

Title: PSI 2015/2016 Quantum Gravity - Lecture 11

Date: Mar 07, 2016 10:15 AM

URL: <http://pirsa.org/16030049>

Abstract:

3rd week

lecture 11 Gauss constraints

lecture 12 Flatness constraints

lecture 13 Geometric operators
+ 4d LQG

lecture 14 Spinfoam models

lecture 15 Guest speakers!



→ fluxes act as derivatives -

• Inner product → Haar measure.

• ONB → matrix elts of irreducible rep of $SU(2)$

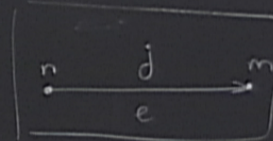
basis states $|j, m, n\rangle$ $\langle g | j, m, n \rangle = D_{mn}^j(g) \sqrt{d_j}$

Left and Right action

Right • $(R_h D_{mn}^j)(g) = D_{mn}^j(gh) = \sum_p D_{mp}^j(g) D_{pn}^j(h)$
 $\Rightarrow R_h |j, m, n\rangle = \sum_p D_{pn}^j(h) |j, m, p\rangle \rightarrow R_h$ acts as D^j on $\{|j, m, p\rangle / p = -j \dots j\}$

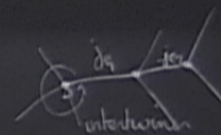
Left • $L_h |j, m, n\rangle = \sum_q D_{mq}^j(h^{-1}) |j, q, n\rangle \rightarrow$ action on the dual representation
 $(D_{mn}^j)^*(g) = D_{nm}^j(g^{-1})$

$\Rightarrow$ Hilbert space associated to one edge $\mathcal{H}_e = \mathcal{L}^2(G, d_h g) = \bigoplus_j (V_j \otimes V_j^*)$



= one edge of a spinnetwork

← Angular momentum associated to both ends of the edge



Left and Right action

$$\text{Right} \cdot (R_h D_{mn}^j)(g) = D_{mn}^j(gh) = \sum_p D_{mp}^j(g) D_{pn}^j(h)$$

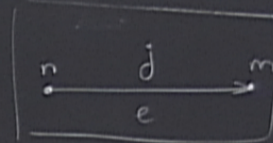
$$\Rightarrow R_h |jmn\rangle = \sum_p D_{pn}^j(h) |jmp\rangle \quad \rightarrow R_h \text{ acts as } D^j \text{ on } \{|j m, p\rangle / p = -j \dots j\}$$

$$\text{Left} \cdot L_h |jmn\rangle = \sum_q D_{mq}^j(h^{-1}) |jqn\rangle \quad \rightarrow \text{action on the dual representation}$$

$$(D_{mn}^j)^*(g) = D_{nm}^j(g^{-1})$$

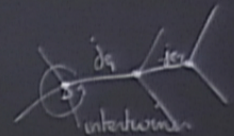
$\Rightarrow$ Hilbert space associated to one edge

$$\mathcal{H}_e = \mathcal{L}^2(G, d\mu(g)) = \bigoplus_j (V_j \otimes V_j^*)$$



= One edge of a spin network

Angular momentum associated to both ends of the edge



$\mathcal{H}_e$

$A \rightarrow h_e$: generalized connection: assignment of $h_e \in SU(2)$ to any path $e \subset \Sigma$
i.e. fundamental variable is the holonomy itself and not its relationship to a smooth connection

Cylindrical functions of generalized connections Cyl

$$Cyl = \bigcup_{\Gamma} Cyl_{\Gamma} \quad \Psi_{\Gamma, f} \in Cyl \quad f \text{ smooth function: } SU(2)^n \rightarrow \mathbb{C}$$

$$\Rightarrow \text{functional of the connection is defined as}$$

$$\Psi_{\Gamma, f}[A] = f(h_{e_1}[A], \dots, h_{e_n}[A])$$

$n = \# \text{ of edges of } \Gamma$

• Hilbert space associated to a graph = tensor product of one edge Hilbert spaces.

states $\Psi_T(g_e \dots g_n) = \langle g_e \dots g_n | \Psi_T \rangle$

• Inner product on the tensor product: $\mathcal{H}_T = \bigotimes_e \mathcal{H}_e$ is given by $\langle \Psi_{T,1} | \Psi_{T,2} \rangle = \int_{G^n} \overline{\Psi_{T,1}(g_e \dots g_n)} \Psi_{T,2}(g_e \dots g_n) dg_e \dots dg_n$

• ONB for $\mathcal{H}_T$ by tensoring the basis functions $|j, m, n\rangle$ for each edge

$$|\{\vec{d}_1, \vec{m}_1, \vec{n}_1\}_T\rangle = |j_1, m_1, n_1\rangle \otimes \dots \otimes |j_n, m_n, n_n\rangle$$

Ku



13-week

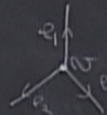
- Lecture 1: Graphs
- Lecture 2: Flows
- Lecture 3: Combinatorics
- Lecture 4: Spectra
- Lecture 5: Graphs

2 edges $\xrightarrow{g_1, g_2}$

$\Psi(g_1, g_2) \rightarrow$ invariant function the type $\Psi(g_1, g_2)$

Generalization
 $S(\hat{\mathcal{G}}_n) =$

Let us a 3 valent vertex



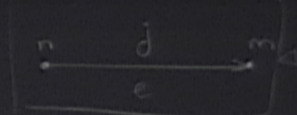
$$S(\hat{\mathcal{G}}_n) \cdot \left(\bigotimes_{i=1,2,3} |d_i, m_i, n_i\rangle \right) = \int dh \, R_h^{e_1} R_h^{e_2} R_h^{e_3} \left(\bigotimes_{i=1}^3 |j_i, m_i, n_i\rangle \right)$$

$$= \int dh \sum_{p_1, p_2, p_3} \underbrace{D_{p_1 m_1}^{j_1}(h) D_{p_2 m_2}^{j_2}(h) D_{p_3 m_3}^{j_3}(h)}_{\text{Hilbert space associated to one edge}} \left(\bigotimes_{i=1}^3 |d_i, m_i, n_i\rangle \right)$$

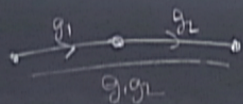
$\Rightarrow$ Hilbert space associated to one edge

$$\left(\mathcal{H}_{n,m}^{j_1} = \mathcal{H}_{n,m}^{j_2} \right)$$

$$= \bigoplus_j (V_j \otimes V_j^*)$$

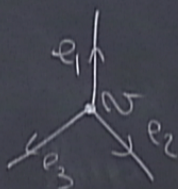


2 edges



$\Psi(g_1, g_2) \rightarrow$ invariant function the type $\Psi(\underbrace{g_1, g_2}_{h, \bar{h}})$

Let us a 3 valent vertex



$$\begin{aligned}
 S(\hat{\mathcal{G}}_{\mathcal{N}} \cdot (\otimes_{i=1,2,3} |j_i, m_i, n_i\rangle)) &= \int d\mathbf{h} R_h^{e_1} R_h^{e_2} R_h^{e_3} \left(\otimes_{i=1}^3 |j_i, m_i, n_i\rangle \right) \\
 &= \int d\mathbf{h} \sum_{p_1, p_2, p_3} D_{p_1 n_1}^{j_1}(\mathbf{h}) D_{p_2 n_2}^{j_2}(\mathbf{h}) D_{p_3 n_3}^{j_3}(\mathbf{h}) (\otimes |j_i, m_i, p_i\rangle)
 \end{aligned}$$

(L)

(m_i, p_i)

Generalization of $\otimes$

$$S(\hat{\mathcal{G}}_N) = \int_G$$

$\overset{\text{in}}{T}$

$\overset{\text{out}}{T}$

$$\left(\bigotimes_{e/N=N_+(e)} L_h^e \right) \otimes \left(\bigotimes_{e'/N=N_-(e)} R_h^{e'} \right) (dh)$$

State $\langle h_e | \Psi_e \rangle = \Psi_e(h_e)$ so state transforms as $\Psi(g(h_s) h_e g^{-1}(h_t)) = \langle h_e | L_{g(h_s)} R_{g^{-1}(h_t)} | \Psi_e \rangle$

$\overset{\text{target}}{h_t}$ $\overset{\text{source}}{h_s}$

→ Gauss constraint: a projector on the subspace of states satisfying the Gauss constraints by averaging over the group $SU(2)$ for both ends of the edge.
(Group averaging)

state based on one edge

$$\delta(\hat{G}) |\Psi_e\rangle = \int L_{g_1} R_{g_2} |\Psi_e\rangle dg_1 dg_2 \stackrel{(*)}{=} \text{const}$$

$$\int_{G \times G} L_{g_1} R_{g_2} \Psi(h_e) dg_1 dg_2 = \int_{G \times G} \Psi(\underbrace{g_1^{-1} h_e g_2}_{g'_e}) dg_1 dg_2 = \int_{G \times G} \Psi(g'_e) dg_1 dg_2 = \text{const}$$

Hilbert space associated to a graph = tensor product of one edge Hilbert spaces.

states $\Psi_r(g_{e_1} \dots g_{e_n}) = \langle g_{e_1} \dots g_{e_n} | \Psi_r \rangle$

Inner product on the tensor product $\mathcal{H}_r = \bigotimes_e \mathcal{H}_e$ is given by $\langle \Psi_{r,1} | \Psi_{r,2} \rangle = \int_{G^n} \overline{\Psi_{r,1}(g_1 \dots g_n)} \Psi_{r,2}(g_1 \dots g_n) dg_1 \dots dg_n$

ONB for $\mathcal{H}_r$ by tensoring the basis functions $|j, m, n\rangle$ for each edge

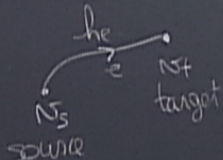
$$|\{\vec{j}, \vec{m}, \vec{n}\}_r\rangle = |j_1, m_1, n_1\rangle \otimes \dots \otimes |j_n, m_n, n_n\rangle$$

Kinematical Hilbert space $\mathcal{H}_r = \bigotimes_e \mathcal{H}_e = \bigotimes_e \mathcal{L}^2(G) \cong \mathcal{L}^2(G^n)$

$(G \sim \partial E + \epsilon A E)$

action $\rightarrow$ generates internal rotation

On a holonomy h_e



$$h_e \rightarrow g(v_s) h_e g^{-1}(v_t)$$

State $\langle h_e | \Psi_e \rangle = \Psi_e(h_e)$ so state transforms as $\Psi(g(v_s) h_e g^{-1}(v_t)) = \langle h_e | L_{g(v_s)} R_{g^{-1}(v_t)} | \Psi_e \rangle$

projector on the subspace of states satisfying the Gauss constraints by averaging over the group $SU(2)$

$$= \frac{(-1)^{j_1 - j_2 - m_3}}{\sqrt{2j_3 + 1}} \langle j_1, m_1, j_2, m_2 | j_3, -m_3 \rangle$$

$$S(\hat{G}) \cdot S(\hat{G}) = S(G)$$

$$S(\hat{G}) \left(\otimes |j_1, m_1, n_1\rangle \right) = \sum_{p_1, p_2, p_3} \begin{pmatrix} j_1 & j_2 & j_3 \\ p_1 & p_2 & p_3 \end{pmatrix} \otimes |j_1, m_1, p_1\rangle$$

$$\sum_{m_1, m_2, m_3} \begin{pmatrix} j_1 & j_2 & j_3 \\ m_1 & m_2 & m_3 \end{pmatrix}^2 = 1$$

$\Rightarrow$ Normalized invariant state

$$\left(\sum_{p_1, p_2, p_3} \begin{pmatrix} j_1 & j_2 & j_3 \\ p_1 & p_2 & p_3 \end{pmatrix} |j_1, m_1, p_1\rangle \right)$$