

Title: PSI 2015/2016 Foundations of Quantum Mechanics - Lecture 6

Date: Jan 11, 2016 11:30 AM

URL: <http://pirsa.org/16010059>

Abstract:

The de Broglie Bohm interpretation.

Take case of - non rel. QM.

with

$\psi(\vec{x}_1, \vec{x}_2, \dots, \vec{x}_N, t)$ evolving

by

$$i\hbar \frac{\partial \psi}{\partial t} = \sum_{n=1}^N \frac{-\hbar^2}{2M_n} \nabla_n^2 \psi + V\psi$$

The ontic state at time t is

$$\left(\psi(\vec{x}_1, \vec{x}_2, \dots, \vec{x}_N, t), (\vec{X}_1, \vec{X}_2, \dots, \vec{X}_N, t) \right)$$

ψ
wave fn on config space

ϕ
actual position in config space.

solving

ψ

The dynamics is given by

$$1) \quad i\hbar \frac{d\psi}{dt} = \sum_{n=1}^N \frac{-\hbar^2}{2M_n} \nabla_n^2 \psi + V\psi$$

Schro.
eqn.

$$2) \quad \left. \frac{dX_n}{dt} = \frac{\hbar}{M_n} \frac{\text{Im}(\psi^* \nabla_n \psi)}{\psi^* \psi} \right|$$

Schro
eqn.

$$x = (\vec{x}_1, \dots, \vec{x}_N), \quad X = (\vec{X}_1, \dots, \vec{X}_N)$$

Sometimes a third assump. is added

$$\rho(\vec{X}_1, \vec{X}_2, \dots, \vec{X}_N, 0) = |\psi(\vec{X}_1, \vec{X}_2, \dots, \vec{X}_N, 0)|^2$$

prob density for finding particles at X

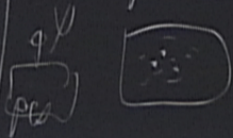
Schro
eqn.

$$x = (\vec{x}_1, \dots, \vec{x}_N), \quad X = (\vec{X}_1, \dots, \vec{X}_N)$$

Sometimes a third assump. is added

$$3) \quad \rho(\vec{X}_1, \vec{X}_2, \dots, \vec{X}_N, 0) = |\psi(\vec{X}_1, \vec{X}_2, \dots, \vec{X}_N, 0)|^2$$

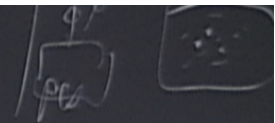
↑
prob density for finding particles at X



= X

t_n

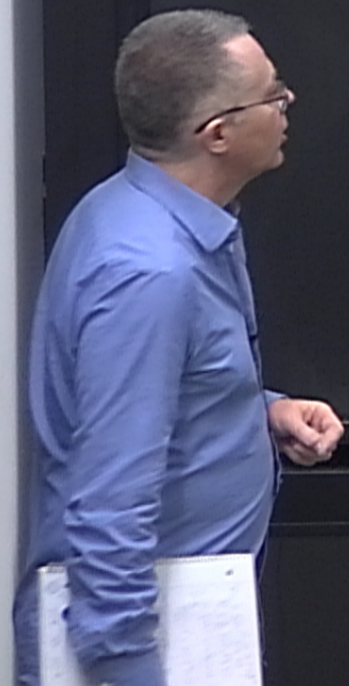
$t_x = X$



What happens when system evolves?

Answer. $\forall t \quad \rho(0) = |\psi(0)\rangle^2$ at $t=0$

then $\rho(t) = |\psi(t)\rangle^2$ for all later t .



t_n (1/2) (2)
What happens when system evolves?

Answer. $\forall t \quad \rho(0) = |\psi(0)|^2 \quad \text{at } t=0$

then $\rho(t) = |\psi(t)|^2 \quad \text{for all later } t$ - quantum equilibrium.

this property is called equivariance. Valentini.

Note that

$$\vec{J}_n^\psi = \frac{\hbar}{M_n} \operatorname{Im}(\psi^* \nabla_k \psi)$$

probability current. It satisfies

$$\frac{\partial |\psi(x,t)|^2}{\partial t} + \operatorname{div} J^\psi(x,t) = 0$$

that

$$= \frac{\hbar}{m} \operatorname{Im}(\psi^* \nabla_k \psi)$$

probability current. It satisfies

$$\frac{\partial |\psi(x,t)|^2}{\partial t} + \operatorname{div} J^y(x,t) = 0$$

where

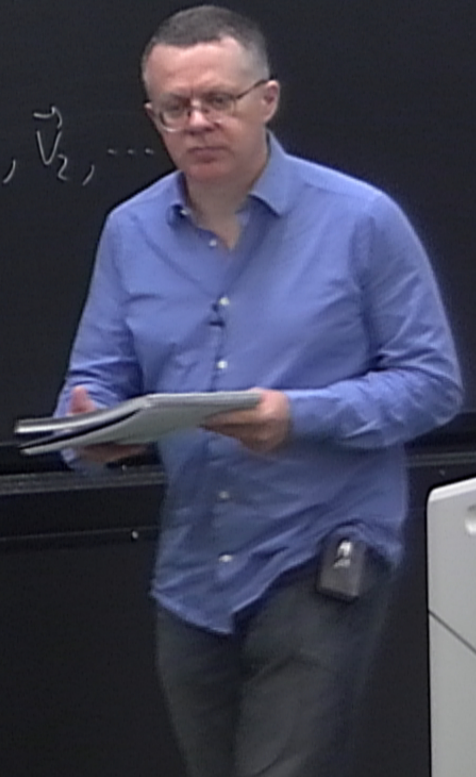
$$J^y(x,t) = (\vec{J}_1^y, \vec{J}_2^y, \dots, \vec{J}_N^y)$$

It we put $|\psi|^2 = \rho$

and $J = \rho v$ ($v = (\vec{v}_1, \vec{v}_2, \dots)$)

from

$$\frac{\partial \rho}{\partial t} + \operatorname{div} \rho v = 0$$

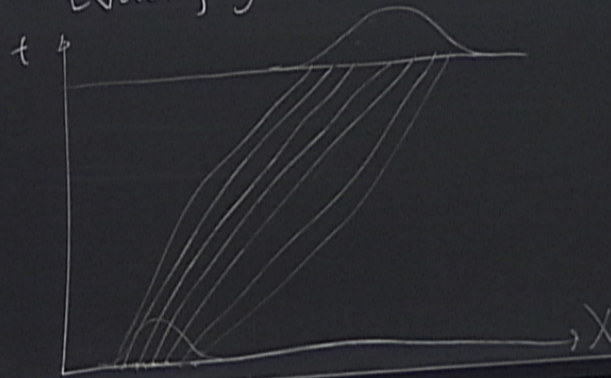


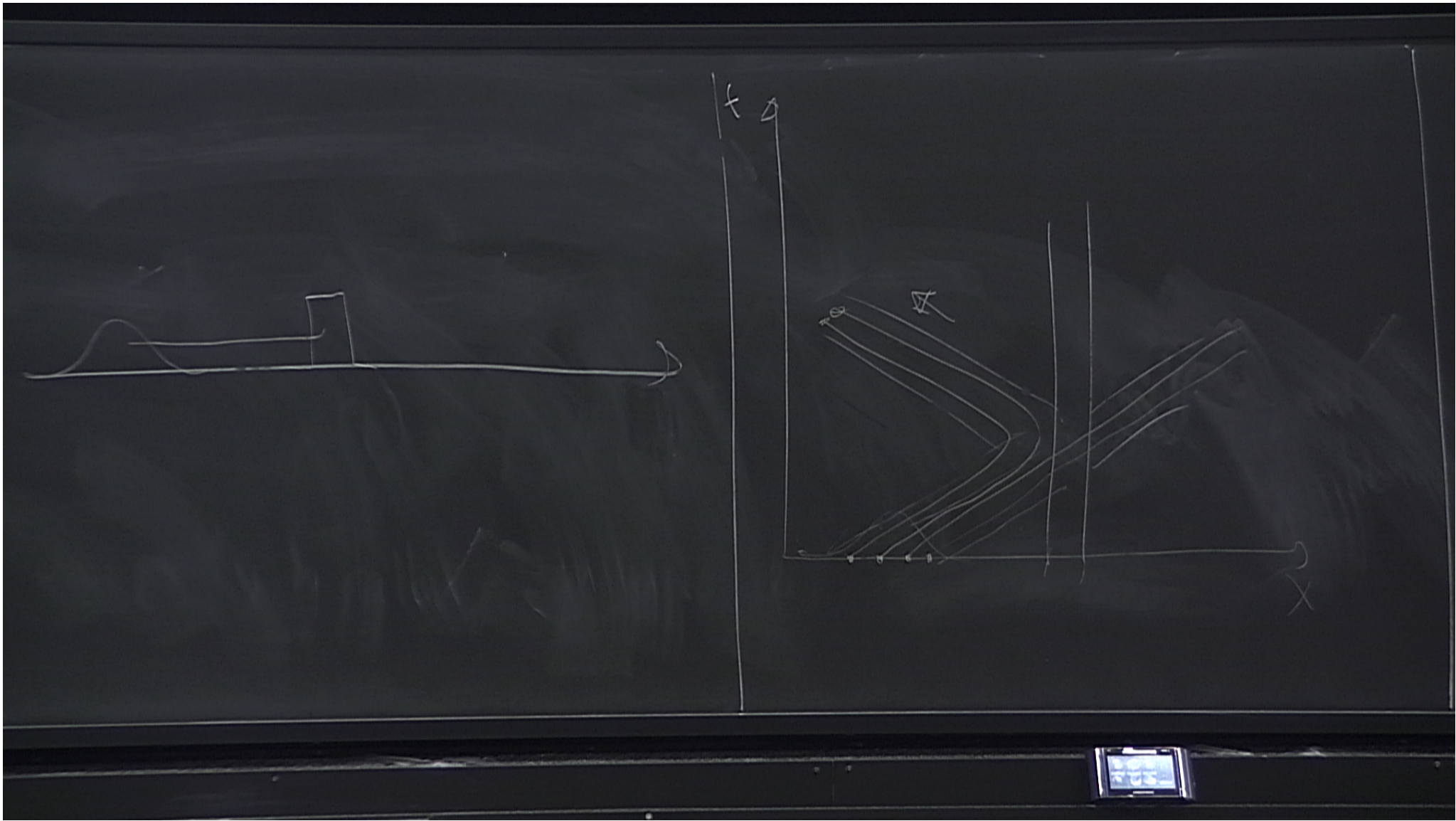
$\psi = \psi(x) = (\psi_1, \psi_2, \dots, \psi_N)$ is a single
valued fn of ψ (and its deriv) evaluated at x



② $\rho = |\psi|^2$ at all times.

fool ex. Evolving gaussian wave packet (1st particle
in 2D)





Measurement in the deBB model

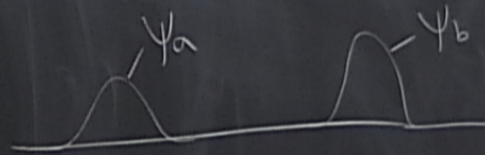
Consider a particle

$$\psi(x) = \alpha \psi_a(x) + \beta \psi_b(x)$$

Measurement in the deBB model.

a particle

$$\psi(x) = \alpha \psi_a(x) + \beta \psi_b(x)$$

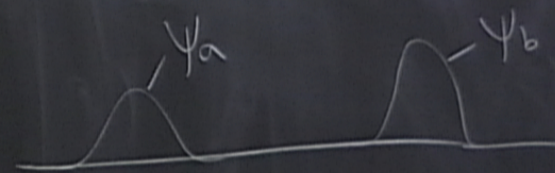


Measurement in the deBB model.

Consider a particle

$$\psi(x) = \alpha \psi_a(x) + \beta \psi_b(x)$$

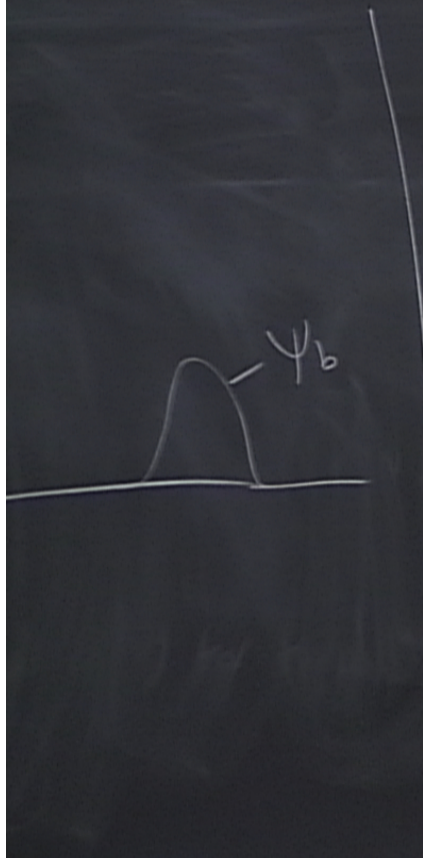
Introduce meas apparatus with pointer (y) s.t.



$$\psi_a(x) \varphi_0(y) \xrightarrow{\text{Schro}} \psi_a(x) \varphi_a(y)$$

$$\psi_b(x) \varphi_0(y) \xrightarrow{\text{Schro}} \psi_b(x) \varphi_b(y)$$

ψ_b



$$\psi_a(x) \varphi_0(y) \xrightarrow{\text{Schro}} \psi_a(x) \varphi_a(y)$$

$$\psi_b(x) \varphi_0(y) \xrightarrow{\text{Schro}} \psi_b(x) \varphi_b(y)$$

$$\psi(x) \varphi_0(y)$$

$$\longrightarrow \alpha \psi_a(x) \varphi_a(y) + \beta \psi_b(x) \varphi_b(y)$$