

Title: PSI 2015/2016 Quantum Field Theory I - Daniel Wohns and Tibra Ali - Lecture 5

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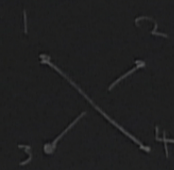
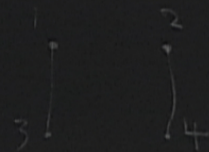
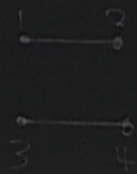
Abstract:

Feynman Diagrams

$$\langle \Phi | S | \Phi \rangle \xrightarrow{\text{LSZ}} \langle \Omega | T \phi_1 \dots \phi_n | \Omega \rangle \xrightarrow{\text{int. pic.}} \langle 0 | T \phi_1 \dots \phi_n | 0 \rangle \xrightarrow{\text{Wick}} \Delta_F$$

Feynman Diagrams

$$\langle 0 | T \phi_1 \phi_2 \phi_3 \phi_4 | 0 \rangle = \Delta_{12} \Delta_{34} + \Delta_{13} \Delta_{24} + \Delta_{14} \Delta_{23}$$



$$\Delta_{ij} = \Delta_F(x_i - x_j)$$

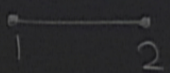
Wick
 $\langle \varphi_1 | \varphi_2 \rangle \xrightarrow{\text{Wick}} \Delta_F$

$\Delta_{ij} = \Delta_F(x_i - x_j)$

$$\langle \Omega | \varphi_1 \varphi_2 | \Omega \rangle = \frac{\langle 0 | T \varphi_1 \varphi_2 \exp[-i \int d^4x H_{int}(x)] | 0 \rangle}{\langle 0 | T \exp[-i \int d^4x H_{int}(x)] | 0 \rangle}$$

$$H_{int}(x) = \int d^3y \frac{\lambda}{4!} \varphi^4(y)$$

Numerator: $\langle 0 | T \varphi_1 \varphi_2 + T \varphi_1 \varphi_2 \left(-i \frac{\lambda}{4!}\right) \int d^4y \varphi_y^4 | 0 \rangle$

$\langle 0 | T \varphi_1 \varphi_2 | 0 \rangle =$ 

$\overbrace{\varphi_y \varphi_y \varphi_y \varphi_y}$ or $\overbrace{\varphi_y \varphi_y \varphi_y \varphi_y}$
 or $\overbrace{\varphi_y \varphi_y \varphi_y \varphi_y}$

$\langle 0 | T \varphi_1 \varphi_2 \left(-i \frac{\lambda}{4!}\right) \int d^4y \varphi_y^4 | 0 \rangle = 3 \frac{-i \lambda}{4!} \Delta_{12} \int d^4y \Delta_{yy} \Delta_{yy} + Q$

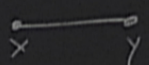
Wick contractions

$$Q = 12 \left(\frac{-i\lambda}{4!} \right) \int d^4y \Delta_{1y} \Delta_{2y} \Delta_{yy}$$

$$\langle 0 | \varphi_1 \varphi_2 \frac{-i\lambda}{4!} \int d^4y \varphi_y | 0 \rangle = \text{diagram 1} + \text{diagram 2}$$

Position

Position Space Feynman Diagrams for n-point functions:

1. For each line:  = $\Delta_F(x-y)$

2. For vertex: $\chi_z = -i\lambda \int d^4z$

Position Space Feynman Diagrams for n-point functions:

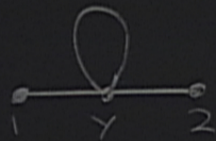
1. For each line: $\text{---} = \Delta_F(x-y)$

2. For vertex: $\times = -i\lambda \int d^4z$ (depends on Hint)

3. For each external point $\bullet = 1$ (will change w/ fermions or gauge bosons)

4. Divide by a symmetry factor S .

Symmetry Factor

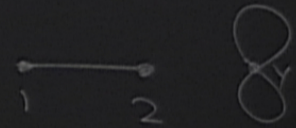


has coefficient $\frac{12}{4!} = \frac{1}{2} = \frac{1}{S}$

$S = \#$ ways to exchange components of a diagram w/o changing the diagram

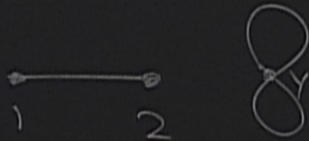
3 types of symme

- exchange en
- exchange two
- exchange two



3 types of symmetries:

- exchange ends of a line
- exchange two lines
- exchange two vertices



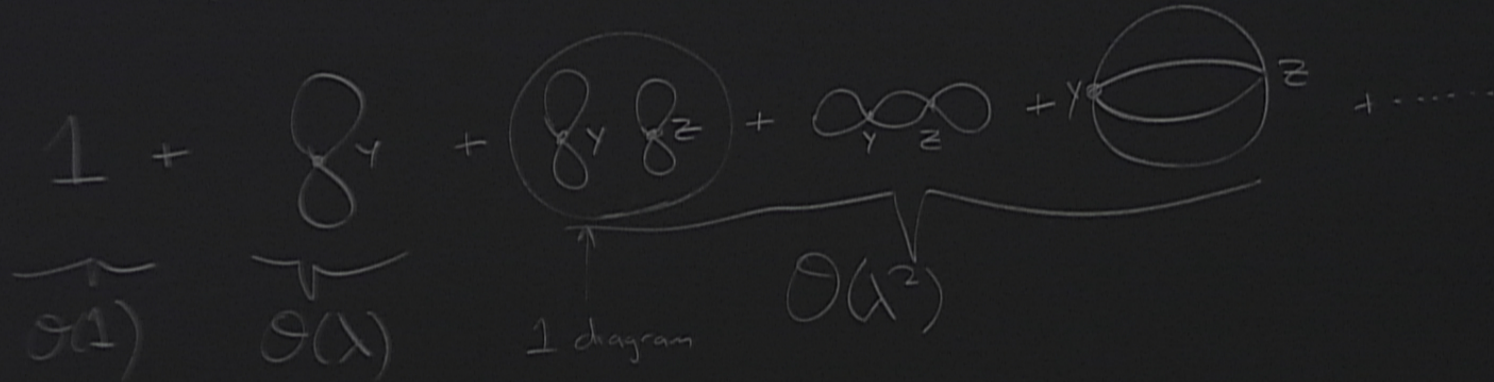
$$\frac{1}{5} = \frac{1}{8} = \frac{3}{4!}$$

exchanging ends of upper line

$$S = 2 \cdot 2 \cdot 2 = 8$$

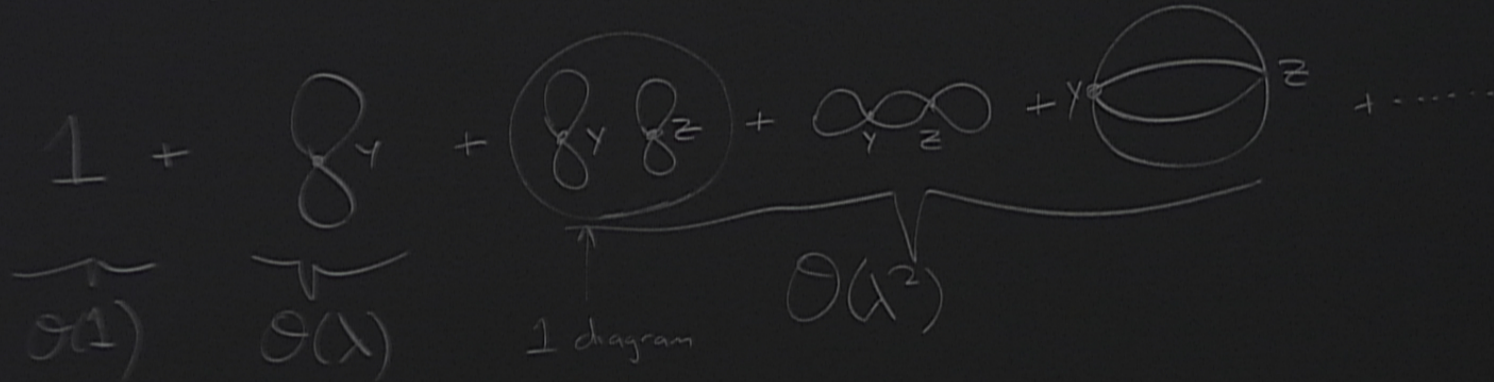
ends of lower line exchanging two lines

$$\langle 0 | T \exp[-i \int dt H_{int}] | 0 \rangle =$$



$$e^{\mathcal{O}} = 1 + \mathcal{O} + \frac{1}{2} (\mathcal{O})^2$$

$$\langle 0 | T \exp[-i \int dt H_{int}] | 0 \rangle =$$




$$e^{\mathcal{G}} = 1 + \mathcal{G} + \frac{1}{2} (\mathcal{G})^2 + \frac{1}{3!} (\mathcal{G})^3$$

Extra symmetry factor from two copies of same diagram

denom

$$\text{denom} = \sum_{\{n_i\}} \prod_i \frac{1}{n_i!} (V_i)^{n_i}$$

$$V_i = \{ \infty, \infty, \text{circle with two dots}, \dots \}$$


$$\text{denom} = \sum_{\{n_i\}} \prod_i \frac{1}{n_i!} (V_i)^{n_i}$$

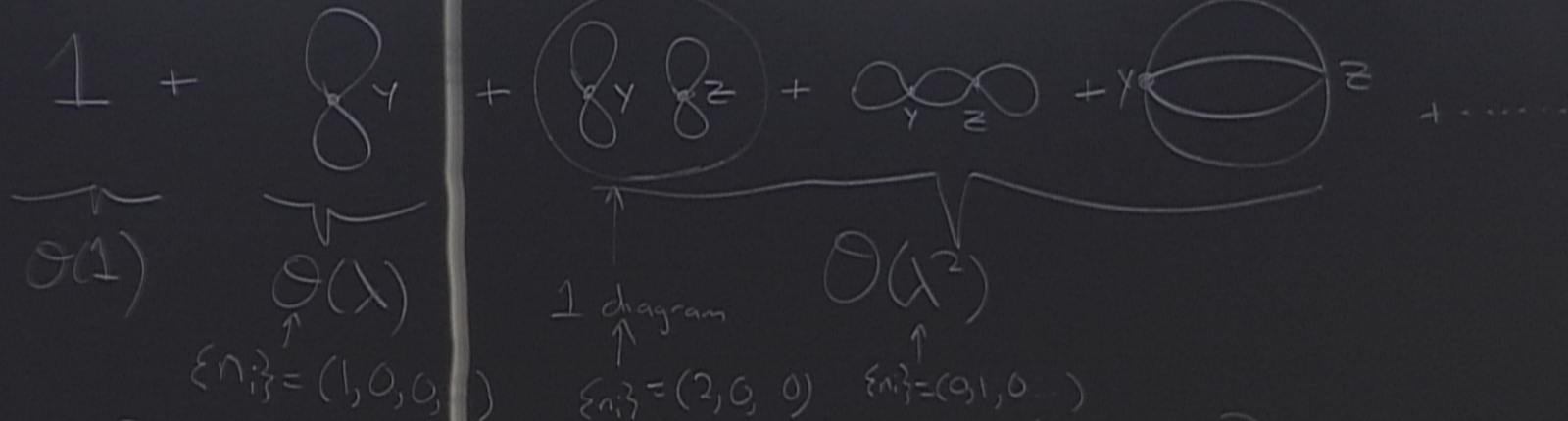
$$V_i = \{\infty, \infty, \ominus, \dots\}$$

$$= \left(\sum_{n_1} \frac{1}{n_1!} V_1^{n_1} \right) \left(\sum_{n_2} \frac{1}{n_2!} V_2^{n_2} \right) \dots \left(\sum_{n_i} \frac{1}{n_i!} V_i^{n_i} \right) \dots$$

$$= \prod_i \left(\sum_{n_i} \frac{1}{n_i!} V_i^{n_i} \right)$$

$$= \prod_i \exp(V_i) = \exp\left(\sum_i V_i\right) = \exp(\infty + \infty + \dots)$$

$$\langle 0 | T \exp \{ -i \int dt H_{int} \} | 0 \rangle =$$



$$e^{\text{loop}} = 1 + \text{loop} + \frac{1}{2} (\text{loop})^2 + \frac{1}{3!} (\text{loop})^3$$

Extra symmetry factor for two copies of same diagram

$$\{C_i\} = \left\{ \begin{array}{c} \text{---} \\ \text{---} \end{array} \right. \left. \begin{array}{c} \text{---} \\ \text{---} \end{array} \right\}$$

$$\text{numerator} = \sum_j \sum_{\{n_i\}} \prod_i C_j \frac{1}{n_i!} (V_i)^{n_i}$$

$$= \sum_j C_j \sum_{\{n_i\}} \prod_i \frac{1}{n_i!} (V_i)^{n_i}$$

$$= \sum_j C_j \exp\left(\sum_i V_i\right)$$

Vacuum diagrams cancel!

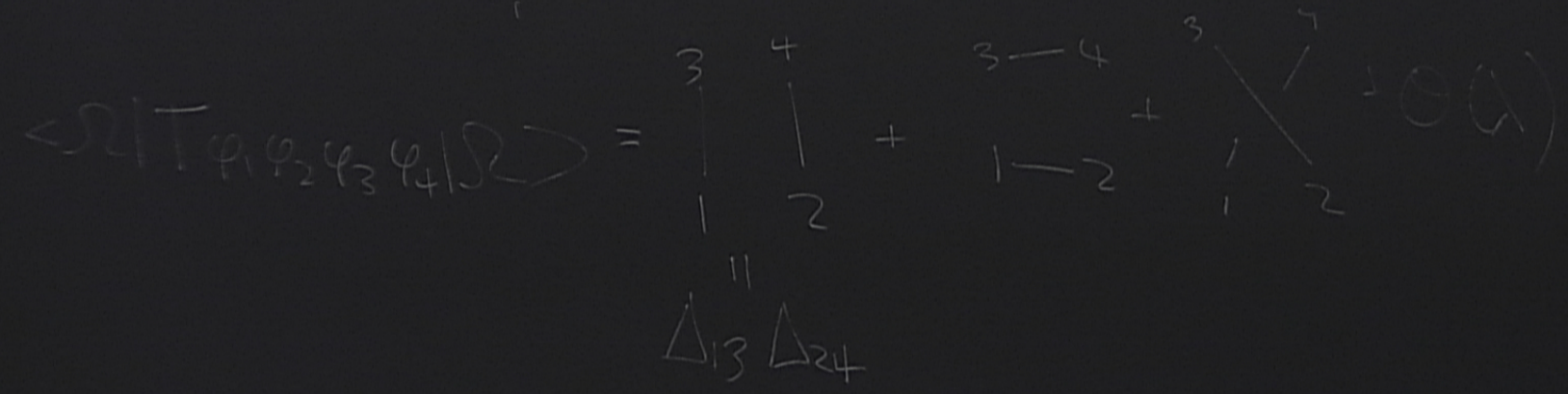
$$= \frac{1}{i^n} \exp(iV_1) = \exp\left(\frac{1}{i} V_1\right) = \exp(-iV_1) = \dots$$

medan

$\langle S | T \varphi_1 \dots \varphi_n | S \rangle =$ sum of all connected n-point functions w/ n external po

Scattering Amplitudes $2 \rightarrow 2$

$$\langle f | S | i \rangle = i^4 \prod_i \int dx_i e^{i \lambda p_i x_i} (\partial_i^2 + m^2) \langle S | T \varphi_1 \dots \varphi_4 | S \rangle$$



Vacuum diagrams cancel!

$$(\not{\partial}_1^2 + m^2)(\not{\partial}_3^2 + m^2) \Delta_F(x_1 - x_3) = F(x_1 - x_3)$$

$$\int d^4 x_1 d^4 x_3 e^{i(p_1 \cdot x_1 - p_3 \cdot x_3)} F(x_1 - x_3) = \frac{1}{2} \int d^4 x_{13} d^4 \bar{x}_{13} e^{i(\bar{p}_{13} \cdot x_{13} + p_{13} \cdot \bar{x}_{13})} F(\bar{x}_{13})$$

$$= \frac{1}{2} \int d^4 x_{13} e^{i\bar{p}_{13} \cdot x_{13}} \hat{F}(p_{13})$$

$$= (2\pi)^4 \delta^4(p_1 - p_3) \hat{F}\left(\frac{p_1 + p_3}{2}\right)$$

$$x_{13} = x_1 + x_3$$

$$\bar{x}_{13} = x_1 - x_3$$

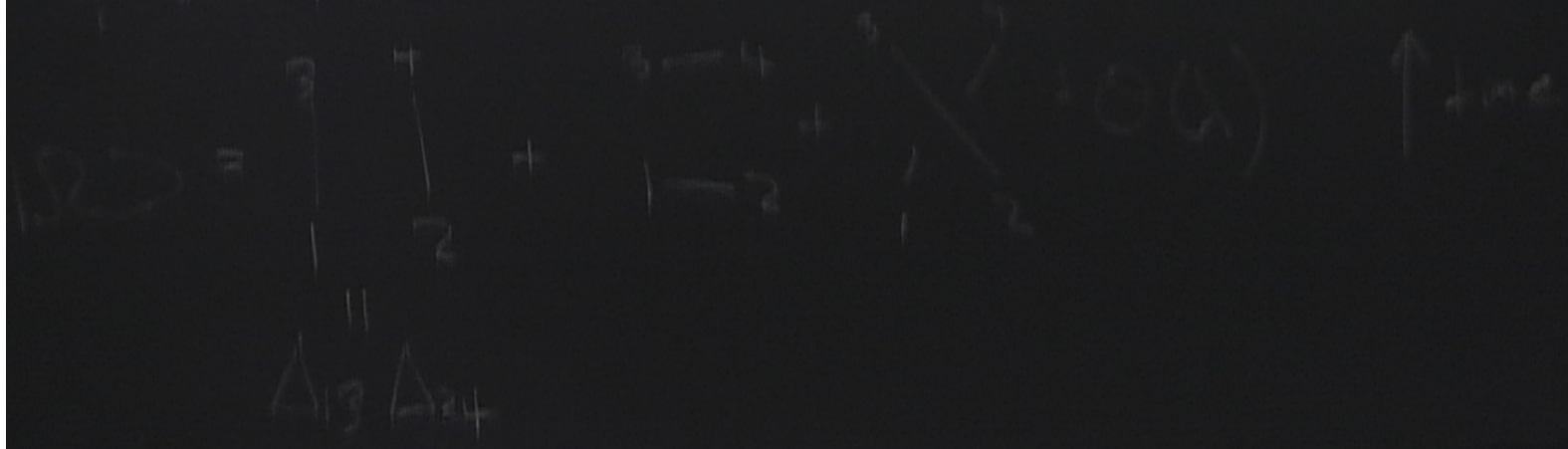
$$p_{13} = \frac{p_1 + p_3}{2}$$

$$\bar{p}_{13} = \frac{p_1 - p_3}{2}$$

Z_n = sum of all connected n-point functions w/ n external points

amplitudes $2 \rightarrow 2$

$$\prod (dx_i e^{i x p_i}) (\partial_i^2 + m^2) \langle \Omega | T \varphi_1 \dots \varphi_n | \Omega \rangle$$



$$\langle f|S|f \rangle_{11} \propto \delta(p_1 - p_3) \delta(p_2 - p_4)$$

same for \bar{X}

$$\langle f|S|f \rangle = \propto \delta^{(4)}(p_1 + p_2) \delta^{(4)}(p_3 + p_4)$$

$$p_1^0 + p_2^0 \geq m + m > 0$$

$$\langle f | S | i \rangle_{11} \propto \delta(p_1 - p_3) \delta(p_2 - p_4)$$

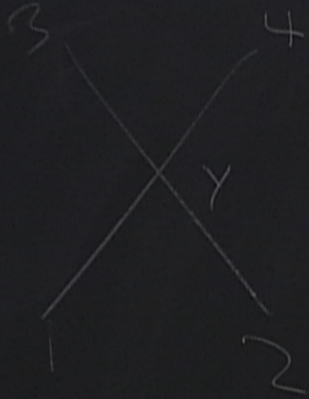
same for \times

$$\langle f | S | i \rangle = \propto \delta^{(+)}(p_1 + p_2) \delta^{(+)}(p_3 + p_4)$$

$$p_1^0 + p_2^0 \geq m + m > 0$$

same for $|k\rangle$ or $|p\rangle$

only fully connected diagrams contribute to scattering



$$= -i\lambda \int d^4y \Delta_{1y} \Delta_{2y} \Delta_{3y} \Delta_{4y}$$

$$(\partial_1^2 + m^2) \Delta_{1y} = -i \delta^{(4)}(x_1 - y)$$

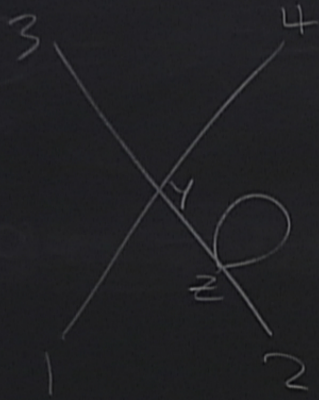
$$\langle f | S | i \rangle_X = -i\lambda \int d^4y e^{i(p_1 + p_2 - p_3 - p_4)y}$$

$$= -i\lambda (2\pi)^4 \delta^4(p_1 + p_2 - p_3 - p_4)$$

$$iM = -i\lambda$$

$$\langle f | iS | i \rangle_{CC} = iM (2\pi)^4 \delta^4(\Sigma p)$$

Completely
Correct



$\langle S | S | i \rangle_{\mathbb{R}}$

$$= \frac{(-i\lambda)^2}{2} \int d^4 y d^4 z \Delta_{2z} \Delta_{zz} \Delta_{1y} \Delta_{3y} \Delta_{1y} \Delta_{yz}$$

$$= \frac{(-i\lambda)^2}{2} \int d^4 y d^4 z e^{i(P_1 - P_3 - P_4) \cdot y} e^{i P_2 \cdot z} \Delta_{zz} \Delta_{yz}$$

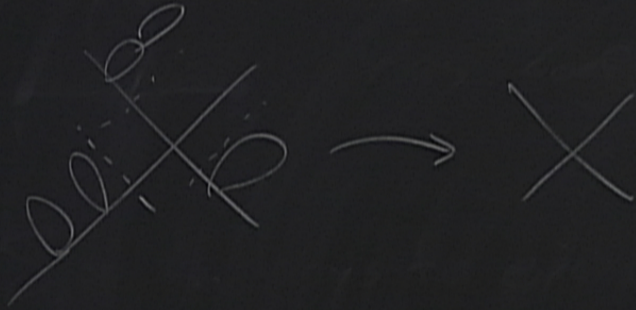
$$\Delta_{yz} = \int \frac{d^4 k}{(2\pi)^4} \frac{i}{k^2 - m^2 + i\epsilon} e^{i k \cdot (y - z)}$$

$$\langle f | S | i \rangle_{\text{tree}} = \frac{(-i\lambda)^2}{2} \int \frac{d^4 k}{(2\pi)^4} \int \frac{d^4 q}{(2\pi)^4} \frac{i}{k^2 - m^2 + i\epsilon} \frac{i}{q^2 - m^2 + i\epsilon}$$

$$(2\pi)^4 \delta^4(p_1 + k - p_3 - p_4) \cdot (2\pi)^4 \delta^4(p_2 - k)$$

$$\int d^4 k \delta(p_2 - k) \frac{1}{k^2 - m^2 + i\epsilon} = \frac{1}{p_2^2 - m^2} = 0$$

= Need to amputate



Momentum Space Fe

Momentum Space Feynman Rules for Scattering Amplitudes

i) \mathcal{M} = sum of all completely connected, amputated diagrams

1. For each internal line $\xrightarrow{\vec{p}} = \frac{i}{p^2 - m^2 + i\epsilon}$

2. For each vertex $\times = -i\lambda$ (depends on \mathcal{H}_{int})

3. For each external line $\rightarrow = 1$ (change for fermions & gauge bosons)

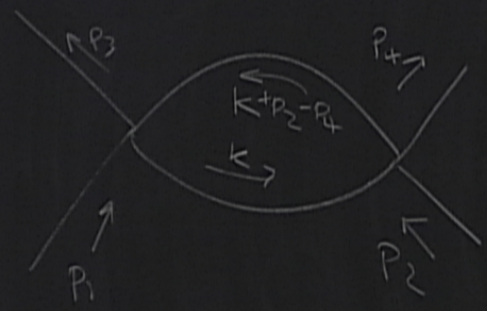
4. Impose momentum conservation at each vertex

vacuum diagrams cancel!

5. Divide by symmetry factor

6. Integrate over any undetermined momenta

denom



$$= \frac{(-i\lambda)^2}{2} \int \frac{d^4 k}{(2\pi)^4} \frac{1}{k^2 - m^2 + i\epsilon} \frac{1}{(k+p_2-p_4)^2 - m^2 + i\epsilon}$$