

Title: Mysterious Magnetars: Maximum Stars

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Abstract:

Magnetars are exceptional neutron stars with the highest magnetic fields (10^{15} gauss) in the universe, an unusual quasi steady X radiation (10^{35} ergs/sec) and also produce flares which are some of the brightest events (10^{46} ergs in one fifth of a second) to be recorded. There is no satisfactory model of magnetars.

The talk will cover neutron stars and a new model for the origin of the magnetic fields in which magnetars arise from a high baryon density (phase transition) magnetized core which forms when they are born. The core magnetic field is initially shielded by the ambient high conductivity plasma. With time the shielding currents dissipate transporting the core field out, first to the crust and then breaking through the crust to the surface of the star. Recent observations provide support for this model which accounts for several properties of magnetars and also enables us to identify new magnetars.

Report

The explosion on the surface of a magnetar, SGR 1806-20, reached Earth from 50,000 light years away.

It was the brightest event known to have impacted this planet from an origin outside our solar system.

The magnetar, which is no more than twenty kilometers (twelve miles) in diameter, released more energy in one-tenth of a second than our sun has released in 100,000 years. A similar blast within ten light years of earth would destroy the ozone layer.

Fortunately, the nearest known magnetar to earth is 1E 2259+586, 13,000 light years (120 exameters) away. SGR 1806-20 is located in the constellation Sagittarius.

Neutron Star Profile

- **Density decreases as one moves outwards from the centre.**
- The outer kilometre or so is the **Crust**. It consists of a lattice of bare nuclei (starting with Fe 56) and a degenerate electron gas. They get neutron rich as we go in, till we go into a neutron drip region and then
- Next to the crust is a **Neutron, Proton Electron..** Plasma in beta equilibrium /
- With POSSIBLY a neutron superfluid and vortices here that contribute to the angular momentum of the star. These vortices also play an important role in the so called glitches whence the star actually speeds up.
- The least understood part is the **core** of the NS. Density here can be in the range of 3-10 times nuclear density. The core can be in one of the many exotic phases of high density baryonic or quark matter.

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Neutron Star Properties

- Radius ~ 10 KM
- Density 10^{15} gm/cc ~ 1baryon / Fermi³
- The electron densities in the core are very high ' 10^{36} /cc and the electrical conductivity is ' 10^{28} (CGS).
- **Birth**
- Pulsar Fields
- The progenitor star (Radius ~ 10^6 km) typically has surface magnetic fields of 1-10 Gauss.
- Collapse : High conductivity -> Flux conservation ($\sim 4 \pi R^2 B$) amplifies surface B ~ 10^{12} (11) Gauss.
- We shall call these fossil or inherited fields.
- Newly born Neutron stars are rapidly rotating with periods of tens of milliseconds.

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Magnetic Dipole Radiation

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$$\frac{2}{3} \omega^4 \frac{B^2 R^6}{c^3}$$

- The rate of rotational energy loss

$$\dot{E}_R = I\omega\dot{\omega}$$

- Spin-down rate

$$\dot{P} = \frac{24\pi^2 B^2 R^6}{3 P c^3 I} \quad \tau_{SD} = \frac{P}{2\dot{P}}$$

- An important relation (using $R = 10 \text{ Km}$, $I = 10^{45} \text{ gmcm}^2$)

$$P\dot{P} = 9.75 \cdot 10^{-38} B^2$$

Navigation icons: back, forward, search, etc.

Magnetars: A New Class of Neutron Stars

- Their surface magnetic fields are much larger than those of pulsars i.e ' 10^{14} – 10^{15} G. These are the largest ever observed fields.
 - They have with periods $P \sim 5\text{ s} - 12\text{ s}$. (>pulsars)
 - Unlike Pulsars they have strong steady X-ray emission :
 - 10^{34} ergs/sec.
 - Unlike Pulsars they exhibit bursts of energy called flares:
 - $10^{41} - 10^{46}$ ergs
- Unlike pulsars (10^6 yr) Magnetar Ages - 10^4 years
- Magnetars progenitors are more massive than Pulsars.
- Anomalous X-ray Pulsars (AXP) have small flares.
 - Soft Gamma Repeaters (SGR) have large flares.

What supplies the energy in flares and steady X-emission?

- Pulsars are powered by Dipole Emission which comes from Rotational energy loss.
- For magnetars with periods of ~ 10 sec the rate of rotational energy loss is too meagre give the steady Xray flux (10^{35} erg/sec).
- Duncan & Thomson suggested that the source is instead the magnetic energy They suggested the Dynamo Mechanism. If stars are born rotating fast ($P \sim 1$ millisec) convection etc can lead to a field amplification from 10^{12} to 10^{15} Gauss

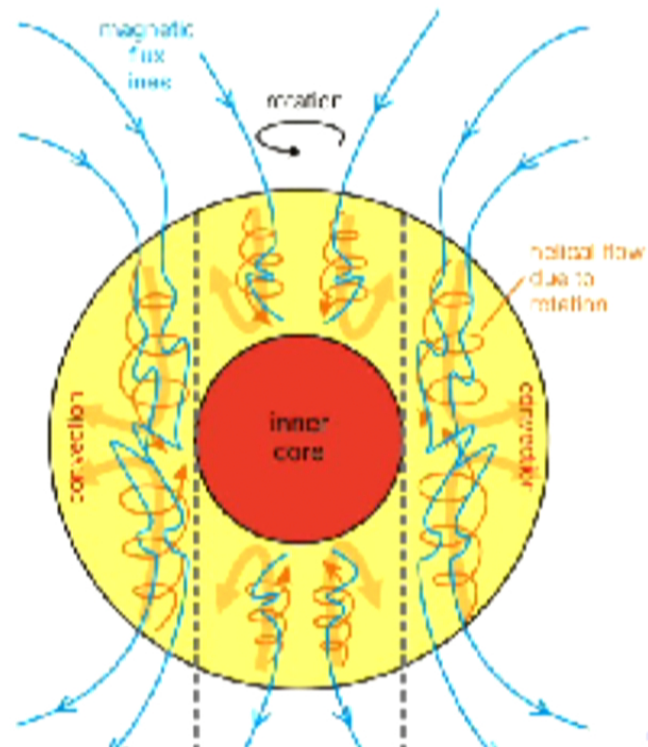
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Dynamos.



N.D. Hari Dass

Magnetars & Pion Condensates

Difficulties with the Dynamo Mechanism

- To amplify pulsar fossil fields to Magnetar fields, P must be of the order of a **1** millisecc.
- This requires **massive progenitors** which supports the fact that magnetars are more massive. But
- **Simulations (Heger et al Ap J 626 2005) and Explosion energies (J. Vink and L. Kuiper, MNRAS 370, 2006))**
- show that even with very massive progenitors its hard to get $P < \sim 5$ millisecc.
- In the dynamo model, magnetars acquire their large fields almost immediately after their birth and magnetic dipole radiation starts promptly. This implies that the spin-down age $T = P / (2dP/df)$ is a good measure of the age of a magnetar.
BUT

Observational Evidence for Age Crisis.

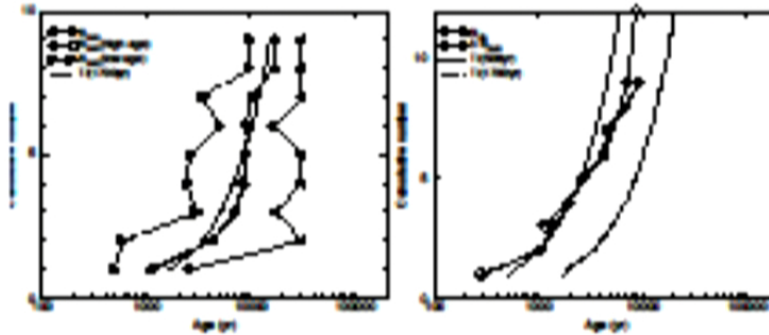


Figure: Leahy Analysis for Age Crisis.

Dynamo Model - More Difficulties

- THE ENERGY CRISIS (A Dar and A De Rijula 2000etc)
- As per the Dynamo Model the **energetics** of the powerful flares as well of the steady X-emissions is **powered by the magnetic field**.
- If so, the surface magnetic fields of magnetars must **decrease** with time.
- So should dP/dt .
- Not only is this not seen in Magnetars, there is even evidence that the opposite happens i.e the surface magnetic fields and P' show phases where they **increase!**
- (Kulkarni etal Ap J)
- This last point has been demonstrated by us by analysing the very precise **timing data of Kaspi et al (arXiv (astro) 0905.3567)** in the post-glitch phase of the **young pulsar J1846-0258**.

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THE MAGNETIZED CORE

This introduces an additional term in the strong interaction hamiltonian
$$- \Sigma \cdot \mathbf{q} \tau_3$$
where \mathbf{q} is the Pion wave-vector, Σ the neutron spin-operator and τ_3 the 3 component of the isospin.

This shows that in the neutral pion condensate phase it becomes energetically favourable for all neutron (or u quark) spins, and

hence all the magnetic moments, to align antiparallel (parallel) to \mathbf{q} .

(F. Dautry and E Nyman , Nuc. Phy. A319 1979 for Nucleon matter
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Note that the minimization yields the a condensate wave vector q , that is proportional to the baryon density, n .

The negative condensate term is proportional to, n^2 , whereas the kinetic energy term for the nucleons goes as, $n^{5/3}$. As the density goes up the negative condensate term dominates leading to a phase transition to the neutral pion condensed groundstate.

The strong interaction Phase Transition to the neutral pion condensate brings down the energy per baryon, E_B , by 20~30 Mev/ baryon

Estimate of the aligned magnetic moments

- At 5 times nuclear density of 10^{15} gm/cm³ there are some 10^{38} neutrons per cc.
- Each has a magnetic moment
 $\mu_n = 2 \cdot eh/2m_n c = 10^{-23}$ (CGS).
- This translates to a magnetic moment density
- $m = 6 \cdot 10^{15}$, and a uniform core field of $5 \cdot 10^{16}$ G.
- We get a similar number for constituent quark matter
- Unlike magnetic fields produced by currents this field is virtually indestructible.

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Time evolution of Magnetic Fields in Our Model – Magnetized Core.

- As the core magnetic moment and the core field builds up, (eg islands in a first order transition) it gets oriented by the inherited/ fossil field
- **The high conductivity plasma generates Lenz currents to shield this field.**
- The core field will remain hidden. Till these currents dissipate there will be no surface magnetic field apart from the fossil field.
- As the shielding currents dissipate the magnetic field in the
- interior starts increasing till it reaches the fully relaxed
- configuration due to the core moment produced by strong
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ENERGY BUDGET

- I) The strong interaction Phase Transition to the neutral pion condensate brings down the energy per baryon, E_B , by ~ 30 MeV/ baryon,
 - So for a core of $R \sim 2\text{km}$ and density 10^{39} baryons/cc
- Net energy reduction is 10^{50} ergs
- II) Magnetic energy of the, $R \sim 2\text{km}$, core with $B \sim 10^{16}$ gauss
 - $E_m \sim 10^{47}(48)$ ergs
- III) Net energy IN SHIELDING CURRENTS $> E_m$
- IV) Net energy from X-ray emission of 10^{36} ergs per second for $\sim 10^4$ yr and 100 flares of 10^{46} ergs $\sim 10^{47}$ ergs
- Neutrino emission $10^{47}(48)$ ergs
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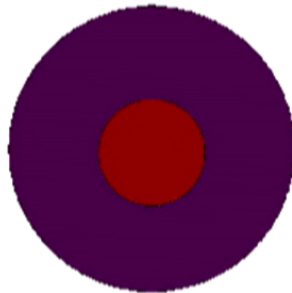


Figure: Early Magnetic Field Configuration.

N.D. Hart Data Magnetars & Plan Condensates

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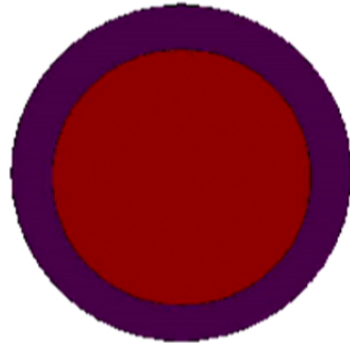


Figure: Late Magnetic Field Configuration.

N.D. Hart Dass Magneters & Pion Condensates

Time Scales of Field Evolution.

- How exactly the shielding currents dissipate is an extremely interesting problem.
- The time scale for **Ohmic Dissipation** is estimated at 10^{11} years and is therefore totally irrelevant.
- There is a dissipation mechanism called **Ambipolar Diffusion** whose time scale is given by

$$t_{amb} \simeq 10^4 B_{16}^{-2} T_{8.5}^{-6}$$

years. This, for the likely parameters of our model, works out to about 10,000 years.

- This nicely matches the offset characterising the age crisis.
- The physics of ambipolar diffusion involves plasma physics, beta equilibrium, condensed matter physics of strongly magnetised media etc...
- During the ambipolar dissipation there will be energy loss from the star in the form of neutrinos, thermal X-rays etc..

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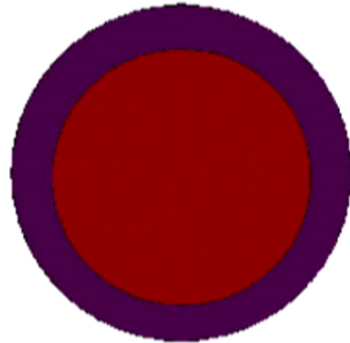


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N.D. Hart Dass Magneters & Pion Condensates

The phases

- Phase I
- Via ambipolar diffusion (P. Goldreich and A Reisenegger, Ap J 395
- 1992) the relaxed magnetic field reaches the inner crust.
- Till then plasma does not allow photon propagation . $T >$ till Neutrino emission is efficient. In this phase the magnetar has no significant EM activity like spin down
- Observational Support II
- In our model THIS IS THE REASON FOR DELAYED AMPLIFICATION AND THE
- RESOLUTION OF THE AGE CRISIS BETWEEN SPIN DOWN AGE AND SNOVA REMNANT (SNR) AGE
- Phase II
- Once the outgoing field reaches the crust the field gradients will be too
- large for the crust to sustain the magnetic stresses.
- Observational Support III
- The crust starts cleaving and field lines will start to leak out giving rise to the phase where **the surface field starts increasing.**
- **X Radiation and Flares**

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Lx and \dot{E}_R

- Magnetic Energy in the crust shell , $\Delta R = 1 \text{ Km}$
- $E_m \sim 4 \pi R^2 \Delta R B^2$ $R \sim 10^6 \text{ cm}$
- Spin down age $\zeta \sim P / 2 (dP/dt)$
- $dP/dt \sim B^2$
- $E_m / \zeta \sim 2 \times 10^{23} B^4$
- $B \sim 4 \times 10^{14} \text{ g}$ **$L_x \sim E_m / \zeta \sim 5 \times 10^{35} \text{ erg/sec}$**
- $\dot{E}_R \sim 4 \pi^2 I (dP/dt) / (P^3)$
- **$L_x / \dot{E}_R \sim 10 \sim (B_{13})^2 P^2 / 1000$**

T.S. Wood and R. Hollerbach arxiv.org/pdf/1501.05149(2015)

K. N. Gourgouliatos and A. Cumming, MNRAS 438, 1618 (2014)

Their model uses ohmic dissipation and hall drift

**$L_x \sim B \sim 0.4 \times 10^{14} \text{ g}$ OUTSIDE CRUST $\rightarrow 0$
INSIDE CRUST**

**GOES TO HALF ITS VALUE IN 0.4 MILLION
YEARS**

$L_x \sim 7.5 \cdot 10^{(32)} \text{ ergs/sec}$

Magnetar periods $< 12\text{sec}$

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II PREDICTING NEW YOUNG MAGNETARS

- **Direct Evidence for Field Increase**
- **The object J1846-0258 has been conventionally considered to be a rotation powered pulsar.**
- **Its $P = 326\text{ms}$.**
- **$B = 5 \cdot 10^{13}\text{ G}$.**
- **$\dot{E}_R > L_x$**
- **Are all counter to magnetars**
- **But it has a high X-ray luminosity more characteristic of magnetars.**
- **Recently **Gavriil et al (Science 319 2008)** et al reported a glitch accompanied by a flare in this object.**

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Post-glitch Timing Analysis of J1846-0258

- Livingstone, Kaspi et al (2010) carried out a very precise timing analysis of this object in the post-glitch period.
- Even after 920 days there was a persistent change in
- $\Delta\dot{\omega} / \dot{\omega} = 5.5 \cdot 10^{-2}$
- which translates via the dipole formula to an increase in magnetic field of about $\Delta B = 0.13 \cdot 10^{13}$ Gauss Or 1% per year.
- **Anomaly**
- **This effect is 10-1000 times in magnitude than in any pulsars**
- More evidence for a potential persistent increase in B is provided by the observed reduction in the braking index for this star, from , **n= 2.65 to 2.16**, after the FLARE, by Livingstone et al (arXiv 2010). Only plausible source is
- **$dB/dt > 0$**