

Title: PSI 2015/2016 Statistical Mechanics - G. Baskaran - Lecture 12

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Abstract:

RENORMALIZATION



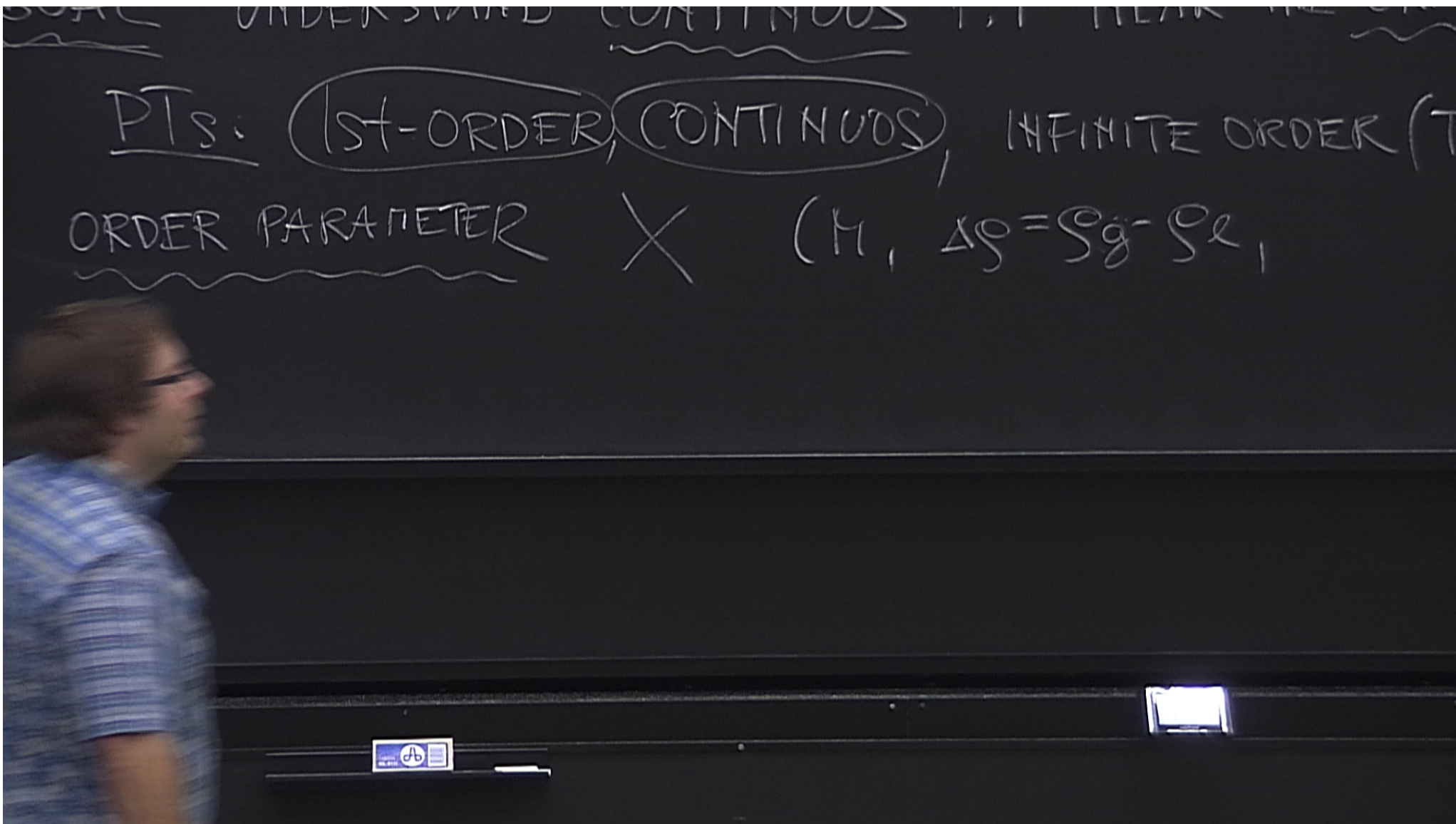
RENORMALIZATION

1) GOAL UNDERSTAND CONTINUOUS P.T NEAR THE CRITICAL
PTS. 1ST-ORDER, CONTINUOUS, INFINITE ORDER (

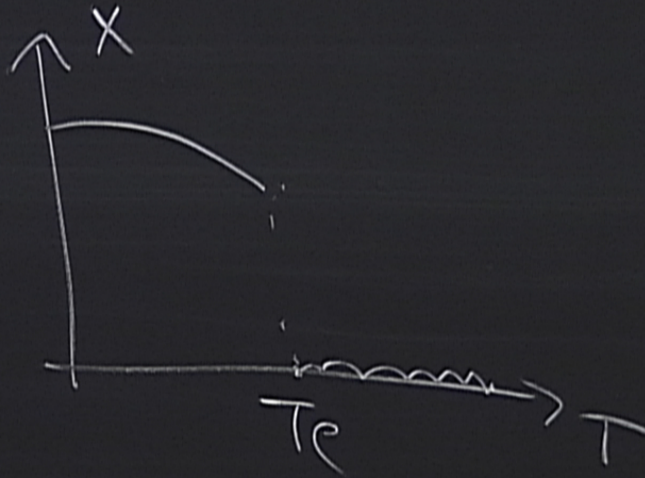


T NEAR THE CRITICAL POINT

INFINITE ORDER (TOPOLOGICAL), LIQUID-GLASS, QUANTUM ($T=0$)

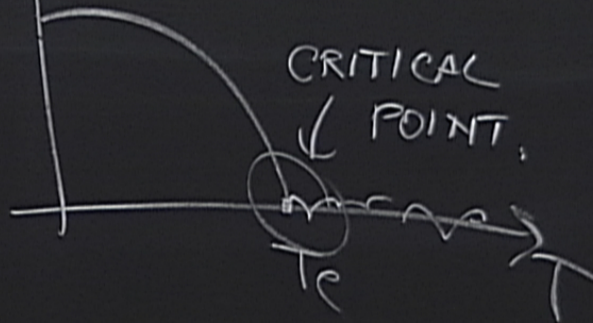


a 1st ORDER

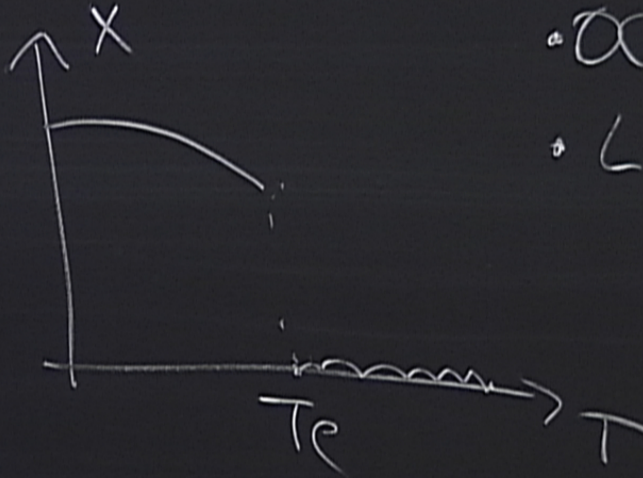


CONTINUOUS

X

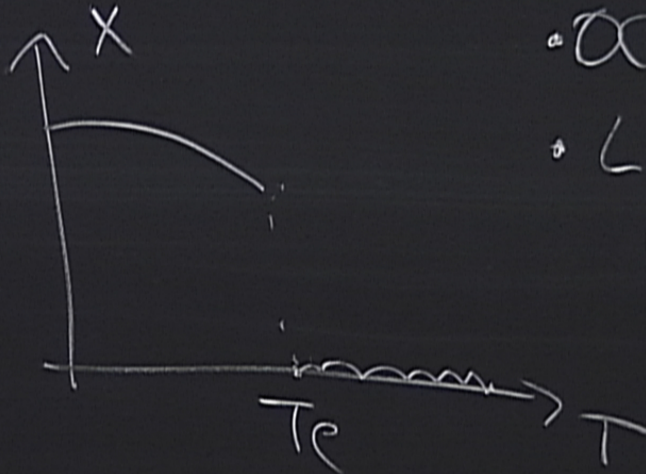


a 1st ORDER

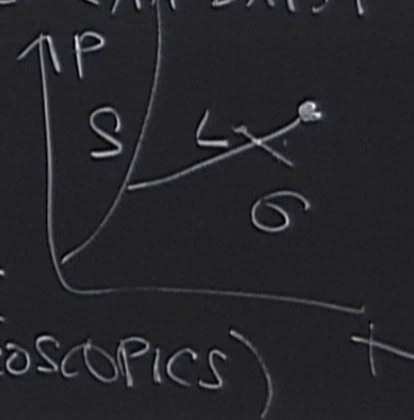


- OCCUR OFTEN IN NATURE
- LATENT HEAT

• 1st ORDER



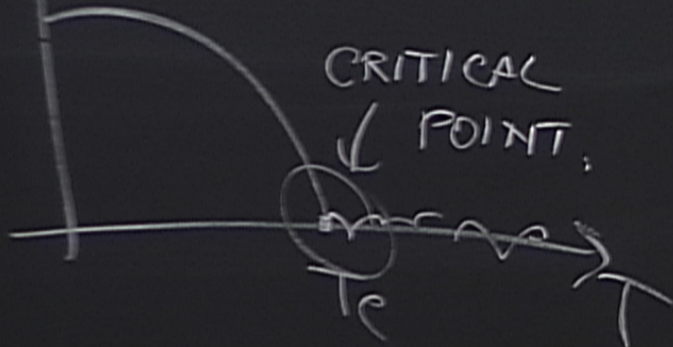
- OCCUR OFTEN IN NATURE
- LATENT HEAT \Rightarrow DISTINCT PHASES CAN EXIST IN EQ.



- DIFFICULT TO DESCRIBE (DEPEND ON MICROSCOPICS)

CONTINUOUS

X



- LESS OFTEN (NEED SYMMETRY)
- DEMONSTRATE UNIVERSALITY

$$t = \frac{T - T_c}{T_c}$$

ORDERED PHASE $t < 0$

$$C \propto |t|^{-d}$$
$$x \propto |t|^{-\alpha}$$

$$1 \propto |t|^\beta$$

μ_0

$$G(n) \propto \frac{2^{-n/3}}{n^{d-2+\gamma}}$$
$$, \quad \xi \propto |t|^{-\nu}$$

$$\begin{aligned}
 C &\propto |t|^{-d} \\
 \chi &\propto |t|^{-\gamma} \\
 M &\propto |t|^{-\beta} \\
 H &\propto M^{\delta}
 \end{aligned}$$

$$G(n) \propto \frac{2^{-n/\beta}}{n^{d-2+\gamma}} \quad , \quad \xi \propto |t|^{-\nu}$$

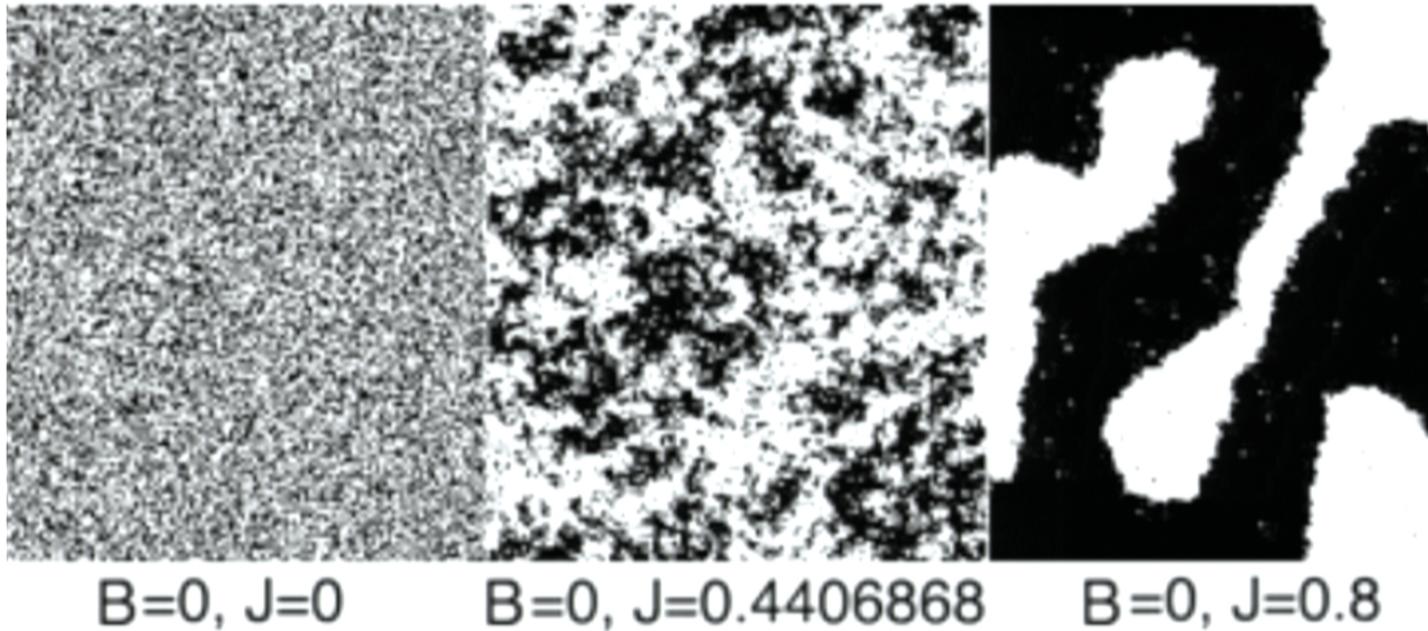
UNIVERSALITY CLASSES (HAVE T)

3 UNIVERSALITY CLASSES (HAVE THE SAME CRIT. EXPTS.)

- MFT .. QUALITATIVELY DESCRIBES C.P.T., BUT GIVES WRONG CRITICAL EXPONENTS.
- AIM: TO INCLUDE FLUCTUATIONS
→ RIGHT CRIT. EX.

$$\alpha + \beta(\delta + 1)$$
$$\gamma = \nu$$
$$\alpha = \nu$$

ISING MODEL: ordering

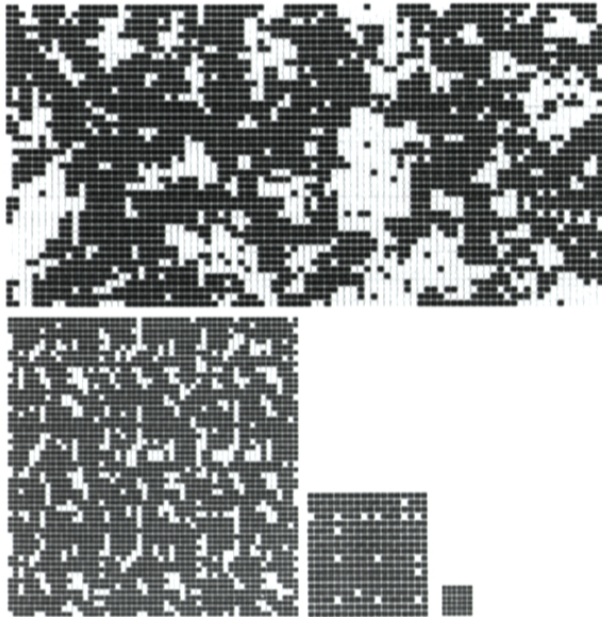


At T_c : scale invariant (fractal) structure

<http://www.dhushara.com/DarkHeart/Ising/Ising.htm>

Block Spin RG

Below T_c ($M=M_0$)



Above T_c ($M=0$)



<https://plus.maths.org/content/going-flow-0>

2) FROM ISING TO FIELD THEORY

• ISING $H = -\frac{1}{2} \sum_{ij} J_{ij} \sigma_i \sigma_j$

$$Z = \sum_{\{\sigma_i\}} e^{\frac{1}{2T} \sum_{ij} J_{ij} \sigma_i \sigma_j}$$

2) FROM ISING TO FIELD THEORY

• ISING $H = -\frac{1}{2} \sum_{ij} J_{ij} \sigma_i \sigma_j$

$$Z = \sum_{\{\sigma_i\}} e^{\frac{1}{2T} \sum_{ij} J_{ij} \sigma_i \sigma_j}$$

a) APPLY HUBBARD-STRATONOVICH TF.

$$\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} d\varphi e^{-\frac{1}{2}a\varphi^2 + b\varphi} = \frac{1}{\sqrt{a}} e^{\frac{b^2}{2a}} \quad a > 0$$

$$e^{\frac{1}{2} \ln A + i \arg b_j} = \frac{1}{\sqrt{A}} (2\pi)^{N/2} \int_{-\infty}^{\infty} d\varphi_1 \dots d\varphi_N e^{-\frac{1}{2} \varphi_i A_{ij}^{-1} \varphi_j + b_i \varphi_i}$$

$$\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} d\varphi e^{-\frac{1}{2}a\varphi^2 + b\varphi} = \frac{1}{\sqrt{a}} e^{\frac{b^2}{2a}} \quad a > 0$$

$$e^{\frac{1}{2} \ln \det A} = \frac{1}{\sqrt{\det A}} \frac{1}{(2\pi)^{N/2}} \int_{-\infty}^{\infty} d\varphi_i \cdot d\varphi_{i+1} \dots e^{-\frac{1}{2} \varphi_i A_{ij}^{-1} \varphi_j + b_i \varphi_i}$$

$N \rightarrow \infty$

$$e^{\frac{1}{2} \ln |A_{ij}|} = \frac{1}{\sqrt{|\det A|}} \frac{1}{(2\pi)^{N/2}} \int_{-\infty}^{\infty} d\phi_i d\phi_{ii} e^{-\frac{1}{2} \phi_i |A_{ij}|^{-1} \phi_j + b_i \phi_i}$$

• IDENTIFY $A_{ij} = J_{ij}/T$, $b_i = C_i$

$\|\alpha\| \rightarrow \infty$

$$e^{\frac{1}{2} \ln |A_{ij} b_j|} = \frac{1}{\sqrt{\det A}} \frac{1}{(2\pi)^{N/2}} \int_{-\infty}^{\infty} d\varphi_i \cdot d\varphi_{ii} \quad e^{-\frac{1}{2} \varphi_i A_{ij}^{-1} \varphi_j + b_i \varphi_i}$$

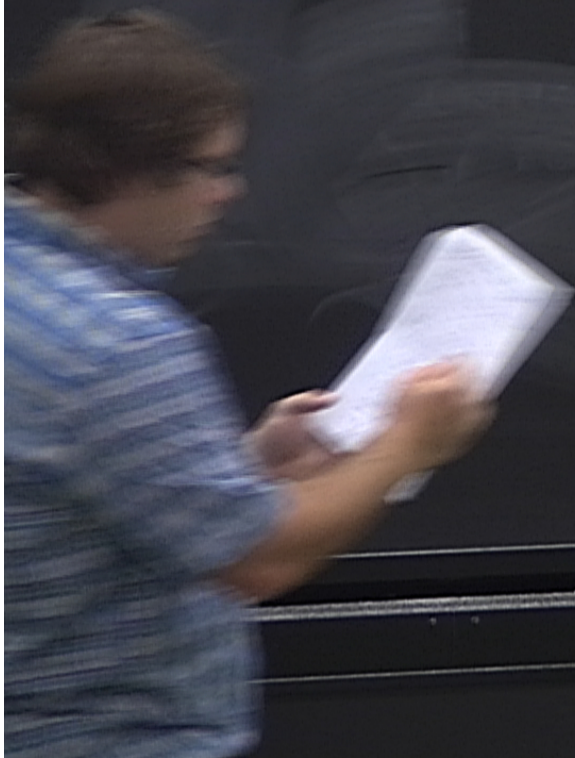
IDENTIFY $A_{ij} = J_{ij}/T$, $b_i = \bar{c}_i$

$$Z = \frac{1}{(2\pi)^{N/2}} \sum_{\{\bar{c}_i\}} \int_{-\infty}^{\infty}$$

$$\varphi_i \rightarrow \varphi_i/T = \frac{1}{\sqrt{d\varphi}} \frac{1}{(2\pi\hbar)^{N/2}} \sum_{\{b_n\}} \int_{-\infty}^{\infty} D\varphi e^{-\frac{1}{2T} \varphi_i J_{ij}^{-1} \varphi_j} + \frac{\varphi_i b_i}{T}$$

$$\frac{\psi_i G_i}{T}$$

INTERACTION OF SPIN G_i WITH SOME FLUCTUATING
"FIELD φ "



$$\frac{\chi_i G_i}{T}$$

INTERACTION OF SPIN G_i WITH SOME FLUCTUATING
"FIELD φ " (MFT $\varphi \ll \langle \varphi \rangle \ll M$)

$$\sum_{\sigma_i = \pm 1} e^{\frac{1}{T} \sum_i \sigma_i \phi_i} = e^{\log [2 \cosh \dots]}$$

$$Z = \int D\varphi e^{-S[\varphi]}$$

$$S[\varphi] = \frac{1}{2T} \varphi_i J_{ij}^{-1} \varphi_j - \sum_i \log [2 \cosh \varphi_i / T]$$

NOTE: $\text{NFT} \Leftrightarrow \text{SADDLE POINT}$
LET $\varphi = \varphi_N$ CONF. MINIMIZES S

$$\frac{\partial S}{\partial \varphi_i} \Big|_{\varphi_i = \bar{\varphi}_i} = 0$$

NOTE: NIFT \Leftrightarrow SADDLE POINT
LET $\overline{\varphi}_1, \dots, \overline{\varphi}_N$ CONF. MINIMIZES \underline{S}

$$\frac{\partial S}{\partial \varphi_i} \Big|_{\varphi_i = \overline{\varphi}_i} = 0 = \frac{1}{T} \sum_j J_{ij}^{-1} \overline{\varphi}_j - \frac{1}{T} \tanh \frac{\overline{\varphi}_i}{T}$$

$\text{MFT} \Leftrightarrow \text{SADDLE POINT}$
 $\bar{\varphi}_1, \dots, \bar{\varphi}_N$ CONF. MINIMIZES \underline{S}

$$\frac{\partial S}{\partial \varphi_i} \Big|_{\varphi_i = \bar{\varphi}_i} = 0 = \frac{1}{T} \sum_j J_{ij} \bar{\varphi}_j - \frac{1}{T} \tanh \frac{\bar{\varphi}_i}{T}$$

$$\Rightarrow \bar{\varphi}_i = \sum_j J_{ij} \tanh \frac{\bar{\varphi}_j}{T}$$

$$\bar{\varphi}_i = \varphi_i$$

MFPT \Leftrightarrow SADDLE POINT

$\varphi_1, \dots, \varphi_N$ CONF. MINIMIZES S

$$\frac{\partial S}{\partial \varphi_i} \Big|_{\varphi_i = \bar{\varphi}_i} = 0 = \frac{1}{T} \sum_j J_{ij} \bar{\varphi}_j - \frac{1}{T} \tanh \frac{\varphi_i}{T}$$

$$\Rightarrow \bar{\varphi}_i = \sum_j J_{ij} \tanh \frac{\bar{\varphi}_j}{T}$$

$$\bar{\varphi} = J \tanh$$

$$\bar{\varphi}_i = \bar{\varphi} \quad T_i$$

b) "CONTINUUM LIMIT" φ_i . DEFINED AT SITE $i \rightarrow \varphi(x)$

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"ZOOM OUT" \leftrightarrow LOOSE INFO ABOUT UNDERLYING LATTICE.

$$Z = \int D\varphi e^{-S[\varphi]}$$

$$S[\varphi] = \frac{1}{2T} \varphi_i J_{ij}^{-1} \varphi_j - \sum_i \log[2 \cosh \varphi_i/T]$$

$$\sum_{\sigma_i = \pm 1} e^{\frac{1}{T} \varphi_i \sigma_i} = e^{\log 2}$$

• LOOK AT $\langle \psi_i | J_{ij}^{-1} | \psi_j \rangle$ USING LATTICE F.T



• LOOK AT $\langle \psi_i | J_{ij}^{-1} | \psi_j \rangle$ USING LATTICE F.T

$$\psi_i = \frac{1}{N} \sum_{\vec{q}} \psi(\vec{q}) e^{i\vec{q} \cdot \vec{r}_i} \quad \frac{1}{N}$$

$$J_{ij} = \frac{1}{N} \sum_{\vec{q}} J(\vec{q}) e^{i\vec{q} \cdot (\vec{r}_i - \vec{r}_j)}$$

USING LATTICE F.T

$$f(\vec{q}) e^{i\vec{q} \cdot \vec{r}_i}$$

$$f(\vec{q}) e^{i\vec{q} \cdot (\vec{r}_i - \vec{r}_j)}$$

$$\frac{1}{N} \sum_i e^{-i(\vec{k} - \vec{k}') \cdot \vec{r}_i} = \delta_{\vec{k}, \vec{k}'}$$

$$\frac{1}{N} \sum_{i,j} e^{i\vec{q} \cdot (\vec{r}_i - \vec{r}_j)} = \delta_{ij}$$

USING LATTICE F.T

$$f(\vec{q}) e^{i\vec{q} \cdot \vec{r}_1}$$

$$f(\vec{q}) e^{i\vec{q} \cdot (\vec{r}_1 - \vec{r}_j)}$$

$$\frac{1}{N} \sum_{\vec{r}} e^{-i(\vec{k} - \vec{k}') \cdot \vec{r}} = \delta_{\vec{k}, \vec{k}'}$$

$$\frac{1}{N} \sum_{\vec{r}_1, \vec{r}_j} e^{i\vec{q} \cdot (\vec{r}_1 - \vec{r}_j)} = \delta_{ij}$$

$$\sum_{ij} \psi_i J_{ij}^{-1} \psi_j = \sum_{ij} \frac{1}{H} \sum_q \psi(q) e^{i\vec{q} \cdot \vec{r}_i} \frac{1}{H} \sum_{q'} \psi^{-1}$$

$$\begin{aligned}
\sum_{ij} \psi_i J_{ij}^{-1} \psi_j &= \sum_{ij} \frac{1}{N} \sum_q \psi(q) e^{i\vec{q} \cdot \vec{r}_i} \frac{1}{N} \sum_{q'} J^{-1}(q') e^{i q' (\vec{r}_i - \vec{r}_j)} \frac{1}{N} \sum_{q''} \psi(q'') \\
&= \frac{1}{N} \sum_{q, q', q''} \psi(q) J^{-1}(q') \psi(q'') \delta_{q, -q'} \delta_{q', q''} \\
&= \frac{1}{N} \sum_q \psi(-q) J^{-1}(q) \psi(q)
\end{aligned}$$

$$\begin{aligned}
\sum_{i,j} \psi_i^{-1} \psi_j &= \sum_{i,j} \frac{1}{N} \sum_q \psi(q) e^{i\vec{q} \cdot \vec{r}_i} \frac{1}{N} \sum_{q'} \psi'(q') e^{i q' (\vec{r}_i - \vec{r}_j)} \frac{1}{N} \sum_{q''} \psi(q'') \\
&= \frac{1}{N} \sum_{q, q', q''} \psi(q) \psi'(q') \psi(q'') \delta_{q, -q'} \delta_{q', q''} \\
&= \frac{1}{N} \sum_q \psi(-q) \psi'(q) \psi(q)
\end{aligned}$$

• WHAT IS $J(q)$?

$$J(q) = \frac{1}{N} \sum_{ij} J_{ij} e^{-iq \cdot (\vec{r}_i - \vec{r}_j)}$$

SMALL q

q

• WHAT IS $J(q)$?

$$J(q) = \frac{1}{N} \sum_{ij} J_{ij} e^{-iq \cdot (\vec{r}_i - \vec{r}_j)}$$

$$\approx \frac{1}{N} \sum_{ij} J_{ij} \left(1 - iq \cdot (\vec{r}_i - \vec{r}_j) - \frac{1}{2} [q \cdot (\vec{r}_i - \vec{r}_j)]^2 + \dots \right)$$

SMALL q $qR \ll 1$
(FOUR OUT)

• WHAT IS $J(q)$?

$$J(q) = \frac{1}{N} \sum_{ij} J_{ij} e^{-iq \cdot (\vec{r}_i - \vec{r}_j)}$$

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NJ

SMALL q $qR \ll 1$
(FOUR OUT) $\frac{qR}{2a}$

• WHAT IS $J(q)$?

$$J(q) = \frac{1}{N} \sum_{ij} J_{ij} e^{-iq \cdot (\vec{r}_i - \vec{r}_j)}$$

SMALL q $qR \ll 1$
(FOUR FOUR)

$$\approx \frac{1}{N} \sum_{ij} J_{ij} \left(1 - iq \cdot (\vec{r}_i - \vec{r}_j) - \frac{1}{2} [q \cdot (\vec{r}_i - \vec{r}_j)]^2 + \dots \right)$$

$$\approx \frac{1}{2N} \sum_{ij} J_{ij} [q \cdot (\vec{r}_i - \vec{r}_j)]^2$$

• WHAT IS $J(q)$?

$$J(q) = \frac{1}{N} \sum_{ij} J_{ij} e^{-i\vec{q} \cdot (\vec{r}_i - \vec{r}_j)}$$

SMALL q $qR \ll 1$
(FOUR FOUR)

$$\approx \frac{1}{N} \sum_{ij} J_{ij} \left(1 - i\vec{q} \cdot (\vec{r}_i - \vec{r}_j) - \frac{1}{2} [\vec{q} \cdot (\vec{r}_i - \vec{r}_j)]^2 + \dots \right)$$

$$\approx \frac{1}{2N} \sum_{ij} \underbrace{J_{ij} (\vec{r}_i - \vec{r}_j)^2}_{q^2 R^2} J_{ij}$$

$\alpha \dots$ SPIN STIFFNESS

$$\alpha \approx \frac{R^2 J}{d}$$

\propto ... SPIN STIFFNESS

$$\partial P \approx \frac{R^2 J}{d}$$

WHAT IS $J(q)$?

$$J(q) = \frac{1}{N} \sum_{ij} J_{ij} e^{-i\vec{q} \cdot (\vec{r}_i - \vec{r}_j)}$$

SMALL q $qR \ll 1$
(FOUR OUT)

$$\approx \frac{1}{N} \sum_{ij} J_{ij} \left(1 - i\vec{q} \cdot (\vec{r}_i - \vec{r}_j) - \frac{1}{2} [\vec{q} \cdot (\vec{r}_i - \vec{r}_j)]^2 + \dots \right)$$

$$\approx J - \frac{1}{2N} \sum_{ij} \underbrace{[\vec{q} \cdot (\vec{r}_i - \vec{r}_j)]^2}_{\frac{q^2 R^2 J \cdot N}{d}} J_{ij} \approx J - \frac{1}{2} \chi(q)$$

$$\bullet S_0[\varphi] = \frac{1}{2T^2} \left[\sum_{ij} T \varphi_i J_{ij}^{-1} \varphi_j - \sum \varphi_i^2 \right]$$

$$\begin{aligned}
 \bullet S_0[\varphi] &= \frac{1}{2T^2} \left[\sum_{ij} T \varphi_i J_{ij}^{-1} \varphi_j - \sum \varphi_i^2 \right] \\
 &= \frac{1}{2T^2 N} \sum_{q_j} \varphi(q_j) \left[\frac{T}{J(q_j)} - 1 \right] \varphi(q_j) \\
 &= \frac{1}{2T^2 N} \sum_{q_j} \left[\frac{T}{J(q_j)} - 1 \right] + \frac{R^2 T}{2d T_c} q_j^2
 \end{aligned}$$

$$P_0[\varphi] = \frac{1}{2T^2} \left[\sum_{ij} T \varphi_i J_{ij}^{-1} \varphi_j - \sum \varphi_i^2 \right] \quad \underline{J = T_c}$$

$$= \frac{1}{2T^2 N} \sum_{q_j} \varphi(q_j) \left[\frac{T}{J(q_j)} - 1 \right] \varphi(q_j)$$

$$\frac{T}{T_c} - 1 + \frac{R^2 T}{2d T_c} q^2 = t + \frac{R^2 q^2}{2d}$$

$$J = Tc$$

$$\varphi \rightarrow \frac{Tc \sqrt{2d}}{R a^{d/2}} \varphi,$$

$$R \equiv \frac{2dt}{R^2}$$

$$t + \frac{R^2 g^2}{2d}$$

$$\left[\sum_{ij} T \psi_i J_{ij}^{-1} \psi_j - \sum \psi_i^2 \right] \quad \underline{J = Tc}$$

$$\psi \rightarrow \frac{Tc \sqrt{2d}}{R a^{d/2}} \psi$$

$$\sum_{q_1} \psi(-q_1) \left[\frac{T}{J(q_1)} - 1 \right] \psi(q_1)$$

$$\frac{T}{Tc} - 1 + \frac{R^2 T}{2d Tc} q^2 = t + \frac{R^2 q^2}{2d}$$

$$\sum_{q_1} \psi(-q_1) (n + q_1^2) \psi(q_1) = \frac{1}{Z} \int \frac{d^d q}{(2\pi)^d} \psi(-q) (n + q^2) \psi(q)$$

- n ... CHANGES SIGN AT THE TRANSITION ($n \sim t$)
- $\mu > 0$

- μ ... CHANGES SIGN AT THE TRANSITION (μt)
- $\mu > 0$ (FOR STABILITY)
- CAN WRITE THIS IMMEDIATELY (S' ... $\delta \rightarrow -\delta$ EVEN

3) DIMENSIONAL ANALYSIS

$$[k] = \dim k = +1$$

$$[x] = -1$$

$$[S] = 0 = [x^d] + \frac{[D^2]}{2}$$

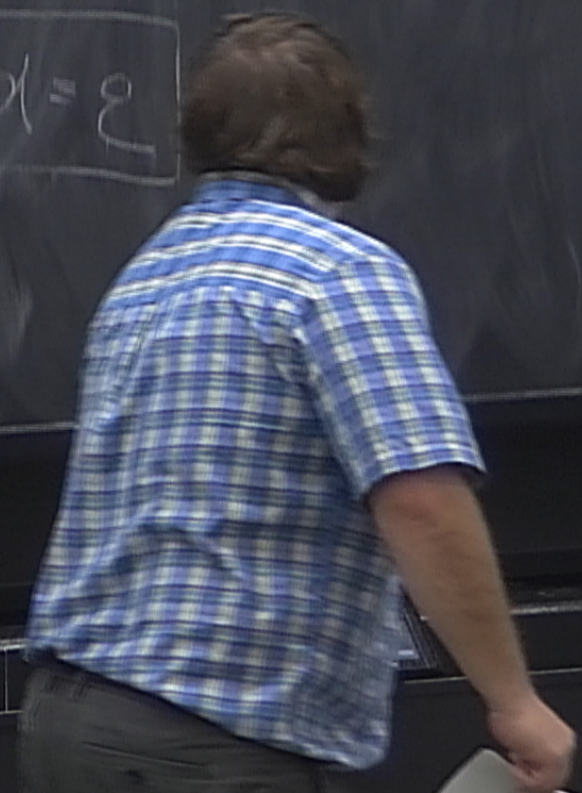
LOCAL ANALYSIS

$$[S] = 0 = \begin{bmatrix} x^2 & 1 \\ -\alpha & 2 \end{bmatrix} + \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} + \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

$$\dim k = +1$$
$$= -1$$

$$\Rightarrow [C] = \frac{\alpha - 2}{2}$$

$$[L] = \alpha, [m] = 4 - \alpha = 2$$



SCALE ANALYSIS

$$= \dim k = +1$$
$$= -1$$

$$[S] = 0 = \underbrace{[x^d]}_{-d} + \underbrace{[D^c]}_2 + \underbrace{[\varphi^c]}$$

• TOT 9 \Rightarrow IN $d=$

$$\Rightarrow [c] = \frac{d-2}{2}$$

$$[n] = d, [m] = 4 - d = 2$$

UPPER CRITICAL DIMENSION
(MFT EXACT)