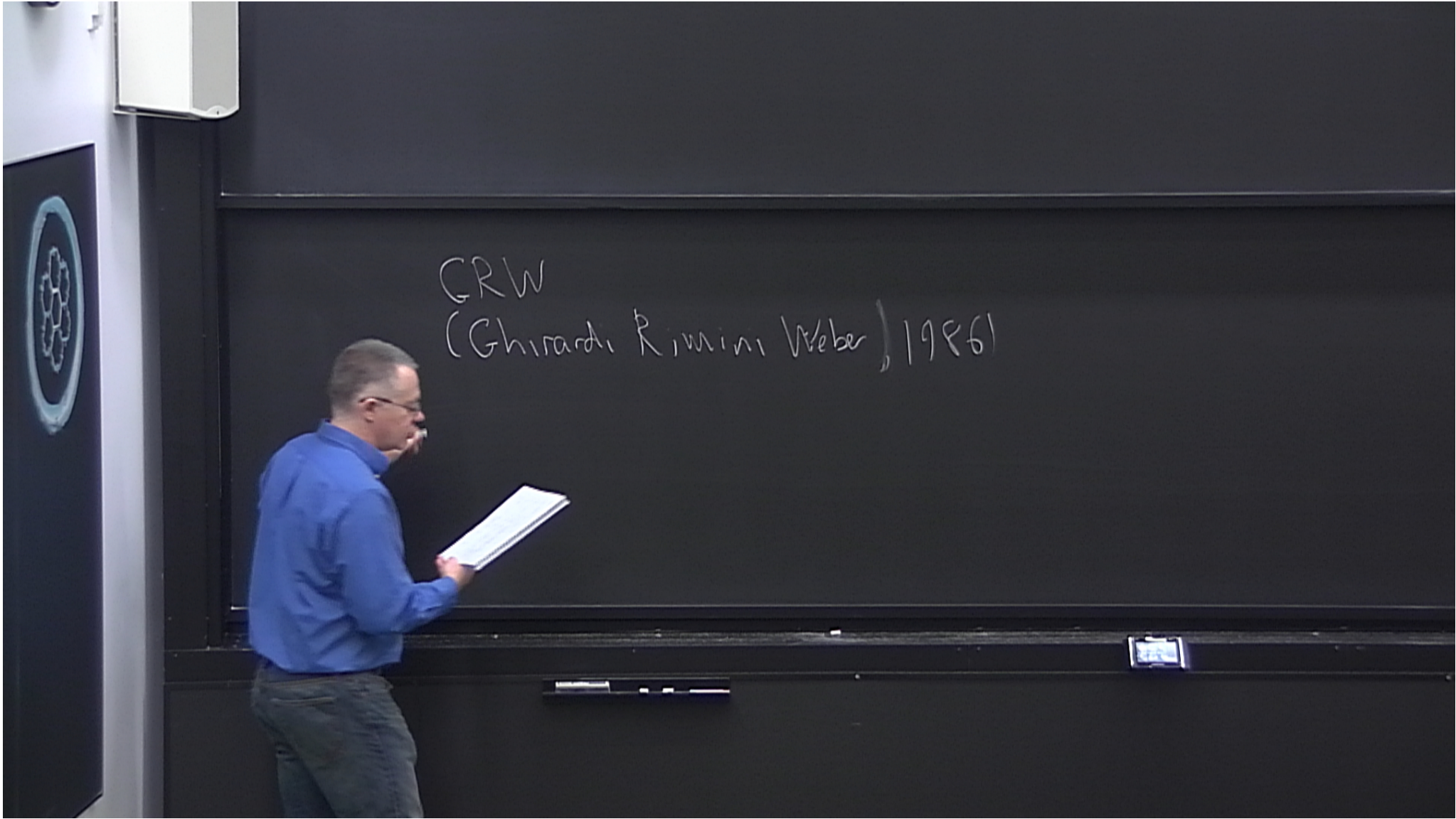


Title: PSI 2015/2016 Foundations of Quantum Mechanics - Lecture 10

Date: Jan 15, 2016 11:30 AM

URL: <http://pirsa.org/16010064>

Abstract:



in Weber, 1986
Are there Q. Jumps?

① The ontology

(a) a wave fn $\psi(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N, t)$

(b) a way to get local beables

all flash
ontology

radi mass density
ontology.

② The dynamics

the Schrödinger eqn except when

in Weber, 1986)

Are there Q. Jumps?

① The ontology

(a) a wave fn $\psi(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N, t)$

(b) a way to get local beables

- Bell flash ontology
- Ghirardi mass density ontology.

② The dynamics

the Schrödinger eqn except when jumps occur.

(2) The dynamics

the Schrödinger equation except when jumps occur

Ghirardi mass density ontology.

(ii) The jump is

$$\psi \rightarrow \frac{j(\vec{x} - \vec{r}_n) \psi(\vec{r}_1, \dots, \vec{r}_N, t)}{R_n(\vec{x})}$$

n is chosen randomly from 1 to N

(2) The dynamics

the Schrödinger equation except when jumps occur

Ghirardi mass density ontology.

(ii) The jump is

$$\rightarrow \frac{j(\vec{x} - \vec{r}_n) \psi(\vec{r}_1, \dots, \vec{r}_N, t)}{R_n(\vec{x})}$$

chosen randomly from 1 to N and

$$\int d^3\vec{x} |j(\vec{x})|^2 = 1$$

$$P(\text{\# jumps in time } T = M) = \frac{e^{-\left(\frac{N}{\tau} T\right)}}{M!}$$

Poisson distribution.

if positions randomly from 1

$$\int d^3 \vec{x} |j(x)|^2 = 1$$

usual choice

$$j(x) = k \exp\left(-x^2 / 2a^2\right)$$

another fundamental constant of nature.

collapse centre \vec{x}

$$P(\text{\# jumps in time } T = M) = \frac{e^{-\left(\frac{N}{\tau} T\right)}}{M!}$$

Poisson distribution.

if chosen randomly from

$$\int d^3 \vec{x} |j(\vec{x})|^2 = 1$$

usual choice

$$j(\vec{x}) = k \exp\left(-\frac{\vec{x}^2}{2a^2}\right)$$

another fundamental constant of nature.

(iii) The collapse centre \vec{x} is chosen with prob

$$d^3 \vec{x} |R_j(\vec{x})|^2$$

... chosen randomly from 1 to N and

$$\int d^3\vec{x} |j(\vec{x})|^2 = 1$$

$$\alpha^2 / 2a^2$$

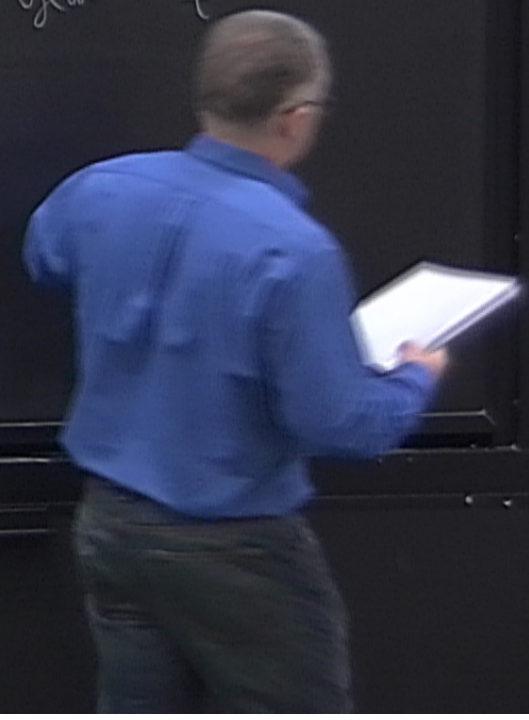
another fundamental constant of nature.

... chosen with prob.

(iv) GRW suggested

$$\tau = 10^{15} \text{ sec} \approx 10^8 \text{ year} \quad \frac{N}{\tau}$$

$$a = 10^{-5} \text{ cm}$$



is chosen randomly from 1 to N and

$$\int d^3\vec{x} |j(\vec{x})|^2 = 1$$

$$\frac{\hbar^2}{2m^2 a^2}$$

another fundamental constant of nature.

chosen with prob.

(iv) GRW suggested

$$\tau = 10^{15} \text{ sec} \approx 10^8 \text{ year} \left(\frac{N}{\tau} \text{ can be big for big } N \right)$$

$$a = 10^{-5} \text{ cm}$$

bigger than an atom
smaller than eye can discern

n chosen randomly from 1 to N and

$$\int d^3 \vec{x} |j(x)|^2 = 1$$

$$\frac{\hbar^2}{2m^2 a^2}$$

another fundamental constant of nature.

chosen with prob.

(iv) GRW suggested

$$\tau = 10^{15} \text{ sec} \approx 10^8 \text{ year} \left(\frac{N}{\tau} \text{ can be big for big } N \right)$$

$$a = 10^{-5} \text{ cm} \quad \begin{array}{l} \text{bigger than an atom} \\ \text{smaller than eye can discern} \end{array}$$

$$\text{and } |R_n(x)|^2 = \int d^3 \vec{r}_1 d^3 \vec{r}_2 \dots d^3 \vec{r}_N |j(\vec{x} - \vec{r}_n) \psi|^2$$

... chosen randomly from 1 to N and

$$\int d^3 \vec{x} |j(x)|^2 = 1$$

$$\frac{\hbar^2}{2m^2}$$

another fundamental constant of nature.

... chosen with prob.

(iv) GRW suggested

$$\tau = 10^{15} \text{ sec} \approx 10^8 \text{ year} \left(\frac{N}{\tau} \text{ can be big for big } N \right)$$

$$a = 10^{-5} \text{ cm} \quad \begin{array}{l} \text{bigger than an atom} \\ \text{smaller than eye can discern} \end{array}$$

$$\text{and } |R_n(x)|^2 = \int d^3 \vec{r}_1 d^3 \vec{r}_2 \dots d^3 \vec{r}_N |j(\vec{x} - \vec{r}_n) \psi|^2$$

N chosen randomly from 1 to N and

$$\int d^3\vec{x} |j(x)|^2 = 1$$

$$\left(\frac{1}{2a}\right)^2$$

another fundamental constant of nature.

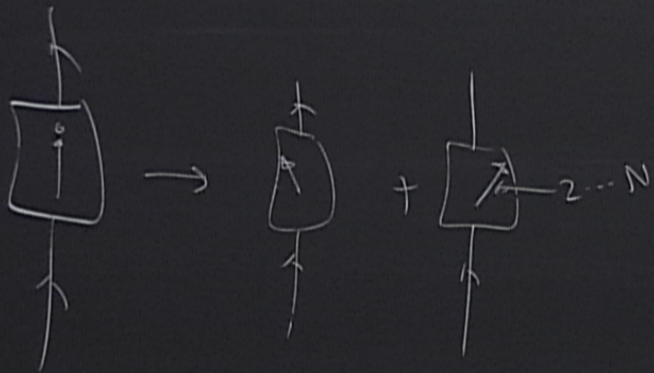
chosen with prob.

(iv) GRW suggested

$$\tau = 10^{15} \text{ sec} \approx 10^8 \text{ years} \left(\frac{N}{\tau} \text{ can be big for big } N \right)$$

$$a = 10^{-5} \text{ cm} \quad \begin{array}{l} \text{bigger than an atom} \\ \text{smaller than eye can discern} \end{array}$$

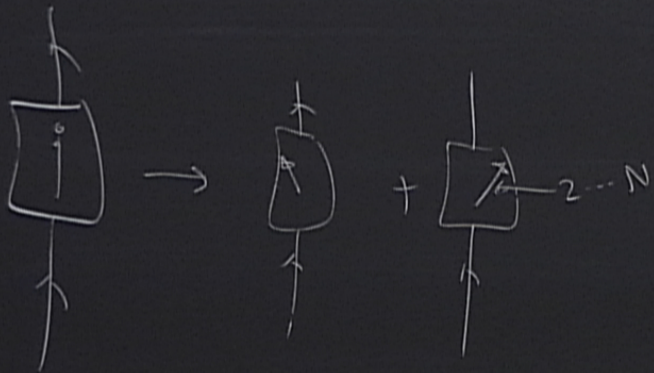
$$\text{and } |R_N(x)|^2 = \int d^3\vec{r}_1 d^3\vec{r}_2 \dots d^3\vec{r}_N |j(\vec{x} - \vec{r}_N) \psi|^2$$



$$|\psi\rangle|A_0\rangle \rightarrow \alpha|a_1\rangle|A_1\rangle + \beta|a_2\rangle|A_2\rangle$$

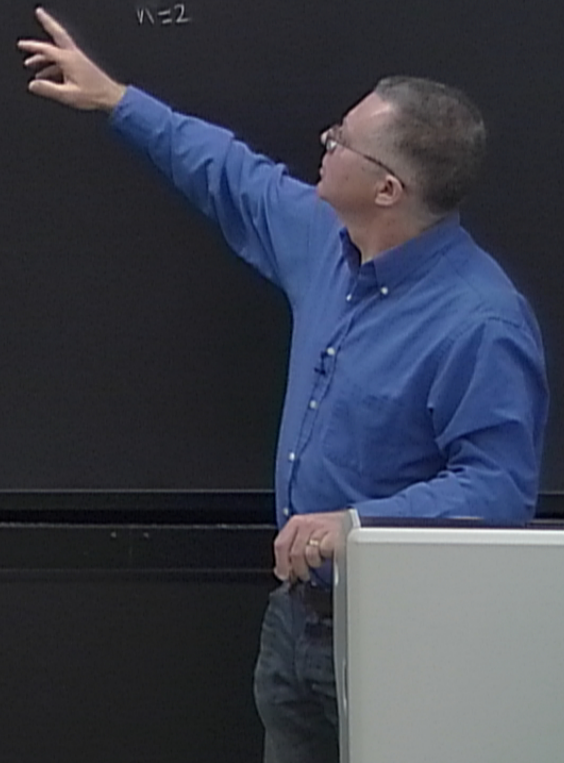
$$\simeq \alpha a(\gamma_1) \prod_{n=2}^N A_1(\gamma_n) + \beta a_2(\gamma_1) \prod_{n=2}^N A_2(\gamma_n)$$

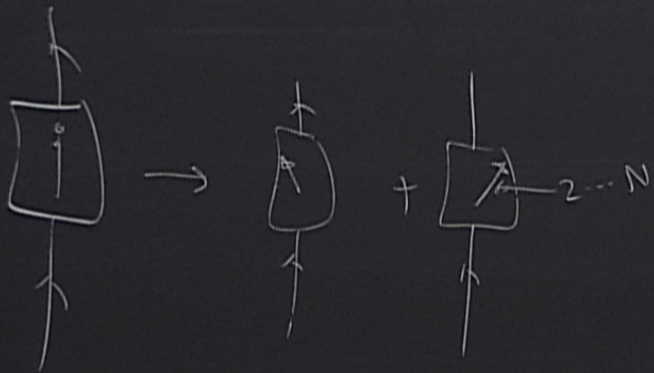




$$|\psi\rangle|A_0\rangle \rightarrow \alpha|a_1\rangle|A_1\rangle + \beta|a_2\rangle|A_2\rangle$$

$$\simeq \alpha a(\gamma_1) \bigotimes_{n=2}^N A_1(\gamma_n) + \beta a_2(\gamma_1) \bigotimes_{n=2}^N A_2(\gamma_n)$$





$$|\psi\rangle|A_0\rangle \rightarrow \alpha|a_1\rangle|A_1\rangle + \beta|a_2\rangle|A_2\rangle$$

$$\simeq \alpha a_1(\gamma_1) \bigotimes_{n=2}^N A_1(\gamma_n) + \beta a_2(\gamma_1) \bigotimes_{n=2}^N A_2(\gamma_n)$$

where $|A_1(\gamma_n)|/|A_2(\gamma_n)| \simeq 0$

$$\approx \alpha a_1(r_1) \prod_{n=2}^N A_1(r_n) + \beta a_2(r_1) \prod_{n=2}^N A_2(r_n)$$

jump \rightarrow

$$a_1(r_1) \prod_{n=2}^N A_1(r_n)$$

where

$$|A_1(r_n)| / |A_2(r_n)| \approx 0$$

2...N

$a_2)/A_2)$

$$|\psi\rangle|A_0\rangle \rightarrow \alpha|a_1\rangle|A_1\rangle + \beta|a_2\rangle|A_2\rangle$$

Ψ

Φ

Ψ Φ



$$|\psi\rangle|A_0\rangle \rightarrow \alpha|a_1\rangle|A_1\rangle + \beta|a_2\rangle|A_2\rangle$$

$$\Psi$$

$$\Phi$$

$$\Psi \quad \Phi$$

The master eqn

$$\rho(\xi, \xi_m^3, \xi, \xi_m^3) = \Psi(\xi, \xi_m^3) \Psi^*(\xi, \xi_m^3)$$

$$\frac{d\rho}{dt} = -\sum_n \frac{1}{\tau} (\rho - T_n(\rho))$$

$$|\Psi\rangle|A_0\rangle \rightarrow \alpha|a_1\rangle|A_1\rangle + \beta|a_2\rangle|A_2\rangle$$

$$\Psi$$

$$\Phi$$

$$\Psi \Phi$$

The master eqn

$$\rho(\xi, \Gamma_m^3, \xi, \Gamma_m^3) = \Psi(\xi, \Gamma_m^3) \Psi^*(\xi, \Gamma_m^3)$$

$$\frac{d\rho}{dt} = -\frac{i}{\hbar} [\hat{H}, \rho] - \sum_n \frac{1}{\tau} (\rho - T_n(\rho))$$

where

$$T_n(\rho) = \int d^3\vec{x} j(x - \Gamma_n) \rho j(x - \Gamma_n')$$

$$|\psi\rangle|A_0\rangle \rightarrow \alpha|a_1\rangle|A_1\rangle + \beta|a_2\rangle|A_2\rangle$$

$$\Psi$$

$$\Phi$$

$$\Psi \quad \Phi$$

The master eqn

$$\rho(\{\Gamma_m\}, \{\Gamma'_m\}) = \Psi(\{\Gamma_m\}) \Psi^*(\{\Gamma'_m\})$$

$$\frac{d\rho}{dt} = \frac{-i}{\hbar} [\hat{H}, \rho] - \sum_n \frac{1}{\tau} (\rho - T_n(\rho))$$

where

$$T_n(\rho) = \int d^3\vec{x} j(x - \Gamma_n) \rho j(x - \Gamma'_n)$$

$$|\psi\rangle|A_0\rangle \rightarrow \alpha|a_1\rangle|A_1\rangle + \beta|a_2\rangle|A_2\rangle$$

$$\Psi$$

$$\Phi$$

$$\Psi \quad \Phi$$

The master eqn

$$\rho(\xi, \Gamma_m^3, \xi, \Gamma_m^3) = \Psi(\xi, \Gamma_m^3) \Psi^*(\xi, \Gamma_m^3)$$

$$\frac{d\rho}{dt} = \frac{-i}{\hbar} [\hat{H}, \rho] - \sum_n \frac{1}{\tau} (\rho - T_n(\rho))$$

where

$$T_n(\rho) = \int d^3\vec{x} j(x - \Gamma_n) \rho j^*(x - \Gamma_n)$$

N ← number of particles
 τ ← a new constant of nature.

$$P(\text{\# jumps in time } T = M) = \frac{e^{-N\tau/T} (N\tau/T)^M}{M!}$$

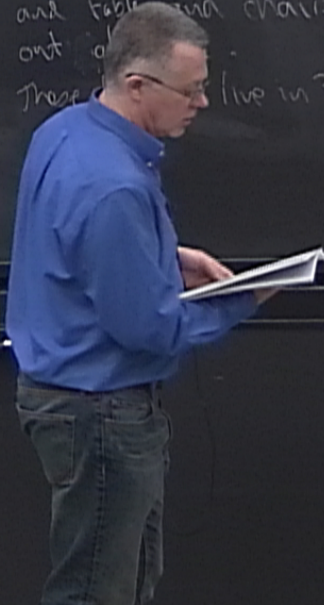
Poisson distribution.

$$\psi \rightarrow \frac{j(\vec{x} - \vec{r}_n) \psi(\vec{r}_1, \dots, \vec{r}_N, t)}{R_n(\vec{x})}$$

n chosen randomly from 1 to N and
 $\int d^3\vec{x} |j(\vec{x})|^2 = 1$

The Hoshino ontology.

the Hoshinos are the jump centres
 and tables and chairs are made
 out of them.
 These things live in 3D.



N ← number of particles
 τ ← a new constant of nature.

$$P(\text{\# jumps in time } T = M) = \frac{e^{-NT/\tau} (NT/\tau)^M}{M!}$$

Poisson distribution.

$$\psi \rightarrow \frac{j(\vec{x} - \vec{r}_n) \psi(\vec{r}_1, \dots, \vec{r}_N, t)}{R_n(\vec{x})}$$

n chosen randomly from 1 to N and
 $\int d^3\vec{x} |j(\vec{x})|^2 = 1$

The Hoshino ontology.

the Hoshinos are the jump centres
 and tables and chairs are made
 out of these.
 These Hoshinos live in 3D.

The mass density ontology - (Ghirardi, Grassano)

$\frac{N}{T}$ ← number of particles
 T ← a new constant of nature.

jumps in time $T = M$

$$= \frac{e^{-\frac{N}{T}T} \left(\frac{N}{T}T\right)^M}{M!}$$

Poisson distribution.

$$\psi \rightarrow \frac{j(\vec{x} - \vec{r}_n) \psi(\vec{r}_1, \dots, \vec{r}_N, t)}{R_n(\vec{x})}$$

n chosen randomly from 1 to N and

$$\int d^3\vec{x} |j(\vec{x})|^2 = 1$$

The flash ontology.

the flashes are the jump centres
 and tables and chairs are made
 out of these.
 These flashes live in 3D.

The mass density ontology - (Ghirardi, Grassi, Benatti, 1995)

Define a mass density op

$$\hat{M}(x) = \sum_k m_k \hat{a}_k(x) \hat{a}_k(x)$$

$m(x, t) = \langle \psi | \hat{M} | \psi \rangle$

k labels types of particles

mass volume

$\frac{N}{T}$ ← number of particles
 T ← a new constant of nature.

$$P(n) = \frac{e^{-N/T} (N/T)^n}{n!}$$

Poisson distribution.

$$\psi \rightarrow \frac{j(\vec{x} - \vec{r}_n) \psi(\vec{r}_1, \dots, \vec{r}_n, t)}{R_n(\vec{x})}$$

n chosen randomly from 1 to N and
 $\int d^3\vec{x} |j(\vec{x})|^2 = 1$

The flash ontology.

the flashes are the jump centres
 and tables and chairs are made
 out of these.

These flashes live in 3D.

The mass density ontology - (Ghirardi, Grassi, Benatti, 1995)

Define a mass density op

$$\hat{M}(x) = \sum_k m_k \underbrace{\hat{a}_k^{\dagger}(x) \hat{a}_k(x)}_{\substack{\# \text{ macro-} \\ \text{volume of size } a}}$$

k labels types of particles

$$m(x,t) = \langle \psi | \hat{M} | \psi \rangle$$