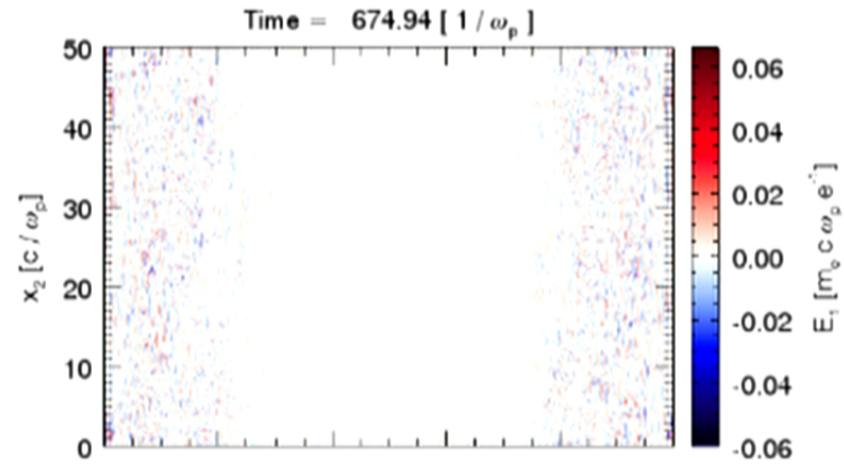


transverse B field

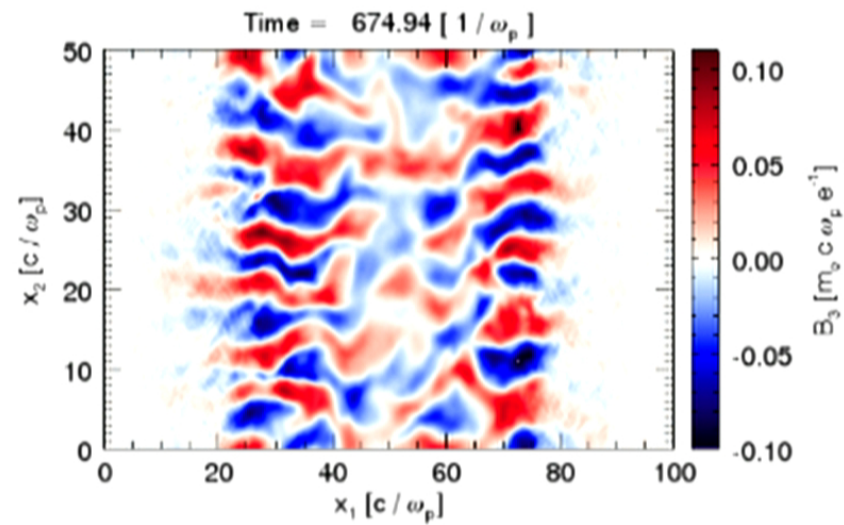


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longitudinal E field



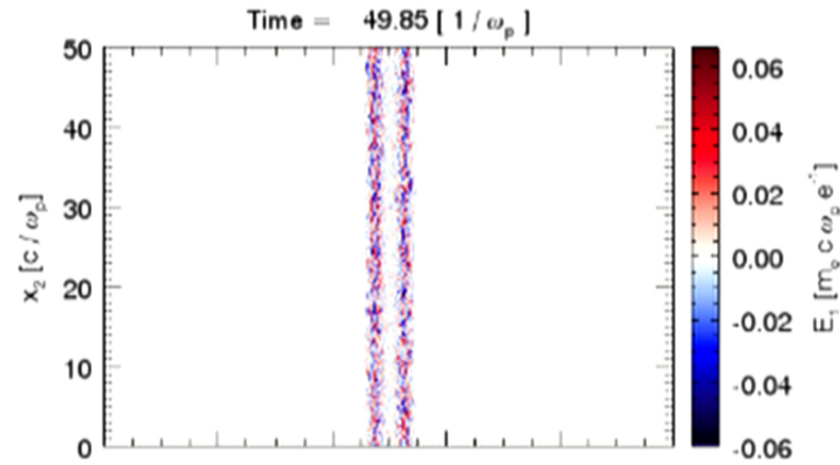
transverse B field



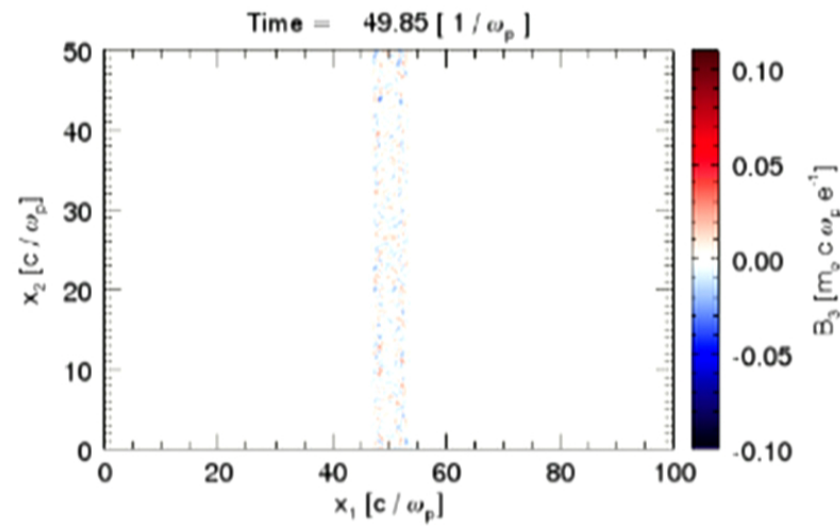
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longitudinal E field

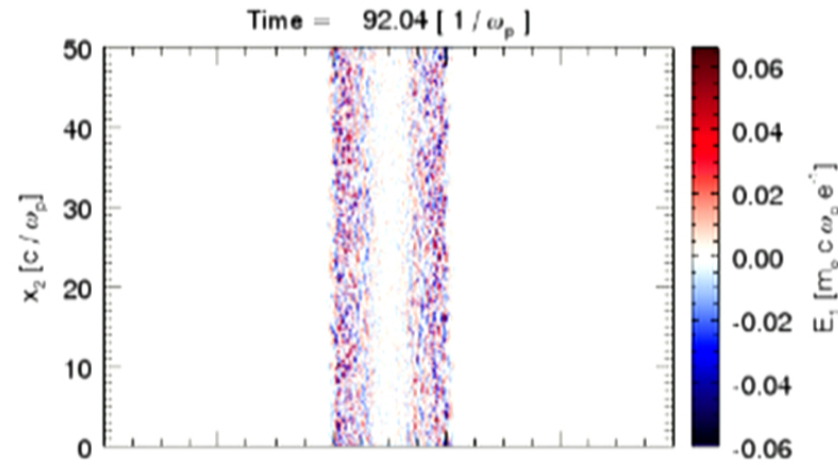


transverse B field

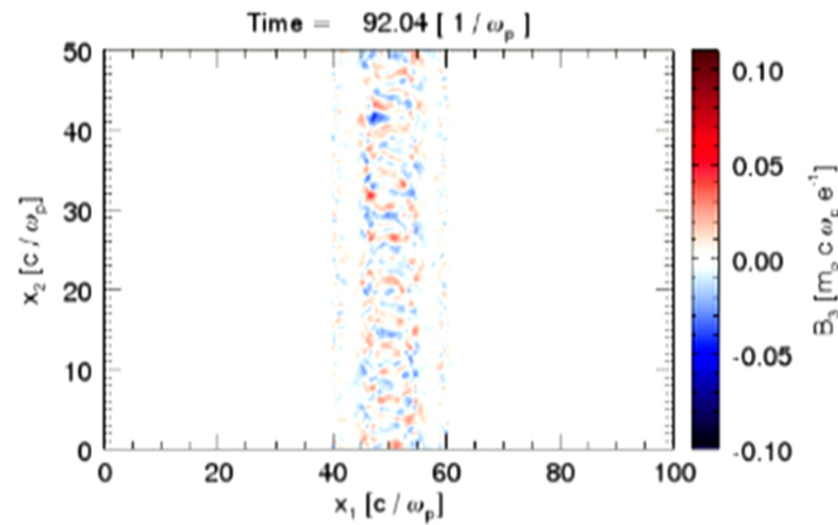


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longitudinal E field

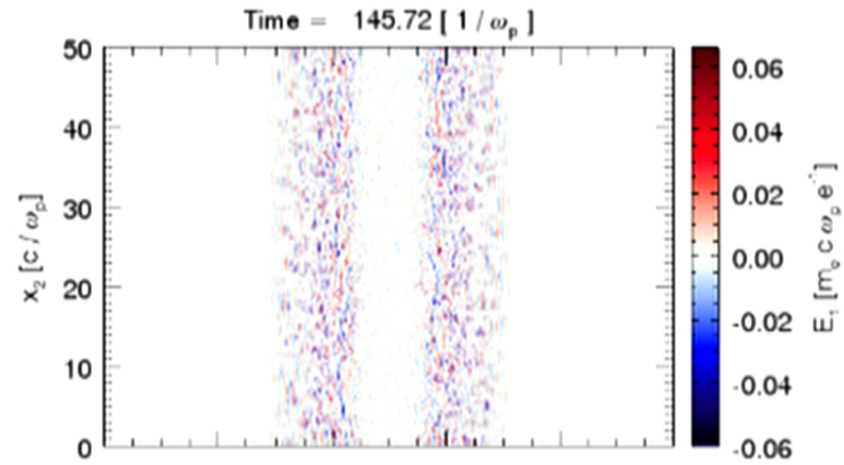


transverse B field

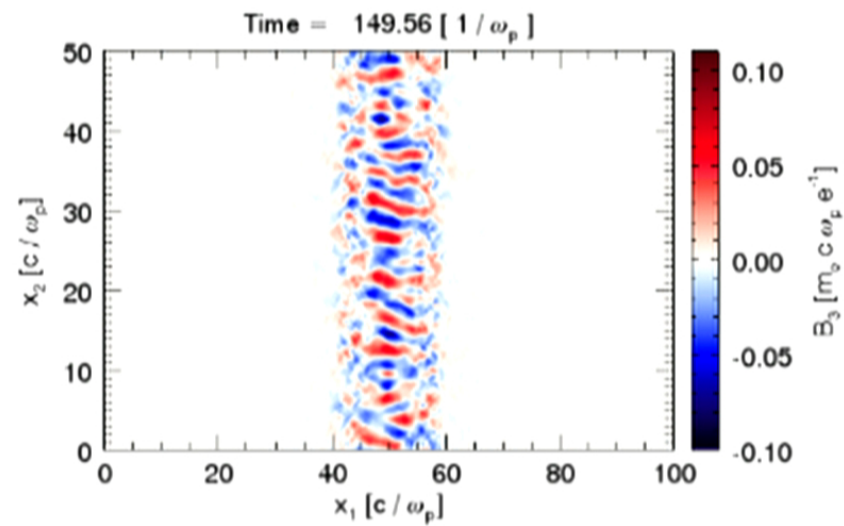


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longitudinal E field

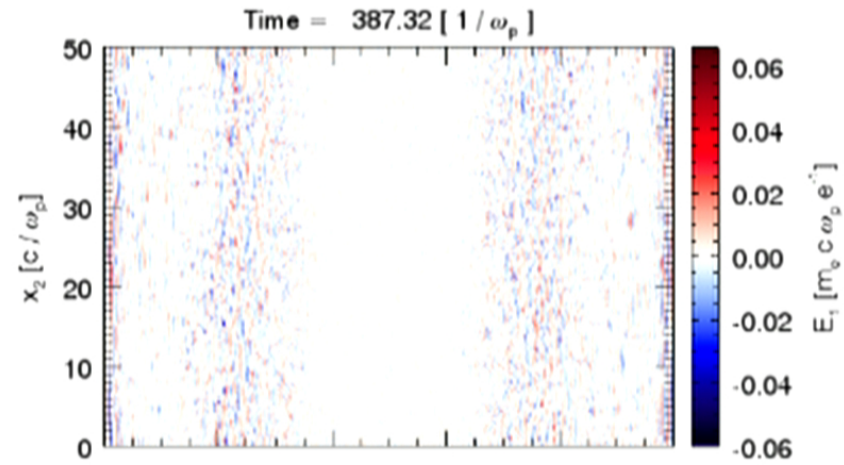


transverse B field

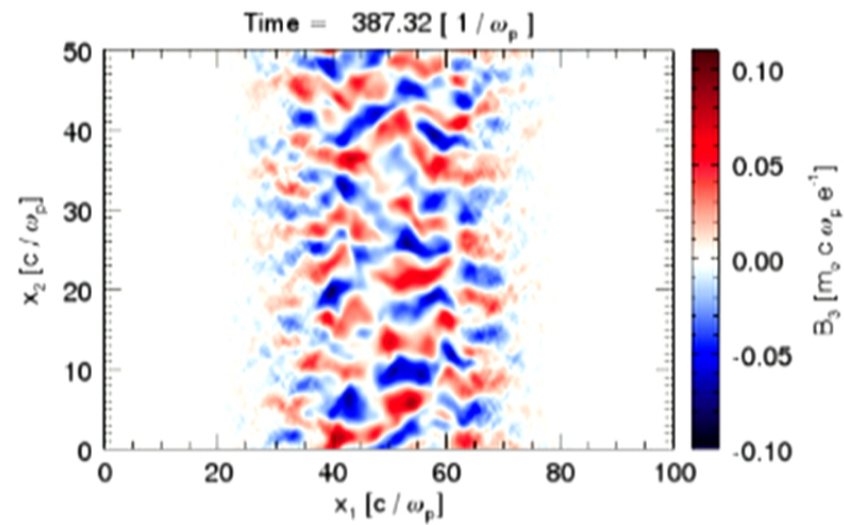


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longitudinal E field



transverse B field



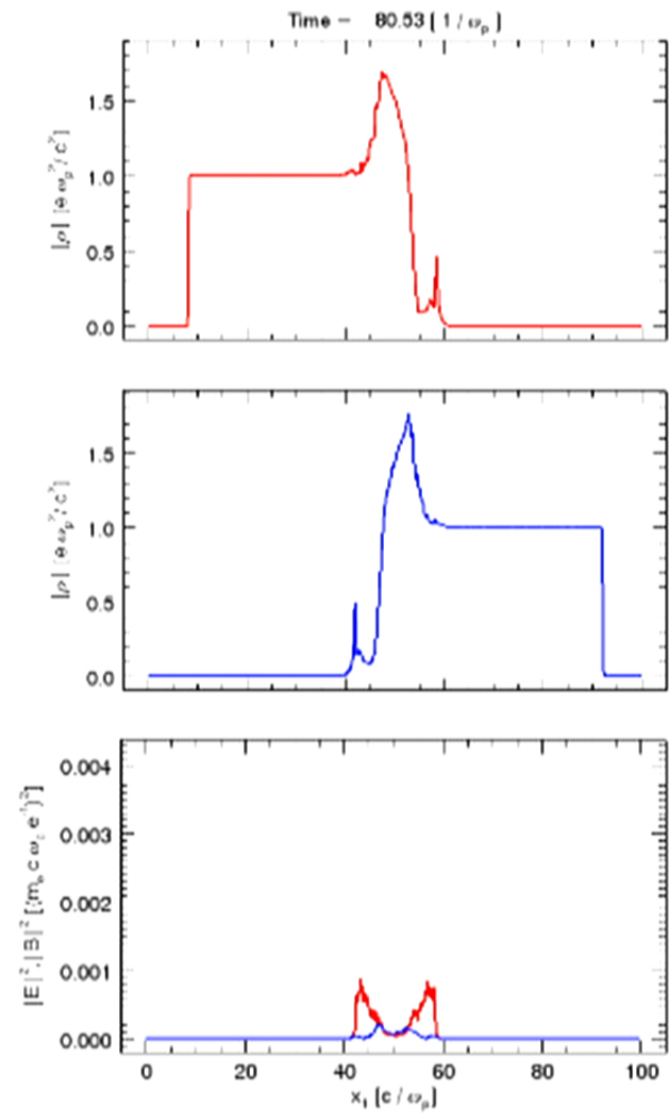
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number density  
(right-movers)

number density  
(left-movers)

longitudinal E field  
transverse B field

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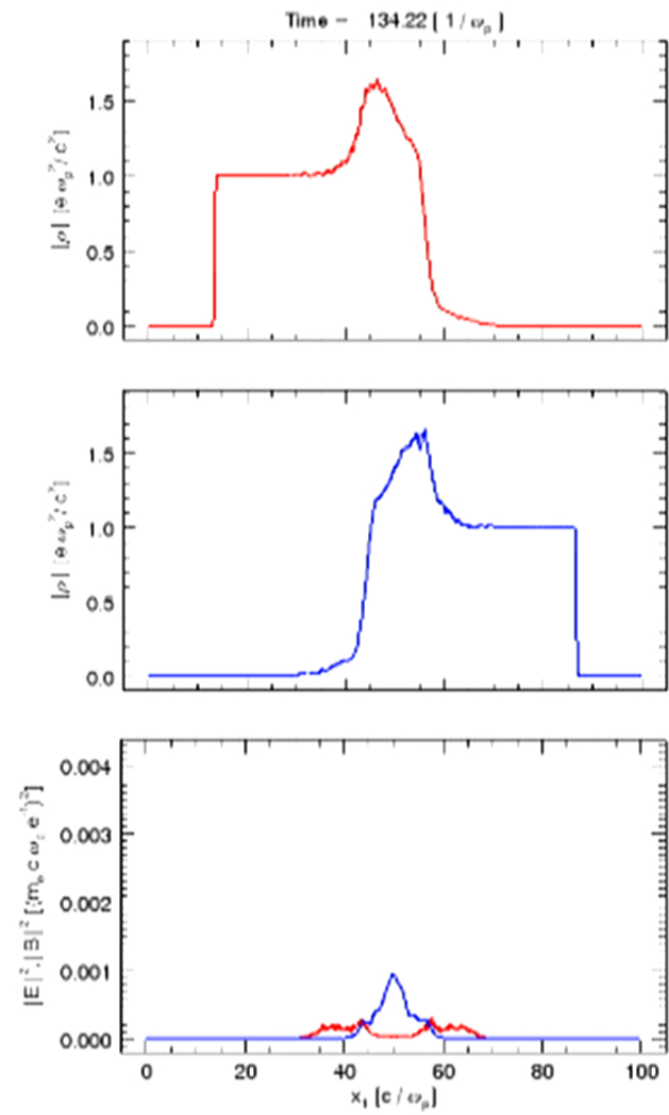


number density  
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longitudinal E field  
transverse B field

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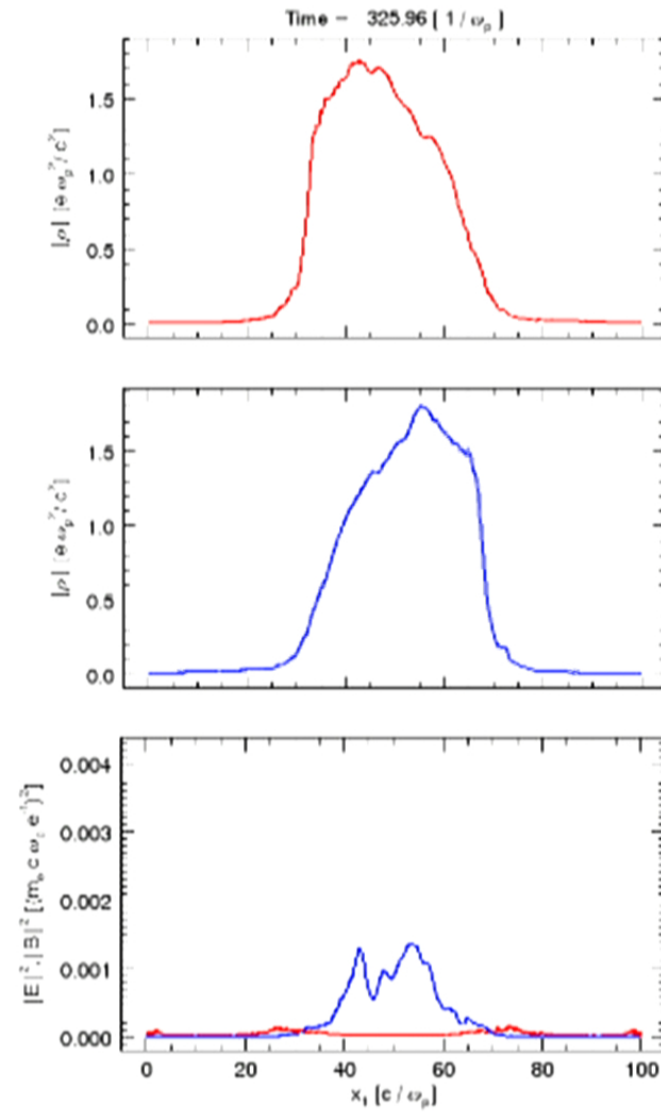


number density  
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longitudinal E field  
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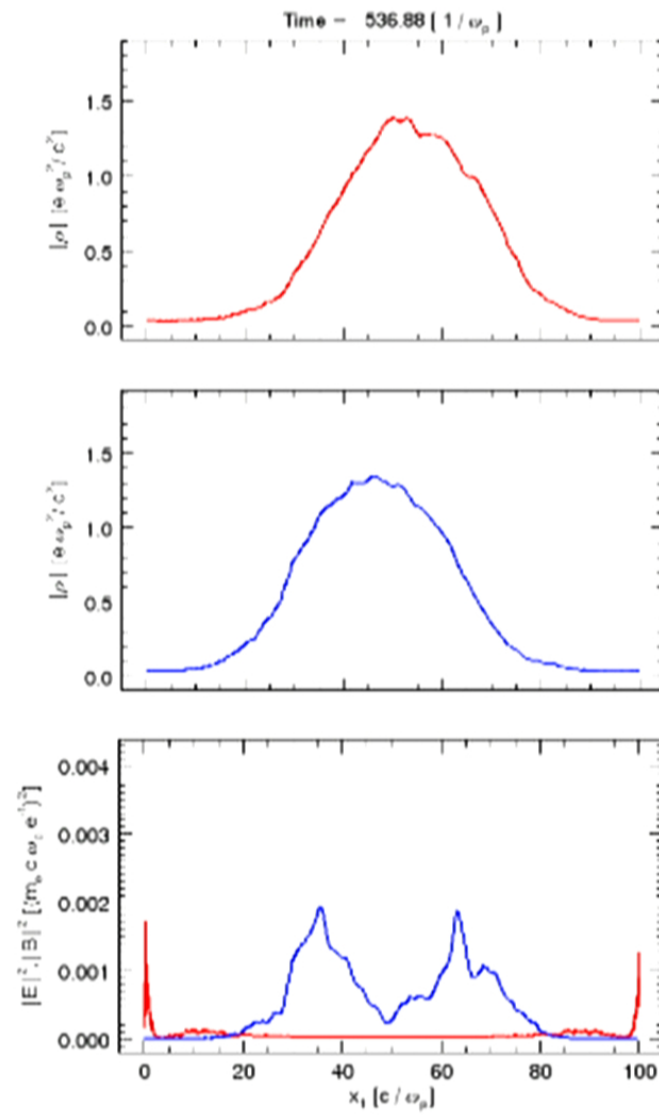


number density  
(right-movers)

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longitudinal E field  
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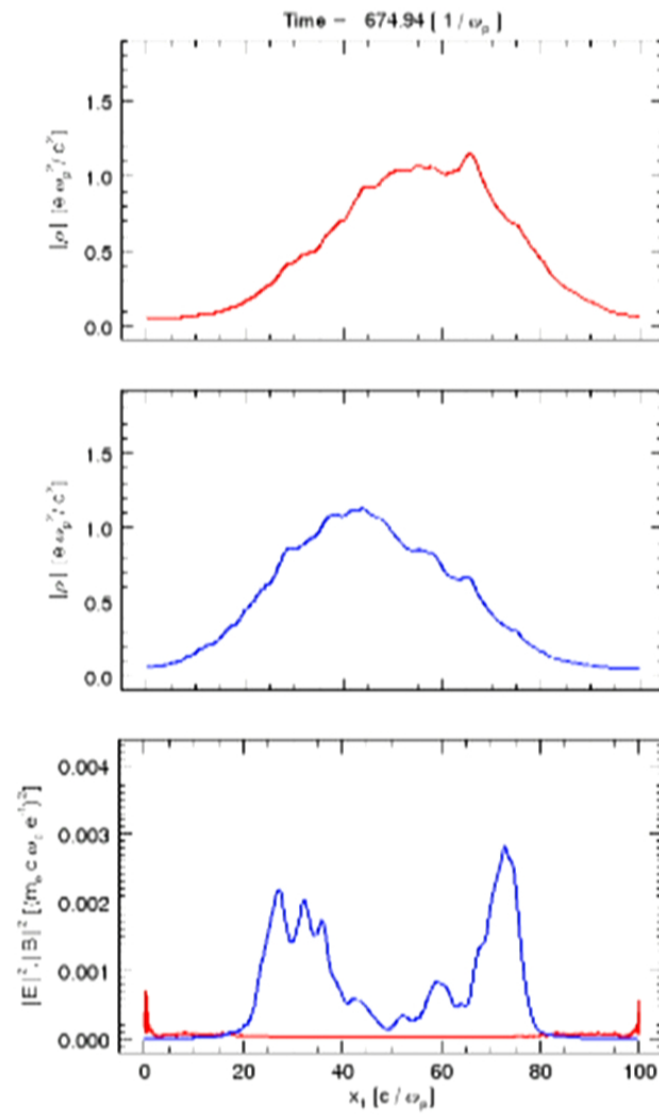


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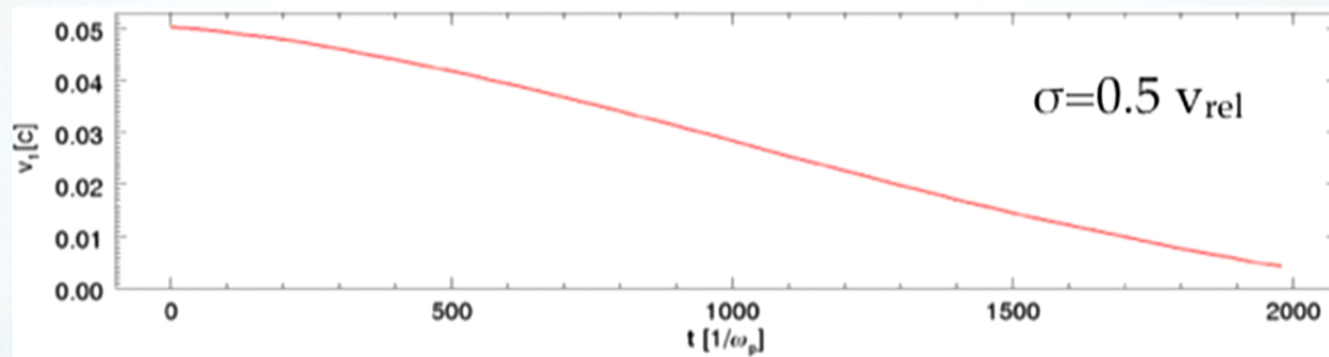
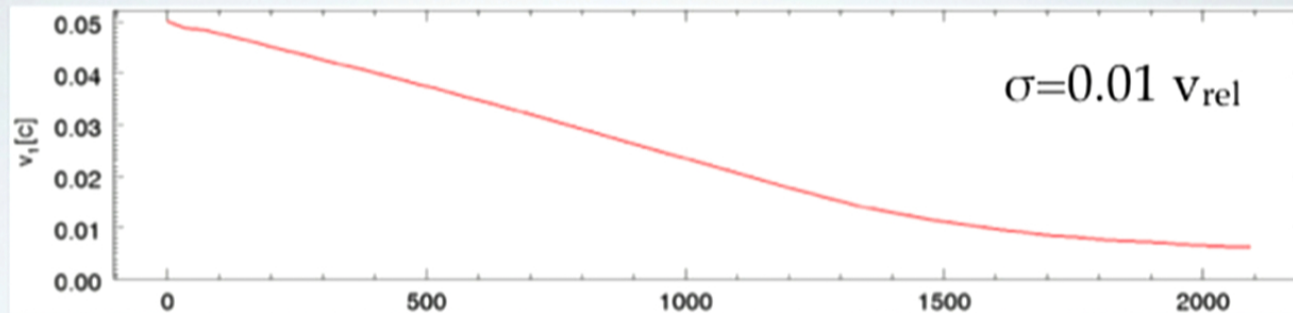
number density  
(left-movers)

longitudinal E field  
transverse B field

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# SLOWDOWN BY TRANSVERSE INST.



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Luis O. Silva

xxxxxx

(unbroken)

$$t_{\text{univ}} \times \sqrt{\quad} \lesssim \sim 100$$

$$\frac{1}{H_0} \approx \frac{M_{\text{pl}}}{\sqrt{\rho}}$$

$\left. \begin{array}{l} 0.3 \omega_{\text{pl}} \\ \omega_{\text{pl}} \end{array} \right\} \begin{array}{l} \text{long} \\ \text{trans} \end{array}$

long:

$$\omega_{\text{pl}} \lesssim 300 \frac{\sqrt{\rho} \sqrt{8\pi}}{M_{\text{pl}}}$$

$$\frac{e}{m} \lesssim 10^3 \frac{\sqrt{\rho}}{\sqrt{\rho}} \frac{1}{M_{\text{pl}}} \sim \frac{1}{M_{\text{pl}}}$$

# CONCLUSIONS

Instabilities are much more effective than collisions in at disrupting counter-streaming plasmas

**DM electromagnetism must be almost weak as gravity**

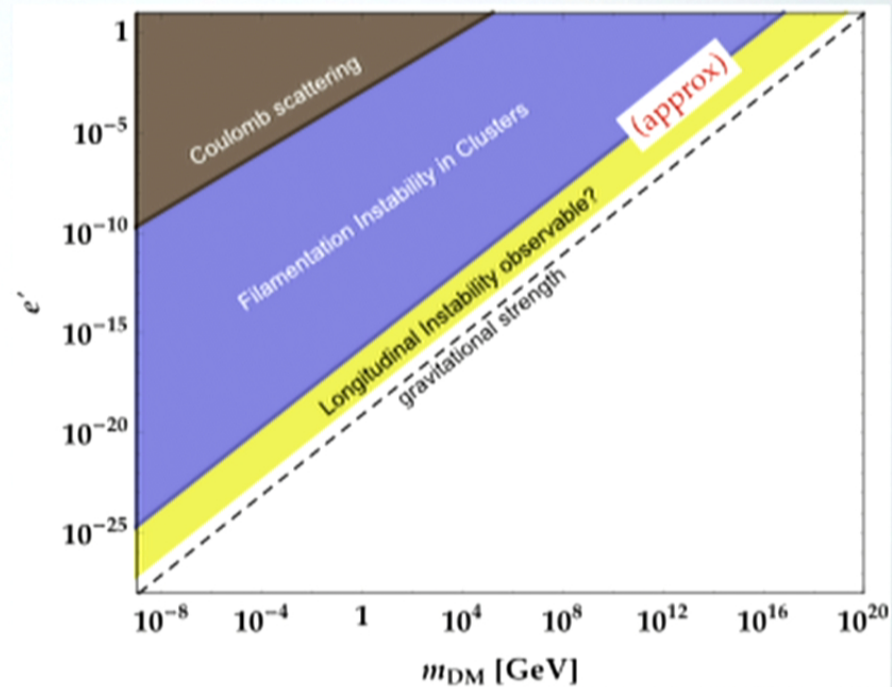
Where is the bound?

Observable signature?

Awaiting sim. results

Can we do better?

- better simulation using methods of T. Abel?
- better astro modeling?
- include gravity?



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