

Title: 2+1D topological orders and braided fusion category

Date: Oct 22, 2015 09:00 AM

URL: <http://pirsa.org/15100105>

Abstract:

# 2+1D topological orders and braided fusion category

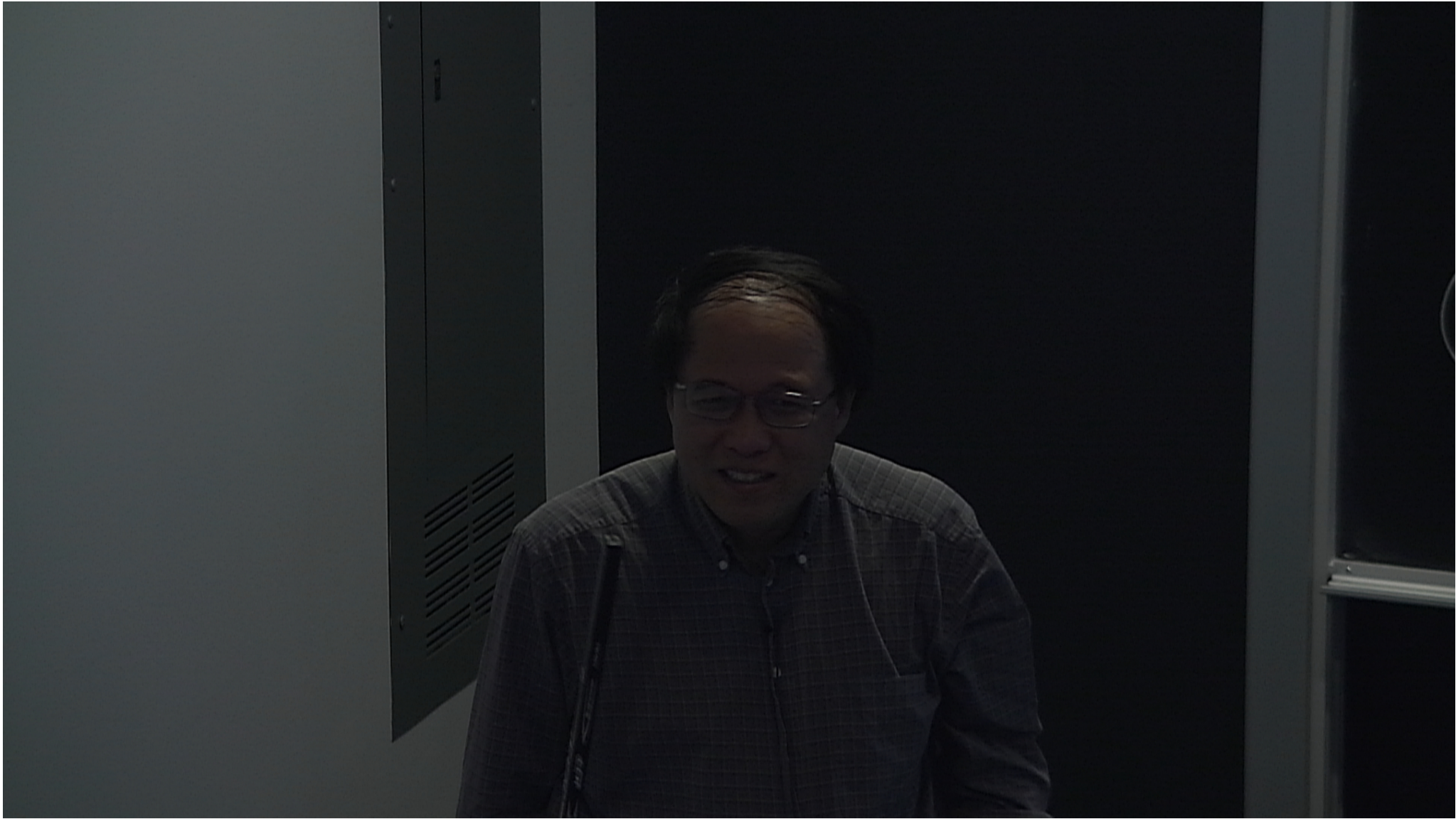
Xiao-Gang Wen  
Oct., 2015

There are gapped phases beyond symmetry breaking order  
→ **topological orders**



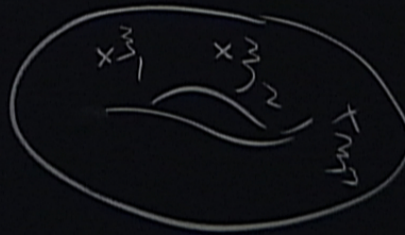
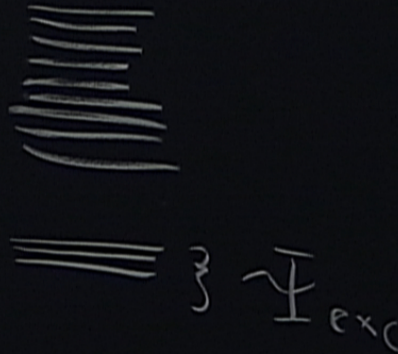
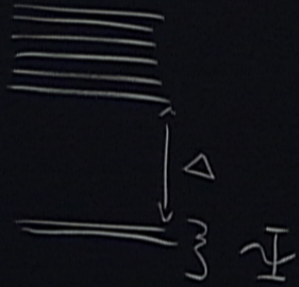
Xiao-Gang Wen Oct., 2015

2+1D topological orders and braided fusion category



H

$$H + \delta H(\vec{z}_1) + \delta H(\vec{z}_2)$$



# 2+1D topological orders and braided fusion category

Xiao-Gang Wen  
Oct., 2015

There are gapped phases beyond symmetry breaking order  
→ **topological orders**



Xiao-Gang Wen Oct., 2015

2+1D topological orders and braided fusion category

## Examples of topo. orders – long range entanglement (LRE)

**Abelian topological order:**  $\rightarrow$  fractional statistics  $N \rightarrow \infty$

- IQH and Laughlin

$$\Psi_{\nu=1}^F = \prod_{1 \leq i < j \leq N} (z_i - z_j) e^{-\frac{1}{4} \sum |z_i|^2}, \quad \Psi_{\nu=1/m}^{F,B} = \prod (z_i - z_j)^m e^{-\frac{m}{4} \sum |z_i|^2}$$

Laughlin PRL 50 1395 (1983)

$$= (\Psi_{\nu=1}^F)^m$$

where  $z_i = x_i + iy_i$ .

## Examples of topo. orders – long range entanglement (LRE)

**Abelian topological order:** → fractional statistics  $N \rightarrow \infty$

- IQH and Laughlin **many-body** state Laughlin PRL 50 1395 (1983)  

$$\Psi_{\nu=1}^F = \prod_{1 \leq i < j \leq N} (z_i - z_j) e^{-\frac{1}{4} \sum |z_i|^2}, \quad \Psi_{\nu=1/m}^{F,B} = \prod (z_i - z_j)^m e^{-\frac{m}{4} \sum |z_i|^2}$$

$$\in \mathcal{V}_{N=0} \oplus \mathcal{V}_{N=1} \cdots = \otimes_z \mathcal{V}_z^{\{|0\rangle, |1\rangle\}} = (\Psi_{\nu=1}^F)^m$$
 where  $z_i = x_i + iy_i$ .  $\Psi_{\nu=n}^{F,IQH}$  has an invertible topological order

- $\Psi^{B/F}$  symmetric/anti-symmetric → bosonic/fermionic topo. order;  
Wen PRB 40, 7387 (89)

**Non-abelian topological order:** → non-abelian statistics

- $SU(N)_2$  state via slave-particle Wen PRL 66 802 (Feb. 1991)  

$$\Psi_{SU(2)_2}^B = (\Psi_{\nu=2}^F)^2, \nu = 1; \quad \Psi_{SU(3)_2}^F = (\Psi_{\nu=2}^F)^3, \nu = \frac{2}{3};$$

→  $SU(N)_2$  Chern-Simons effective theory → non-abelian statistics

- Pfaffian state via CFT Moore-Read NPB 360 362 (Aug. 1991)

$$\Psi_{Pfa}^B = \mathcal{A} \left[ \frac{1}{z_1 - z_2} \frac{1}{z_3 - z_4} \cdots \right] \prod (z_i - z_j) e^{-\frac{1}{4} \sum |z_i|^2}, \quad \nu = 1$$

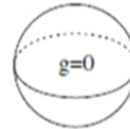
- The Pfaffian and  $SU(2)_2$  have the same non-abelian statistics
- The  $SU(3)_2$  state has the Fibonacci non-abelian statistics

# Topological invariants that define LRE and topo. orders

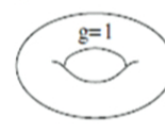
(1)  $\Psi$  = space of **locally indistinguishable (LI)** vectors  $\subset \otimes \mathcal{V}_z$

- Given vector  $\Psi_1$ ,  $\exists$  other LI  $\Psi_2, \dots$ .

- Topo. degeneracy  $D_g \equiv \dim \Psi$ ,  
depends on topology of space



Deg.=1



Deg.=D<sub>1</sub>



Wen PRB 40, 7387 (89), Wen-Niu PRB 41, 9377 (90)



## Examples of topo. orders – long range entanglement (LRE)

**Abelian topological order:** → fractional statistics  $N \rightarrow \infty$

- IQH and Laughlin **many-body** state Laughlin PRL 50 1395 (1983)  

$$\Psi_{\nu=1}^F = \prod_{1 \leq i < j \leq N} (z_i - z_j) e^{-\frac{1}{4} \sum |z_i|^2}, \quad \Psi_{\nu=1/m}^{F,B} = \prod (z_i - z_j)^m e^{-\frac{m}{4} \sum |z_i|^2}$$

$$\in \mathcal{V}_{N=0} \oplus \mathcal{V}_{N=1} \cdots = \otimes_z \mathcal{V}_z^{\{|0\rangle, |1\rangle\}} = (\Psi_{\nu=1}^F)^m$$

where  $z_i = x_i + iy_i$ .  $\Psi_{\nu=n}^{F,IQH}$  has an invertible topological order

- $\Psi^{B/F}$  symmetric/anti-symmetric → bosonic/fermionic topo. order;  
Wen PRB 40, 7387 (89)

**Non-abelian topological order:** → non-abelian statistics

- $SU(N)_2$  state via slave-particle Wen PRL 66 802 (Feb. 1991)  

$$\Psi_{SU(2)_2}^B = (\Psi_{\nu=2}^F)^2, \nu = 1; \quad \Psi_{SU(3)_2}^F = (\Psi_{\nu=2}^F)^3, \nu = \frac{2}{3};$$

→  $SU(N)_2$  Chern-Simons effective theory → non-abelian statistics

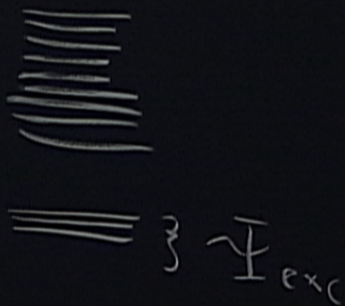
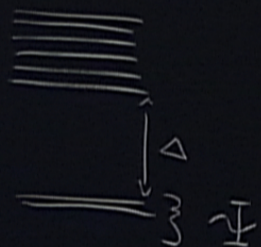
- Pfaffian state via CFT Moore-Read NPB 360 362 (Aug. 1991)

$$\Psi_{Pfa}^B = \mathcal{A} \left[ \frac{1}{z_1 - z_2} \frac{1}{z_3 - z_4} \cdots \right] \prod (z_i - z_j) e^{-\frac{1}{4} \sum |z_i|^2}, \quad \nu = 1$$

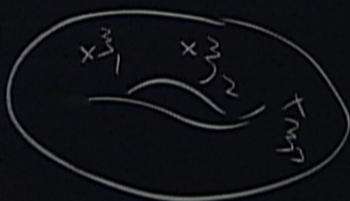
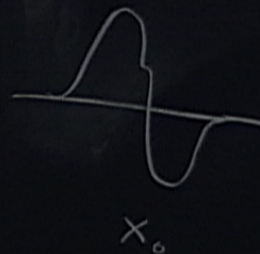
- The Pfaffian and  $SU(2)_2$  have the same non-abelian statistics
- The  $SU(3)_2$  state has the Fibonacci non-abelian statistics

H

$$H + \delta H(x_1) + \delta H(x_2)$$



$$\sum_i f(x_i - x_0, y_i - y_0)$$

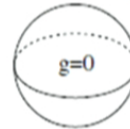


# Topological invariants that define LRE and topo. orders

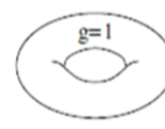
(1)  $\Psi$  = space of **locally indistinguishable (LI)** vectors  $\subset \otimes \mathcal{V}_Z$

- Given vector  $\Psi_1$ ,  $\exists$  other LI  $\Psi_2, \dots$ .

- Topo. degeneracy  $D_g \equiv \dim \Psi$ ,  
depends on topology of space



Deg.=1



Deg.=D<sub>1</sub>



Wen PRB 40, 7387 (89), Wen-Niu PRB 41, 9377 (90)

(2) **Vector bundle on the moduli space**

**i.** Consider a torus  $\Sigma_1$  w/ metrics  $g_{ij}$ . **ii.** Different metrics  $g_{ij}$  form the moduli space  $\mathcal{M} = \{g_{ij}\}$ . **iii.** The LI states depend on spacial metrics:  $\Psi_\alpha(g_{ij}) \rightarrow$  a vector bundle over  $\mathcal{M}$  with fiber  $\Psi_\alpha(g_{ij})$ .

## Examples of topo. orders – long range entanglement (LRE)

**Abelian topological order:** → fractional statistics  $N \rightarrow \infty$

- IQH and Laughlin **many-body** state Laughlin PRL 50 1395 (1983)  

$$\Psi_{\nu=1}^F = \prod_{1 \leq i < j \leq N} (z_i - z_j) e^{-\frac{1}{4} \sum |z_i|^2}, \quad \Psi_{\nu=1/m}^{F,B} = \prod (z_i - z_j)^m e^{-\frac{m}{4} \sum |z_i|^2}$$

$$\in \mathcal{V}_{N=0} \oplus \mathcal{V}_{N=1} \cdots = \otimes_z \mathcal{V}_z^{\{|0\rangle, |1\rangle\}} = (\Psi_{\nu=1}^F)^m$$

where  $z_i = x_i + iy_i$ .  $\Psi_{\nu=n}^{F,IQH}$  has an invertible topological order

- $\Psi^{B/F}$  symmetric/anti-symmetric → bosonic/fermionic topo. order;  
Wen PRB 40, 7387 (89)

**Non-abelian topological order:** → non-abelian statistics

- $SU(N)_2$  state via slave-particle Wen PRL 66 802 (Feb. 1991)  

$$\Psi_{SU(2)_2}^B = (\Psi_{\nu=2}^F)^2, \nu = 1; \quad \Psi_{SU(3)_2}^F = (\Psi_{\nu=2}^F)^3, \nu = \frac{2}{3};$$

→  $SU(N)_2$  Chern-Simons effective theory → non-abelian statistics

- Pfaffian state via CFT Moore-Read NPB 360 362 (Aug. 1991)

$$\Psi_{Pfa}^B = \mathcal{A} \left[ \frac{1}{z_1 - z_2} \frac{1}{z_3 - z_4} \cdots \right] \prod (z_i - z_j) e^{-\frac{1}{4} \sum |z_i|^2}, \quad \nu = 1$$

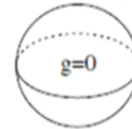
- The Pfaffian and  $SU(2)_2$  have the same non-abelian statistics
- The  $SU(3)_2$  state has the Fibonacci non-abelian statistics

# Topological invariants that define LRE and topo. orders

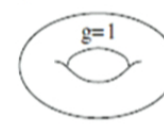
(1)  $\Psi$  = space of **locally indistinguishable (LI)** vectors  $\subset \otimes \mathcal{V}_Z$

- Given vector  $\Psi_1$ ,  $\exists$  other LI  $\Psi_2, \dots$ .

- Topo. degeneracy  $D_g \equiv \dim \Psi$ ,  
depends on topology of space



Deg.=1



Deg.=D<sub>1</sub>



Wen PRB 40, 7387 (89), Wen-Niu PRB 41, 9377 (90)

(2) **Vector bundle on the moduli space**

i. Consider a torus  $\Sigma_1$  w/ metrics  $g_{ij}$ . ii. Different metrics  $g_{ij}$  form the moduli space  $\mathcal{M} = \{g_{ij}\}$ . iii. The LI states depend on spacial metrics:  $\Psi_\alpha(g_{ij}) \rightarrow$  a vector bundle over  $\mathcal{M}$  with fiber  $\Psi_\alpha(g_{ij})$ .

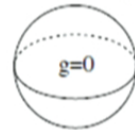
- Local curvature detects grav. Chern-Simons term  $e^{i \frac{2\pi c}{24} \int_{M^2 \times S^1} \omega_3}$

# Topological invariants that define LRE and topo. orders

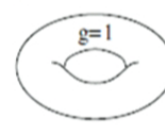
(1)  $\Psi$  = space of **locally indistinguishable (LI)** vectors  $\subset \otimes \mathcal{V}_Z$

- Given vector  $\Psi_1$ ,  $\exists$  other LI  $\Psi_2, \dots$ .

- Topo. degeneracy  $D_g \equiv \dim \Psi$ ,  
depends on topology of space



Deg.=1



Deg.=D<sub>1</sub>



Wen PRB 40, 7387 (89), Wen-Niu PRB 41, 9377 (90)

(2) **Vector bundle on the moduli space**

i. Consider a torus  $\Sigma_1$  w/ metrics  $g_{ij}$ . ii. Different metrics  $g_{ij}$  form the moduli space  $\mathcal{M} = \{g_{ij}\}$ . iii. The LI states depend on spacial metrics:  $\Psi_\alpha(g_{ij}) \rightarrow$  a vector bundle over  $\mathcal{M}$  with fiber  $\Psi_\alpha(g_{ij})$ .

• Local curvature detects grav. Chern-Simons term  $e^{i \frac{2\pi c}{24} \int_{M^2 \times S^1} \omega_3}$

• Loops  $\pi_1(\mathcal{M}) = SL(2, \mathbb{Z})$ :  $90^\circ$  rotation  $|\Psi_\alpha\rangle \rightarrow |\Psi'_\alpha\rangle = S_{\alpha\beta} |\Psi_\beta\rangle$

Dehn twist:  $|\Psi_\alpha\rangle \rightarrow |\Psi'_\alpha\rangle = T_{\alpha\beta} |\Psi_\beta\rangle$  =

$S, T$  generate a rep. of modular group:  $S^2 = (ST)^3 = C, C^2 = 1$

Wen IJMPB 4, 239 (90); KeskiVakkuri-Wen IJMPB 7, 4227 (93)

## Classify 2+1D topo. orders (ie patterns of entanglement)

via the topological invariants  $(S, T, c)$

- A 2+1D topological order  $\rightarrow$  a  $(S, T, c)$
- An arbitrary  $(S, T, c) \not\rightarrow$  a 2+1D topological order
  
- $(S, T, c)$ 's satisfying **a set of conditions**  $\leftrightarrow$  2+1D topo. orders

# Bosonic invertible topological orders

**invertible topological orders** have no topological bulk excitation (particle-like, string-like, ...), but have non-trivial boundary with **gauge/gravitational anomalies**.  $(S, T, c) = (1, 1, c)$ .

- SPT orders  $\subset$  invertible topological orders with symmetry  $G$ .
- Construct SPT orders or more general invertible topological orders:

$$S = \int d^d x \lambda |(\partial + iA + i\Gamma)g|^2 + i2\pi \int W[(\partial + iA + i\Gamma)g]$$

$$g \in G \times SO(\infty)$$

→ SPT orders are labeled by  $\mathcal{H}^d[G \times SO(\infty), U(1)] / \sim$ .

$G \setminus d =$	0+1	1+1	2+1	3+1	4+1	5+1	6+1
$iTO^d$	0	0	$\mathbb{Z}$	0	$\mathbb{Z}_2$	0	0
$\mathbb{Z}_n$	$\mathbb{Z}_n$	0	$\mathbb{Z}_n$	0	$\mathbb{Z}_n \oplus \mathbb{Z}_n$	$\mathbb{Z}_{\langle n,2 \rangle}$	$\mathbb{Z}_n \oplus \mathbb{Z}_n \oplus \mathbb{Z}_{\langle n,2 \rangle}$
$\mathbb{Z}_2^T$	0	$\mathbb{Z}_2$	0	$\mathbb{Z}_2 \oplus \mathbb{Z}_2$	0	$\mathbb{Z}_2 \oplus 2\mathbb{Z}_2$	$\mathbb{Z}_2$
$U(1)$	$\mathbb{Z}$	0	$\mathbb{Z}$	0	$\mathbb{Z} \oplus \mathbb{Z}$	0	$\mathbb{Z} \oplus \mathbb{Z} \oplus \mathbb{Z}_2$
$U(1) \times \mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z} \oplus \mathbb{Z}_2$	$\mathbb{Z}_2$	$2\mathbb{Z}_2 \oplus \mathbb{Z}_2$	$2\mathbb{Z}_2 \oplus 2\mathbb{Z}_2$	$\mathbb{Z} \oplus 2\mathbb{Z}_2 \oplus \mathbb{Z} \oplus 2\mathbb{Z}_2$
$U(1) \times \mathbb{Z}_2^T$	0	$2\mathbb{Z}_2$	0	$3\mathbb{Z}_2 \oplus \mathbb{Z}_2$	0	$4\mathbb{Z}_2 \oplus 3\mathbb{Z}_2$	$2\mathbb{Z}_2 \oplus \mathbb{Z}_2$
$U(1) \times \mathbb{Z}_2^T$	$\mathbb{Z}$	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$2\mathbb{Z}_2 \oplus \mathbb{Z}_2$	$\mathbb{Z} \oplus \mathbb{Z}_2 \oplus \mathbb{Z}$	$2\mathbb{Z}_2 \oplus 2\mathbb{Z}_2$	$2\mathbb{Z}_2 \oplus 3\mathbb{Z}_2 \oplus \mathbb{Z}_2$

Wen arXiv:1410.8477



## Classify 2+1D topo. orders (ie patterns of entanglement)

via the topological invariants  $(S, T, c)$

- A 2+1D topological order  $\rightarrow$  a  $(S, T, c)$
- An arbitrary  $(S, T, c) \not\rightarrow$  a 2+1D topological order
  
- $(S, T, c)$ 's satisfying **a set of conditions**  $\leftrightarrow$  2+1D topo. orders

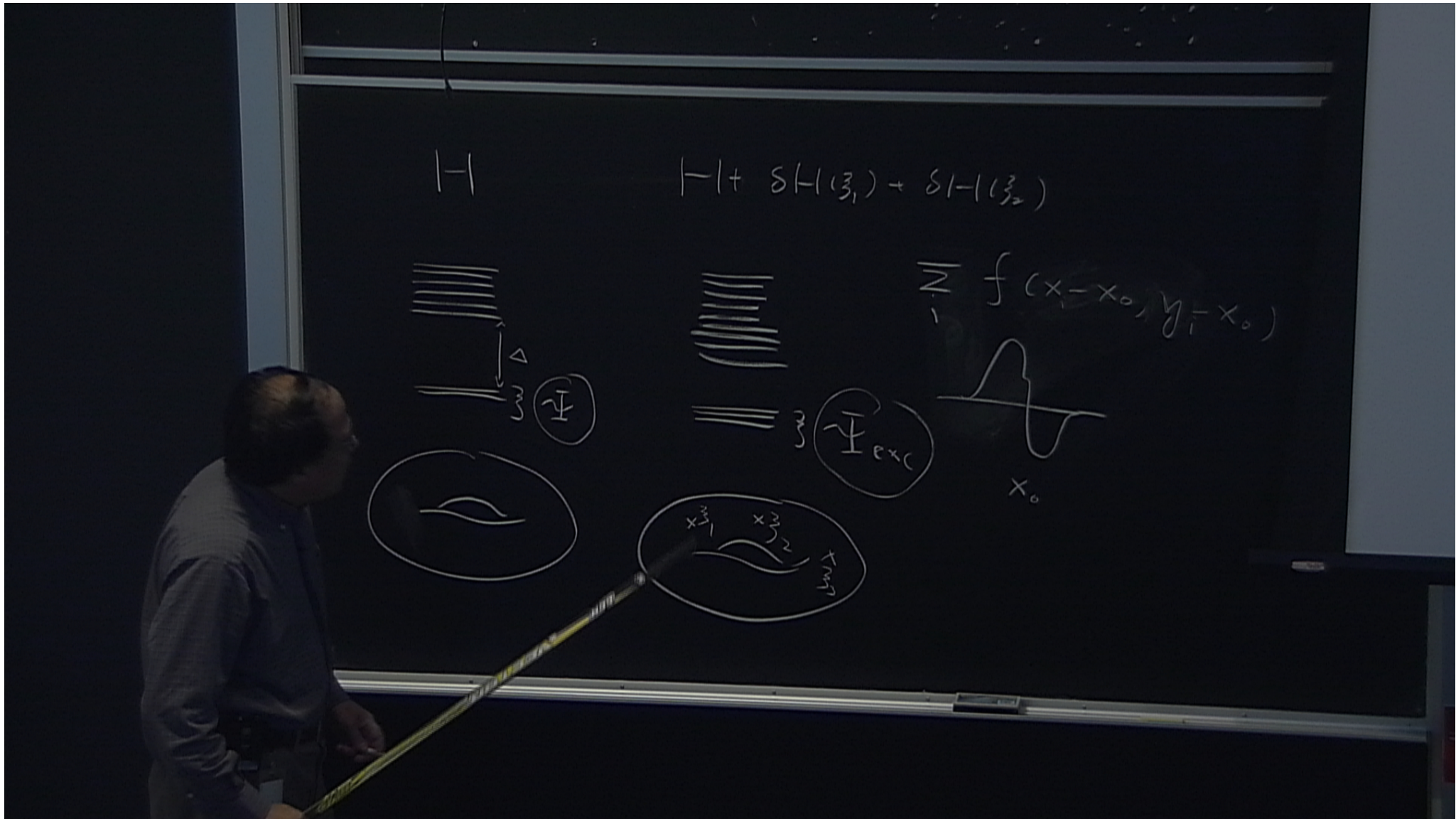
## Classify 2+1D topo. orders (ie patterns of entanglement)

via the topological invariants  $(S, T, c)$

- A 2+1D topological order  $\rightarrow$  a  $(S, T, c)$
- An arbitrary  $(S, T, c) \not\rightarrow$  a 2+1D topological order
- $(S, T, c)$ 's satisfying **a set of conditions**  $\leftrightarrow$  2+1D topo. orders assuming each  $(S, T, c) \rightarrow$  one topological order, otherwise  $(S, T, c)$ 's satisfying **a set of conditions**  $\leftrightarrow$  several topo. orders
- How to find the conditions?  
Study topological excitations above the ground states  
 $\rightarrow$  unitary modular tensor category theory (UMTC)

## Theory of topological excitations = category theory

- Local excitations: 1)  $\Psi_{\text{exc}}$  and  $\Psi$  are LI except near a few points  $\xi_I$ .  
2)  $\Psi_{\text{exc}}$  = ground-state subspace of  $H_{\text{trap}} = H + \delta H_{\xi_1} + \delta H_{\xi_2} \dots$ .  
Ex.  $\Psi = \prod (z_i - z_j)^m$ ,  $\Psi_{\text{exc}} = \prod_{i,I} (z_i - \xi_I) \prod (z_i - z_j)^m \in \otimes_z \mathcal{V}_z$



## Theory of topological excitations = category theory

- Local excitations: 1)  $\Psi_{\text{exc}}$  and  $\Psi$  are LI except near a few points  $\xi_l$ .  
 2)  $\Psi_{\text{exc}}$  = ground-state subspace of  $H_{\text{trap}} = H + \delta H_{\xi_1} + \delta H_{\xi_2} \dots$ .  
 Ex.  $\Psi = \prod (z_i - z_j)^m$ ,  $\Psi_{\text{exc}} = \prod_{i,l} (z_i - \xi_l) \prod (z_i - z_j)^m \in \otimes_z \mathcal{V}_z$
- **Trivial excitation**: can be created by local operators  
 $O(\xi)\Psi \subset \Psi_{\text{exc}} : \Psi \rightarrow \Psi_{\text{exc}}$ .  $O'(\xi)\Psi_{\text{exc}} \subset \Psi : \Psi_{\text{exc}} \rightarrow \Psi$ .  $O(\xi)$  acts on  $\mathcal{V}_\xi$   
*" $\rightarrow$ " (contained in) = morphism in category.*  $\Psi_{\text{trivial exc}} \leftrightarrow \Psi$
- **Topological excitation** if cannot be created by local operators  
 (or more precisely,  $\Psi_{\text{topo. exc}}(\xi_1, \xi_2) \not\rightarrow \Psi$ ,  $\Psi \not\rightarrow \Psi_{\text{topo. exc}}(\xi_1, \xi_2)$  )

# Theory of topological excitations = category theory

- Local excitations: 1)  $\Psi_{\text{exc}}$  and  $\Psi$  are LI except near a few points  $\xi_I$ .  
 2)  $\Psi_{\text{exc}}$  = ground-state subspace of  $H_{\text{trap}} = H + \delta H_{\xi_1} + \delta H_{\xi_2} \dots$ .  
 Ex.  $\Psi = \prod (z_i - z_j)^m$ ,  $\Psi_{\text{exc}} = \prod_{i,I} (z_i - \xi_I) \prod (z_i - z_j)^m \in \otimes_z \mathcal{V}_z$
- **Trivial excitation**: can be created by local operators  
 $O(\xi)\Psi \subset \Psi_{\text{exc}} : \Psi \rightarrow \Psi_{\text{exc}}$ .  $O'(\xi)\Psi_{\text{exc}} \subset \Psi : \Psi_{\text{exc}} \rightarrow \Psi$ .  $O(\xi)$  acts on  $\mathcal{V}_\xi$   
*" $\rightarrow$ " (contained in) = morphism in category.*  $\Psi_{\text{trivial exc}} \leftrightarrow \Psi$
- **Topological excitation** if cannot be created by local operators  
 (or more precisely,  $\Psi_{\text{topo. exc}}(\xi_1, \xi_2) \not\rightarrow \Psi$ ,  $\Psi \not\rightarrow \Psi_{\text{topo. exc}}(\xi_1, \xi_2)$  )
- **Topological type** = equivalence class of  $\Psi_{\text{exc}}$ :  $\Psi_{\text{exc}} \sim \Psi'_{\text{exc}}$  iff  
 $\Psi_{\text{exc}} \rightarrow \Psi'_{\text{exc}}$  and  $\Psi'_{\text{exc}} \rightarrow \Psi_{\text{exc}}$  *isomorphic in category*
- **simple type**:  $\Psi_{\text{exc}} \rightarrow \Psi_{\text{exc}}^{\text{simple}}$  implies  $\Psi_{\text{exc}}^{\text{simple}} \rightarrow \Psi_{\text{exc}}$
- **composite type**:  $k = i \oplus j$ ,  $i \rightarrow k$ ,  $j \rightarrow k$ . *Accidental degeneracy*
- Topo. excitations (topo. types)  $\leftrightarrow$  objects in category  $\mathcal{C}$ :  
**A category  $\mathcal{C}$  = a set of topo. exc.:**  $\mathcal{C} = \{i\}_{\text{simple}} + \{i \oplus j, \dots\}$   
*Example:*  $\mathcal{C} = \{\text{spin-0, spin-1, ..}\}_{\text{simple}} + \{\text{spin-1} \oplus \text{spin-2}, \dots\}$

# Theory of topological excitations = tensor category theory

- Topological excitations can fuse (form bound states)  $\rightarrow$   $\{ \text{Topological excitations} \} \leftrightarrow$  A tensor (fusion) category  $\mathcal{C}$ .

## Data to describe fusion of simple types:

- $i \otimes j = k \rightarrow$  maps  $\{i\}_{\text{simple}} \times \{j\}_{\text{simple}} \rightarrow \{k\}_{\text{simple}}$   
 $\rightarrow k_{ij} \in \{i\}_{\text{simple}}, \forall i, j \in \{i\}_{\text{simple}}$ . **Wrong!**
- Bound state of simple types  $i, j$  may correspond to several  $k$ 's with accidental degeneracy:  $i \otimes j = k_1 \oplus k_2 \oplus k_3 \oplus \dots = \oplus_k N_k^{ij} k$   
 Ex. for  $H$  with  $SO(3)$ :  $\text{spin-1} \otimes \text{spin-1} = \text{spin-0} \oplus \text{spin-1} \oplus \text{spin-2}$
- Associativity condition:  
 $(i \otimes j) \times k = i \otimes (j \otimes k) \rightarrow \sum_m N_m^{ij} N_l^{mk} = \sum_m N_l^{im} N_m^{jk}$

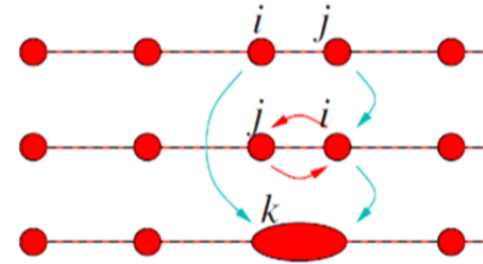
$N_k^{ij}$  are the data to describe fusion of the tensor category.

$N_k^{ij}$  = topological inv.  $\rightarrow$  fusion ring of a tensor category.

$N_k^{ij}$   $\rightarrow$  quantum dimension  $d_i$  for topological type- $i$ .

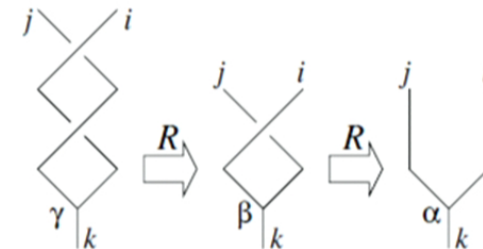
# Theory of topological excitations = braided fusion category

- Particles can also braid → unitary braided fusion category
- Braiding requires that  $N_k^{ij} = N_k^{ji}$ .



- Braiding → **mutual statistics**  $e^{i\theta_{ij}^{(k)}}$  and non-trivial **spin**  $s_i$   
 $2\pi$  rotation of  $(i, j) = 2\pi$  rotation of  $k$   
 $2\pi$  rotation of  $(i, j) = 2\pi$  rotation of  $i$  and  $j$  and exchange  $i, j$  twice

$$e^{i2\pi s_i} e^{i2\pi s_j} e^{i\theta_{ij}^{(k)}} = e^{i2\pi s_k}$$



A unitary braided fusion category (UBFC) is a set of quasiparticles with fusion and braiding, which is described by data  $(N_k^{ij}, s_i)$



## Relation between $(S, T, c)$ and $(N_k^{ij}, s_i, c)$

**Conjecture:** A bosonic topological order [ie a non-degenerate UBFC  $\equiv$  an unitary modular tensor category (UMTC)] is fully characterized by data  $(S, T, c)$  or by data  $(N_k^{ij}, s_i, c)$ .

- From  $(S, T, c)$  to  $(N_k^{ij}, s_i, c)$ : Verlinde formula

$$N_k^{ij} = \sum_l \frac{S_{li} S_{lj} (S_{lk})^*}{S_{1l}}, \quad e^{i2\pi s_i} e^{-i2\pi \frac{c}{24}} = T_{ii}.$$

- From  $(N_k^{ij}, s_i, c)$  to  $(S, T, c)$ :

$$S_{ij} = \frac{1}{\sqrt{\sum_i d_i^2}} \sum_k N_k^{ij} e^{2\pi i (s_i + s_j - s_k)} d_k, \quad T_{ii} = e^{i2\pi s_i} e^{-i2\pi \frac{c}{24}}$$

**Conditions on  $(N_k^{ij}, s_i, c) \leftrightarrow$  Conditions on  $(S, T, c)$**

**$\rightarrow$  A theory of unitary modular tensor category (UMTC)**

## Relation between $(S, T, c)$ and $(N_k^{ij}, s_i, c)$

**Conjecture:** A bosonic topological order [ie a non-degenerate UBFC  $\equiv$  an unitary modular tensor category (UMTC)] is fully characterized by data  $(S, T, c)$  or by data  $(N_k^{ij}, s_i, c)$ .

- From  $(S, T, c)$  to  $(N_k^{ij}, s_i, c)$ : Verlinde formula

$$N_k^{ij} = \sum_l \frac{S_{li} S_{lj} (S_{lk})^*}{S_{1l}}, \quad e^{i2\pi s_i} e^{-i2\pi \frac{c}{24}} = T_{ii}.$$

- From  $(N_k^{ij}, s_i, c)$  to  $(S, T, c)$ :

$$S_{ij} = \frac{1}{\sqrt{\sum_i d_i^2}} \sum_k N_k^{ij} e^{2\pi i (s_i + s_j - s_k)} d_k, \quad T_{ii} = e^{i2\pi s_i} e^{-i2\pi \frac{c}{24}}$$

**Conditions on  $(N_k^{ij}, s_i, c) \leftrightarrow$  Conditions on  $(S, T, c)$**

**$\rightarrow$  A theory of unitary modular tensor category (UMTC)**

*simplified theory of UMTC*

Rowell-Stong-Wang arXiv:0712.1377

- The standard point of view:

UMTC's are fully characterized by  $(N_k^{ij}, F_{klm}^{ijm;\alpha\beta}, R_{k;\beta}^{ij;\alpha})$  (but not one-to-one). Conditions on those data + the equivalent relations

$\rightarrow$  a theory of UMTC.

*hard to work with*

# A simplified theory of UMTC based on $(N_k^{ij}, s_i, c)$

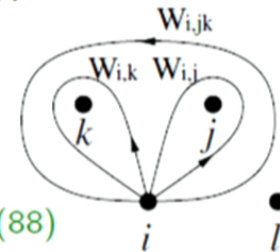
- **Fusion ring:**  $N_k^{ij}$  are non-negative integers that satisfy

$$N_k^{ij} = N_k^{ji}, \quad N_j^{1i} = \delta_{ij}, \quad \sum_{k=1}^N N_1^{ik} N_1^{kj} = \delta_{ij},$$

$$\sum_{m=1}^N N_m^{ij} N_m^{mk} = \sum_{m=1}^N N_m^{im} N_m^{jk} \text{ or } \mathbf{N}_k \mathbf{N}_i = \mathbf{N}_i \mathbf{N}_k$$

where  $i, j, \dots = 1, 2, \dots, N$ , and the matrix  $\mathbf{N}_i$  is given by  $(\mathbf{N}_i)_{kj} = N_k^{ij}$ .  $N_1^{ij}$  defines a charge conjugation  $i \rightarrow \bar{i}$ :

$$N_1^{ij} = \delta_{\bar{i}\bar{j}}.$$



We refer  $N$  as the rank.

- **Vafa's theorem:**  $N_k^{ij}$  and  $s_i$  satisfy Vafa PLB 206, 421 (88)

$$\det(W_{i,j}) \det(W_{i,k}) = \det(W_{i,jk}) \rightarrow \sum_r V_{ijkl}^r s_r = 0 \text{ mod } 1$$

$$V_{ijkl}^r = N_r^{ij} N_r^{kl} + N_r^{il} N_r^{jk} + N_r^{ik} N_r^{jl} - (\delta_{ir} + \delta_{jr} + \delta_{kr} + \delta_{lr}) \sum_m N_m^{ij} N_m^{kl}$$

# 2+1D bosonic topo. orders (up to $E_8$ -states) via $(N_k^{ij}, s_i, c)$

$$\zeta_n^m = \frac{\sin(\pi(m+1)/(n+2))}{\sin(\pi/(n+2))}$$

Rowell-Stong-Wang arXiv:0712.1377; Wen arXiv:1506.05768

$N_c^B$	$d_1, d_2, \dots$	$s_1, s_2, \dots$	wave func.	$N_c^B$	$d_1, d_2, \dots$	$s_1, s_2, \dots$	wave func.
$1_1^B$	1	0					
$2_1^B$	1, 1	$0, \frac{1}{4}$	$\prod(z_i - z_j)^2$	$2_{-1}^B$	1, 1	$0, -\frac{1}{4}$	$\prod(z_i^* - z_j^*)^2$
$2_{14/5}^B$	$1, \zeta_3^1$	$0, \frac{2}{5}$		$2_{-14/5}^B$	$1, \zeta_3^1$	$0, -\frac{2}{5}$	
$3_2^B$	1, 1, 1	$0, \frac{1}{3}, \frac{1}{3}$	(221) double-layer	$3_{-2}^B$	1, 1, 1	$0, -\frac{1}{3}, -\frac{1}{3}$	
$3_{8/7}^B$	$1, \zeta_5^1, \zeta_5^2$	$0, -\frac{1}{7}, \frac{2}{7}$		$3_{-8/7}^B$	$1, \zeta_5^1, \zeta_5^2$	$0, \frac{1}{7}, -\frac{2}{7}$	
$3_{1/2}^B$	$1, 1, \zeta_2^1$	$0, \frac{1}{2}, \frac{1}{16}$	$\Psi_{\text{Pfaffian}}$ $\Psi_{\nu=2}^2 SU(2)_2$	$3_{-1/2}^B$	$1, 1, \zeta_2^1$	$0, \frac{1}{2}, -\frac{1}{16}$	
$3_{3/2}^B$	$1, 1, \zeta_2^1$	$0, \frac{1}{2}, \frac{3}{16}$		$3_{-3/2}^B$	$1, 1, \zeta_2^1$	$0, \frac{1}{2}, -\frac{3}{16}$	
$3_{5/2}^B$	$1, 1, \zeta_2^1$	$0, \frac{1}{2}, \frac{5}{16}$		$3_{-5/2}^B$	$1, 1, \zeta_2^1$	$0, \frac{1}{2}, -\frac{5}{16}$	
$3_{7/2}^B$	$1, 1, \zeta_2^1$	$0, \frac{1}{2}, \frac{7}{16}$		$3_{-7/2}^B$	$1, 1, \zeta_2^1$	$0, \frac{1}{2}, -\frac{7}{16}$	
$4_0^{B,a}$	1, 1, 1, 1	$0, 0, 0, \frac{1}{2}$	$Z_2$ gauge	$4_4^B$	1, 1, 1, 1	$0, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}$	
$4_1^B$	1, 1, 1, 1	$0, \frac{1}{8}, \frac{1}{8}, \frac{1}{2}$	$\prod(z_i - z_j)^4$	$4_{-1}^B$	1, 1, 1, 1	$0, -\frac{1}{8}, -\frac{1}{8}, \frac{1}{2}$	
$4_2^B$	1, 1, 1, 1	$0, \frac{1}{4}, \frac{1}{4}, \frac{1}{2}$	(220) double-layer	$4_{-2}^B$	1, 1, 1, 1	$0, -\frac{1}{4}, -\frac{1}{4}, \frac{1}{2}$	
$4_3^B$	1, 1, 1, 1	$0, \frac{3}{8}, \frac{3}{8}, \frac{1}{2}$		$4_{-3}^B$	1, 1, 1, 1	$0, -\frac{3}{8}, -\frac{3}{8}, \frac{1}{2}$	
$4_0^{B,b}$	1, 1, 1, 1	$0, 0, \frac{1}{4}, -\frac{1}{4}$	double semion	$4_{9/5}^B$	$1, 1, \zeta_3^1, \zeta_3^1$	$0, -\frac{1}{4}, \frac{3}{20}, \frac{2}{5}$	
$4_{-9/5}^B$	$1, 1, \zeta_3^1, \zeta_3^1$	$0, \frac{1}{4}, -\frac{3}{20}, -\frac{2}{5}$		$4_{19/5}^B$	$1, 1, \zeta_3^1, \zeta_3^1$	$0, \frac{1}{4}, -\frac{7}{20}, \frac{2}{5}$	
$4_{-19/5}^B$	$1, 1, \zeta_3^1, \zeta_3^1$	$0, -\frac{1}{4}, \frac{7}{20}, -\frac{2}{5}$	$\Psi_{\nu=3}^2 SU(2)_3$	$4_0^{B,c}$	$1, \zeta_3^1, \zeta_3^1, \zeta_3^1 \zeta_3^1$	$0, \frac{2}{5}, -\frac{2}{5}, 0$	Fibonacci <sup>2</sup>
$4_{12/5}^B$	$1, \zeta_3^1, \zeta_3^1, \zeta_3^1 \zeta_3^1$	$0, -\frac{2}{5}, -\frac{2}{5}, \frac{1}{5}$		$4_{-12/5}^B$	$1, \zeta_3^1, \zeta_3^1, \zeta_3^1 \zeta_3^1$	$0, \frac{2}{5}, \frac{2}{5}, -\frac{1}{5}$	
$4_{10/3}^B$	$1, \zeta_7^1, \zeta_7^2, \zeta_7^3$	$0, \frac{1}{3}, \frac{2}{3}, -\frac{1}{3}$		$4_{-10/3}^B$	$1, \zeta_7^1, \zeta_7^2, \zeta_7^3$	$0, -\frac{1}{3}, -\frac{2}{3}, \frac{1}{3}$	
$5_0^B$	1, 1, 1, 1, 1	$0, \frac{1}{5}, \frac{1}{5}, -\frac{1}{5}, -\frac{1}{5}$	(223) DL	$5_4^B$	1, 1, 1, 1, 1	$0, \frac{2}{5}, \frac{2}{5}, -\frac{2}{5}, -\frac{2}{5}$	
$5_2^{B,a}$	$1, 1, \zeta_4^1, \zeta_4^1, 2$	$0, 0, \frac{1}{8}, -\frac{1}{8}, \frac{1}{3}$		$5_2^{B,b}$	$1, 1, \zeta_4^1, \zeta_4^1, 2$	$0, 0, -\frac{1}{8}, \frac{1}{8}, \frac{1}{3}$	
$5_{-2}^B$	$1, 1, \zeta_4^1, \zeta_4^1, 2$	$0, 0, \frac{1}{8}, -\frac{1}{8}, -\frac{1}{3}$		$5_{-2}^{B,a}$	$1, 1, \zeta_4^1, \zeta_4^1, 2$	$0, 0, -\frac{1}{8}, \frac{1}{8}, -\frac{1}{3}$	
$5_{16/11}^B$	$1, \zeta_9^1, \zeta_9^2, \zeta_9^3, \zeta_9^4$	$0, -\frac{2}{11}, \frac{2}{11}, \frac{1}{11}, -\frac{5}{11}$		$5_{-16/11}^B$	$1, \zeta_9^1, \zeta_9^2, \zeta_9^3, \zeta_9^4$	$0, \frac{2}{11}, -\frac{2}{11}, -\frac{1}{11}, \frac{5}{11}$	
$5_{18/7}^B$	$1, \zeta_5^2, \zeta_5^2, \zeta_{12}^2, \zeta_{12}^4$	$0, -\frac{1}{7}, -\frac{1}{7}, \frac{1}{7}, \frac{3}{7}$		$5_{-18/7}^B$	$1, \zeta_5^2, \zeta_5^2, \zeta_{12}^2, \zeta_{12}^4$	$0, \frac{1}{7}, \frac{1}{7}, -\frac{1}{7}, -\frac{3}{7}$	

Xiao-Gang Wen Oct., 2015

2+1D topological orders and braided fusion category

## Remote detectability: why those $(N_k^{ij}, s_i, c)$ are realizable

- The list cover all the 2+1D bosonic topological orders. But the list might contain fake entries that are not realizable. [Schoutens-Wen arXiv:1508.01111](#) used simple current algebra to construct many-body wave functions for all the entries in the list.



**All the topological order in the table can be realized in multilayer FQH systems**

# A simplified theory of UMTC based on $(N_k^{ij}, s_i, c)$

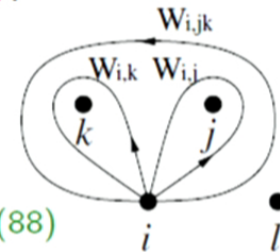
- **Fusion ring:**  $N_k^{ij}$  are non-negative integers that satisfy

$$N_k^{ij} = N_k^{ji}, \quad N_j^{1i} = \delta_{ij}, \quad \sum_{k=1}^N N_1^{ik} N_1^{kj} = \delta_{ij},$$

$$\sum_{m=1}^N N_m^{ij} N_m^{mk} = \sum_{m=1}^N N_m^{im} N_m^{jk} \text{ or } \mathbf{N}_k \mathbf{N}_i = \mathbf{N}_i \mathbf{N}_k$$

where  $i, j, \dots = 1, 2, \dots, N$ , and the matrix  $\mathbf{N}_i$  is given by  $(\mathbf{N}_i)_{kj} = N_k^{ij}$ .  $N_1^{ij}$  defines a charge conjugation  $i \rightarrow \bar{i}$ :

$$N_1^{ij} = \delta_{\bar{i}\bar{j}}.$$



We refer  $N$  as the rank.

- **Vafa's theorem:**  $N_k^{ij}$  and  $s_i$  satisfy Vafa PLB 206, 421 (88)

$$\det(W_{i,j}) \det(W_{i,k}) = \det(W_{i,jk}) \rightarrow \sum_r V_{ijkl}^r s_r = 0 \text{ mod } 1$$

$$V_{ijkl}^r = N_r^{ij} N_r^{kl} + N_r^{il} N_r^{jk} + N_r^{ik} N_r^{jl} - (\delta_{ir} + \delta_{jr} + \delta_{kr} + \delta_{lr}) \sum_m N_m^{ij} N_m^{kl}$$

## A simplified theory of UMTC based on $(N_k^{ij}, s_i, c)$

From  $(N_k^{ij}, s_i, c) \rightarrow (S, T)$

- Let  $d_i$  be the largest eigenvalue of the matrix  $\mathbf{N}_i$ . Let

$$S_{ij} = \frac{1}{D} \sum_k N_k^{ij} e^{2\pi i(s_i + s_j - s_k)} d_k, \quad D^2 = \sum_i d_i^2.$$

Then,  $S$  satisfies

$$S_{11} > 0, \quad \sum_k S_{kl} N_k^{ij} = \frac{S_{li} S_{lj}}{S_{1l}}, \quad S = S^\dagger C, \quad C_{ij} \equiv N_1^{ij}.$$

- Let  $T_{ij} = e^{i2\pi s_i} e^{-i2\pi \frac{c}{24} \delta_{ij}}$  then (rep. of modular group  $SL(2, \mathbb{Z})$ )

$$S^2 = (ST)^3 = C.$$

- Let  $\nu_i = \frac{1}{D^2} \sum_{jk} N_i^{jk} d_j d_k e^{4\pi i(s_j - s_k)}$ . Then  $\nu_i = 0$  if  $i \neq \bar{i}$ , and  $\nu_i = \pm 1$  if  $i = \bar{i}$ .

Rowell-Stong-Wang arXiv:0712.1377

## Remote detectability: why those $(N_k^{ij}, s_i, c)$ are realizable

- The list cover all the 2+1D bosonic topological orders. But the list might contain fake entries that are not realizable. [Schoutens-Wen arXiv:1508.01111](#) used simple current algebra to construct many-body wave functions for all the entries in the list.



**All the topological order in the table can be realized in multilayer FQH systems**



## Remote detectability: why those $(N_k^{ij}, s_i, c)$ are realizable

- The list cover all the 2+1D bosonic topological orders. But the list might contain fake entries that are not realizable. Schoutens-Wen arXiv:1508.01111 used simple current algebra to construct many-body wave functions for all the entries in the list.

**All the topological order in the table can be realized in multilayer FQH systems**

Levin arXiv:1301.7355, Kong-Wen arXiv:1405.5858

- **Remote detectable = Realizable (anomaly-free):**

Every non-trivial topo. excitation  $i$  can be remotely detected by at least one other topo. excitation  $j$  via the non-zero mutual braiding  $\theta_{ij}^{(k)} \neq 0 \rightarrow S_{ij} = \frac{1}{D} \sum_k N_k^{ij} e^{-i\theta_{ij}^{(k)}} d_k$  is unitary (one of conditions)  $\rightarrow$  the topological order is realizable in the same dimension.

- The M-center of BFC  $\mathcal{C} =$  the set of particles with trivial mutual statistics respecting to all others:  $Z_M(\mathcal{C}) \equiv \{i \mid \theta_{ij}^{(k)} = 0, \forall j, k\}$ .  
Remote detectable  $\leftrightarrow Z_M(\mathcal{C}) = \{1\} \leftrightarrow$  Realizable (anomaly-free)



## A simplified theory of UMTC based on $(N_k^{ij}, s_i, c)$

From  $(N_k^{ij}, s_i, c) \rightarrow (S, T)$

- Let  $d_i$  be the largest eigenvalue of the matrix  $\mathbf{N}_i$ . Let

$$S_{ij} = \frac{1}{D} \sum_k N_k^{ij} e^{2\pi i(s_i + s_j - s_k)} d_k, \quad D^2 = \sum_i d_i^2.$$

Then,  $S$  satisfies

$$S_{11} > 0, \quad \sum_k S_{kl} N_k^{ij} = \frac{S_{li} S_{lj}}{S_{1l}}, \quad S = S^\dagger C, \quad C_{ij} \equiv N_1^{ij}.$$

- Let  $T_{ij} = e^{i2\pi s_i} e^{-i2\pi \frac{c}{24} \delta_{ij}}$  then (rep. of modular group  $SL(2, \mathbb{Z})$ )

$$S^2 = (ST)^3 = C.$$

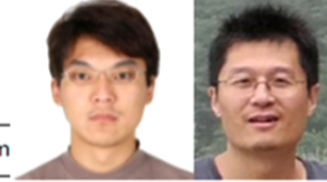
- Let  $\nu_i = \frac{1}{D^2} \sum_{jk} N_i^{jk} d_j d_k e^{4\pi i(s_j - s_k)}$ . Then  $\nu_i = 0$  if  $i \neq \bar{i}$ , and  $\nu_i = \pm 1$  if  $i = \bar{i}$ .

Rowell-Stong-Wang arXiv:0712.1377

# 2+1D fermionic topo. orders (up to $p + ip$ ) via $(N_k^{ij}, s_i, c)$

Classified by **UBFC's with M-center  $\{1, f\}$** .

Lan-Kong-Wen arXiv:1507.04673



$N_c^F(\frac{ \Theta_2 }{\angle\Theta_2/2\pi})$	$D^2$	$d_1, d_2, \dots$	$s_1, s_2, \dots$	com
$2_0^F(\zeta_2^1)$	2	1, 1	$0, \frac{1}{2}$	trivial $\mathcal{F}_0$
$4_0^F(\zeta_2^1)$	4	1, 1, 1, 1	$0, \frac{1}{2}, \frac{1}{4}, -\frac{1}{4}$	$\mathcal{F}_0 \boxtimes 2_1^B(\zeta_2^1)$ , $K = \begin{pmatrix} 2 & 2 \\ 2 & 1 \end{pmatrix}$
$4_{1/5}^F(\zeta_2^1, \zeta_3^1)$	7.2360	$1, 1, \zeta_3^1, \zeta_3^1$	$0, \frac{1}{2}, \frac{1}{10}, -\frac{2}{5}$	$\mathcal{F}_0 \boxtimes 2_{-14/5}^B(\zeta_3^1)$
$4_{-1/5}^F(\zeta_2^1, \zeta_3^1)$	7.2360	$1, 1, \zeta_3^1, \zeta_3^1$	$0, \frac{1}{2}, -\frac{1}{10}, \frac{2}{5}$	$\mathcal{F}_0 \boxtimes 2_{14/5}^B(\zeta_3^1)$
$4_{1/4}^F(\zeta_6^3)$	13.656	$1, 1, \zeta_6^2, \zeta_6^2 = 1 + \sqrt{2}$	$0, \frac{1}{2}, \frac{1}{4}, -\frac{1}{4}$	$\mathcal{F}_{(A_1, 6)}$
$6_0^F(\zeta_2^1)$	6	1, 1, 1, 1, 1, 1	$0, \frac{1}{2}, \frac{1}{6}, -\frac{1}{3}, \frac{1}{6}, -\frac{1}{3}$	$\mathcal{F}_0 \boxtimes 3_{-2}^B(\zeta_2^1)$ , $K = (3)$ , $\Psi_{1/3}(z_i)$
$6_0^F(\zeta_2^1)$	6	1, 1, 1, 1, 1, 1	$0, \frac{1}{2}, -\frac{1}{6}, \frac{1}{3}, -\frac{1}{6}, \frac{1}{3}$	$\mathcal{F}_0 \boxtimes 3_2^B(\zeta_2^1)$ , $K = (-3)$ , $\Psi_{1/3}^*(z_i)$
$6_0^F(\zeta_6^3)$	8	$1, 1, 1, 1, \zeta_2^1, \zeta_2^1 = \sqrt{2}$	$0, \frac{1}{2}, 0, \frac{1}{2}, \frac{1}{16}, -\frac{7}{16}$	$\mathcal{F}_0 \boxtimes 3_{1/2}^B(\zeta_6^1)$ , $\mathcal{F}_{U(1)_2/\mathbb{Z}_2}$
$6_0^F(\zeta_6^3)$	8	$1, 1, 1, 1, \zeta_2^1, \zeta_2^1$	$0, \frac{1}{2}, 0, \frac{1}{2}, -\frac{1}{16}, \frac{7}{16}$	$\mathcal{F}_0 \boxtimes 3_{-1/2}^B(\zeta_6^1)$
$6_0^F(\frac{1.0823}{3/16})$	8	$1, 1, 1, 1, \zeta_2^1, \zeta_2^1$	$0, \frac{1}{2}, 0, \frac{1}{2}, \frac{3}{16}, -\frac{5}{16}$	$\mathcal{F}_0 \boxtimes 3_{3/2}^B(\frac{0.7653}{3/16})$
$6_0^F(\frac{1.0823}{-3/16})$	8	$1, 1, 1, 1, \zeta_2^1, \zeta_2^1$	$0, \frac{1}{2}, 0, \frac{1}{2}, -\frac{3}{16}, \frac{5}{16}$	$\mathcal{F}_0 \boxtimes 3_{-3/2}^B(\frac{0.7653}{-3/16})$
$6_{1/7}^F(\zeta_2^1, \zeta_5^2)$	18.591	$1, 1, \zeta_5^1, \zeta_5^1, \zeta_5^2, \zeta_5^2$	$0, \frac{1}{2}, \frac{5}{14}, -\frac{1}{7}, -\frac{3}{14}, \frac{2}{7}$	$\mathcal{F}_0 \boxtimes 3_{8/7}^B(\zeta_5^2)$
$6_{-1/7}^F(\zeta_2^1, \zeta_5^2)$	18.591	$1, 1, \zeta_5^1, \zeta_5^1, \zeta_5^2, \zeta_5^2$	$0, \frac{1}{2}, -\frac{5}{14}, \frac{1}{7}, \frac{3}{14}, -\frac{2}{7}$	$\mathcal{F}_0 \boxtimes 3_{-8/7}^B(\zeta_5^2)$
$6_0^F(\zeta_{10}^2)$	44.784	$1, 1, \zeta_{10}^2, \zeta_{10}^2, \zeta_{10}^4, \zeta_{10}^4$	$0, \frac{1}{2}, \frac{1}{3}, -\frac{1}{6}, 0, \frac{1}{2}$	$\mathcal{F}_{(A_1, -10)}$
$6_0^F(\zeta_{10}^2)$	44.784	$1, 1, \zeta_{10}^2, \zeta_{10}^2, \zeta_{10}^4, \zeta_{10}^4$	$0, \frac{1}{2}, -\frac{1}{3}, \frac{1}{6}, 0, \frac{1}{2}$	$\mathcal{F}_{(A_1, 10)}$

Xiao-Gang Wen Oct., 2015

2+1D topological orders and braided fusion category