

Title: PSI 2015/2016 Quantum Gravity - Lecture 8

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URL: <http://pirsa.org/16030046>

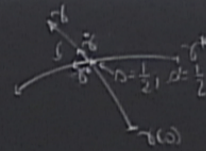
Abstract:

Basic variables

holonomy $h_\gamma[A]$

$$\text{flux } E_\alpha = \int_{\sigma^*} h_{\alpha\sigma^*} e_a(\gamma^*(s)) \dot{\gamma}^a(s) (h_{\sigma\sigma^*})^{-1} ds$$

↳ // transport from $\sigma^*(s)$ to $\sigma(s)$.



tangent vectors at the intersection positively oriented
 $\Rightarrow \det(\dot{\gamma}^* \dot{\gamma}^a) = \det \begin{pmatrix} \dot{\gamma}^1 & \dot{\gamma}^{a1} \\ \dot{\gamma}^2 & \dot{\gamma}^{a2} \end{pmatrix} > 0$

Basic variables

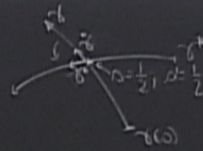
→ holonomy $h_\gamma[A]$

→ flux $E_\gamma = \int_{\gamma^*} h_{\gamma^*} e_a(\gamma^*(s)) \dot{\gamma}^*(s) (h_{\gamma^*})^{-1} ds$

↳ sits at $\gamma(s)$ ↳ transport from $\gamma^*(s)$ to $\gamma(s)$.

* $E_\gamma \rightarrow g(\gamma(s)) E_0 g(\gamma(s))^{-1}$

* $h_{\gamma^{-1}} = (h_\gamma)^{-1}$; $E_{\gamma^{-1}} = -$



tangent vectors at the intersection positively oriented $\Rightarrow \det(\dot{\gamma}^* \dot{\gamma}^*) = \det \begin{pmatrix} \dot{\gamma}^1 & \dot{\gamma}^{*2} \\ \dot{\gamma}^2 & \dot{\gamma}^{*1} \end{pmatrix} > 0$

Poisson algebra

$T_j = -\frac{i}{2} \sigma_j$ metric on the Lie algebra $\text{tr}(T_i T_j) = -\frac{1}{2} \delta_{ij}$

\Rightarrow extract component: $E_j = -2 \text{tr} \left(\underbrace{E}_{E_k T^k} T_j \right)$

$e_a = \tilde{E}_{ba} E^{bj} T_j$

so we can define the flux component.

$(E_b)_j = \int_{\mathcal{D}^4} -2 \text{tr} \left(T_j \underbrace{h_{\alpha\beta} T_k h_{\alpha\beta}^{-1}}_{(h_{\alpha\beta} T^k)} \right) \tilde{E}_{ba} E^{bk} \delta^{x^a} d\sigma$



Poisson algebra

$T_j = -\frac{i}{2} \sigma_j$ metric on the Lie algebra $\text{tr}(T_i T_j) = -\frac{1}{2} \delta_{ij}$

\Rightarrow extract component: $E_j = -2 \text{tr}(E T_j)$

$e_a = \tilde{E}_{ba} E^{bj} T_j$

so we can define the flux component.

$$(E_b)_j = \int_{\mathcal{S}^2} -2 \text{tr} \left(T_j h_{rs}(\omega) T_k h_{rs}^{-1}(\omega) \right) \tilde{E}_{ba} E^{bk}(\omega) \dot{\gamma}^{ra} d\omega = \int_{\mathcal{S}^2} R_{jk} E^{bk}(\omega) \dot{\gamma}^a \tilde{E}_{ba} d\omega$$

$(h_{rs}(\omega) E^b h_{rs}^{-1})^d = R_{jk} E^{bk}$

on the Lie algebra $\text{tr}(T_i T_j) = -\frac{1}{2} \delta_{ij}$

component: $E_j = -2 \text{tr} \left(\frac{E T_j}{E_k T^k} \right)$

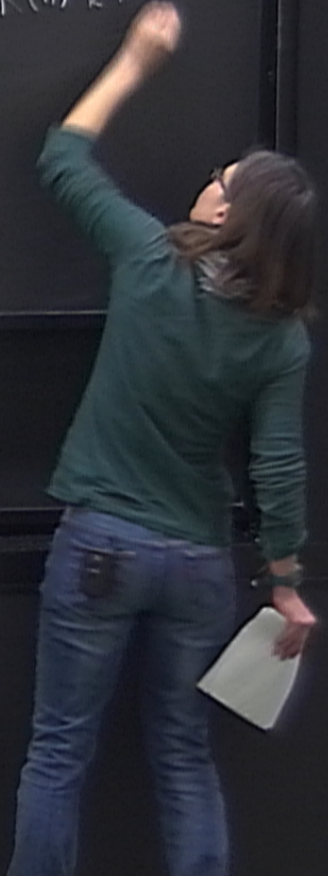
the flux component.

$$= \int_{\mathcal{D}^*} -2 \text{tr} \left(T_j h_{\mathcal{D}^*} \omega T_k h_{\mathcal{D}^*}^{-1} \right) \tilde{E}_{ba} E^{bk}(\mathcal{D}^*) \delta^{*a} d\mathcal{D}^* = \int_{\mathcal{D}^*} R_{jk} E^{bk}(\mathcal{D}^*) \delta^a \tilde{E}_{ba} d\mathcal{D}^*.$$

$(h_{\mathcal{D}^*} \omega h_{\mathcal{D}^*}^{-1})^d = R_{jk}^{(h)} E^{bk}$

$$e_a = \tilde{E}_{ba} E^{bj} T_j$$

$$h \cdot T_j h^{-1} = R(h)^j_k T^k$$



on the Lie algebra $\text{tr}(T_i T_j) = -\frac{1}{2} \delta_{ij}$

component: $E_j = -2 \text{tr} \left(\frac{E}{E^k T^k} T_j \right)$

the flux component.

$$= \int_{\mathcal{D}^4} -2 \text{tr} \left(T_j h_{\alpha\beta}(\omega) T_k h^{\alpha\beta}(\omega) \right) \tilde{E}_{ba} E^{bk}(\mathcal{D}^4(\omega)) \delta^{*a} d\omega = \int_{\mathcal{D}^4} R_{jk} E^{bk}(\mathcal{D}^4(\omega)) \delta^a \tilde{E}_{ba} d\omega.$$

$(h_{\alpha\beta}(\omega) E^b h^{\alpha\beta}(\omega))^\delta = R_{jk}^{(h)} E^{bk}$

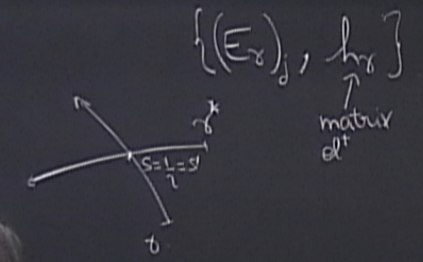
$$e_a = \tilde{E}_{ba} E^{bj} T_j$$

$$h \cdot T^j h^{-1} = R^{(h)j}_k T^k$$

$E_j = -\alpha \partial_k (E^k T_j)$ $\partial_a = \partial_{ba} L$
 so we can define the flux component.

$$(E_\sigma)_j = \int_{\sigma^*} -2 \text{tr} \left(T_j h_{\sigma\sigma^*} T_k h_{\sigma\sigma^*}^{-1} \right) \tilde{E}_{ba} E^{bk}(\sigma^*) \dot{\sigma}^{xa} d\sigma = \int_{\sigma^*} R_{jk} E^{bl}(\sigma^*) \dot{\sigma}^a \tilde{E}_{ba} d\sigma$$

$$(h_{\sigma\sigma^*}) E^b h_{\sigma\sigma^*}^{-1} = R_{jk}^{(h)} E^{bk}$$

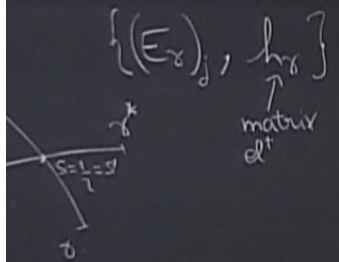


\rightarrow only the part near the intersection of σ and σ^*
 so splitting $h_\sigma = h_\sigma \left(\frac{1+\epsilon}{2}, \frac{1-\epsilon}{2} \right) h_\sigma \left(\frac{1+\epsilon}{2}, \frac{1-\epsilon}{2} \right) h_\sigma \left(\frac{1-\epsilon}{2}, 0 \right)$

so we can define the flux component.

$$(E_\sigma)_j = \int_{\sigma^*} -2 \text{tr} \left(T_j h_{\sigma\sigma^*} T_k h_{\sigma\sigma^*}^{-1} \right) \tilde{E}_{ba} E^{bk}(\sigma^*(\omega)) \dot{\gamma}^a d\omega = \int_{\sigma^*} R_{jk} E^{bl}(\sigma^*(\omega)) \dot{\gamma}^a \tilde{E}_{ba} d\omega.$$

$$(h_{\sigma\sigma^*}(\omega) E^b h_{\sigma\sigma^*}^{-1})^\delta = R_{jk}^{(h)} E^{bk}$$



→ only the part near the intersection of σ and σ^*

so splitting $h_\sigma = h_\sigma \left(\frac{1}{2}, \frac{1}{2}, \frac{\epsilon}{2} \right) h_\sigma \left(\frac{1}{2}, \frac{1}{2}, \frac{\epsilon}{2} \right) h_\sigma \left(\frac{1}{2}, \frac{1}{2}, 0 \right)$

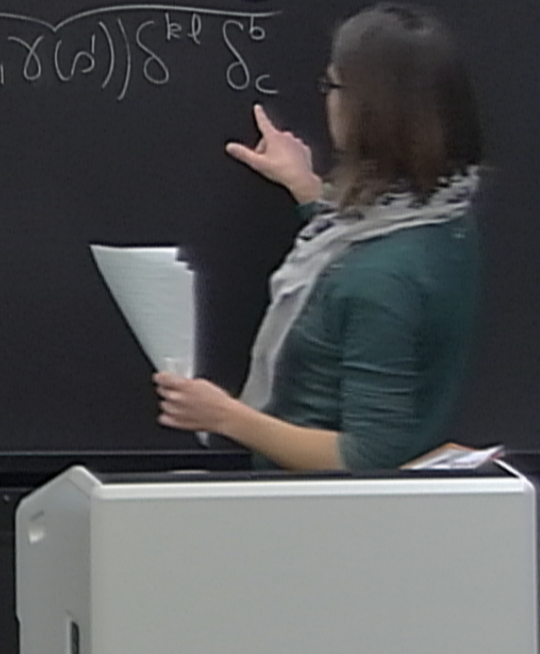
commute with the flux approximation

$$= \int_{\frac{1}{2}-\frac{\epsilon}{2}}^{\frac{1}{2}+\frac{\epsilon}{2}} \frac{1}{T_\ell} \dot{\gamma}^c(\sigma') A_c^l(\sigma(s')) d\omega' + O(\epsilon^2)$$

$$\{ (E_{\alpha_j}, h_{\alpha}) = - h_{\alpha} (1, \frac{1+\epsilon}{2}) h_{\alpha} (\frac{1-\epsilon}{2}, 0) \int_{\delta \frac{1-\epsilon}{2}}^{\frac{1+\epsilon}{2}} R_{jk}(\rho) \bar{E}_{ba} \dot{\gamma}^a(\rho) \dot{\gamma}^c(\rho) \{ E^{bk}(\dot{\gamma}(\rho)) \} ,$$

$$= -h_{\gamma(1, \frac{1+\epsilon}{2})} h_{\gamma(\frac{1-\epsilon}{2}, 0)} \int_{\delta^* \frac{1-\epsilon}{2}}^{\frac{1+\epsilon}{2}} R_{jk}(\rho) \bar{E}_{ba} \dot{\gamma}^a(\rho) \dot{\gamma}^c(\rho) \left\{ E^{bk}(\gamma^*(\rho)), A_c(\gamma(\rho')) \right\} d\rho d\rho' + O(\epsilon^2)$$

① $\overbrace{\delta(\gamma^*(\rho), \gamma(\rho'))} \delta^{kl} \delta_c^b$



$$= -h_{\gamma(1, \frac{1+\epsilon}{2})} h_{\gamma(\frac{1-\epsilon}{2}, 0)} \int_{\frac{1-\epsilon}{2}}^{\frac{1+\epsilon}{2}} R_{jk}(\rho) \bar{\epsilon}_{ba} \ddot{\gamma}^{*a}(\omega) \ddot{\gamma}^c(\omega') \left\{ E^{bk}(\ddot{\gamma}^*(\omega)), A_c(\ddot{\gamma}(\omega')) \right\} d\omega d\omega' + O(\epsilon^2)$$

$$\textcircled{1} \overbrace{\delta(\ddot{\gamma}^*(\omega), \ddot{\gamma}(\omega'))} \delta^{kl} \delta_c^b$$

$$\textcircled{2} \bar{\epsilon}_{ca} \ddot{\gamma}^{*a}(\omega) \ddot{\gamma}^c(\omega') = \det \begin{pmatrix} \ddot{\gamma}^1(\omega') & \ddot{\gamma}^2(\omega') \\ \ddot{\gamma}^{*1}(\omega) & \ddot{\gamma}^{*2}(\omega) \end{pmatrix}$$



$$\textcircled{4} \lim_{\varepsilon \rightarrow 0}$$

$$\textcircled{5} R_{jk}(\frac{1}{2}) T^k = h_{\sigma}(\frac{0, \frac{1}{2}})^{-1} \cdot T_j \cdot h_{\sigma}(\frac{0, \frac{1}{2}})$$

$$\Rightarrow \left\{ (E_{\sigma})_j, h_{\sigma} \right\} = h_{\sigma} T_j \rightarrow \text{hold matrix component with } \left\{ (h_{\sigma})_{MN} \right\} =$$

$\lim_{\epsilon \rightarrow 0}$

$$R_{j\frac{1}{2}}(\frac{1}{2}) T^k = h_{\sigma}(0, \frac{1}{2})^{-1} \cdot T_j \cdot h_{\sigma}(0, \frac{1}{2})$$

$$\boxed{\{(E_{\sigma})_j, h_{\sigma}\} = h_{\sigma} T_j} \rightarrow \text{hold matrix component wise} \quad \{(E_{\sigma})_j, (h_{\sigma})_{MN}\} = (h_{\sigma} T_j)_{MN}$$

$$\Rightarrow \left\{ (h_{\alpha})_{MN}, (h_{\alpha'})_{M'N'} \right\} = 0$$

$\Rightarrow \{ , \}$ has to satisfy the Jacobi identity.

$$\left\{ \left\{ E_{\alpha}^j, E_{\alpha}^k \right\}, h_{\alpha} \right\} = \left\{ \left\{ E_{\alpha}^j, h_{\alpha} \right\}, E_{\alpha}^k \right\} - \left\{ \left\{ E_{\alpha}^k, h_{\alpha} \right\}, E_{\alpha}^j \right\}$$

$$\Rightarrow \left\{ (h_{\alpha})_{MN}, (h_{\alpha'})_{M'N'} \right\} = 0$$

$\Rightarrow \{, \}$ has to satisfy the Jacobi identity.

$$\begin{aligned} \left\{ \left\{ E_{\alpha}^i, E_{\alpha}^k \right\}, h_{\alpha} \right\} &= \left\{ \left\{ E_{\alpha}^i, h_{\alpha} \right\}, E_{\alpha}^k \right\} - \left\{ \left\{ E_{\alpha}^k, h_{\alpha} \right\}, E_{\alpha}^i \right\} \\ &= -h_{\alpha} T^k T^i + h_{\alpha} T^i T^k \\ &= -h_{\alpha} \varepsilon^{ki} T^l \end{aligned}$$



$$\Rightarrow \{ (h_{\alpha})_{MN}, (h_{\alpha'})_{M'N'} \} = 0$$

$\Rightarrow \{ , \}$ has to satisfy the Jacobi identity.

$$\begin{aligned} \{ \{ E_{\alpha}^i, E_{\alpha}^k \}, h_{\alpha} \} &= \{ \{ E_{\alpha}^j, h_{\alpha} \}, E_{\alpha}^k \} - \{ \{ E_{\alpha}^k, h_{\alpha} \}, E_{\alpha}^j \} \\ &= -h_{\alpha} T^k T^j + h_{\alpha} T^j T^k \\ &= -h_{\alpha} \varepsilon^{kj} T^l \end{aligned} \rightarrow \{ E$$

$$\left\{ \left[E_{\gamma}^k, h_{\alpha} \right], E_{\gamma}^j \right\} \rightarrow \left\{ E_{\gamma}^j, E_{\gamma}^k \right\} = \varepsilon^{jk} E_{\gamma}^l$$

Closed Algebra

$$\left\{ (h_{\alpha})_{MN}, (h_{\alpha})_{M'N'} \right\} = 0$$

$$\left\{ E_{\gamma}^j, h_{\alpha} \right\} = h_{\alpha} T^j$$

$$\left\{ E_{\gamma}^j, E_{\gamma}^k \right\} = \varepsilon^{jk} E_{\gamma}^l$$

→ length

$$l_{e^*} = |E_e|$$

dihedral angles $\phi_{ee'}$ between (e^*, e'^*) of a triangle.

is contained in the dot product $E_e, E_{e'}$

$$\cos \phi_{ee'} = - \epsilon_{ee'} \frac{E_e \cdot E_{e'}}{|E_e| |E_{e'}|}$$

$\epsilon_{ee'} = +1$ if e, e'