

Title: The Unnatural (or Split) Composite Higgs

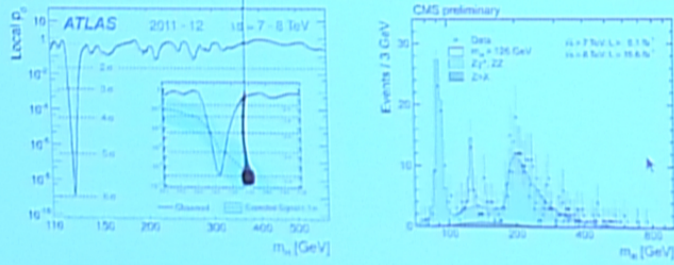
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URL: <http://pirsa.org/16020052>

Abstract:

A simple way to trivially satisfy precision-electroweak and flavor constraints in composite Higgs models is to require a large global symmetry breaking scale, $f > 10$ TeV. This leads to a tuning of order 10^{-4} to obtain the observed Higgs mass, but gives rise to a 'split' spectrum where the strong-sector resonances with masses greater than 10 TeV are separated from the pseudo Nambu-Goldstone bosons, which remain near the electroweak scale. To preserve gauge-coupling unification (due to a composite top quark), the symmetry breaking scale satisfies an upper bound $f < 100$ - 1000 TeV, which implies that the resonances are not arbitrarily heavy and may be accessible at future colliders. Furthermore, by identifying dark matter with a pseudo Nambu-Goldstone boson, the smallest coset space containing a stable, scalar singlet and an unbroken $SU(5)$ symmetry is $SU(7) / SU(6) \times U(1)$. Interestingly, this coset space also contains a metastable color-triplet pseudo Nambu-Goldstone boson that can decay via a displaced vertex when produced at colliders, leading to a distinctive signal of unnaturalness.

Higgs discovery - LHC Run I



Higgs potential: $V(h) = -\mu_h^2 |H|^2 + \lambda_h |H|^4$ $(H) = \frac{1}{\sqrt{2}}(v+h)$

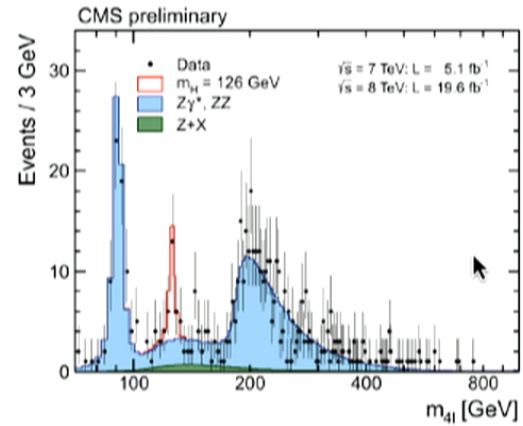
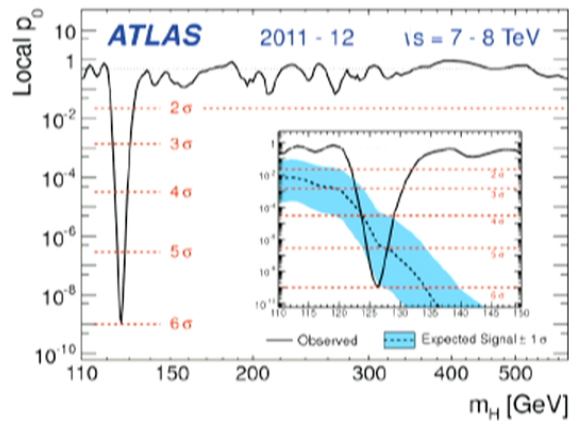
$v^2 = \frac{\mu_h^2}{\lambda_h} \simeq (246 \text{ GeV})^2$ $m_h^2 = 2\lambda_h v^2 \simeq (126 \text{ GeV})^2$

→ $\mu_h^2 \simeq (89 \text{ GeV})^2$ $\lambda_h \simeq 0.13$

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Higgs discovery - LHC Run I



Higgs potential: $V(h) = -\mu_h^2 |H|^2 + \lambda_h |H|^4$ $\langle H \rangle = \frac{1}{\sqrt{2}}(v + h)$

$$v^2 = \frac{\mu_h^2}{\lambda_h} \simeq (246 \text{ GeV})^2$$

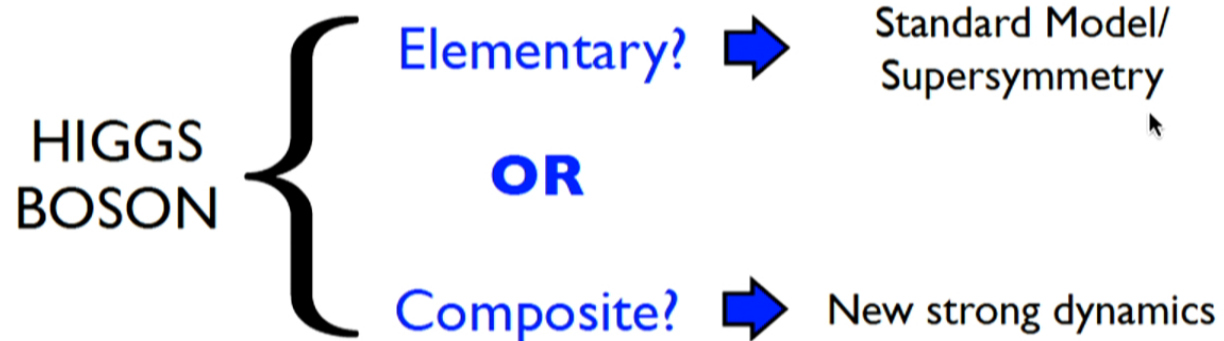
$$m_h^2 = 2\lambda_h v^2 \simeq (126 \text{ GeV})^2$$



$$\mu_h^2 \simeq (89 \text{ GeV})^2$$

$$\lambda_h \simeq 0.13$$

What is the nature of the Higgs boson?

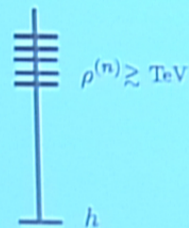


Understanding why $m_h \ll M_p$ can help
address shortcomings in the SM

Composite Higgs

Higgs as a pseudo Nambu-Goldstone boson (Georgi-Kaplan '84)

Global symmetry G spontaneously broken to subgroup H at scale f



Resonance mass: $m_{\rho} \sim g_{\rho} f \quad 1 \lesssim g_{\rho}^2 \lesssim 4\pi$

coset $G/H \ni h$

Higgs mass protected by G/H symmetry
-- like pions in QCD

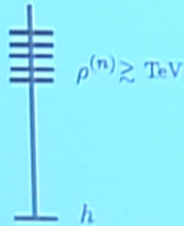
BUT global symmetry must be explicitly broken to generate $V(h) \neq 0$

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Composite Higgs

Higgs as a pseudo Nambu-Goldstone boson (Georgi-Kaplan '84)

Global symmetry G spontaneously broken to subgroup H at scale f



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coset $G/H \supset h$

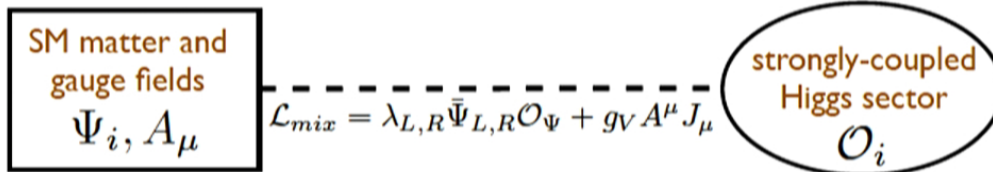
Higgs mass protected by shift symmetry
— like pions in QCD

BUT global symmetry must be explicitly broken to generate $V(h) \neq 0$

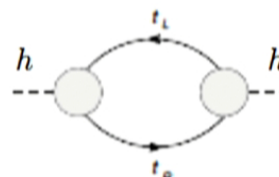
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Global symmetry broken by mixing with elementary sector

[Contino, Nomura, Pomarol '03; Agashe, Contino, Pomarol '04]



Higgs potential:

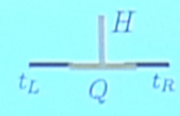


$$V(h) = -\mu_h^2 |H|^2 + \lambda_h |H|^4$$

where $\mu_h^2 \sim \frac{g_{SM}^2}{16\pi^2} g_\rho^2 f^2$ $\lambda_h \sim \frac{g_{SM}^2}{16\pi^2} g_\rho^2$

EWSB ($\langle H \rangle = \frac{v}{\sqrt{2}}$) $v^2 = \frac{\mu_h^2}{\lambda_h}$ \rightarrow Prefers: $f \sim v$

Higgs mass: $m_h^2 \approx \frac{N_c}{\pi^2} m_t^2 \left(\frac{m_Q^2}{f^2} \right) = g_Q^2$



m_Q = fermion resonance mass

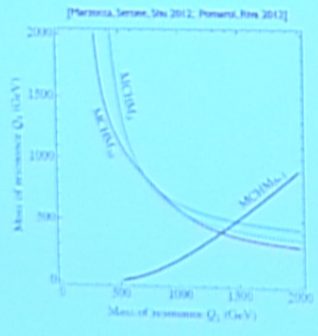
$m_Q \sim m_p \gtrsim 2.5 \text{ TeV} \quad (g_Q \sim g_p \gtrsim 3) \quad \Rightarrow \quad m_h \gtrsim m_t$

But, no need for $m_Q \sim m_p$

$m_h \sim 125 \text{ GeV}$

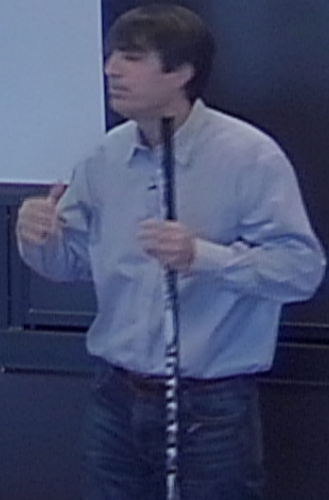
$\Rightarrow m_Q < m_p$

light fermion resonances!



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HOWEVER, precision electroweak, flavor constraints

$$\text{EWPT: } \frac{s}{16\pi^2 v^2} H^\dagger \tau^a H B^{\mu\nu} W_{a\mu\nu} \quad S = \frac{s}{2\pi} \sim \frac{m_W^2}{m_\rho^2} \Rightarrow f \gtrsim \frac{2.5 \text{ TeV}}{g_\rho}$$

$$\frac{-t}{16\pi^2 v^2} ((D^\mu H)^\dagger H) (H^\dagger D_\mu H) \quad T = \frac{t}{8\pi c^2} \sim \frac{v^2}{f^2} \Rightarrow f \gtrsim 5.5 \text{ TeV}$$

$$\text{e.g. FCNC } c_q^i c_q^j c_q^k c_q^l \frac{g_\rho^2}{m_\rho^2} \bar{q}^i \bar{q}^j \bar{q}^k q^l \quad c_q^i \sim \frac{g_i}{g_\rho} \Rightarrow f \gtrsim 10 \text{ TeV}$$

[DeFazio, Csaki, Serra, 1401.2417]
[Pomar, Wilczek 1506.01961]

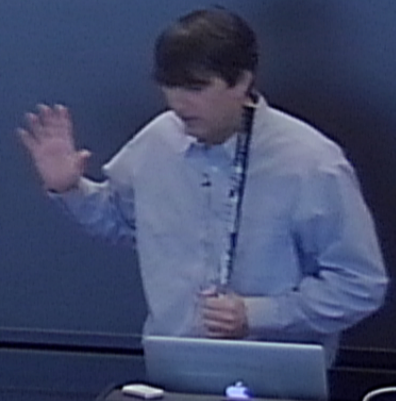
$$\Rightarrow \boxed{f \gg v} \quad \text{"Little" hierarchy}$$

Tension partly alleviated by complicating minimal models

e.g. custodial symmetry, flavor symmetry...

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HOWEVER, precision electroweak, flavor constraints

EWPT: $\frac{s}{16\pi^2 v^2} H^\dagger \tau^a H B^{\mu\nu} W_{a\mu\nu}$ $S = \frac{s}{2\pi} \sim \frac{m_W^2}{m_\rho^2} \Rightarrow f \gtrsim \frac{2.5 \text{ TeV}}{g_\rho}$

$\frac{-t}{16\pi^2 v^2} ((D^\mu H)^\dagger H) (H^\dagger D_\mu H)$ $T = \frac{t}{8\pi e^2} \sim \frac{v^2}{f^2} \Rightarrow f \gtrsim 5.5 \text{ TeV}$

e.g. FCNC $e_q^L e_q^L e_q^L \frac{g_\rho^2}{m_\rho^2} q^i q^j q^k q^l$ $e_q^L \sim \frac{g_L}{g_\rho} \Rightarrow f \gtrsim 10 \text{ TeV}$

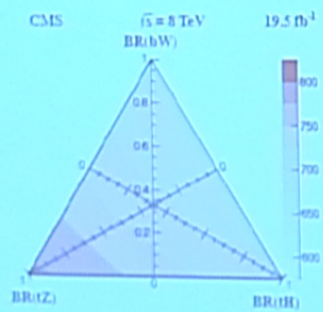
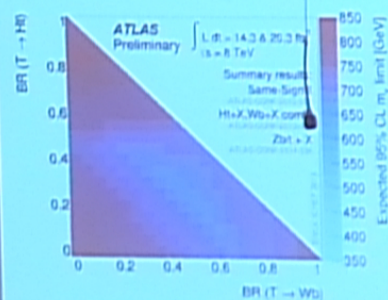
[Deffainsi, Ciuchci, Serra 1401.2457]
[Pomar, Valzer 1506.01961]

\Rightarrow $f \gg v$ "little" hierarchy

Tension partly alleviated by complicating minimal models

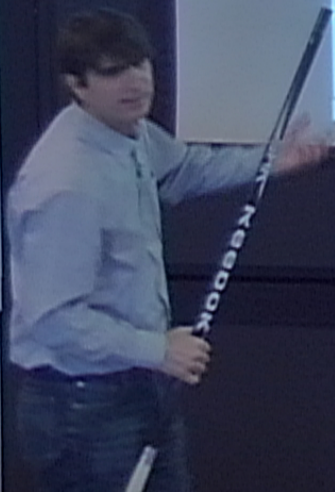
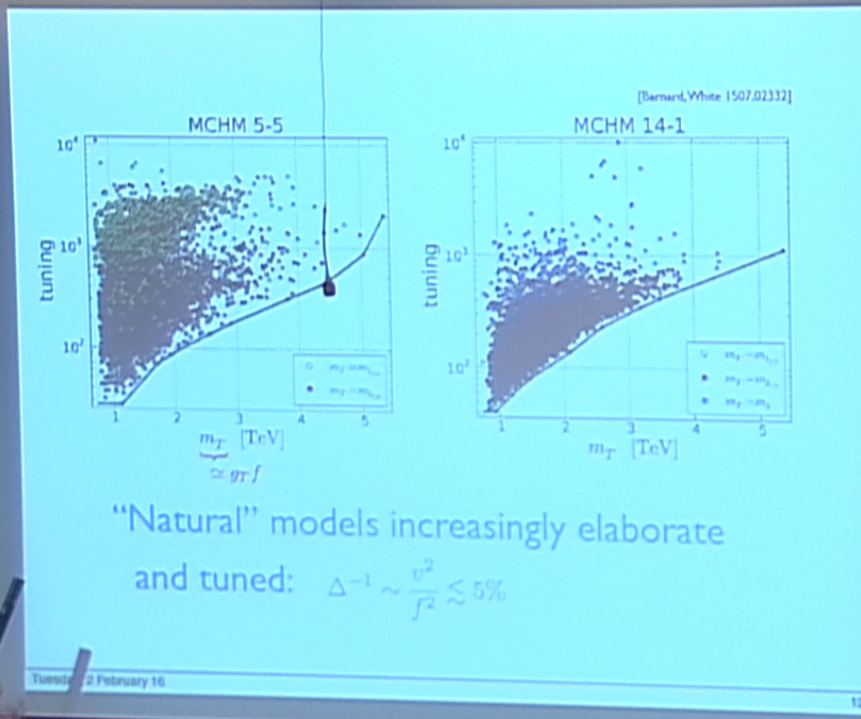
e.g. custodial symmetry, flavor symmetry...

LHC Limits: The Missing Resonances Problem



➔ $m_T \gtrsim 600 - 800 \text{ GeV}$

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Simple solution:

Assume $f \gtrsim 10 \text{ TeV}$ – no need for custodial or flavor symmetries!


Tuned Higgs potential

$$V \sim c_2 f^2 |H|^2 + c_4 |H|^4 \quad \text{tuning} \quad \frac{v^2}{f^2} \lesssim 10^{-4}$$

This compares to $\sim 10^{-28}$ in SM!

e.g. QCD - sensitivity in neutron, proton mass $\frac{m_{u,d}}{m_{\text{nucleon}}} \sim 10^{-3}$

Is there a motivated upper bound for f ?

Yes! 

Gauge coupling unification [Aguiar, Coimbra, Sundrum '05]

Assume composite t_R and coset \mathcal{G}/\mathcal{H}

$(t_R, \chi^c) =$ complete \mathcal{H} multiplet

Decoupled with top "companions" χ Dirac mass: $m_\chi \sim \lambda_\chi f$

New contribution to the running of SM gauge couplings

$$\alpha_i(\mu) - \alpha_j(\mu) = \text{SM} - \underbrace{\left\{ \underbrace{H, t^c}_{\text{composite Higgs, top}}, \overbrace{t^c}^{\text{top "companions" contribution}} \right\}}_{\text{composite Higgs, top}}$$

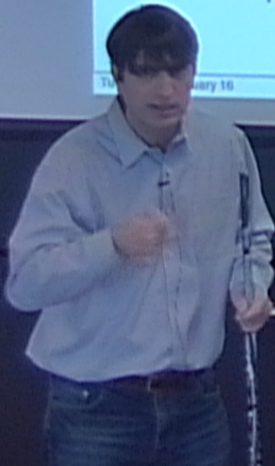
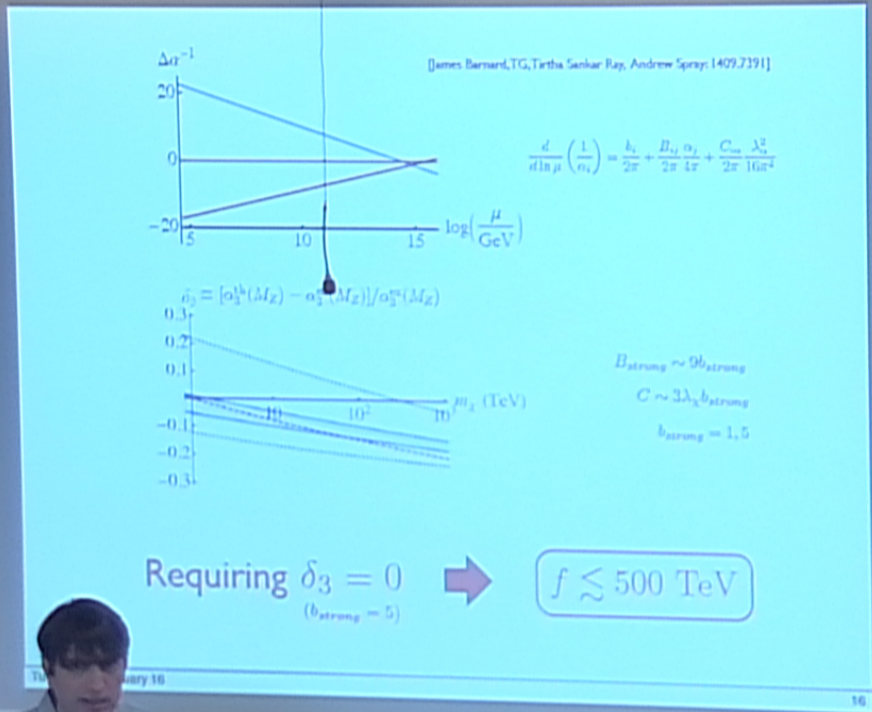
One-loop beta function coefficients:

$$b_1 - b_2 = \frac{94}{15} \quad b_2 - b_3 = \frac{13}{3} \quad \Rightarrow \quad \frac{b_2 - b_3}{b_1 - b_2} \approx 0.69$$

c.f. MSSM value = 0.71

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Effective Lagrangian

Integrate out strong sector

$$\begin{aligned}
 \mathcal{L}_{\text{eff}} \supset & (\bar{q}^c, \bar{e})_{i_4 i_2} \not{p} (\tilde{q}^c, \tilde{e})^{j_4 j_2} \left[\Pi^{\chi\chi} (\lambda_\chi^{10*})_{IJK}^{i_4 i_2} (\lambda_\chi^{10})_{j_4 j_2}^{IJL} \right] \Sigma_L^K \\
 & + (\bar{q}^c, \bar{e})_{i_4 i_2} \not{p} (\tilde{d}^c, \tilde{l})^{j_5} \left[\Pi^{\chi\chi} (\lambda_\chi^{10*})_{IJK}^{i_4 i_2} (\lambda_\chi^5)_{j_5}^{IJL} \right] \Sigma_L^K + \text{h.c.} \\
 & + (\bar{d}^c, \bar{l})_{i_5} \not{p} (\tilde{d}^c, \tilde{l})^{j_5} \left[\Pi^{\chi\chi} (\lambda_\chi^{5*})_{IJK}^{i_5} (\lambda_\chi^5)_{j_5}^{IJL} \right] \Sigma_L^K \\
 & + \bar{q}^{i_3 i_2} \not{p} q_{j_3 j_2} \left[\Pi^{tt} (\lambda_t^*)_{i_3 i_2, IJK} (\lambda_t)^{j_3 j_2, IJL} + \Pi^{bb} (\lambda_b^*)_{i_3 i_2}^{IJL} (\lambda_b)^{j_3 j_2}_{IJK} \right] \Sigma_L^K \\
 & + \bar{b}_{i_3}^c \not{p} b^{c j_3} \left[\Pi^{b^c b^c} (\lambda_{b^c}^*)_{IJK}^{i_3} (\lambda_{b^c})_{j_3}^{IJL} \right] \Sigma_L^K \\
 & + (\bar{q}^c, \bar{e})_{i_4 i_2} \not{p} q_{j_3 j_2} \left[\Pi^{\chi t} (\lambda_\chi^{10*})_{IJK}^{i_4 i_2} (\lambda_t)^{j_3 j_2, IJL} \right] \Sigma_L^K + \text{h.c.} \\
 & + (\bar{d}^c, \bar{l})_{i_5} \not{p} q_{j_3 j_2} \left[\Pi^{\chi t} (\lambda_\chi^{5*})_{IJK}^{i_5} (\lambda_t)^{j_3 j_2, IJL} \right] \Sigma_L^K + \text{h.c.} \\
 & + q_{i_3 i_2} b^{c j_3} \left[M^{b^c b^c} (\lambda_b)_{IJK}^{i_3 i_2} (\lambda_{b^c})_{j_3}^{IJL} \right] \Sigma_L^K + \text{h.c.}
 \end{aligned}$$

where $\Sigma = w^\dagger w$ = adjoint spurion (contains D, T and S)

$\Pi, M^{b^c b^c}$ = momentum dependent form factors

Obtain:

$$V(|D|) = -\frac{\alpha}{f^2}|D|^2 + \frac{\beta}{f^4}|D|^4$$

HIGGS POTENTIAL

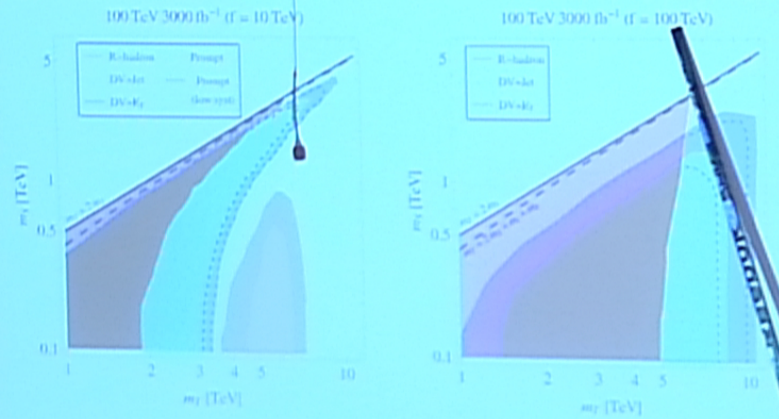
Electroweak VEV: $v = f \sqrt{\frac{\alpha}{\beta}}$ *must be tuned*

Higgs mass: $m_h^2 = \frac{2\beta v^2}{f^4} = \frac{3c_2^A g_\rho^2}{8\pi^2} M_W^2$ *W-boson mass* *Requires: $c_2^A \sim \frac{64}{g_\rho^2} \sim 0.5 - 4$*

where $\alpha = \frac{g_\rho^2}{16\pi^2} f^4 \left(\frac{14}{3} c_1^X |\lambda_X|^2 - 2c_1^t |\lambda_t|^2 - 2c_1^b |\lambda_b|^2 + 2c_1^{bc} |\lambda_{bc}|^2 - \frac{9}{4} c_1^A g_2^2 \right)$
 $\beta = \frac{g_\rho^2}{16\pi^2} f^4 \left(\frac{3}{4} c_2^A g_2^2 \right)$

Future 100 TeV collider:

[Bernard, Cox, TG, Spray: 1510.06405]



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Summary

- $f \gtrsim 10$ TeV simply eliminates all precision electroweak and flavor constraints
 - Higgs potential is tuned at 10^{-4} level
 - “Unnatural” or “split” composite Higgs
- SU(7)/SU(6)xU(1) minimal model
 - Improves gauge coupling unification
 - Explains fermion mass hierarchy
- Higgs partners: S = dark matter, T = color triplet
- Long-lived T decays = sign of unnaturalness!