

Title: BMS Symmetry in Three Dimensions

Date: Dec 08, 2015 03:30 PM

URL: <http://pirsa.org/15120022>

Abstract: <p>In this talk I will sketch the relation between unitary representations of the BMS3 group and three-dimensional, asymptotically flat gravity. More precisely, after giving an exact definition of the BMS group in three dimensions, I will argue that its unitary representations are classified by orbits of CFT stress tensors under conformal transformations. These stress tensors, in turn, can be interpreted as Bondi mass aspects for asymptotically flat metrics. I will also show how one can compute characters of the BMS3 group, which coincide with suitable gravitational one-loop partition functions</p>

BMS Symmetry in Three Dimensions

Blagoje Oblak

ULB (Brussels, Belgium) & DAMTP (Cambridge, UK)

December 2015

MOTIVATION

Flat space-time

MOTIVATION

Flat space-time

- ▶ Symmetries : Poincaré \sim Lorentz \times Translations

But space-time is not flat...

- ▶ “Asymptotically flat” space-times

MOTIVATION

Flat space-time

- ▶ Symmetries : Poincaré \sim Lorentz \times Translations

But space-time is not flat...

- ▶ “Asymptotically flat” space-times
- ▶ Symmetries : BMS \sim Superrotations \times Supertranslations

[Bondi *et al.* 1962, Barnich *et al.* 2009]

PLAN OF THE TALK

- 1. Unitary reps of semi-direct products**
2. Representations of BMS_3
3. Relation to gravity

PLAN OF THE TALK

- 1. Unitary reps of semi-direct products**
2. Representations of BMS_3
3. Relation to gravity

1. Unitary representations of semi-direct products

SEMI-DIRECT PRODUCTS

$$P = \text{SO}(2, 1) \ltimes \mathbb{R}^3$$

- ▶ Elements of $P =$ pairs (f, α)
- ▶ Group operation $(f, \alpha) \cdot (g, \beta) \equiv (f \cdot g, \alpha + f \cdot \beta)$

SEMI-DIRECT PRODUCTS

$$P = \text{SO}(2, 1) \ltimes \mathbb{R}^3$$

- ▶ Elements of P = pairs (f, α)
- ▶ Group operation $(f, \alpha) \cdot (g, \beta) \equiv (f \cdot g, \alpha + f \cdot \beta)$
- ▶ P is a **semi-direct product** :

$$P = G \ltimes_{\sigma} A$$

SEMI-DIRECT PRODUCTS

UIRREPS of $P = SO(2, 1) \ltimes \mathbb{R}^3$?

[Wigner 1939]

SEMI-DIRECT PRODUCTS

UIRREPS of $P = G \ltimes_{\sigma} A$?

[Wigner 1939]

- ▶ Start with Abelian group A
- ▶ UIRREPS are one-dimensional :

$$\alpha \longmapsto e^{i \langle p, \alpha \rangle}, \quad p \in A^*$$

- ▶ $p =$ **momentum**

SEMI-DIRECT PRODUCTS

UIRREPS of $P = G \ltimes_{\sigma} A$?

[Wigner 1939]

- ▶ Start with Abelian group A
- ▶ UIRREPS are one-dimensional :

$$\alpha \longmapsto e^{i \langle p, \alpha \rangle}, \quad p \in A^*$$

- ▶ $p =$ **momentum**

A UIRREP of P contains many momenta

ORBITS

Let $\mathcal{T} = \text{UIRREP of } G \ltimes_{\sigma} A$

ORBITS

Let $\mathcal{T} = \text{UIRREP}$ of $G \ltimes_{\sigma} A$ in a Hilbert space \mathcal{H} , with

$$\mathcal{T}[(e, \alpha)] \Big|_{\mathbb{V}} = e^{i\langle p, \alpha \rangle} \quad \text{on } \mathbb{V} \leq \mathcal{H}.$$

Pick $f \in G$.

ORBITS

Let $\mathcal{T} = \text{UIRREP}$ of $G \ltimes_{\sigma} A$ in a Hilbert space \mathcal{H} , with

$$\mathcal{T}[(e, \alpha)] \Big|_{\mathbb{V}} = e^{i\langle p, \alpha \rangle} \quad \text{on } \mathbb{V} \leq \mathcal{H}.$$

Pick $f \in G$.

Then, $\exists \mathbb{V}' \leq \mathcal{H}$ s.t.

$$\mathcal{T}[(e, \alpha)] \Big|_{\mathbb{V}'} = e^{i\langle p, \sigma_{f^{-1}} \alpha \rangle}$$

ORBITS

Let $\mathcal{T} = \text{UIRREP}$ of $G \ltimes_{\sigma} A$ in a Hilbert space \mathcal{H} , with

$$\mathcal{T}[(e, \alpha)] \Big|_{\mathbb{V}} = e^{i\langle p, \alpha \rangle} \quad \text{on } \mathbb{V} \leq \mathcal{H}.$$

Pick $f \in G$.

Then, $\exists \mathbb{V}' \leq \mathcal{H}$ s.t.

$$\mathcal{T}[(e, \alpha)] \Big|_{\mathbb{V}'} = e^{i\langle p, \sigma_{f^{-1}} \alpha \rangle}$$

Define $\langle f \cdot p, \alpha \rangle \equiv \langle p, \sigma_{f^{-1}} \alpha \rangle$

► **Orbit** of p : $\mathcal{O}_p \equiv \{f \cdot p \mid f \in G\}$

UNITARY REPS

Let's build UIRREPS of P !

UNITARY REPS

Let's build UIRREPS of P !

- ▶ Pick an orbit \mathcal{O}_p
- ▶ Let $\mathcal{H} =$ space of **wavefunctions** $\Psi : \mathcal{O}_p \rightarrow \mathbb{C} : q \mapsto \Psi(q)$

UNITARY REPS

Let's build UIRREPS of P !

- ▶ Pick an orbit \mathcal{O}_p
- ▶ Let $\mathcal{H} =$ space of **wavefunctions** $\Psi : \mathcal{O}_p \rightarrow \mathbb{C} : q \mapsto \Psi(q)$
- ▶ Define $(\mathcal{T}[(f, \alpha)]\Psi)(q) = e^{i\langle q, \alpha \rangle} \Psi(f^{-1} \cdot q)$

UNITARY REPS

Let's build UIRREPS of P !

- ▶ Pick an orbit \mathcal{O}_p
- ▶ Let $\mathcal{H} =$ space of **wavefunctions** $\Psi : \mathcal{O}_p \rightarrow \mathbb{C} : q \mapsto \Psi(q)$
- ▶ Define $(\mathcal{T}[(f, \alpha)]\Psi)(q) = e^{i\langle q, \alpha \rangle} \Psi(f^{-1} \cdot q)$
- ▶ This is an IRREP of P !

UNITARY REPS

Let's build UIRREPS of P !

- ▶ Pick an orbit \mathcal{O}_p
- ▶ Let $\mathcal{H} =$ space of **wavefunctions** $\Psi : \mathcal{O}_p \rightarrow \mathbb{C} : q \mapsto \Psi(q)$
- ▶ Define $(\mathcal{T}[(f, \alpha)]\Psi)(q) = e^{i\langle q, \alpha \rangle} \Psi(f^{-1} \cdot q)$
- ▶ This is an IRREP of P !

How to make \mathcal{T} unitary?

- ▶ Pick a G -invariant measure μ on \mathcal{O}_p
- ▶ Scalar product $\langle \Phi | \Psi \rangle \equiv \int_{\mathcal{O}_p} d\mu(q) (\Phi(q))^* \Psi(q)$

UNITARY REPS

Let's build UIRREPS of P !

- ▶ Pick an orbit \mathcal{O}_p
- ▶ Let $\mathcal{H} =$ space of **wavefunctions** $\Psi : \mathcal{O}_p \rightarrow \mathbb{C} : q \mapsto \Psi(q)$
- ▶ Define $(\mathcal{T}[(f, \alpha)]\Psi)(q) = e^{i\langle q, \alpha \rangle} \Psi(f^{-1} \cdot q)$
- ▶ This is an IRREP of P !

How to make \mathcal{T} unitary?

- ▶ Pick a G -invariant measure μ on \mathcal{O}_p
- ▶ Scalar product $\langle \Phi | \Psi \rangle \equiv \int_{\mathcal{O}_p} d\mu(q) (\Phi(q))^* \Psi(q)$

All UIRREPS of P are of this form !

[Mackey ~1950]

UNITARY REPS

Let's build UIRREPS of P !

- ▶ Pick an orbit \mathcal{O}_p
- ▶ Let $\mathcal{H} =$ space of **wavefunctions** $\Psi : \mathcal{O}_p \rightarrow \mathbb{C} : q \mapsto \Psi(q)$
- ▶ Define $(\mathcal{T}[(f, \alpha)]\Psi)(q) = e^{i\langle q, \alpha \rangle} \Psi(f^{-1} \cdot q)$
- ▶ This is an IRREP of P !

How to make \mathcal{T} unitary?

- ▶ Pick a G -invariant measure μ on \mathcal{O}_p
- ▶ Scalar product $\langle \Phi | \Psi \rangle \equiv \int_{\mathcal{O}_p} d\mu(q) (\Phi(q))^* \Psi(q)$

All UIRREPS of P are of this form !

[Mackey ~1950]

DEFINITION OF BMS_3

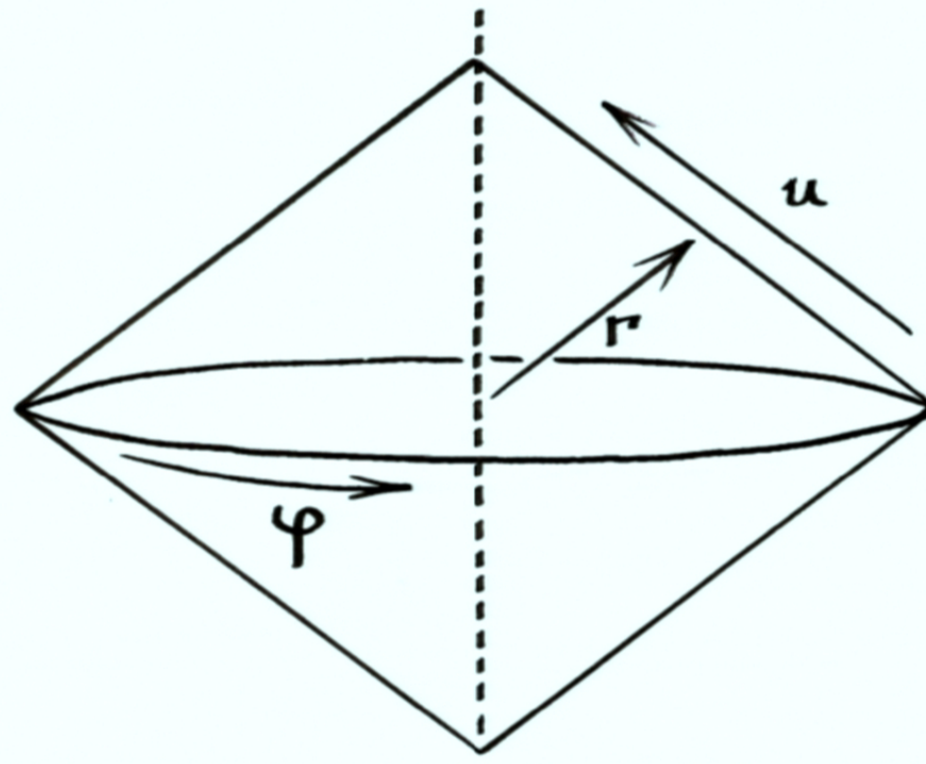
BMS_3 = Aspt. symmetry group of 3D aspt. flat space-times

[Ashtekar *et al.* 1996, Barnich *et al.* 2009]

DEFINITION OF BMS_3

BMS_3 = Aspt. symmetry group of 3D aspt. flat space-times

[Ashtekar *et al.* 1996, Barnich *et al.* 2009]



DEFINITION OF BMS_3

BMS_3 = Aspt. symmetry group of 3D aspt. flat space-times

[Ashtekar *et al.* 1996, Barnich *et al.* 2009]

- ▶ Aspt. flat metrics in 3D :

$$ds^2 \stackrel{r \rightarrow +\infty}{\sim} -du^2 - 2dudr + r^2 d\varphi^2 + \text{subleading terms}$$

- ▶ Aspt. symmetry transformations :

$$\varphi \mapsto f(\varphi), \quad u \mapsto [u + \alpha(\varphi)] f'(\varphi)$$

↙
Superrotations

↓
Supertranslations

DEFINITION OF BMS_3

BMS_3 = Aspt. symmetry group of 3D aspt. flat space-times

[Ashtekar *et al.* 1996, Barnich *et al.* 2009]

- ▶ Aspt. flat metrics in 3D :

$$ds^2 \stackrel{r \rightarrow +\infty}{\sim} -du^2 - 2dudr + r^2 d\varphi^2 + \text{subleading terms}$$

- ▶ Aspt. symmetry transformations :

$$\varphi \mapsto f(\varphi), \quad u \mapsto [u + \alpha(\varphi)] f'(\varphi)$$

↙
Superrotations

↓
Supertranslations

- ▶ Infinite-dimensional extension of Poincaré

DEFINITION OF BMS_3

BMS_3 = Aspt. symmetry group of 3D aspt. flat space-times

[Ashtekar *et al.* 1996, Barnich *et al.* 2009]

- ▶ Aspt. flat metrics in 3D :

$$ds^2 \stackrel{r \rightarrow +\infty}{\sim} -du^2 - 2dudr + r^2 d\varphi^2 + \text{subleading terms}$$

- ▶ Aspt. symmetry transformations :

$$\varphi \mapsto f(\varphi), \quad u \mapsto [u + \alpha(\varphi)] f'(\varphi)$$

↙
Superrotations

↓
Supertranslations

- ▶ Infinite-dimensional extension of Poincaré

DEFINITION OF BMS_3

BMS_3 transformations :

$$\varphi \mapsto f(\varphi), \quad u \mapsto [u + \alpha(\varphi)] f'(\varphi)$$

DEFINITION OF BMS_3

BMS_3 transformations :

$$\varphi \mapsto f(\varphi), \quad u \mapsto [u + \alpha(\varphi)] f'(\varphi)$$

▶ Elements of BMS_3 = pairs (f, α)

▶ Group operation :

$$(f, \alpha) \cdot (g, \beta) = (f \circ g, \alpha + \sigma_f \beta), \quad \sigma_f \beta|_{f(\varphi)} = f'(\varphi) \beta(\varphi)$$

▶ σ = tsf. law of vector fields under diffeos !

▶ $\alpha = \alpha(\varphi) \frac{\partial}{\partial \varphi}$

$$BMS_3 \equiv \text{Diff}(S^1) \ltimes \text{Vect}(S^1)$$

[Barnich & BO 2014]

Unitary reps of BMS_3 ...

- ▶ What should we expect ?

Poincaré : exact space-time symmetry

- ▶ $UIRREP = \mathcal{H}_{\text{Poinc}} = \text{Particle}$

BMS_3 : aspt. space-time symmetry

- ▶ $UIRREP = \mathcal{H}_{BMS} = \text{Particle dressed w/ soft gravitons}$
 \equiv **BMS_3 particle**

- ▶ $\mathcal{H}_{BMS} = \mathcal{H}_{\text{Poinc}} \otimes \mathcal{H}_{\text{Soft grav}}$

ORBITS AND UNITARY REPS

$$\text{Unitary reps of } BMS_3 = \underbrace{\text{Diff}(S^1) \ltimes \text{Vect}(S^1)}_{G \ltimes A} ?$$

- ▶ $\text{Vect}(S^1)^* =$ space of **supermomenta** $p(\varphi)$:

$$\underbrace{\langle p, \alpha \rangle}_{\in \mathbb{R}} = \frac{1}{2\pi} \int d\varphi p(\varphi) \alpha(\varphi)$$

$$p(\varphi) = \sum_{n \in \mathbb{Z}} p_n e^{-in\varphi}$$

- ▶ $p_0 =$ **energy**

ORBITS AND UNITARY REPS

Fix a supermomentum $p(\varphi)$

- ▶ Find all $f \cdot p$, where $f \in \text{Diff}(S^1)$.

$$\langle f \cdot p, \alpha \rangle = \langle p, \sigma_{f^{-1}} \alpha \rangle$$

- ▶ $f \cdot p|_{f(\varphi)} = \frac{1}{(f'(\varphi))^2} \left[p(\varphi) + \frac{c}{12} \{f; \varphi\} \right]$

- ▶ $p(\varphi) \sim$ CFT stress tensor on S^1

ORBITS AND UNITARY REPS

Fix a supermomentum $p(\varphi)$

- ▶ Find all $f \cdot p$, where $f \in \text{Diff}(S^1)$.

$$\langle f \cdot p, \alpha \rangle = \langle p, \sigma_{f^{-1}} \alpha \rangle$$

- ▶ $f \cdot p|_{f(\varphi)} = \frac{1}{(f'(\varphi))^2} \left[p(\varphi) + \frac{c}{12} \{f; \varphi\} \right]$

- ▶ $p(\varphi) \sim$ CFT stress tensor on S^1

- ▶ BMS_3 orbits = orbits of stress tensors under conf. tsfs !
= **coadjoint orbits of the Virasoro group**

[Lazutkin & Pankratova 1975, Witten 1988, Balog *et al.* 1997]

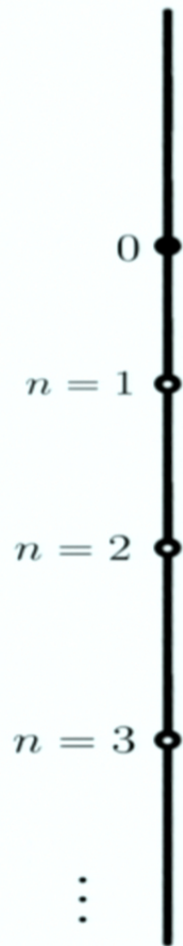
$$f \cdot p|_{f(\varphi)} = \frac{1}{(f'(\varphi))^2} \left[p_0 + \frac{c}{12} \{f; \varphi\} \right]$$

- ▶ $\mathcal{O}_p = \{f \cdot p | f \in \text{Diff}(S^1)\} \rightarrow \text{complicated !}$
- ▶ Let's make it simple :
 1. Constant supermomentum $p(\varphi) = p_0$ (\sim particle at rest)
 2. Look for stabilizer G_p

$$\frac{1}{(f'(\varphi))^2} \left[p_0 + \frac{c}{12} \{f; \varphi\} \right] = p_0$$

- ▶ $\mathcal{O}_p = \{f \cdot p \mid f \in \text{Diff}(S^1)\} \rightarrow \text{complicated !}$
- ▶ Let's make it simple :
 1. Constant supermomentum $p(\varphi) = p_0$ (\sim particle at rest)
 2. Look for stabilizer G_p
 - ▶ $\mathcal{O}_p \cong \text{Diff}(S^1)/G_p$

BMS_3 orbits = coadjoint orbits of Virasoro group :



$\uparrow p_0$

Generic p_0 : $\mathcal{O}_p \cong \text{Diff}(S^1)/S^1$

$p_0 = -n^2c/24$: $\mathcal{O}_p \cong \text{Diff}(S^1)/\text{SO}^{(n)}(2, 1)$

BMS_3 orbits = coadjoint orbits of Virasoro group :



Generic p_0 : $\mathcal{O}_p \cong \text{Diff}(S^1)/S^1$

$p_0 = -n^2c/24$: $\mathcal{O}_p \cong \text{Diff}(S^1)/\text{SO}^{(n)}(2, 1)$

- ▶ What else ?
- ▶ Perturbations $-\frac{n^2c}{24} + \delta p(\varphi)$
- ▶ $\delta p = 3$ -momentum under $\text{SO}^{(n)}(2, 1)$
- ▶ “Poincaré orbits for each n ” !

SUPERMOMENTUM & BONDI MASS

On-shell aspt. flat metrics

SUPERMOMENTUM & BONDI MASS

On-shell aspt. flat metrics :

[Barnich & Troessaert 2010]

$$ds^2 = 8G p(\varphi) du^2 - 2dudr + r^2 d\varphi^2 + \dots$$

$p(\varphi)$ = Bondi mass aspect

▶ Action of BMS_3 on $p(\varphi)$:

$$f \cdot p|_{f(\varphi)} = \frac{1}{f'^2} \left[p + \frac{c}{12} \{f; \varphi\} \right] \quad \text{with } c = 3/G$$

CHARACTERS & PARTITION FUNCTIONS

BMS_3 particle = Particle \otimes Soft gravitons

- ▶ Vacuum BMS_3 character \leftrightarrow graviton partition function ?

Characters of unitary reps of semi-direct products :

- ▶ Orbit \mathcal{O}_p
- ▶ Character :

$$\chi[(f, \alpha)] = \text{Tr} (\mathcal{T}[(f, \alpha)]) = \int_{\mathcal{O}_p} d\mu(q) e^{i\langle q, \alpha \rangle}$$

CHARACTERS & PARTITION FUNCTIONS

Massive BMS_3 particle

▶ $p = p_0$



CHARACTERS & PARTITION FUNCTIONS

Massive BMS_3 particle

- ▶ $p = p_0 \rightarrow \mathcal{O}_p = \text{Diff}(S^1)/S^1$

Take $f(\varphi) = \varphi + \theta$ (rotation by θ)

- ▶ Character :

$$\chi[(\text{rot}_\theta, \alpha)] = \int_{\mathcal{O}_p} d\mu(q) \delta(q, \text{rot}_\theta \cdot q) e^{i\langle q, \alpha \rangle}$$

$$(\text{rot}_\theta \cdot q)(\varphi) = q(\varphi - \theta)$$

- ▶ The integral “localizes” to a point !

CHARACTERS & PARTITION FUNCTIONS

$$\chi_{p_0}[(\mathbf{rot}_\theta, \alpha)] = e^{ip_0\alpha^0} \frac{1}{\prod_{n=1}^{+\infty} |1 - e^{in(\theta+i\epsilon)}|^2}$$

[BO 2015]

CHARACTERS & PARTITION FUNCTIONS

$$\chi_{\text{vac}}[(\text{rot}_\theta, \alpha)] = e^{\beta c/24} \frac{1}{\prod_{n=2}^{+\infty} |1 - e^{in(\theta+i\epsilon)}|^2} \quad [\text{BO 2015}]$$

- ▶ One-loop partition fct of gravitons on thermal flat space !
[Barnich, González, Maloney, BO 2015]

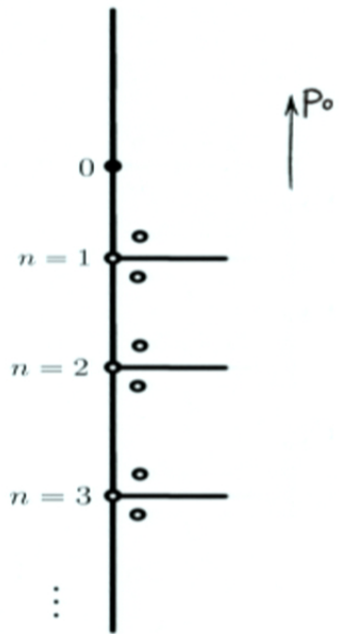
CONCLUSION

$$BMS_3 = \text{Superrotations} \times \text{Supertranslations}$$

CONCLUSION

$BMS_3 = \text{Superrotations} \times \text{Supertranslations}$

- ▶ “Supermomenta”
- ▶ UIRREPs classified by supermomentum orbits



Thank you !

