

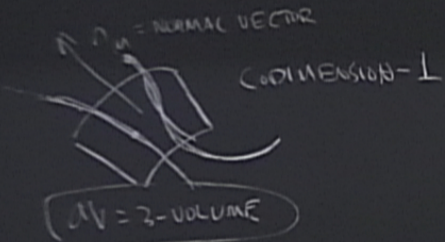
Title: PSI 2015/2016 Explorations in Cosmology - Kendrick Smith - 12

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Abstract:

GR



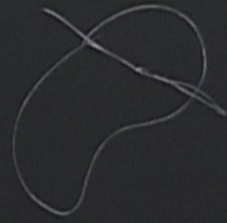
(SIGNED NUMBER OF WORLDELINE CROSSINGS)

$$= N^M n_M dV$$

LET  $X_i^M(\tau)$  BE A SET OF PARAMETRIZED WORLDELINES

$$N^M(x) = \sum_i \int d\tau \frac{1}{\sqrt{-g}} \delta^4(x - X_i^M(\tau)) \frac{dX_i^M}{d\tau}$$

$$\nabla_M N^M = 0$$



$$\int d^4x \sqrt{-g} \frac{1}{\sqrt{-g}} \delta^4(x - x_0) = 1$$

$$\frac{1}{\sqrt{-g}} \delta^4(x - x_0)$$

$$\oint_{\Sigma} n_M N^M = 0$$



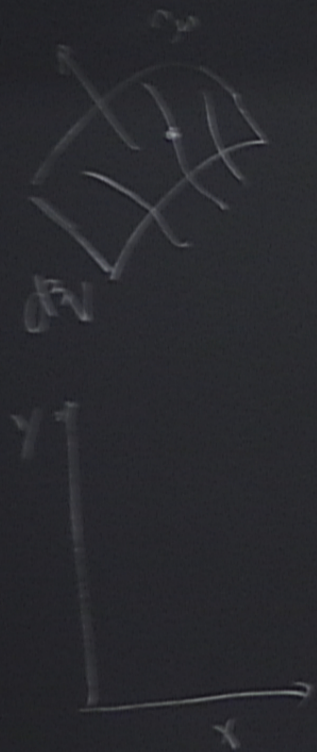
SR

ENERGY  $E \longleftrightarrow$  4-MOMENTUM  $p^\mu = (E, \vec{p})$

NUMBER DENSITY  $n \longleftrightarrow$  NUMBER CURRENT  $N^\mu = (n, n\vec{v})$

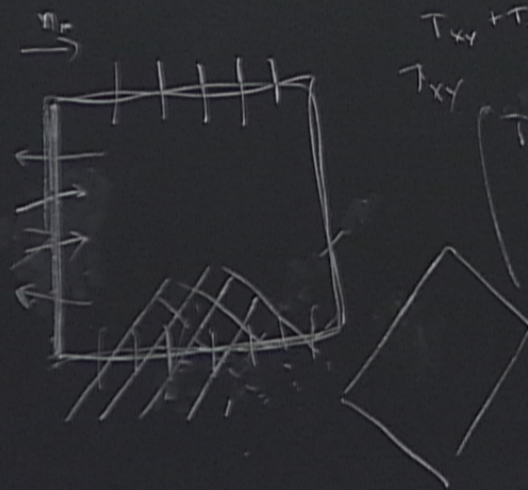
ENERGY DENSITY  $\rho \longleftrightarrow$  STRESS ENERGY  $T^{\mu\nu}$





4-MOMENTUM FLUX =  $\sum_{\text{WORLDLINES}} \dot{p}^\mu$   
 $F^\mu$

$T_{\mu\nu} = T^{\mu\nu} n_\nu$   
 $T_{xx} - T_{yy}$   
 $T_{xy} + T_{yx}$   
 $T_{xx}$   
 $T_{yy}$



$$\begin{pmatrix} T_{00} & T_{i0} \\ T_{0i} & \rho \dot{\phi}_i + \Pi_i \end{pmatrix}$$

$$\begin{pmatrix} \frac{1}{2} \dot{\phi}^2 \\ -V(\phi) \end{pmatrix}$$



COVARIANT FORMULA:  $X_i^{\mu}(\tau) = \text{PARAMETERIZED WORLDLINES } (i=1, \dots, N)$   
 PROPER TIME  $\frac{dx^{\mu}}{dt} = p^{\mu}$

$$\nabla_{\mu} T^{\mu\nu} = 0 \quad \Rightarrow \quad T^{\mu\nu} = 2 \int d\tau \frac{1}{\sqrt{g}} \delta^4(x - X_i^{\mu}(\tau)) \frac{dX_i^{\mu}}{d\tau} \frac{dX_i^{\nu}}{d\tau}$$

$$\frac{2}{\sqrt{g}} \frac{\delta S^m}{\delta g^{\mu\nu}} = T_{\mu\nu}$$

$$\nabla_{\mu} T^{\mu\nu} = 0 \quad \text{IF TRAJECTORIES ARE GEODESICS}$$



COVARIANT FORMULA:

$X_i^M(\tau) = \text{PARAMETERIZED WORLDLINES } i=1, \dots, N$   
↑ PROPER TIME  $\frac{dx^\mu}{dt} = p^\mu$

$$\nabla_\mu T^{\mu\nu} = 0$$

$$\Rightarrow T^{\mu\nu} = \sum_i \int d\tau \frac{1}{\sqrt{g}} \delta^4(x - X_i^\mu(\tau)) \frac{dx_i^\mu}{d\tau} \frac{dx_i^\nu}{d\tau}$$

$$\frac{2}{\sqrt{g}} \frac{\delta S^M}{\delta g^{\mu\nu}} = T_{\mu\nu}$$

$$\nabla_\mu T^{\mu\nu} = 0$$

IF TRAJECTORIES ARE GEODESICS

$$= \sum_i \int d\tau \frac{1}{\sqrt{g}} \delta^4(x - X_i^\mu(\tau)) \frac{D^2 X_i^\nu}{D\tau^2} = 0 \text{ FOR GEODESIC}$$



NEXT FEW LECTURES: MODEL UNIVERSE AS SUM OF 3

→ FLUIDS (BARYONS, CDM, PHOTONS) IN FIRST  
ORDER PERTURBATION THEORY

PERFECT

- AN APPROXIMATION WHICH NEGLECTS HIGHER MULTIPOLES
- + ANISOTROPIC STRESS
- NEUTRINOS OMITTED



• EACH SPECIES HAS A BACKGROUND DENSITY  $\bar{\rho}_J$  AND PRESSURE  $\bar{p}_J$  ( $J \in \{b, c, \gamma\}$ )  
 + FLUCTUATIONS  $\{\delta\rho^J, v_i^J, \delta p^J\}$

• FOR  $J \in \{b, c\}$  WE HAVE  $\delta\rho^J = 0$  (ORDER  $\frac{v^2}{c^2}$ )

• FOR PHOTONS,  $\nabla_{\mu\nu} T^{\mu\nu}_\gamma = 0$   
 ↑  
 MASSLESS

FOR PERFECT FLUID  $T^{\mu\nu}_\gamma = (\rho + p)U^\mu U^\nu + pg^{\mu\nu}$

WHERE  $U^\mu U_\mu = -1$  AND  $p = \bar{p} + \delta p$   
 $\rho = \bar{\rho} + \delta\rho$

$$0 = \nabla_{\mu\nu} T^{\mu\nu}_\gamma = -(\rho + p) + 4p = -\rho + 3p$$

$$= (\bar{\rho} + 3\bar{p}) + (-\delta\rho + 3\delta p) \Rightarrow$$

↳ 0

$$\delta\rho_J = \frac{1}{3}\delta p_J$$



• EACH SPECIES HAS A BACKGROUND DENSITY  $\bar{\rho}_J$  AND PRESSURE  $\bar{p}_J$  ( $J \in \{b, c, \gamma\}$ )  
 + FLUCTUATIONS  $\{\delta\rho^J, v_i^J, \delta p^J\}$

• FOR  $J \in \{b, c\}$  WE HAVE  $\delta\rho^J = 0$  [ORDER  $\frac{v^2}{c^2}$ ]

• FOR PHOTONS,  $g_{\mu\nu} T_{\gamma}^{\mu\nu} = 0$   
 ↑  
 MASSLESS

FOR PERFECT FLUID  $T_{\gamma}^{\mu\nu} = (\rho + p)U^{\mu}U^{\nu} + pg^{\mu\nu}$

WHERE  $U^{\mu}U_{\mu} = -1$  AND  $p = \bar{p} + \delta p$   
 $\rho = \bar{\rho} + \delta\rho$

$$0 = g_{\mu\nu} T_{\gamma}^{\mu\nu} = -(\rho + p) + 4p = -\rho + 3p$$

$$= \underbrace{(-\bar{\rho} + 3\bar{p})}_{\rightarrow 0} + (-\delta\rho + 3\delta p) \Rightarrow$$

$$\delta\rho_J = \frac{1}{3}\delta p_J$$

$$\bar{p}_J = \frac{1}{3}\bar{\rho}_J$$



NOT GENERALLY TRUE THAT  $\bar{p}_j = w_j \bar{p}_j \Rightarrow \delta p_j = w_j \delta p_j$

[COUNTEREXAMPLE: SINGLE FIELD INFLATION]  $\frac{\bar{p}}{\bar{p}} \neq \frac{\delta p}{\delta p}$

TOTAL STRESS ENERGY APPEARING IN EINSTEIN EQS

$$\bar{p} = \sum_j \bar{p}_j \quad \bar{p} = \sum_j \bar{p}_j$$

$$\delta p = \sum_j \delta p_j \quad \delta p = \sum_j \delta p_j$$

$$(\bar{p} + \bar{p})v = \sum_j (\bar{p}_j + \bar{p}_j)v_j$$

← CONSISTENT WITH

$$T_{\mu\nu} = \sum_j T_{\mu\nu}^j$$

AND  $T_{0i} = (\bar{p} + \bar{p})v_i$



→ 0

FINALLY WE DEFINE

$$\delta_J = \frac{\delta p_J}{\bar{p}_J} \quad \text{"FRACTIONAL OVERDENSITY"}$$

DYNAMICAL QUANTITIES:  $\{ \delta p_{Ji}, v_{Ji}^i, \text{METRIC} \}$



EQS MOTION FOR THE FLUIDS

$$\nabla_M T_J^{M\nu} = (\text{FORCE}) \quad \text{SCHEMATICALLY}$$

START WITH LHS:

$$\begin{aligned} \nabla_M T_J^{M0} &= \frac{\partial}{\partial t} (\delta p_J) + 3H (\delta p_J + \overset{w_J}{\delta p_J}) + (\bar{p}_J + \bar{p}_J) (a^{-1} \partial_i v_i^J + \dot{h}_{ii}) \\ &= \frac{\partial}{\partial t} (\bar{p}_J \delta_J) + 3H (1 + w_J) \bar{p}_J \delta_J + (1 + w_J) \bar{p}_J (a^{-1} \partial_i v_i^J + \dot{h}_{ii}) \\ &= \end{aligned}$$



MATICALLY

$$\bar{p}_j \propto \alpha^{-3(1+w_j)}$$

$$\Rightarrow \frac{\partial \bar{p}_j}{\partial \alpha} = -3(1+w_j) \bar{p}_j$$

$$v_i^j + h_{ii}$$

$$(\alpha^{-1} \partial_i v_i^j + h_{ii})$$



$$\nabla_{\mu} N^{\mu} = 0$$

EQS MOTION FOR THE FLUIDS

$$\nabla_{\mu} T^{\mu\nu} = (\text{FORCE}) \quad \text{SCHEMATICALLY}$$

START WITH LHS:

$$\begin{aligned} \nabla_{\mu} T^{\mu 0} &= \frac{\partial}{\partial t} (\delta\rho_J) + 3H(\delta\rho_J + \overset{''}{\delta\rho_J}) + (\bar{\rho}_J + \bar{p}_J)(a^{-1} \partial_i v_i^J + \dot{h}_{ii}) \\ &= \frac{\partial}{\partial t} (\bar{\rho}_J \delta_J) + 3H(1+w_J)\bar{\rho}_J \delta_J + (1+w_J)\bar{p}_J(a^{-1} \partial_i v_i^J + \dot{h}_{ii}) \\ &= -3(1+w_J)H\bar{\rho}_J \delta_J + \bar{\rho}_J \dot{\delta}_J + 3H(1+w_J)\bar{\rho}_J \delta_J + (1+w_J)\bar{p}_J(a^{-1} \partial_i v_i^J + \dot{h}_{ii}) \\ &= \bar{\rho}_J \left[ \dot{\delta}_J + (1+w_J)(a^{-1} \partial_i v_i^J + \dot{h}_{ii}) \right] \end{aligned}$$



$$= P_j \left( \delta_j + (1 + w_j) (a^{-1} \partial_i V_i + h_{ii}) \right)$$

SIMILARLY  $\nabla^M T_{Mi}^j = a(1+w_j) \bar{P}_j \left[ \left( \frac{\partial}{\partial t} + H(1-3w_j) \right) (V_i^j - B_i) + a^{-1} \partial_i A \right] + w_j \partial_i (\delta_j)$

BARYONS & PHOTONS INTERACT BY THOMSON SCATTERING



IN REST FRAME OF BARYONS WE HAVE

$$\frac{d\sigma}{d\Omega} = \frac{3}{8\pi} \sigma_T \left( \frac{1 + \cos^2 \theta}{2} \right)$$

CAN WE TRANSLATE THIS TO EXPRESSIONS FOR  $\nabla_M T_b^{M\nu}$ ,  $\nabla_M T_c^{M\nu}$ ?



ANALYZE IN NORMAL COORDINATES, IN REST FRAME OF  $e^-$

$$U_b^M = (1, 0, 0, 0) \quad g_{\mu\nu} = (-1, 1, 1, 1)$$

$$U_\gamma^M = (1, v_i^\gamma)$$

$n_e =$  FREE ELECTRON DENSITY

$P_\gamma =$  PHOTON ENERGY DENSITY

$V_\gamma^i =$  PHOTON BULK VELOCITY

$$\Pi_\gamma^i = \text{PHOTON MOMENTUM DENSITY} = \frac{4}{3} P_\gamma V_\gamma^i = (P_\gamma + P_\gamma) V_\gamma^i$$

$\uparrow$   
 $\frac{1}{3} P_\gamma$

$$T^{\mu\nu} = (p+p) U^\mu U^\nu + p g^{\mu\nu}$$

$$= \begin{pmatrix} p & (p|p|u) \\ (p|p|u) & p \delta_{ij} \end{pmatrix}$$

WHY  $\frac{4}{3}$ ?

$\omega_\gamma \delta_i(\delta_j)$