Title: AMATH 875/PHYS 786 - Fall 2015 - Lecture 12

Date: Oct 26, 2015 01:30 PM

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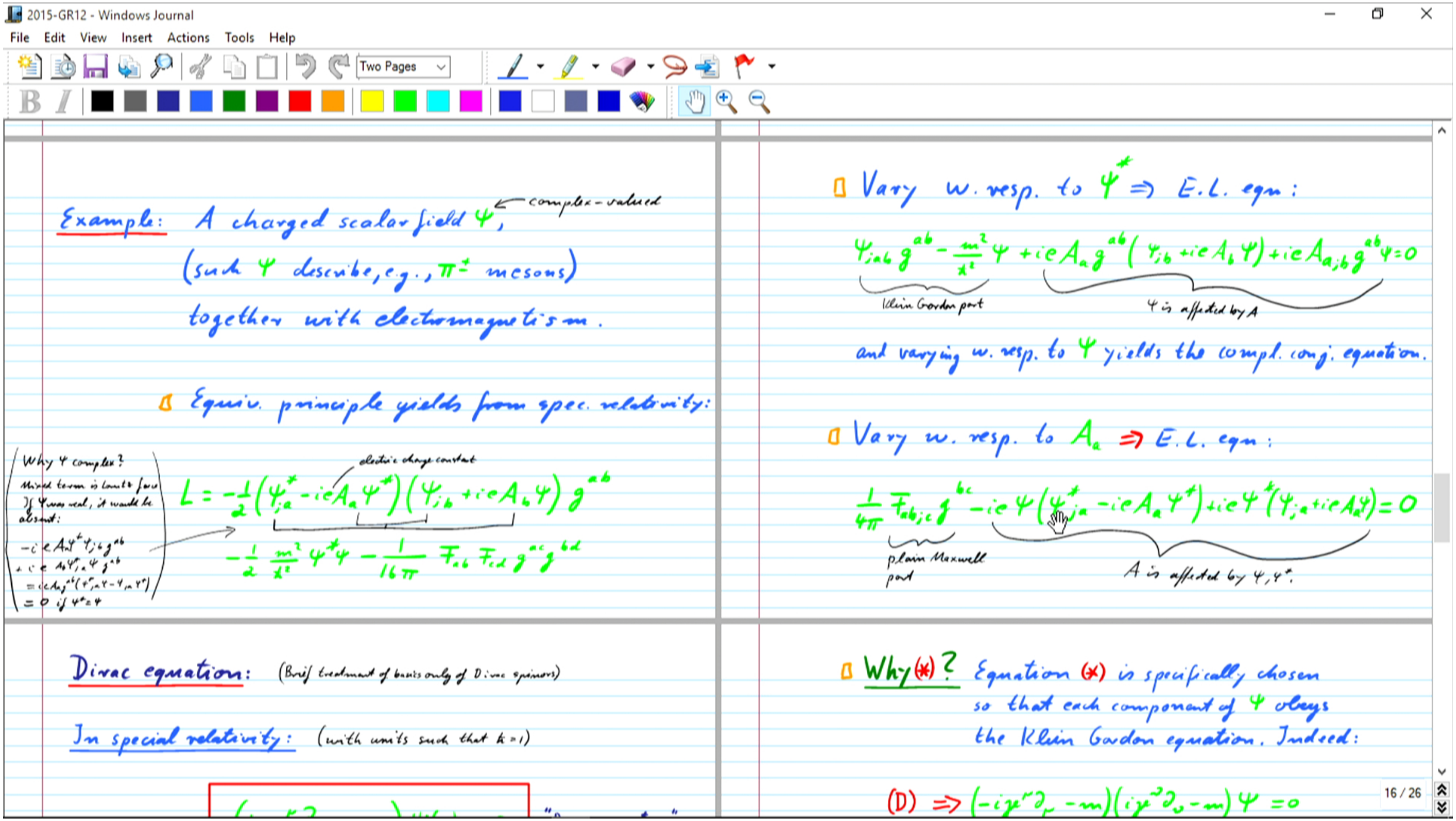
Abstract: <p>Course Description coming soon.</p>



Example: A charged scalar field ψ, ← complex-valued

(such ψ describe, e.g., $\pi^\pm$ mesons)

together with electromagnetism.

Equiv. principle yields from spec. relativity:

$$L = -\frac{1}{2}\left(\psi^*_{;a} - i\overset{\text{electric charge constant}}{e}A_a\psi^*\right)\left(\psi_{;b} + ieA_b\psi\right)g^{ab} - \frac{1}{2}\frac{m^2}{\hbar^2}\psi^*\psi - \frac{1}{16\pi}F_{ab}F_{cd}g^{ac}g^{bd}$$

Why ψ complex?
Mixed term is hard to form
If ψ was real, it would be absent:

$$-ieA_a\psi^*\psi_{;b}g^{ab} + ieA_b\psi^*_{;a}\psi g^{ab} = ieA_b g^{ab}(\psi^*_{;a}\psi - \psi_{;a}\psi^*) = 0 \text{ if } \psi^* = \psi$$

Vary w. resp. to ψ^* ⇒ E.L. eqn:

$$\underbrace{\psi_{;ab}g^{ab} - \frac{m^2}{\hbar^2}\psi}_{\text{Klein Gordon part}} \underbrace{+ieA_a g^{ab}(\psi_{;b} + ieA_b\psi) + ieA_{a;b}g^{ab}\psi}_{\psi \text{ is affected by } A} = 0$$

and varying w. resp. to ψ yields the compl. conj. equation.

Vary w. resp. to A_a ⇒ E.L. eqn:

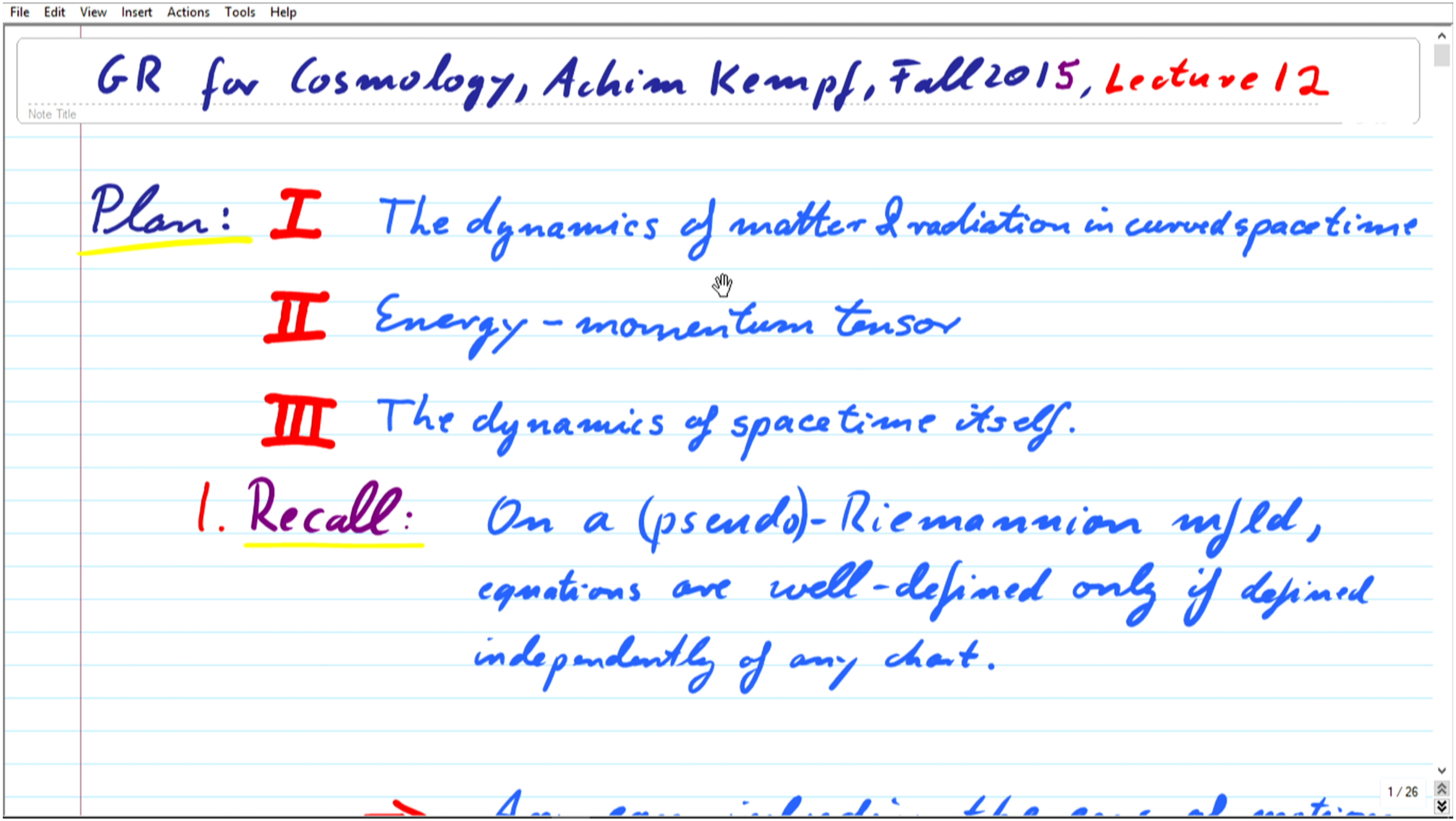
$$\underbrace{\frac{1}{4\pi}F_{ab;c}g^{bc}}_{\text{plain Maxwell part}} \underbrace{- ie\psi(\psi^*_{;a} - ieA_a\psi^*) + ie\psi^*(\psi_{;a} + ieA_a\psi)}_{A \text{ is affected by } \psi, \psi^*} = 0$$

Dirac equation: (Brief treatment of basics only of Dirac spinors)

In special relativity: (with units such that $\hbar = 1$)

Why (*)? Equation (*) is specifically chosen so that each component of ψ obeys the Klein Gordon equation. Indeed:

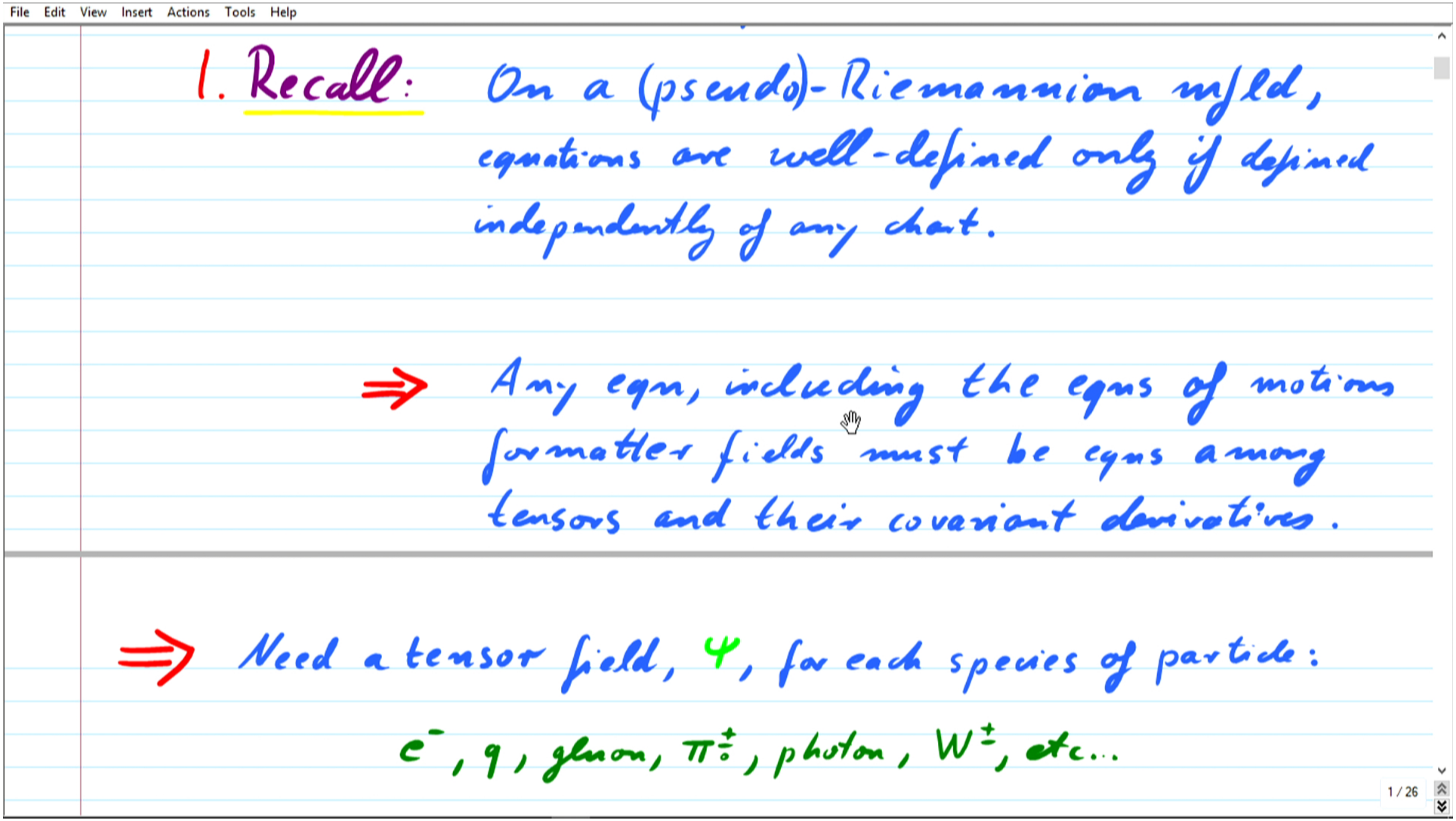
$$(D) \Rightarrow (-i\gamma^\mu\partial_\mu - m)(i\gamma^\nu\partial_\nu - m)\psi = 0$$



GR for Cosmology, Achim Kempf, Fall 2015, Lecture 12

Plan:
- **I** The dynamics of matter & radiation in curved spacetime
- **II** Energy-momentum tensor
- **III** The dynamics of spacetime itself.

1. Recall: On a (pseudo)-Riemannian mfld, equations are well-defined only if defined independently of any chart.

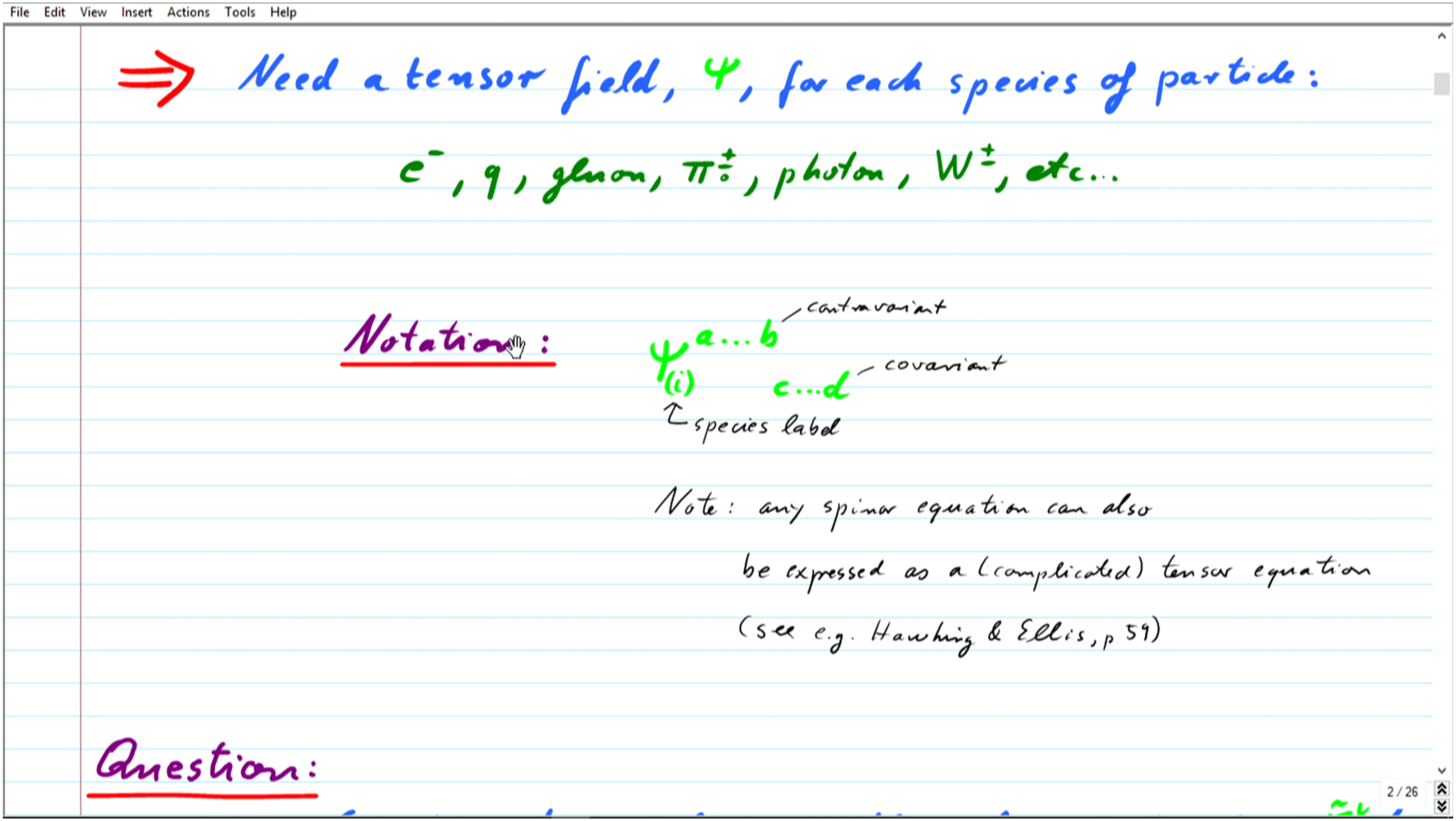


1. Recall: On a (pseudo)-Riemannian mfld, equations are well-defined only if defined independently of any chart.

⇒ Any eqn, including the eqns of motions for matter fields must be eqns among tensors and their covariant derivatives.

⇒ Need a tensor field, Ψ, for each species of particle:

e^-, q, gluon, $\pi^{\substack{\pm\\0}}$, photon, $W^\pm$, etc...



$\Rightarrow$ Need a tensor field, Ψ, for each species of particle:

e^-, q, gluon, $\pi^{\substack{\pm\\0}}$, photon, $W^{\pm}$, etc...

Notation:

$$\Psi_{(i)}{}^{a\ldots b}{}_{c\ldots d}$$

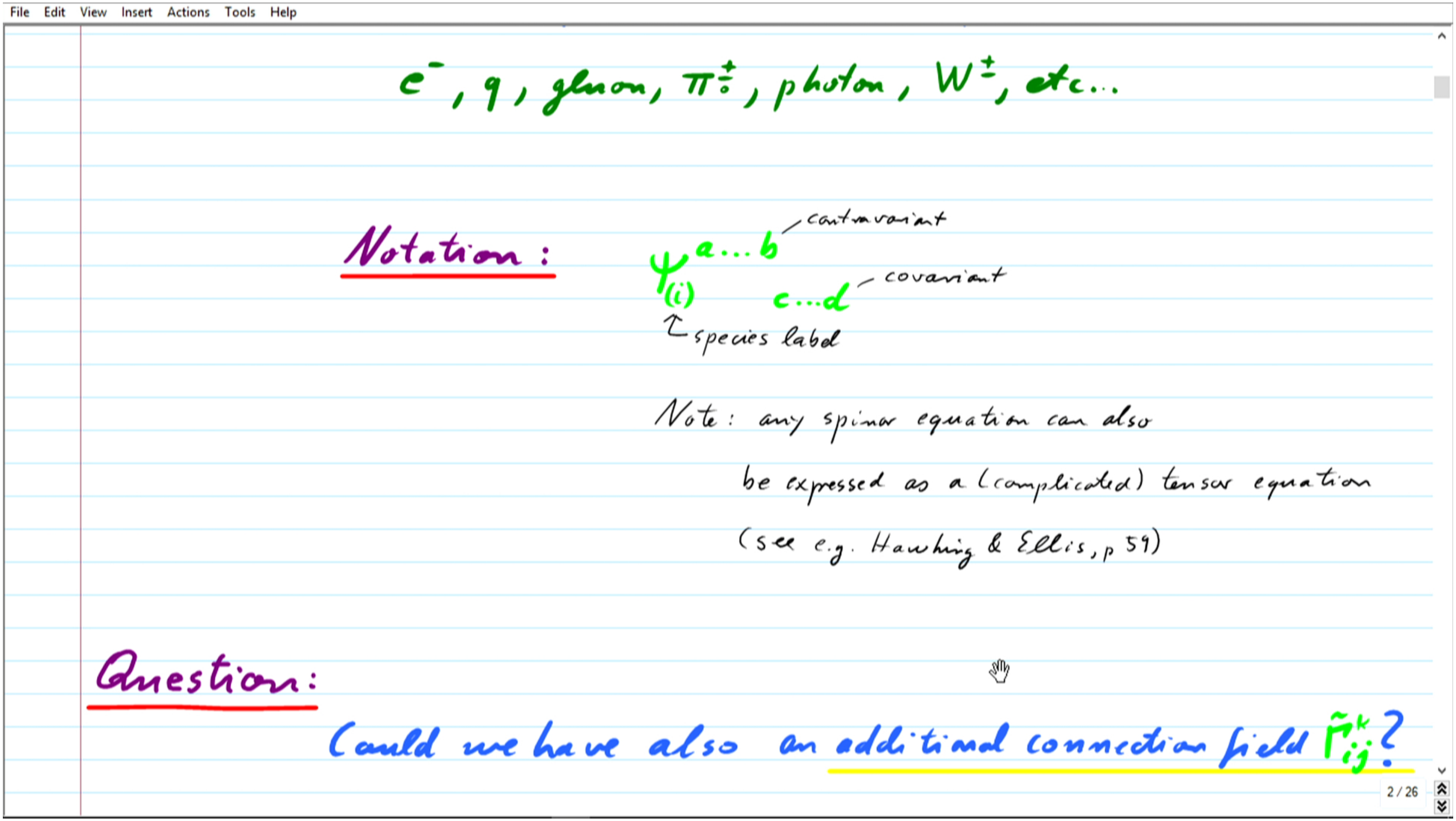
contravariant

covariant

species label

Note: any spinor equation can also be expressed as a (complicated) tensor equation (see e.g. Hawking & Ellis, p 59)

Question:



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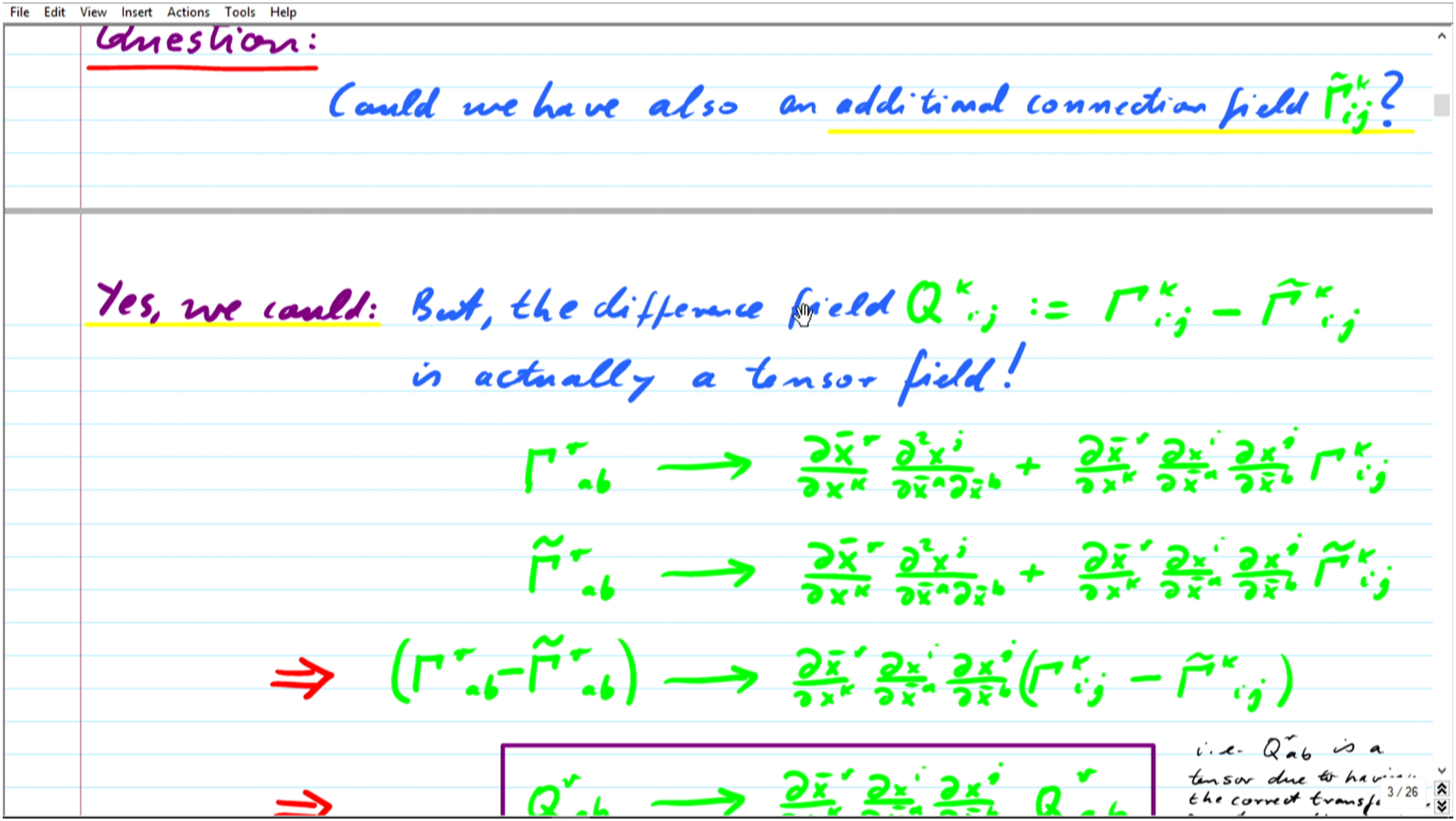
Question:

Could we have also an additional connection field $\tilde{\Gamma}^k_{ij}$?



Question:

Could we have also an additional connection field $\tilde{\Gamma}^{k}_{ij}$?

Yes, we could: But, the difference field $Q^{k}{}_{ij} := \Gamma^{k}{}_{ij} - \tilde{\Gamma}^{k}{}_{ij}$ is actually a tensor field!

$$\Gamma^{r}{}_{ab} \longrightarrow \frac{\partial \bar{x}^{r}}{\partial x^{k}} \frac{\partial^2 x^{j}}{\partial \bar{x}^{a} \partial \bar{x}^{b}} + \frac{\partial \bar{x}^{r}}{\partial x^{k}} \frac{\partial x^{i}}{\partial \bar{x}^{a}} \frac{\partial x^{j}}{\partial \bar{x}^{b}} \Gamma^{k}{}_{ij}$$

$$\tilde{\Gamma}^{r}{}_{ab} \longrightarrow \frac{\partial \bar{x}^{r}}{\partial x^{k}} \frac{\partial^2 x^{j}}{\partial \bar{x}^{a} \partial \bar{x}^{b}} + \frac{\partial \bar{x}^{r}}{\partial x^{k}} \frac{\partial x^{i}}{\partial \bar{x}^{a}} \frac{\partial x^{j}}{\partial \bar{x}^{b}} \tilde{\Gamma}^{k}{}_{ij}$$

$$\Rightarrow \left(\Gamma^{r}{}_{ab} - \tilde{\Gamma}^{r}{}_{ab}\right) \longrightarrow \frac{\partial \bar{x}^{r}}{\partial x^{k}} \frac{\partial x^{i}}{\partial \bar{x}^{a}} \frac{\partial x^{j}}{\partial \bar{x}^{b}} \left(\Gamma^{k}{}_{ij} - \tilde{\Gamma}^{k}{}_{ij}\right)$$

$$\Rightarrow \boxed{Q^{r}{}_{ab} \longrightarrow \frac{\partial \bar{x}^{r}}{\partial x^{k}} \frac{\partial x^{i}}{\partial \bar{x}^{a}} \frac{\partial x^{j}}{\partial \bar{x}^{b}} Q^{r}{}_{ab}}$$

i.e. $Q^{r}{}_{ab}$ is a tensor due to havi
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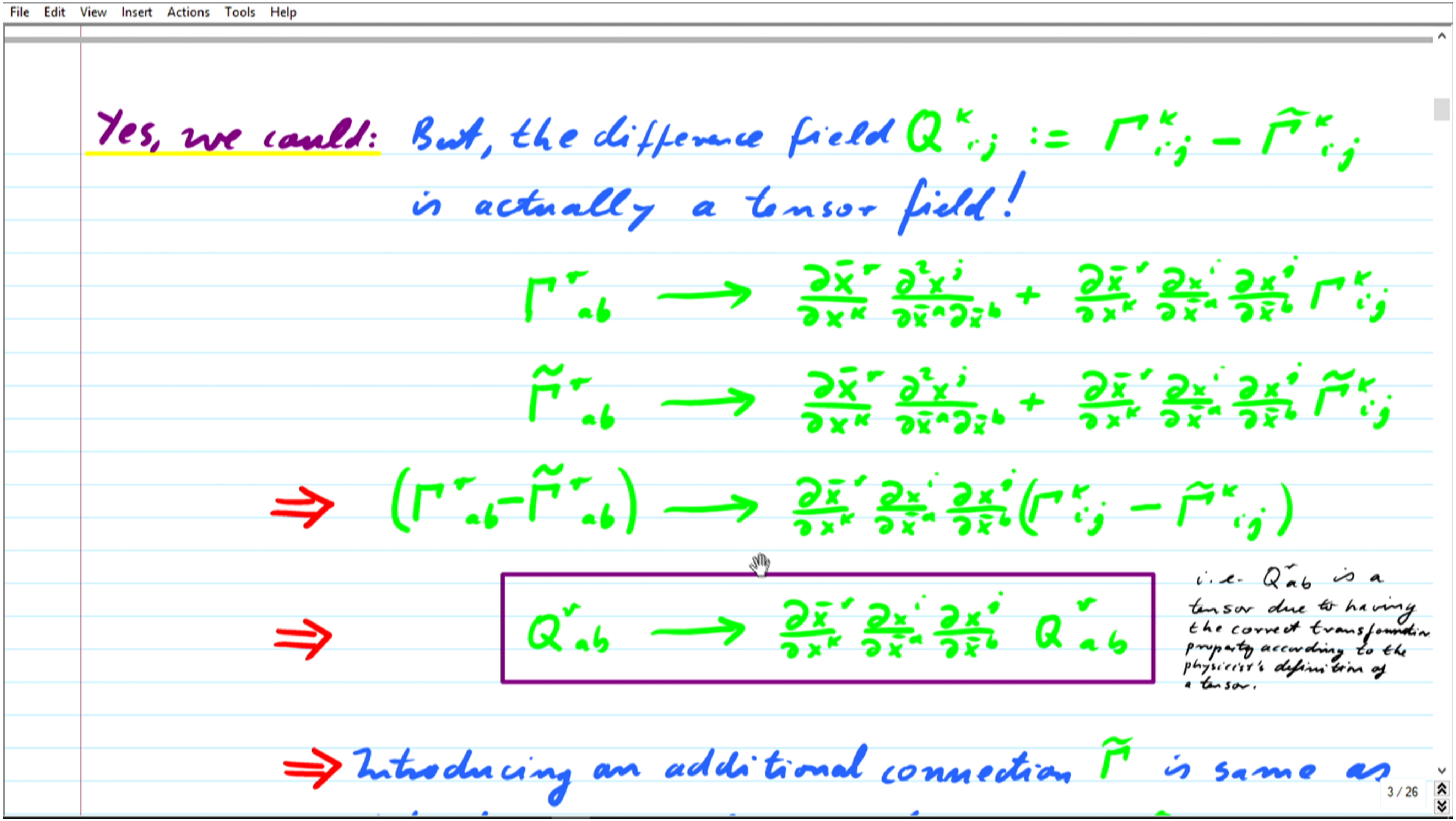

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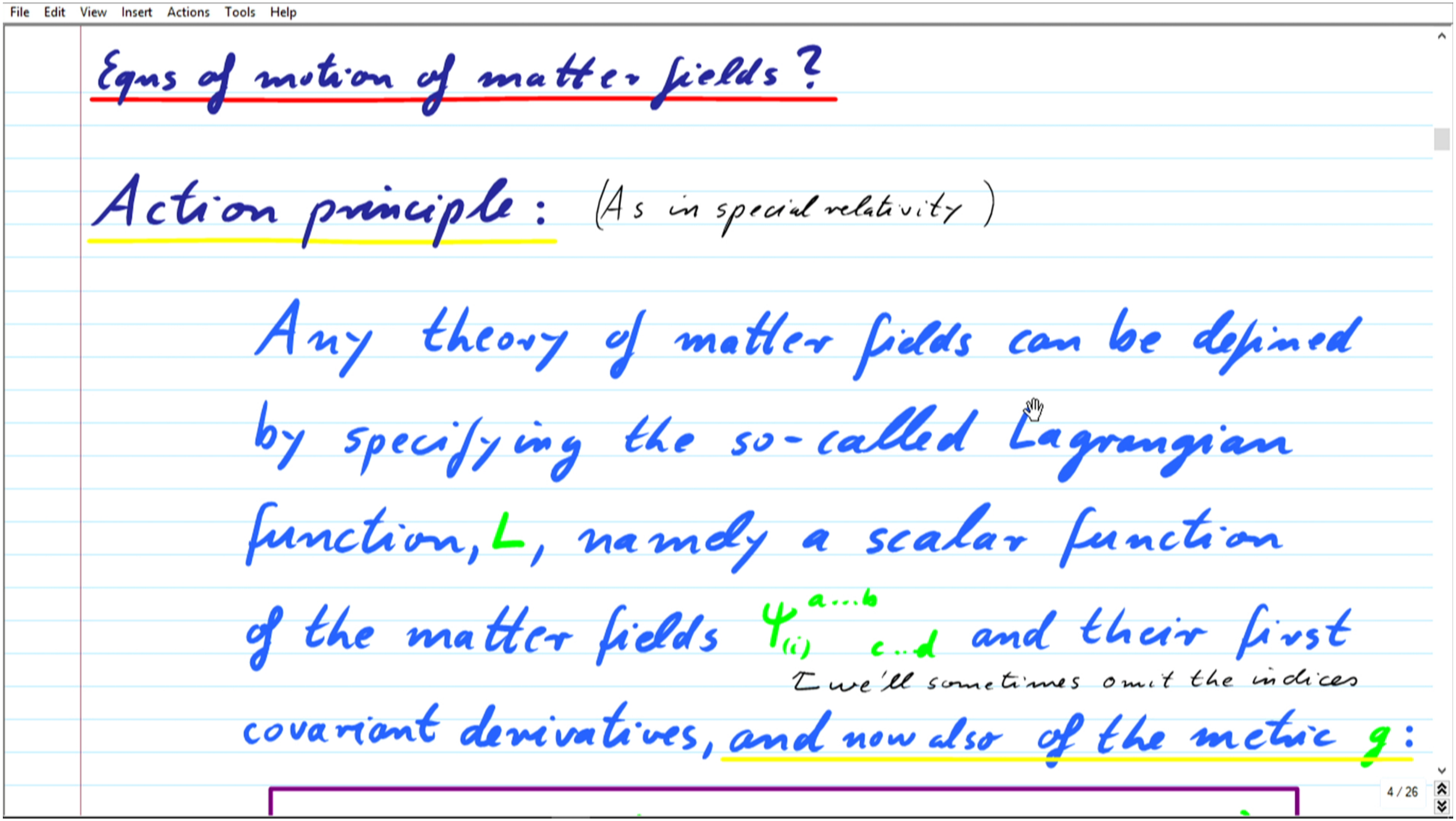
i.e. $Q^{r}{}_{ab}$ is a tensor due to having the correct transformation property according to the physicist's definition of a tensor.

$\Rightarrow$ Introducing an additional connection $\tilde{\Gamma}$ is same as

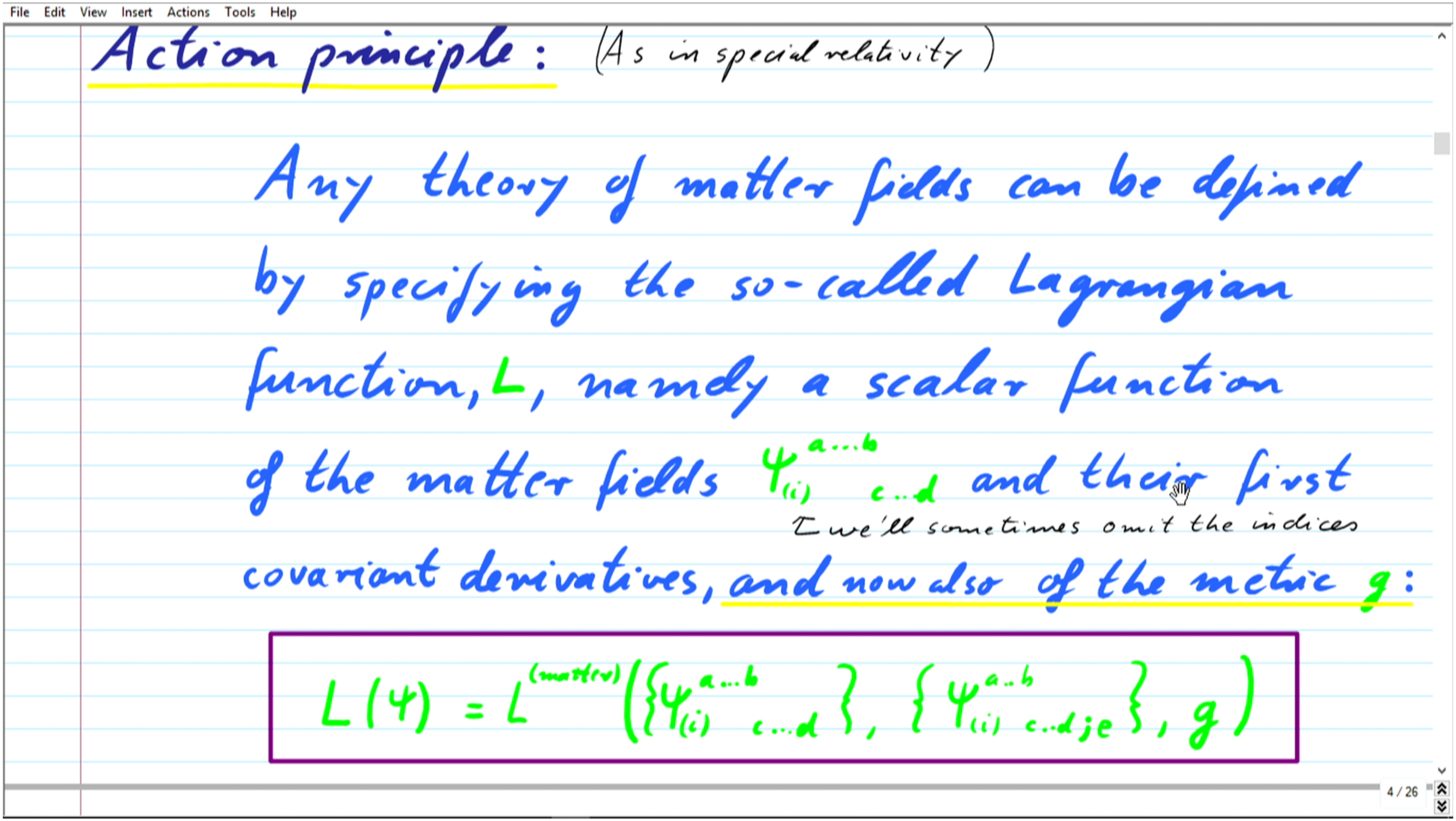


Equs of motion of matter fields?

Action principle: (As in special relativity)

Any theory of matter fields can be defined by specifying the so-called Lagrangian function, L, namely a scalar function of the matter fields $\psi_{(i)}{}^{a\dots b}{}_{c\dots d}$ and their first covariant derivatives, and now also of the metric g:

[we'll sometimes omit the indices



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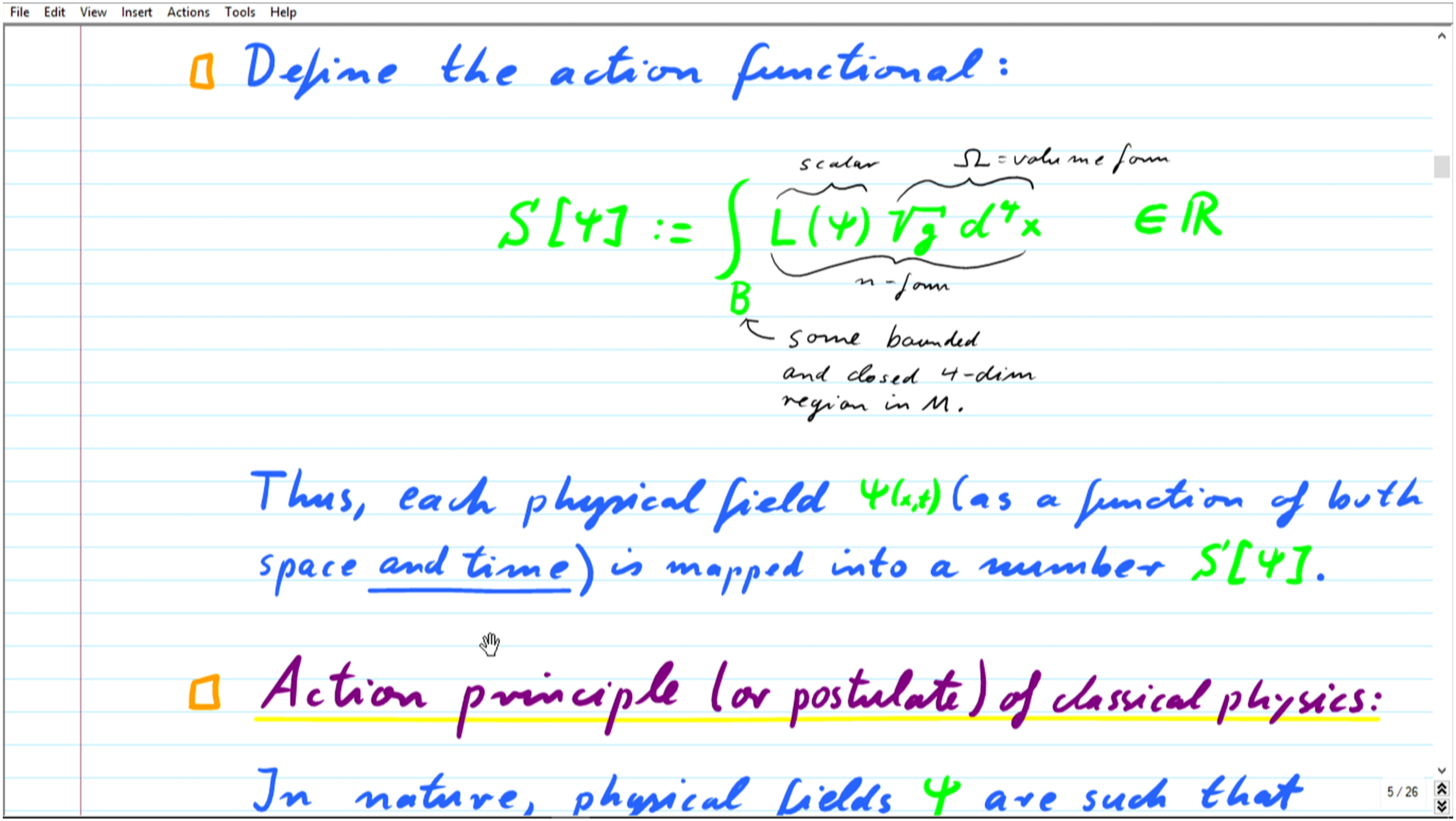
$$\boxed{L(\Psi) = L^{(matter)}\left(\left\{\Psi_{(i)}{}^{a\dots b}{}_{c\dots d}\right\}, \left\{\Psi_{(i)}{}^{a..b}{}_{c..d;e}\right\}, g\right)}$$



☐ Define the action functional:

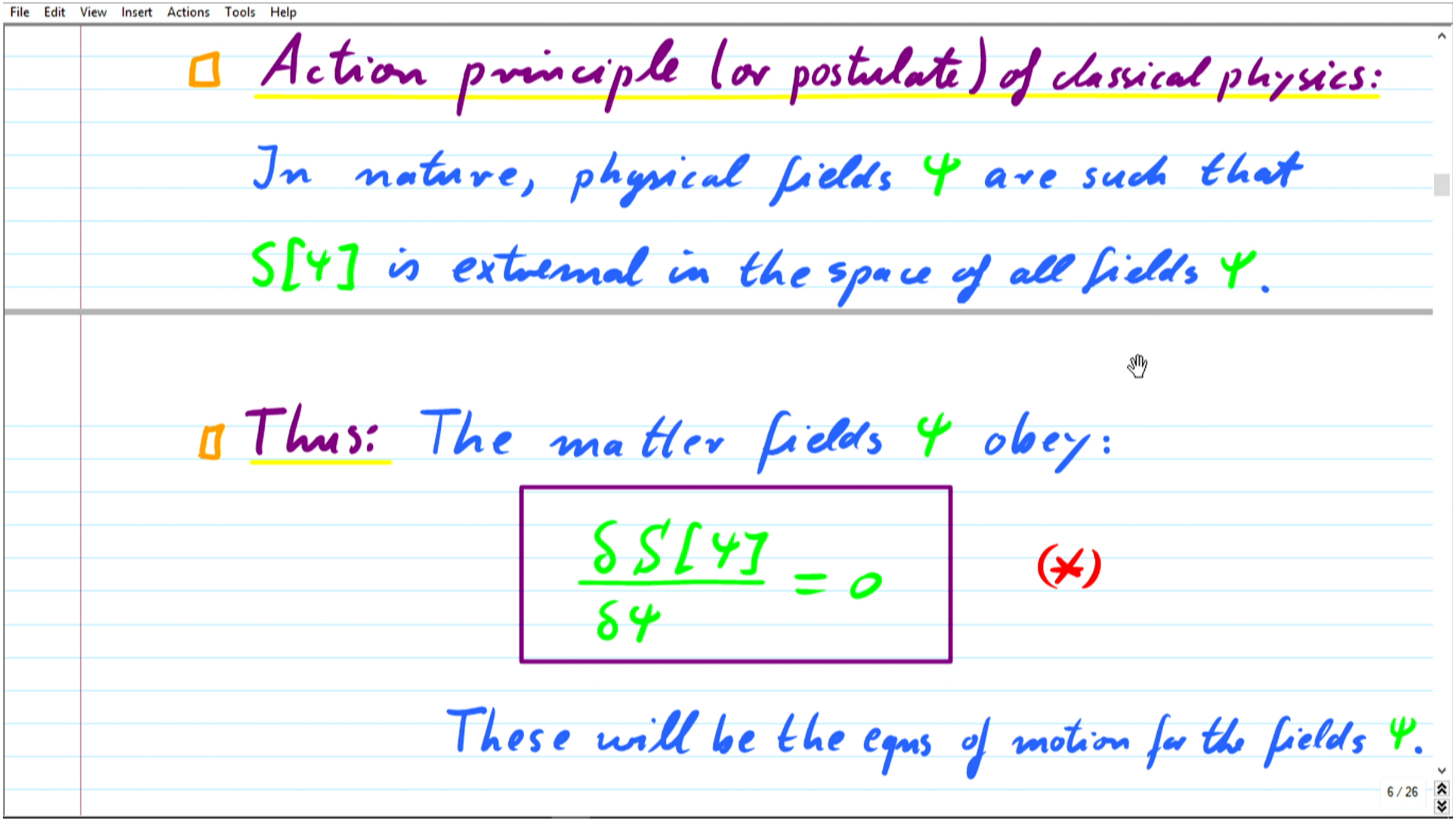
$$S[\psi] := \int_B \underbrace{\overbrace{L(\psi)}^{\text{scalar}}\ \overbrace{\sqrt{g}\, d^4x}^{\Omega = \text{volume form}}}_{n\text{-form}} \quad \in \mathbb{R}$$

B ← some bounded and closed 4-dim region in M.

Thus, each physical field $\psi(x,t)$ (as a function of both space <u>and time</u>) is mapped into a number $S[\psi]$.

☐ **Action principle (or postulate) of classical physics:**

In nature, physical fields ψ are such that



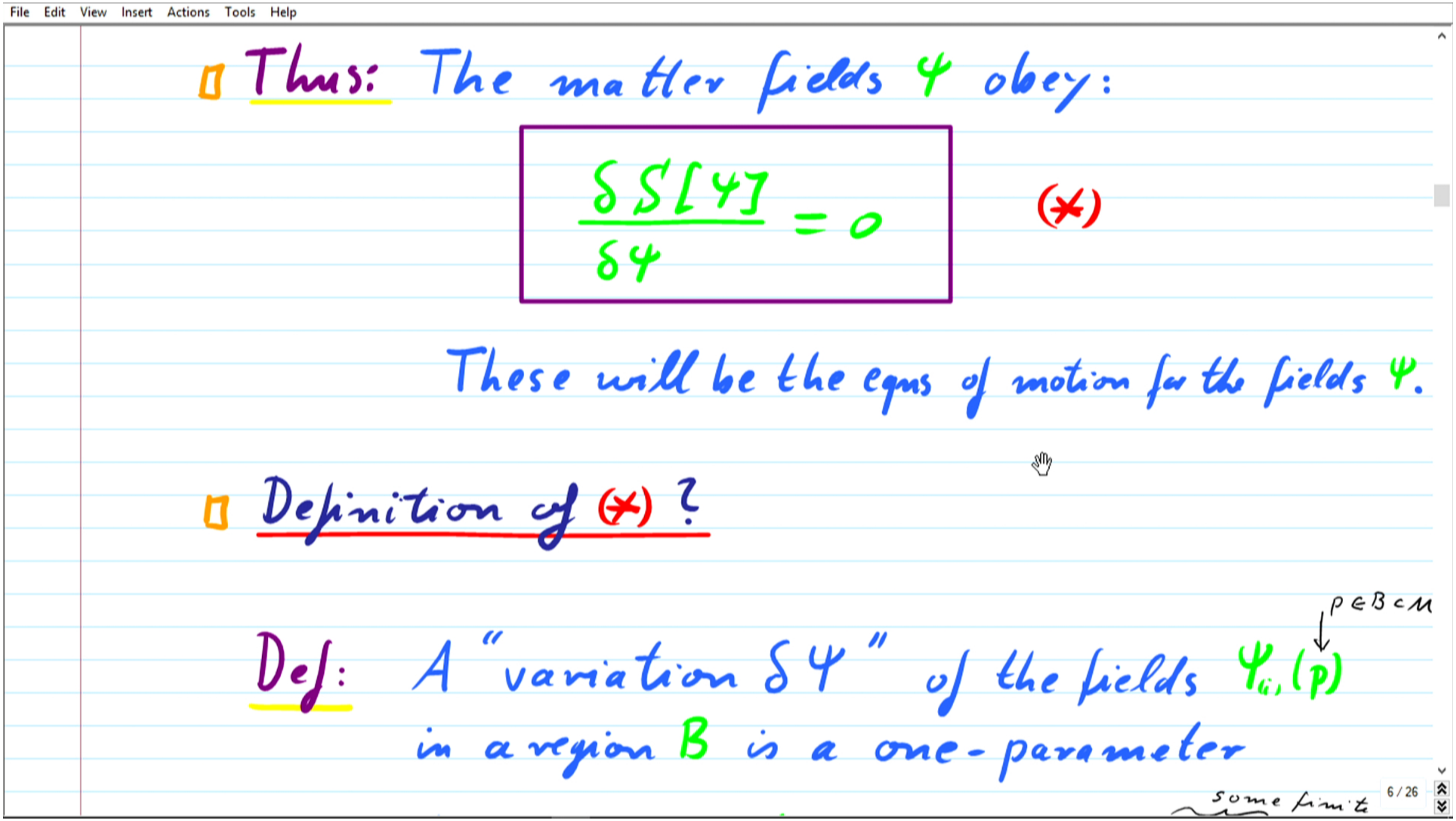
☐ **Action principle (or postulate) of classical physics:**

In nature, physical fields ψ are such that $S[\psi]$ is extremal in the space of all fields ψ.

☐ **Thus:** The matter fields ψ obey:

$$\boxed{\frac{\delta S[\psi]}{\delta \psi} = 0} \qquad (*)$$

These will be the eqns of motion for the fields ψ.



- **Thus:** The matter fields ψ obey:

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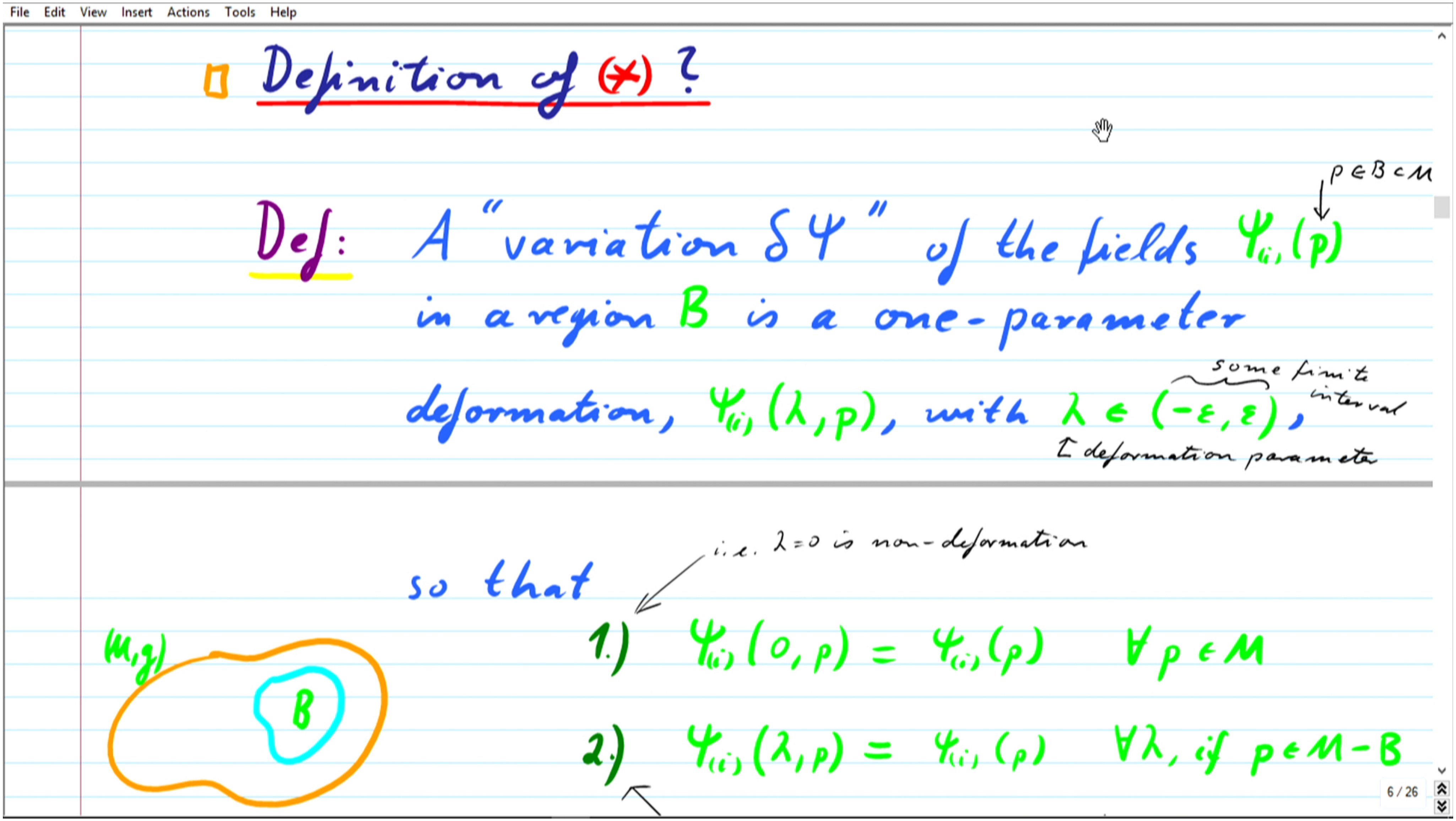
These will be the eqns of motion for the fields ψ.

- Definition of $(*)$?

Def: A "variation $\delta\psi$" of the fields $\psi_{(i)}(\overset{p \in B \subset M}{\downarrow}p)$ in a region B is a one-parameter

some finite



☐ Definition of (*) ?

p ∈ B ⊂ M

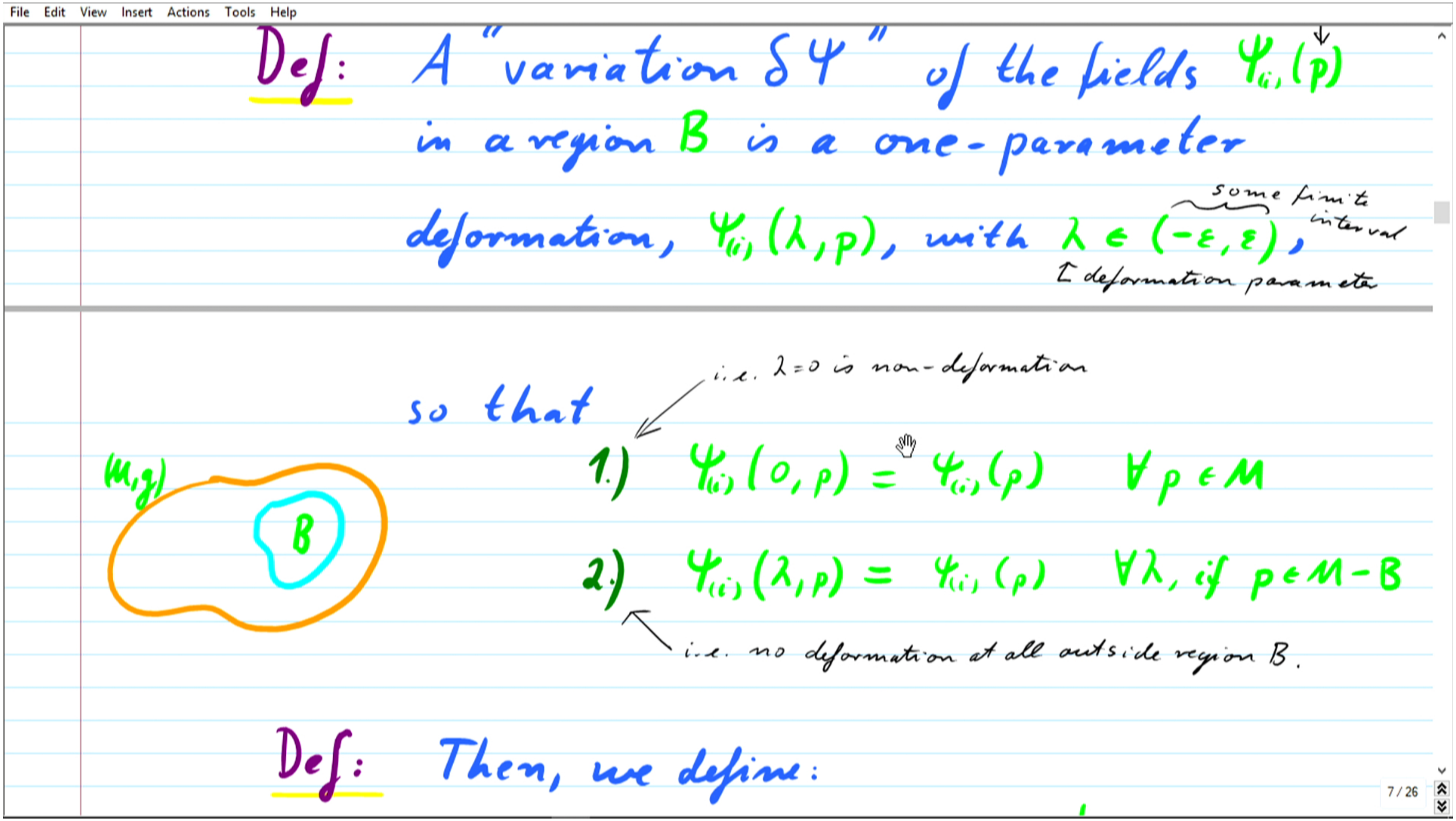
Def: A "variation $\delta\Psi$" of the fields $\Psi_{(i)}(p)$ in a region B is a one-parameter deformation, $\Psi_{(i)}(\lambda, p)$, with $\underset{\uparrow\ \text{deformation parameter}}{\lambda} \in \overbrace{(-\varepsilon, \varepsilon)}^{\text{some finite interval}}$,

so that

i.e. λ=0 is non-deformation

1) $\Psi_{(i)}(0,p) = \Psi_{(i)}(p) \quad \forall\, p \in M$

2) $\Psi_{(i)}(\lambda,p) = \Psi_{(i)}(p) \quad \forall \lambda, \text{ if } p \in M - B$



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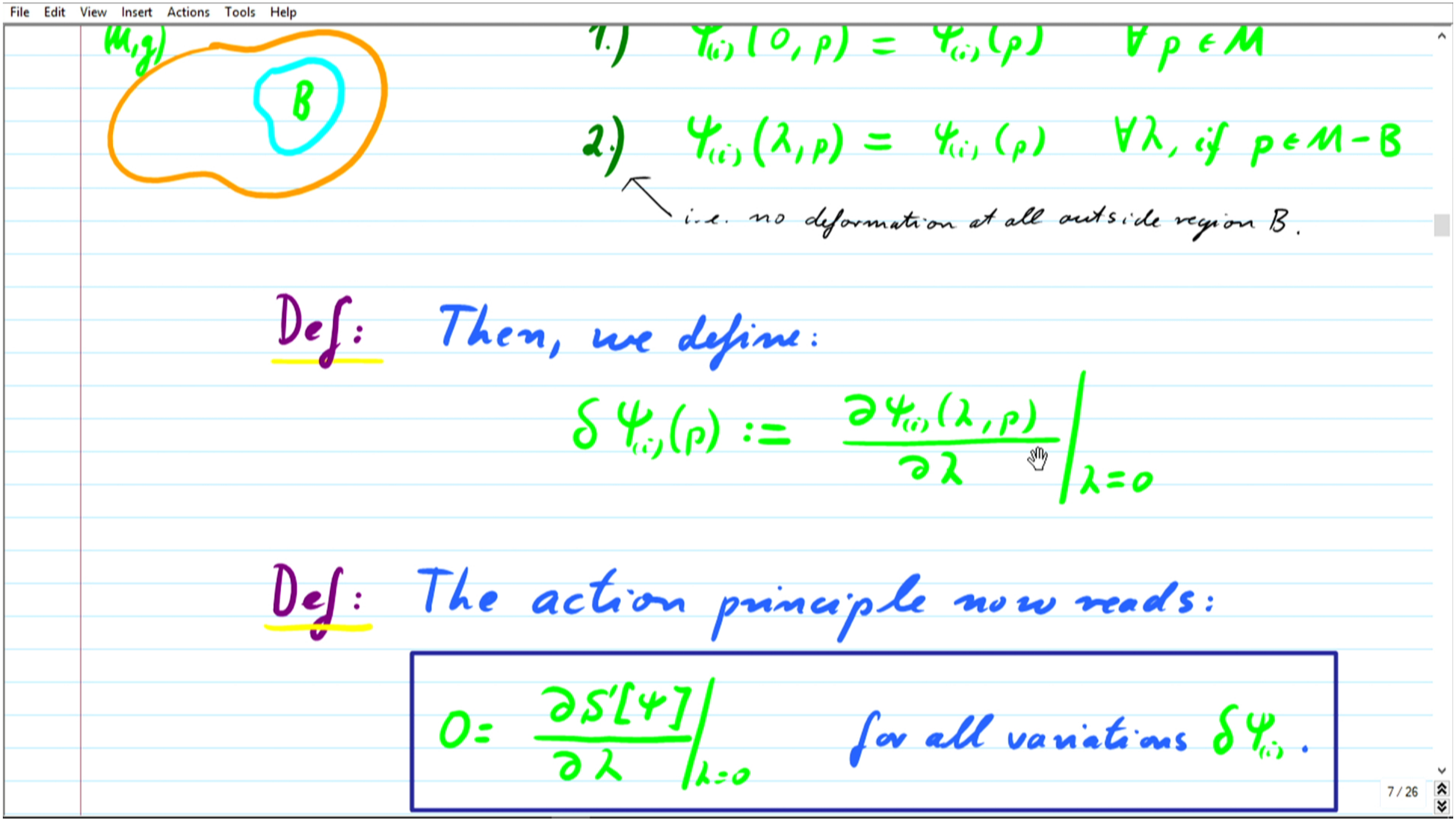
i.e. no deformation at all outside region B.

Def: Then, we define:



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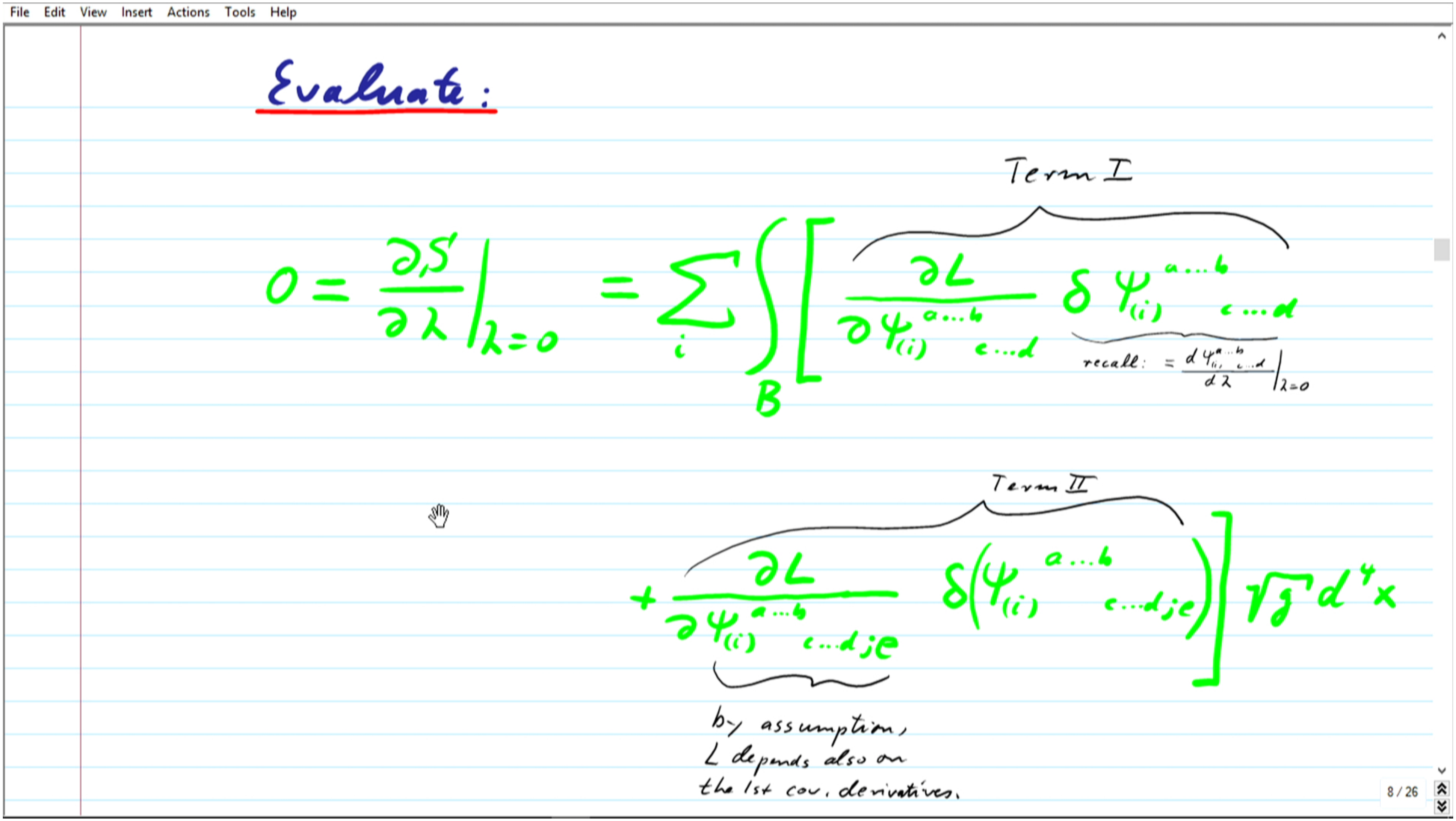
i.e. no deformation at all outside region B.

Def: Then, we define:

$$\delta\Psi_{(i)}(p) := \left.\frac{\partial \Psi_{(i)}(\lambda,p)}{\partial \lambda}\right|_{\lambda=0}$$

Def: The action principle now reads:

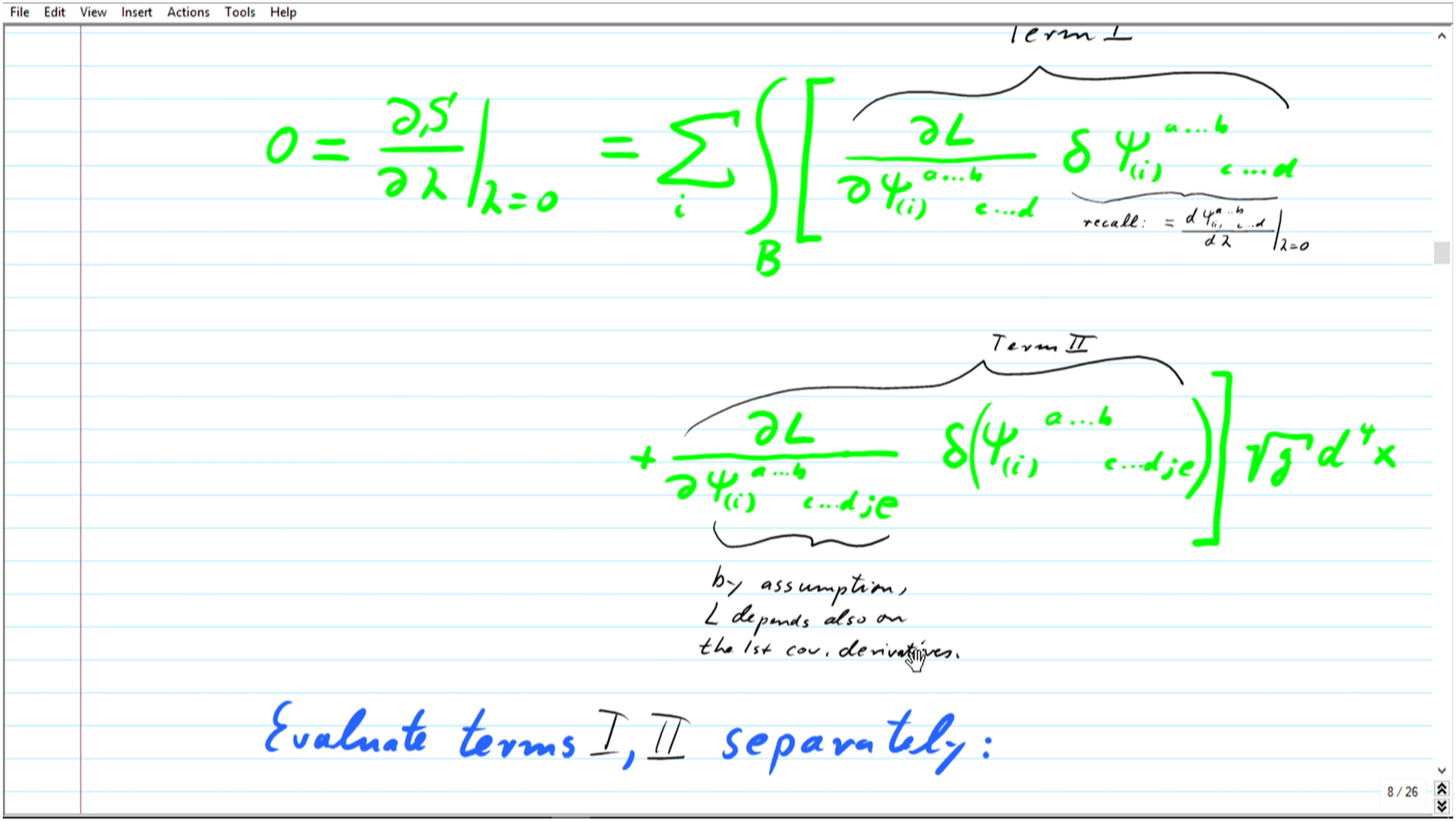
$$\boxed{0 = \left.\frac{\partial S'[\Psi]}{\partial \lambda}\right|_{\lambda=0} \quad \text{for all variations } \delta\Psi_{(i)}.}$$



Evaluate:

$$0 = \frac{\partial S}{\partial \lambda}\Big|_{\lambda=0} = \sum_i \int_B \Bigg[\overbrace{\frac{\partial L}{\partial \psi_{(i)}{}^{a\dots b}{}_{c\dots d}} \, \underbrace{\delta \psi_{(i)}{}^{a\dots b}{}_{c\dots d}}_{\text{recall: } = \frac{d\psi_{(i)}{}^{a\dots b}{}_{c\dots d}}{d\lambda}\big|_{\lambda=0}}}^{\text{Term I}}$$

$$+ \overbrace{\frac{\partial L}{\underbrace{\partial \psi_{(i)}{}^{a\dots b}{}_{c\dots d;e}}_{\substack{\text{by assumption,}\\ L \text{ depends also on}\\ \text{the 1st cov. derivatives.}}}} \, \delta\left(\psi_{(i)}{}^{a\dots b}{}_{c\dots d;e}\right)}^{\text{Term II}} \Bigg] \sqrt{g}\, d^4x$$



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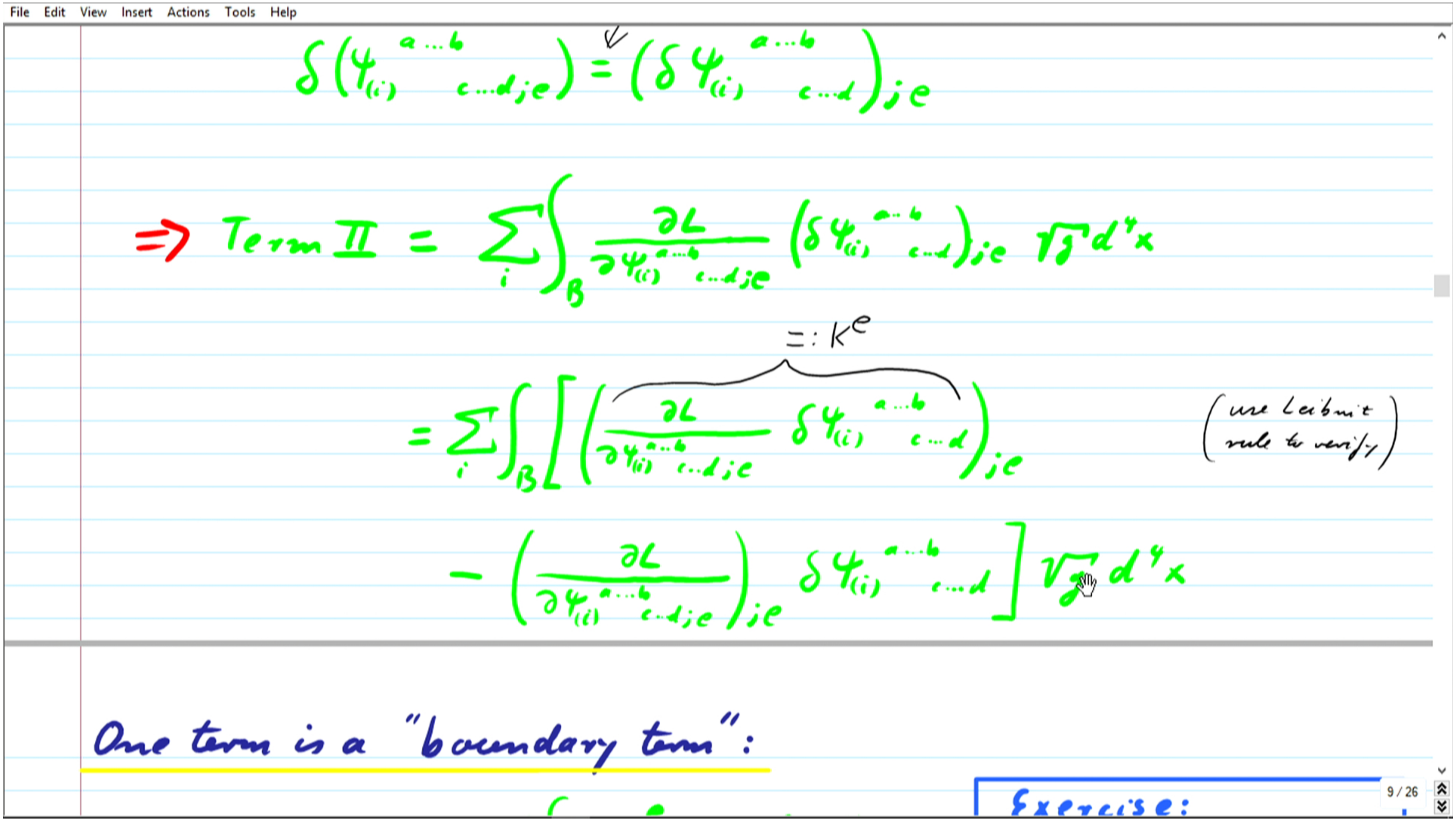
Evaluate terms $\underline{I}$, $\underline{II}$ separately:



$$\delta\left(\psi_{(i)}{}^{a\dots b}{}_{c\dots d;e}\right) \overset{\checkmark}{=} \left(\delta\psi_{(i)}{}^{a\dots b}{}_{c\dots d}\right)_{;e}$$

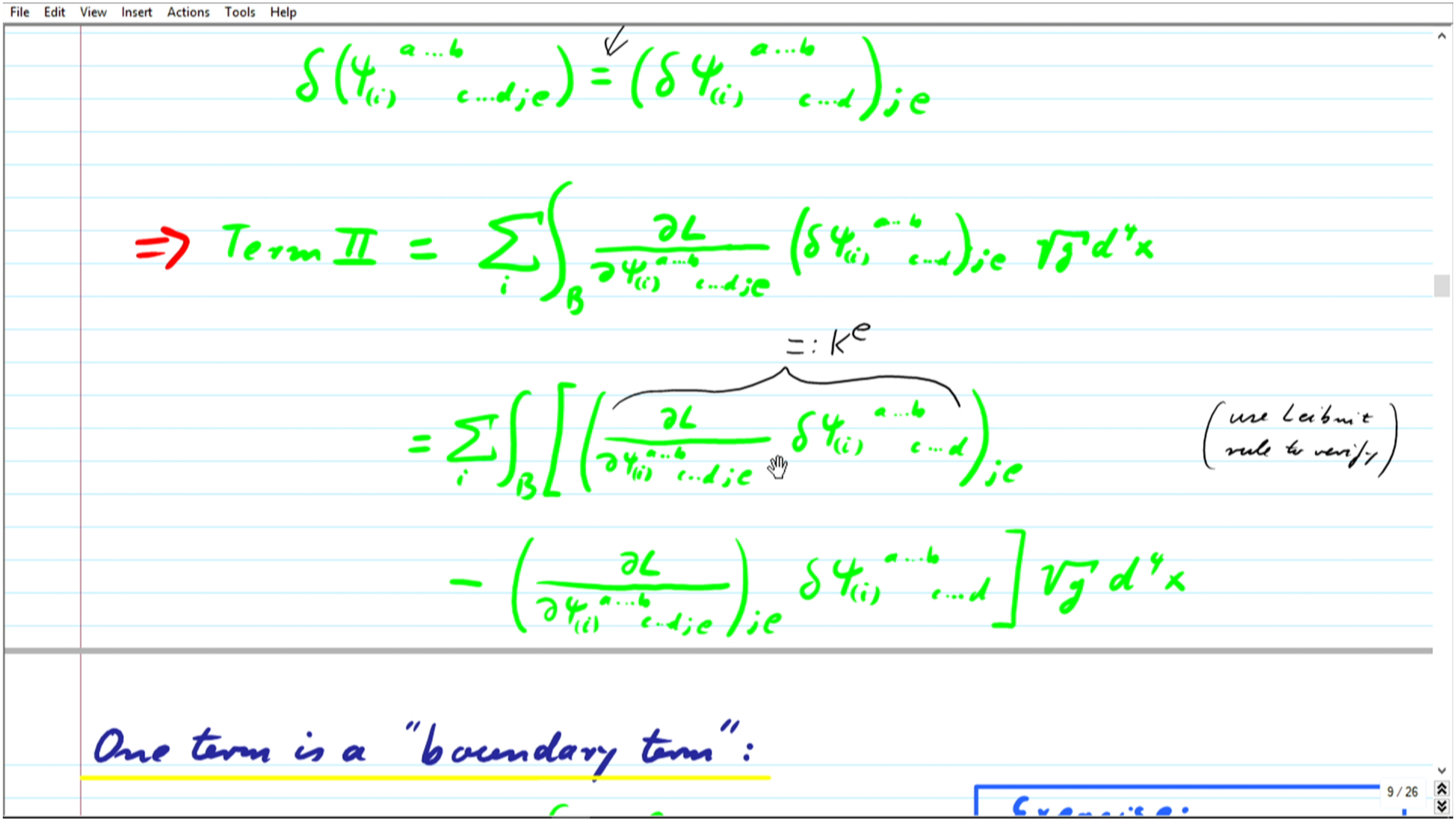
$$\Rightarrow \text{Term II} = \sum_i \int_B \frac{\partial L}{\partial \psi_{(i)}{}^{a\dots b}{}_{c\dots d;e}} \left(\delta\psi_{(i)}{}^{a\dots b}{}_{c\dots d}\right)_{;e} \sqrt{g}\, d^4x$$

$$= \sum_i \int_B \Bigg[\Bigg(\overbrace{\frac{\partial L}{\partial \psi_{(i)}{}^{a\dots b}{}_{c\dots d;e}} \delta\psi_{(i)}{}^{a\dots b}{}_{c\dots d}}^{=:K^e} \Bigg)_{;e} - \left(\frac{\partial L}{\partial \psi_{(i)}{}^{a\dots b}{}_{c\dots d;e}}\right)_{;e} \delta\psi_{(i)}{}^{a\dots b}{}_{c\dots d} \Bigg] \sqrt{g}\, d^4x$$

(use Leibniz rule to verify)

One term is a "boundary term":

Exercise:



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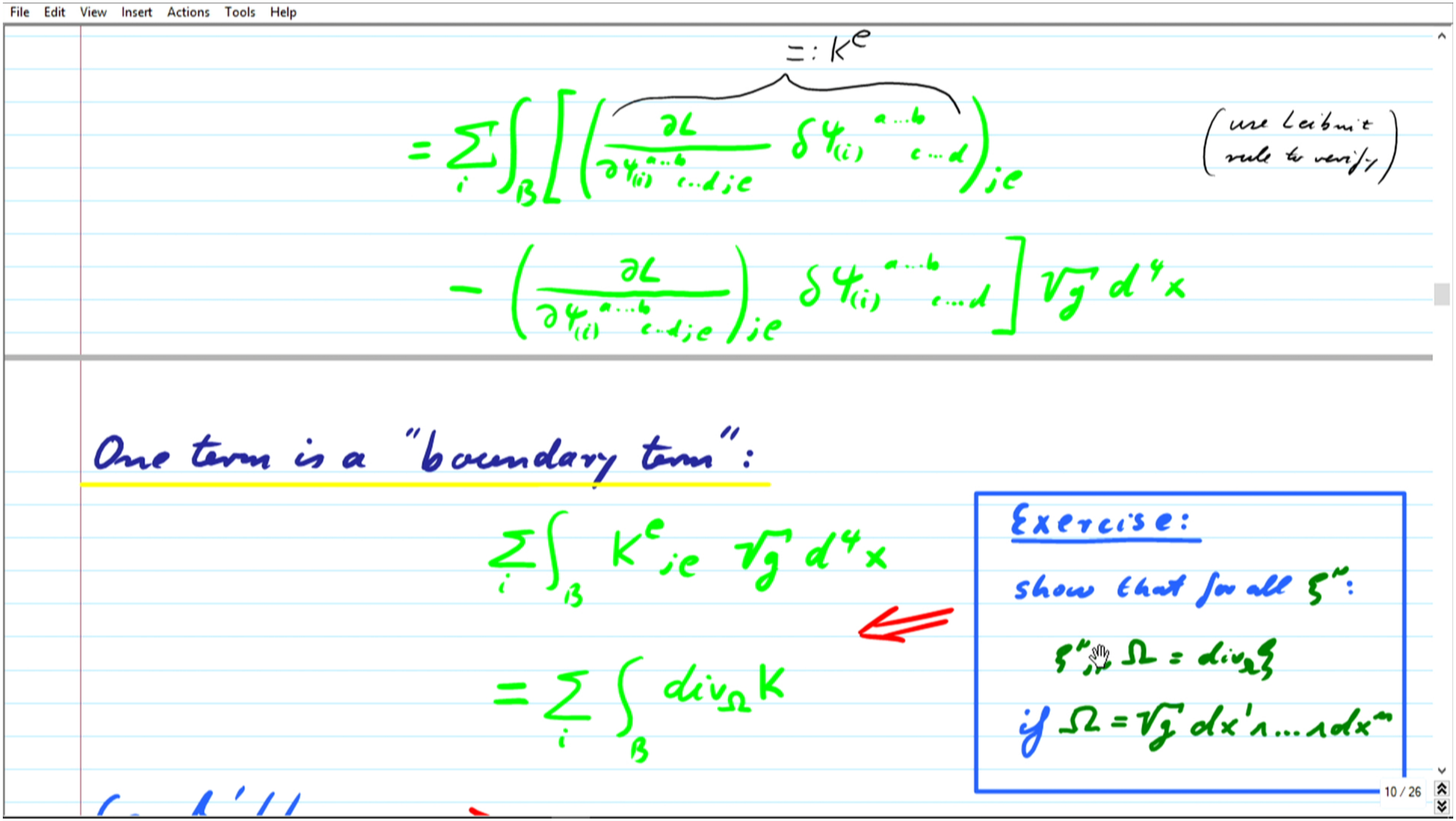


$$= \sum_i \int_B \Big[\Big(\overbrace{\frac{\partial L}{\partial \psi_{(i)}{}^{a\dots b}{}_{c\dots d;e}} \, \delta \psi_{(i)}{}^{a\dots b}{}_{c\dots d}}^{=: K^e} \Big)_{;e}$$

$$- \Big(\frac{\partial L}{\partial \psi_{(i)}{}^{a\dots b}{}_{c\dots d;e}}\Big)_{;e} \, \delta\psi_{(i)}{}^{a\dots b}{}_{c\dots d} \Big] \sqrt{g}\, d^4x$$

(use Leibniz rule to verify)

One term is a "boundary term":

$$\sum_i \int_B K^e{}_{;e} \sqrt{g}\, d^4x$$

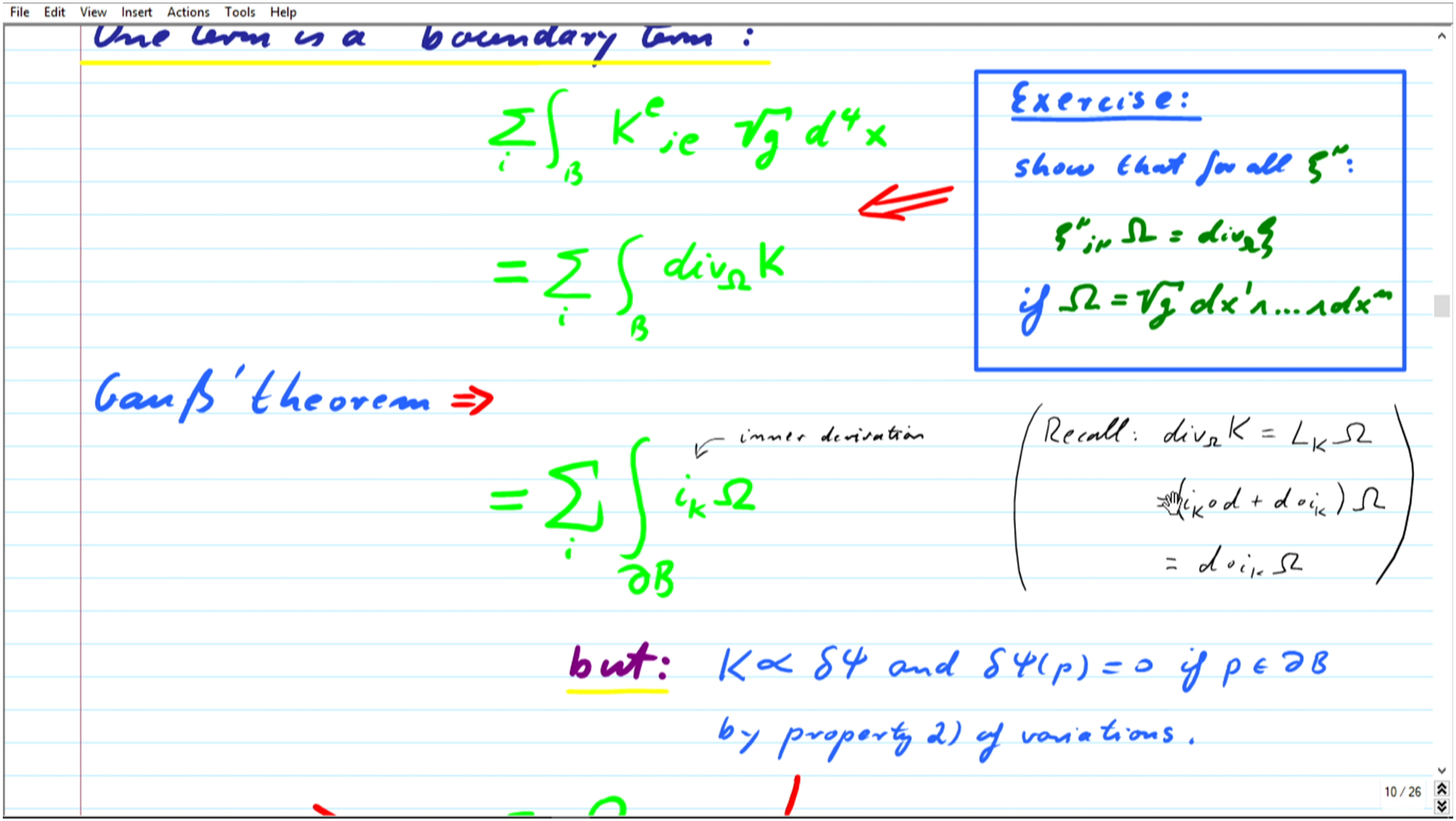
$$= \sum_i \int_B \mathrm{div}_\Omega K$$

Exercise:

show that for all ξ^μ:

$$\xi^\mu{}_{;\mu}\, \Omega = \mathrm{div}_\Omega \xi$$

if $\Omega = \sqrt{g}\, dx^1 \wedge \dots \wedge dx^n$



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Gauß' theorem $\Rightarrow$

$$= \sum_i \int_{\partial B} i_K \Omega$$

inner derivation

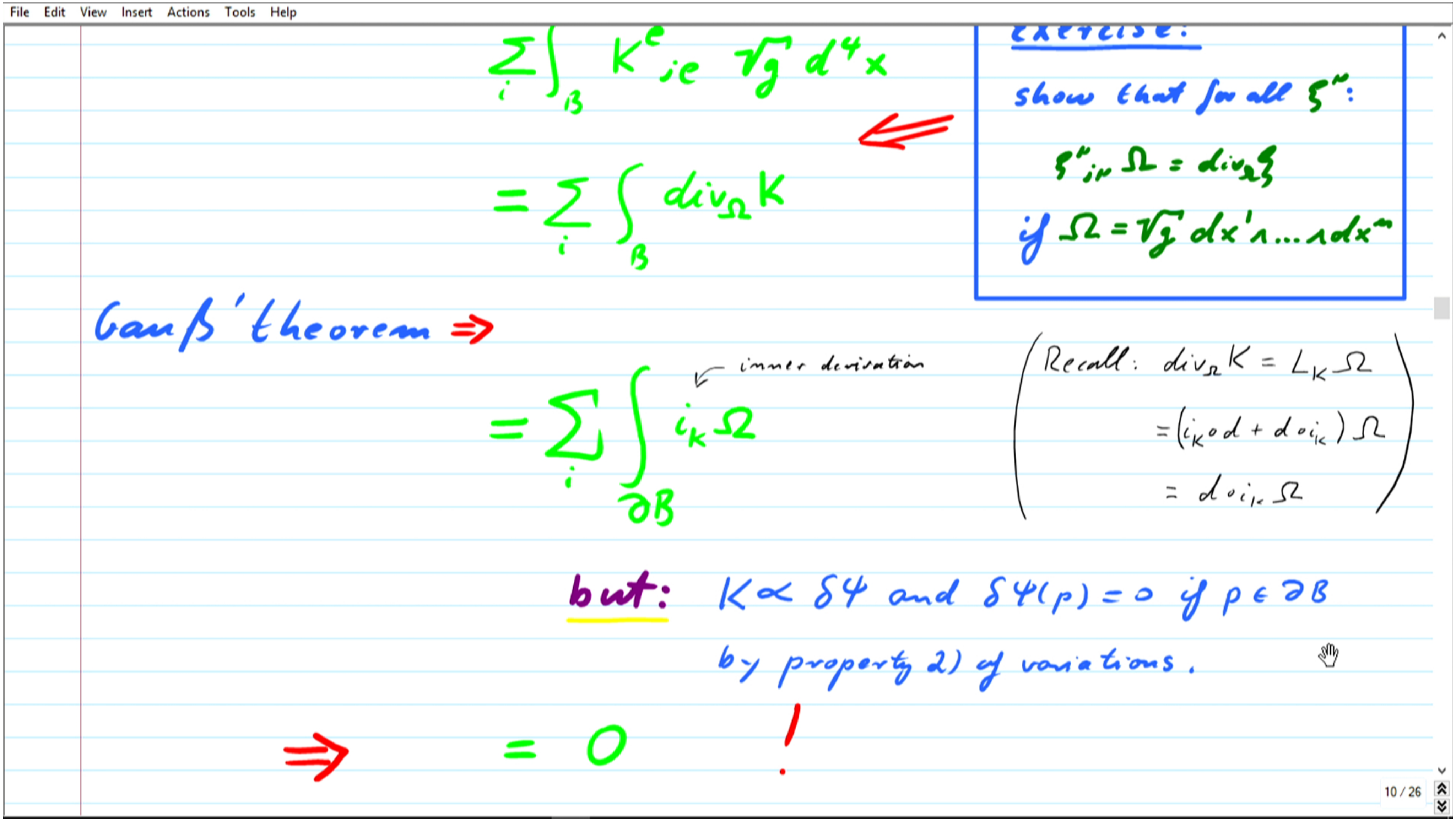
Recall: $\mathrm{div}_\Omega K = L_K \Omega$

$= (i_K \circ d + d \circ i_K)\, \Omega$

$= d \circ i_K \Omega$

but: $K \propto \delta\psi$ and $\delta\psi(p) = 0$ if $p \in \partial B$

by property 2) of variations.



$$\sum_i \int_B K^e{}_{;e} \sqrt{g}\, d^4x \Leftarrow$$

$$= \sum_i \int_B \mathrm{div}_\Omega K$$

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Gauß' theorem $\Rightarrow$

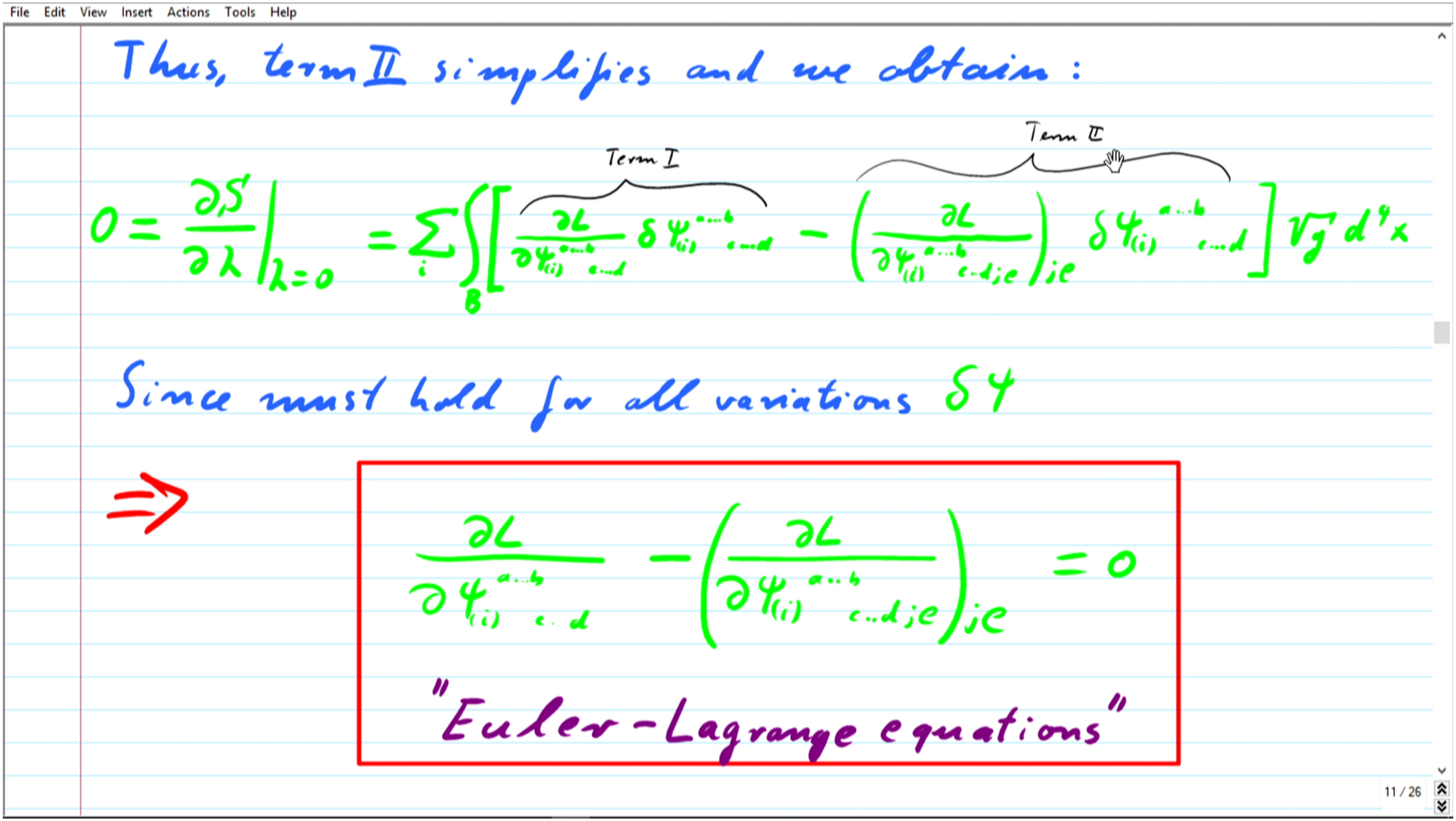
inner derivation

$$= \sum_i \int_{\partial B} i_K \Omega$$

$$\left(\begin{aligned} \text{Recall: } \mathrm{div}_\Omega K &= L_K \Omega \\ &= (i_K \circ d + d \circ i_K)\Omega \\ &= d \circ i_K \Omega \end{aligned} \right)$$

but: $K \propto \delta\psi$ and $\delta\psi(p) = 0$ if $p \in \partial B$ by property 2) of variations.

$\Rightarrow \quad = 0$!



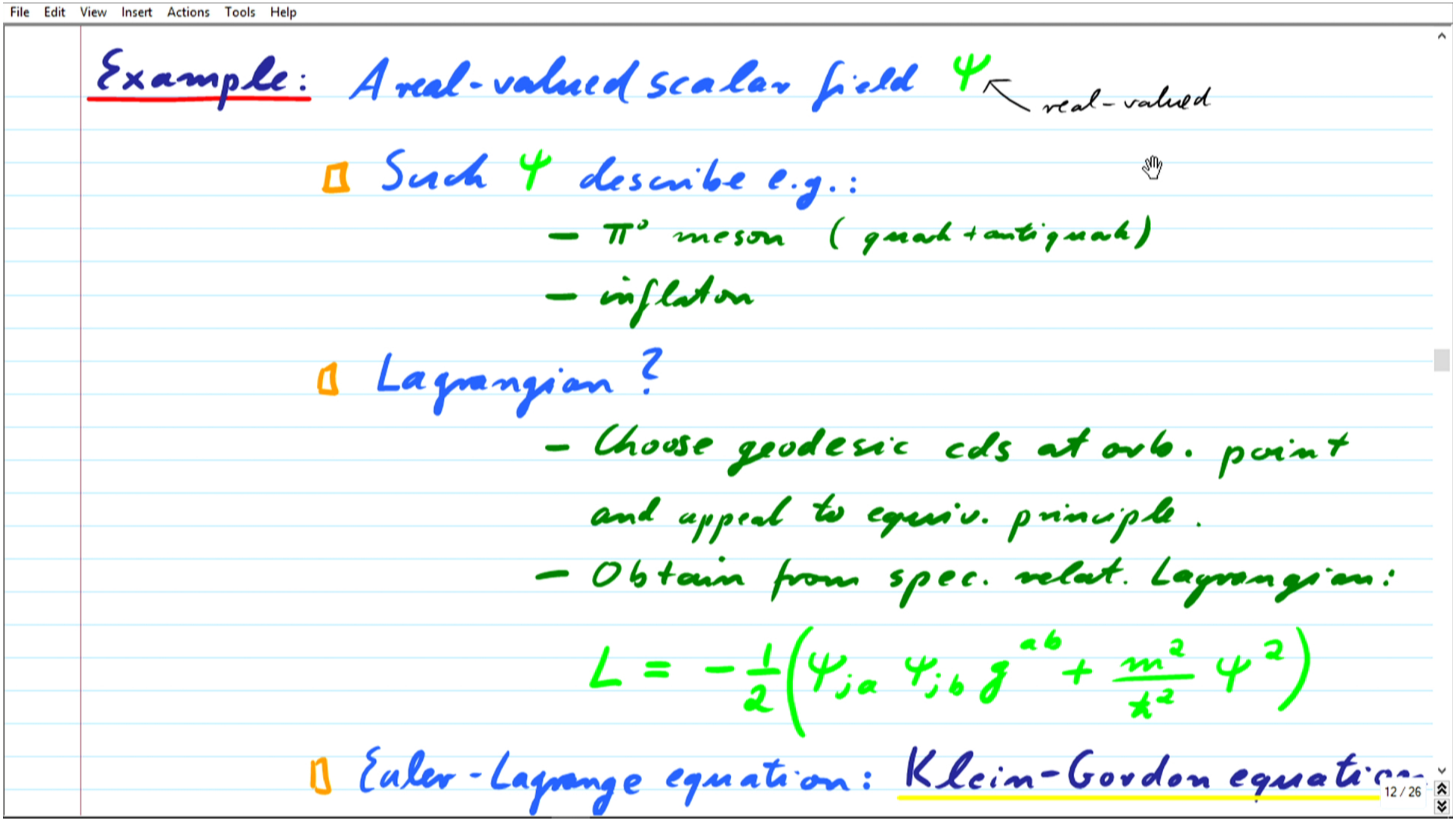
Thus, term II simplifies and we obtain:

$$0 = \frac{\partial S}{\partial \lambda}\Big|_{\lambda=0} = \sum_i \int_B \Big[\overbrace{\frac{\partial L}{\partial \psi_{(i)}{}^{a\ldots b}{}_{c\ldots d}} \delta \psi_{(i)}{}^{a\ldots b}{}_{c\ldots d}}^{\text{Term I}} - \overbrace{\left(\frac{\partial L}{\partial \psi_{(i)}{}^{a\ldots b}{}_{c\ldots d;e}}\right)_{;e} \delta \psi_{(i)}{}^{a\ldots b}{}_{c\ldots d}}^{\text{Term II}} \Big] \sqrt{g}\, d^4x$$

Since must hold for all variations $\delta\psi$

$$\Rightarrow \quad \boxed{\frac{\partial L}{\partial \psi_{(i)}{}^{a\ldots b}{}_{c\ldots d}} - \left(\frac{\partial L}{\partial \psi_{(i)}{}^{a\ldots b}{}_{c\ldots d;e}}\right)_{;e} = 0}$$

"Euler-Lagrange equations"



Example: A real-valued scalar field ϕ ← real-valued

☐ Such ϕ describe e.g.:

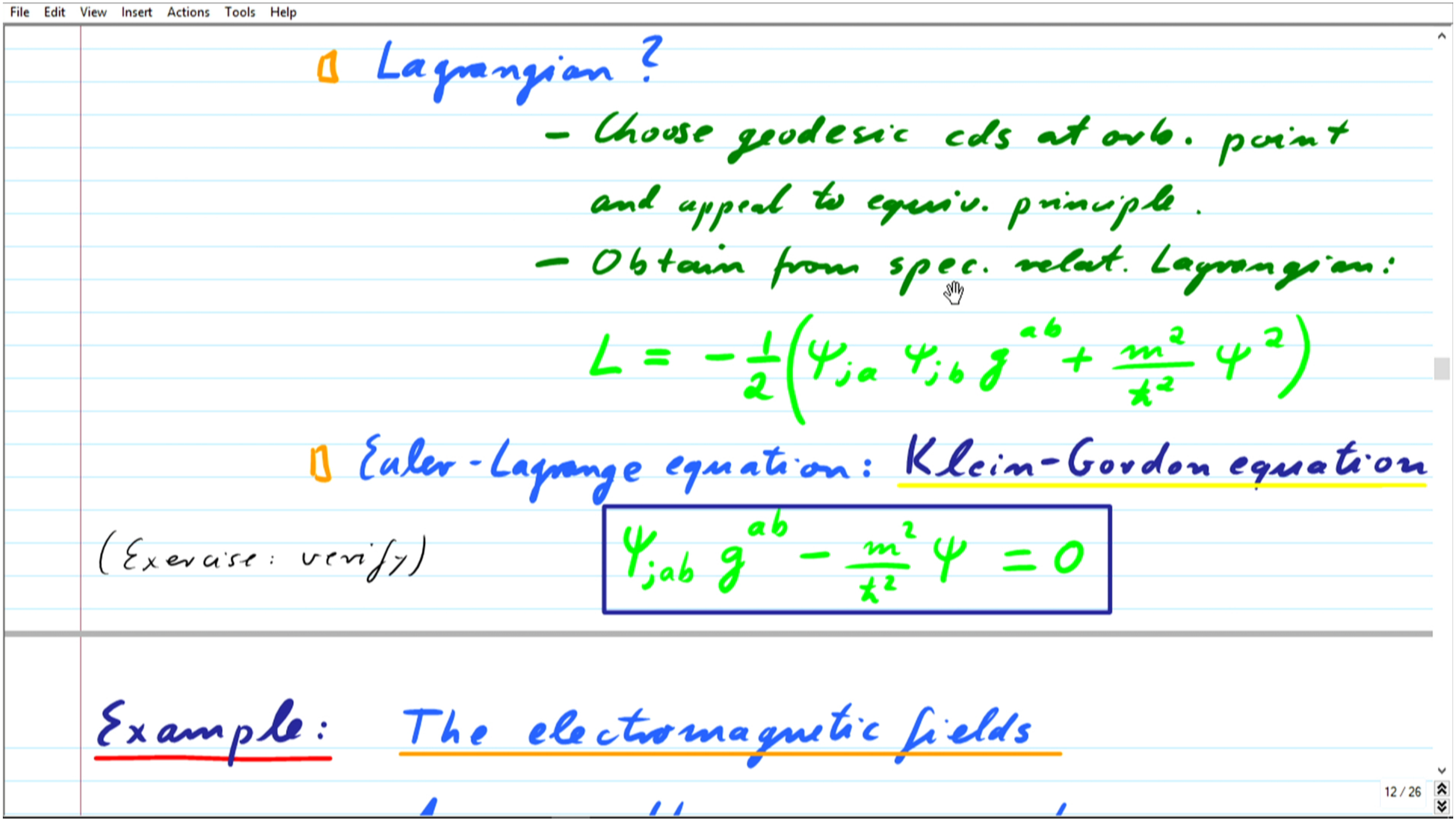
- π^0 meson (quark + antiquark)
- inflaton

☐ Lagrangian?

- Choose geodesic cds at arb. point and appeal to equiv. principle.
- Obtain from spec. relat. Lagrangian:

$$L = -\frac{1}{2}\left(\phi_{;a}\,\phi_{;b}\,g^{ab} + \frac{m^2}{\hbar^2}\phi^2\right)$$

☐ Euler-Lagrange equation: Klein-Gordon equatio



☐ Lagrangian?

- Choose geodesic cds at arb. point and appeal to equiv. principle.
- Obtain from spec. relat. Lagrangian:

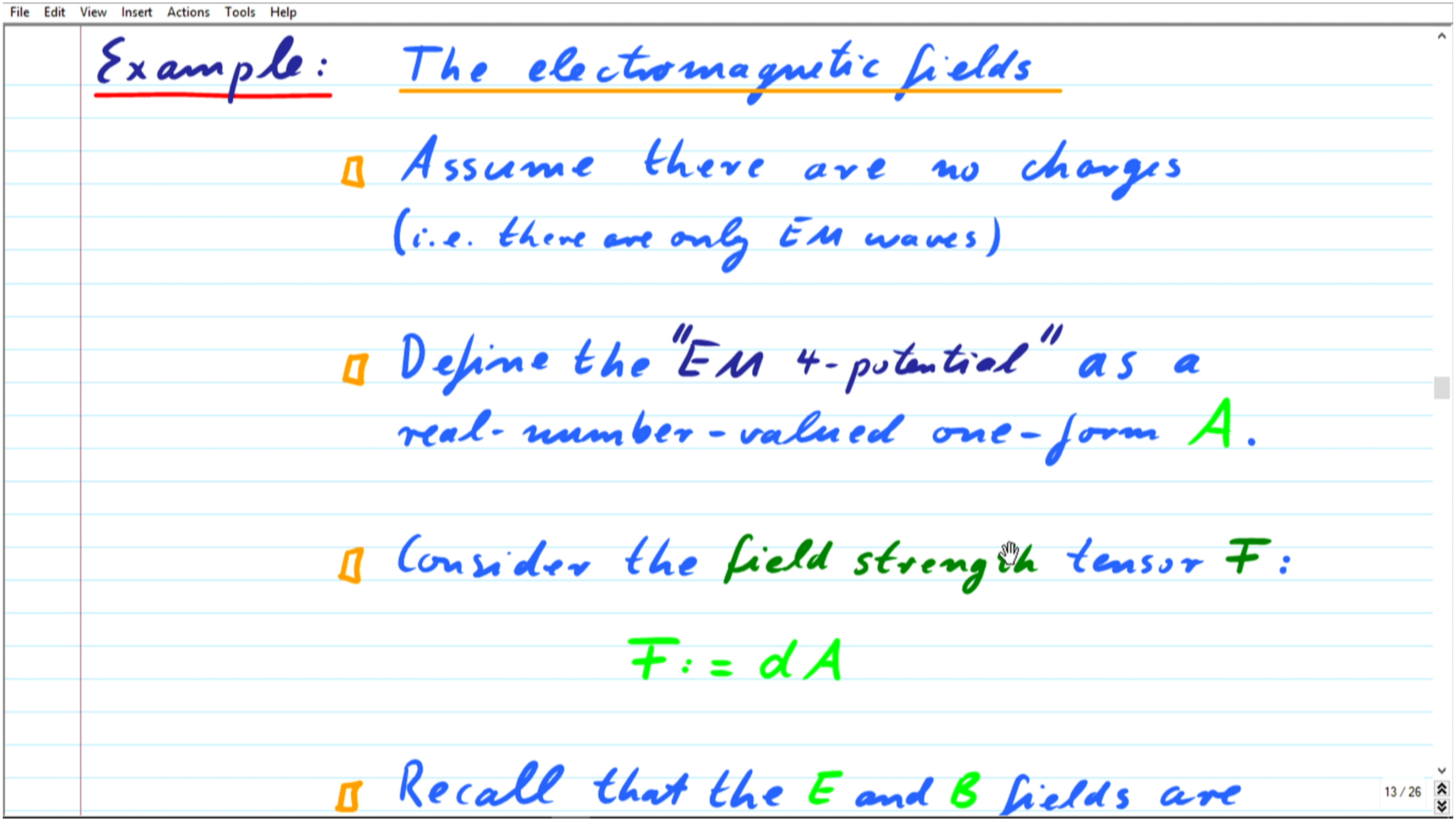
$$L = -\frac{1}{2}\left(\psi_{;a}\,\psi_{;b}\,g^{ab} + \frac{m^2}{\hbar^2}\psi^2\right)$$

☐ Euler-Lagrange equation: Klein-Gordon equation

(Exercise: verify)

$$\boxed{\psi_{;ab}\,g^{ab} - \frac{m^2}{\hbar^2}\psi = 0}$$

Example: The electromagnetic fields

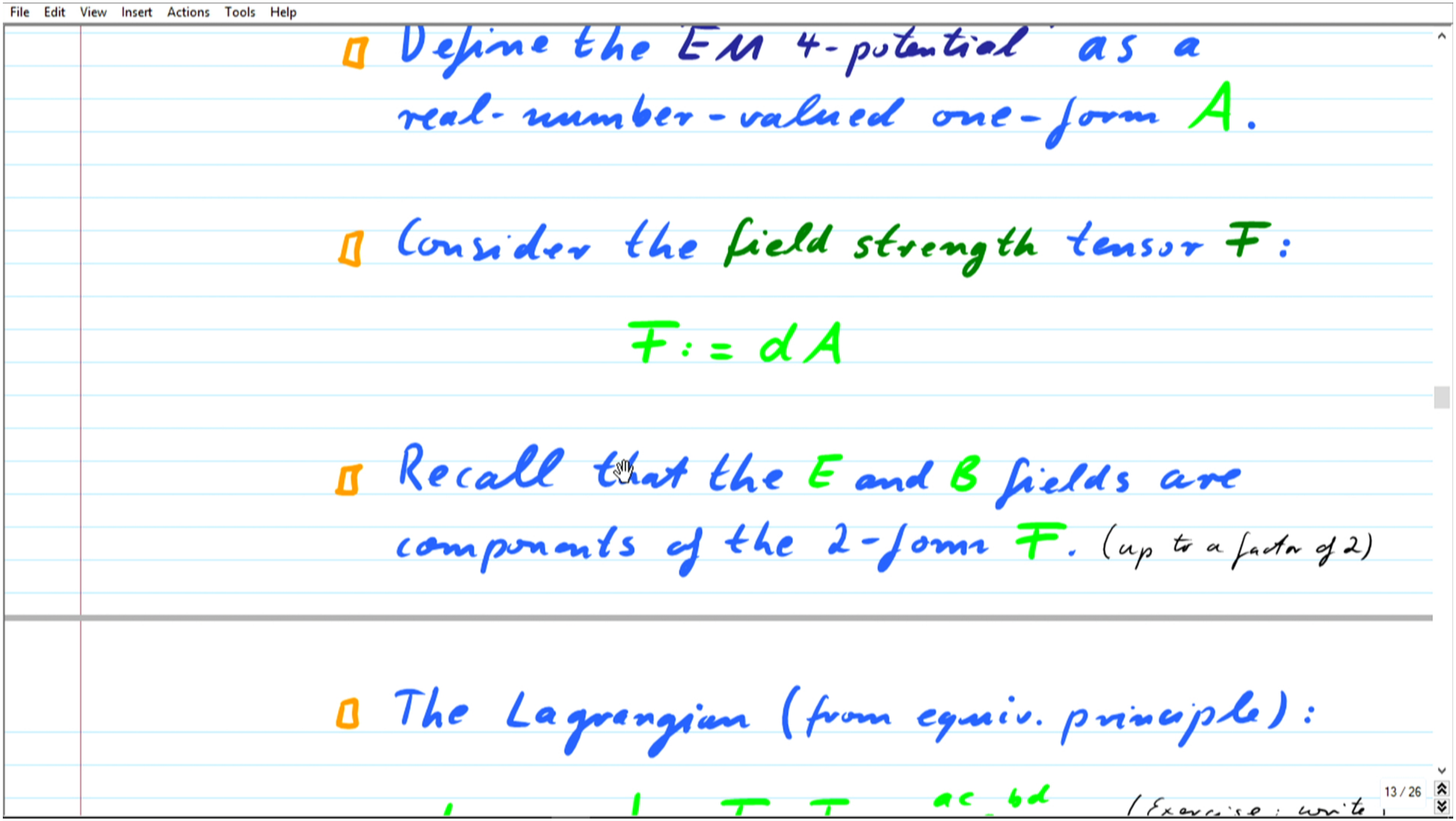


Example: The electromagnetic fields

- Assume there are no charges (i.e. there are only EM waves)
- Define the "EM 4-potential" as a real-number-valued one-form A.
- Consider the field strength tensor F:

$$F := dA$$

- Recall that the E and B fields are



- Define the EM 4-potential as a real-number-valued one-form A.

- Consider the field strength tensor F:

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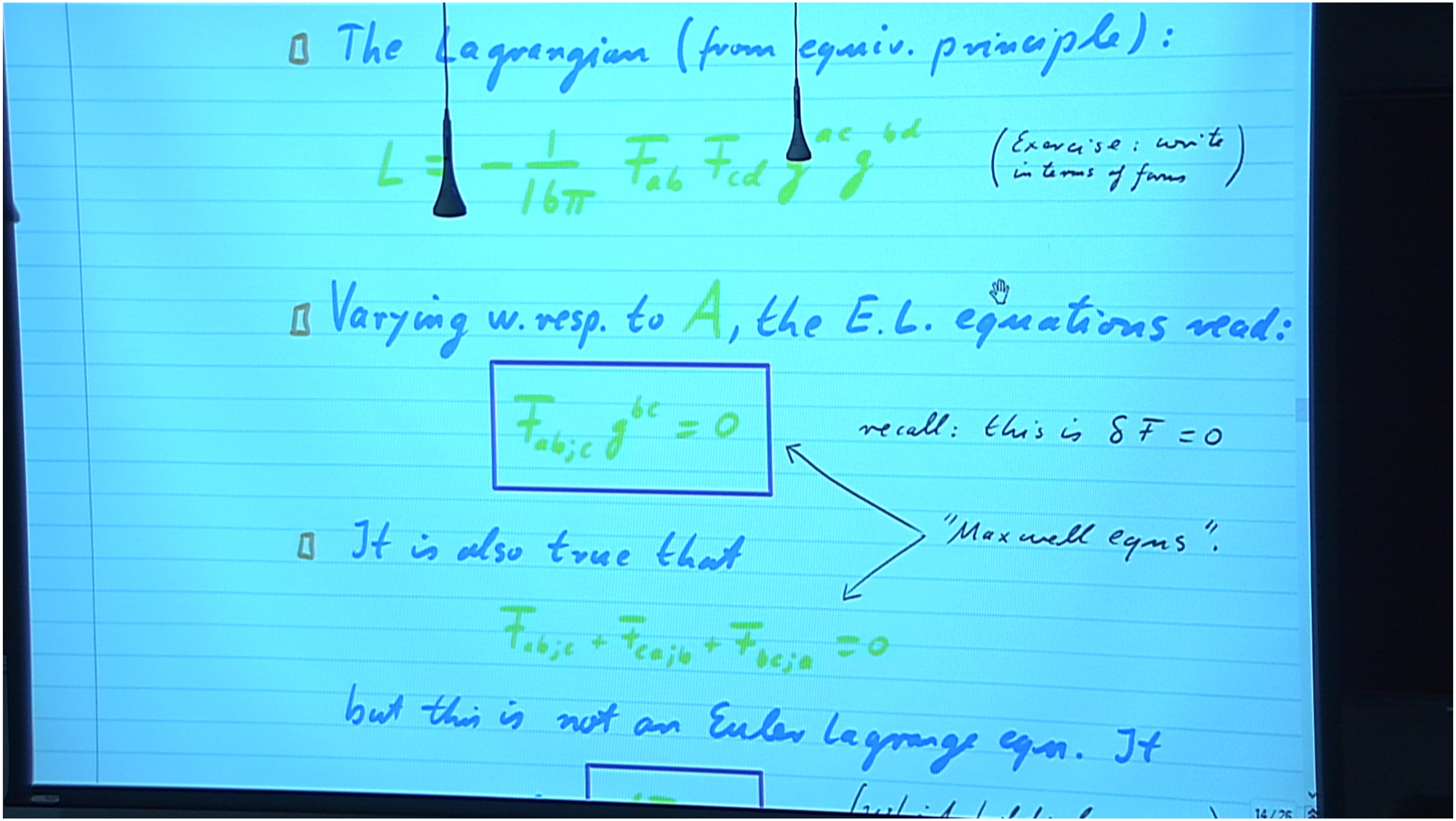
- Recall that the E and B fields are components of the 2-form F. (up to a factor of 2)

- The Lagrangian (from equiv. principle):

(Exercise: write



- The Lagrangian (from equiv. principle):

$$L = -\frac{1}{16\pi} F_{ab} F_{cd} g^{ac} g^{bd}$$

(Exercise: write in terms of forms)

- Varying w. resp. to A, the E.L. equations read:

$$\boxed{F_{ab;c}\, g^{bc} = 0}$$

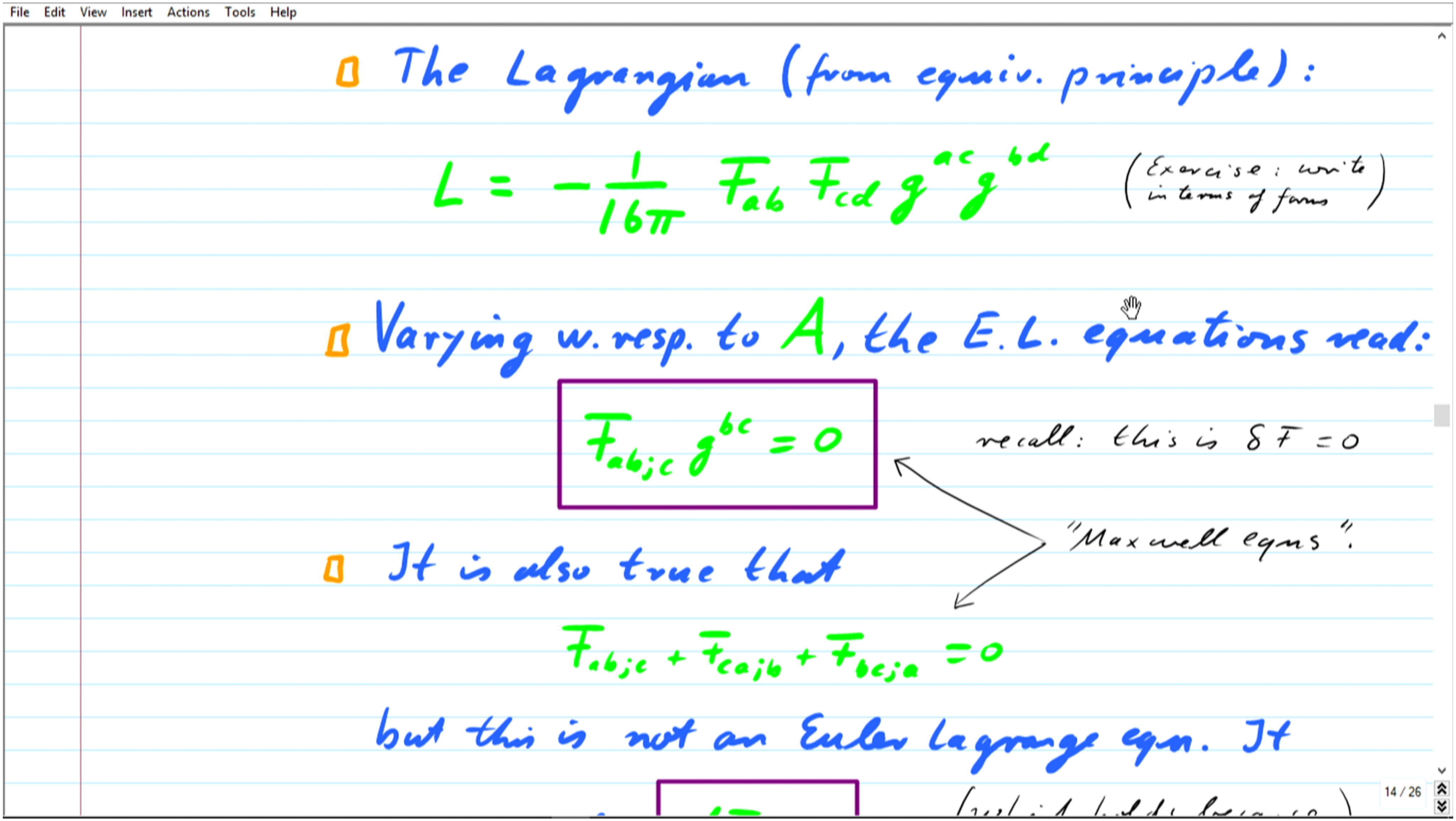
recall: this is $\delta F = 0$

"Maxwell eqns".

- It is also true that

$$F_{ab;c} + F_{ca;b} + F_{bc;a} = 0$$

but this is not an Euler Lagrange eqn. It



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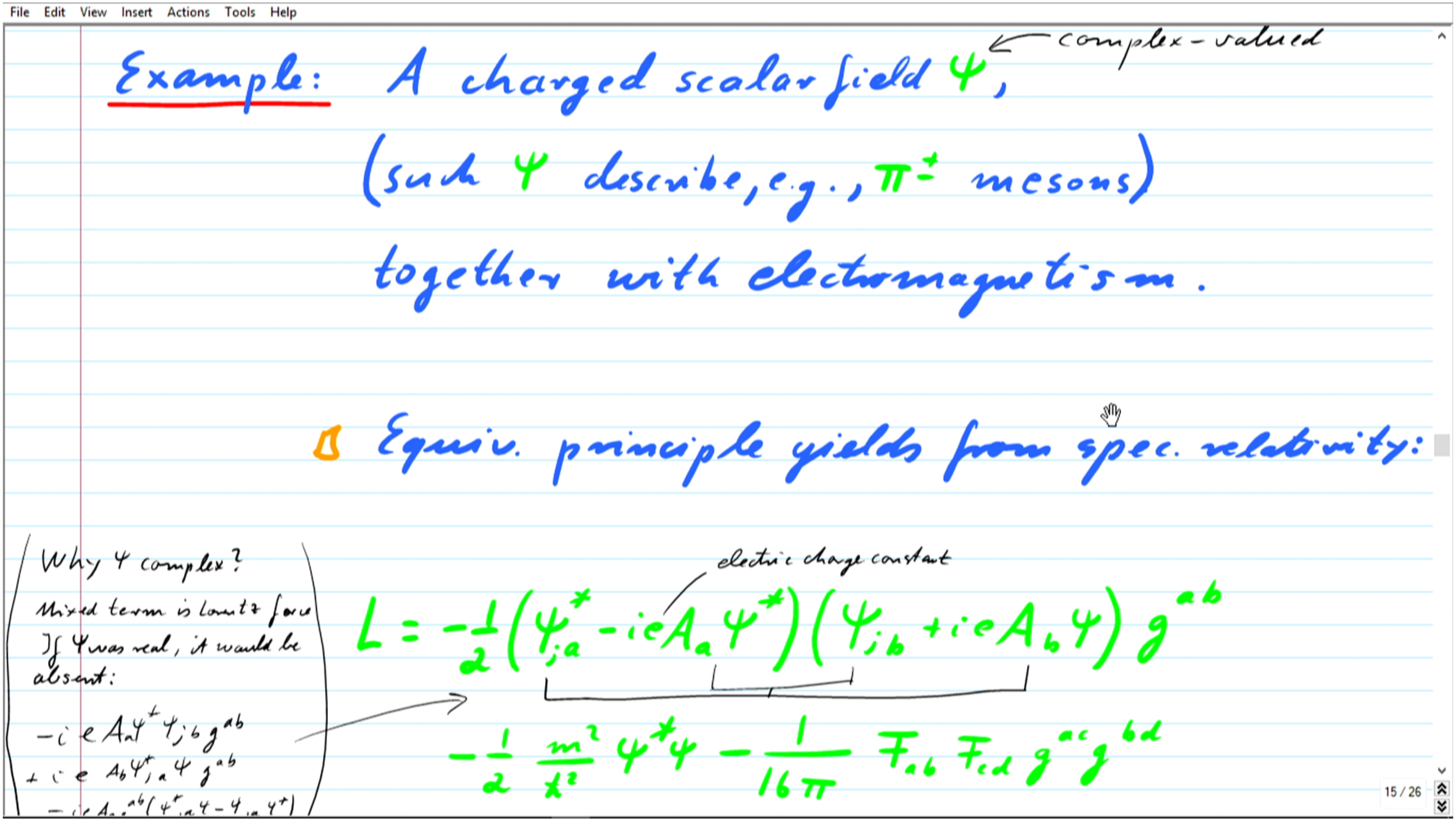


Example: A charged scalar field ψ, ← complex-valued

(such ψ describe, e.g., $\pi^{\pm}$ mesons)

together with electromagnetism.

Equiv. principle yields from spec. relativity:

Why ψ complex?

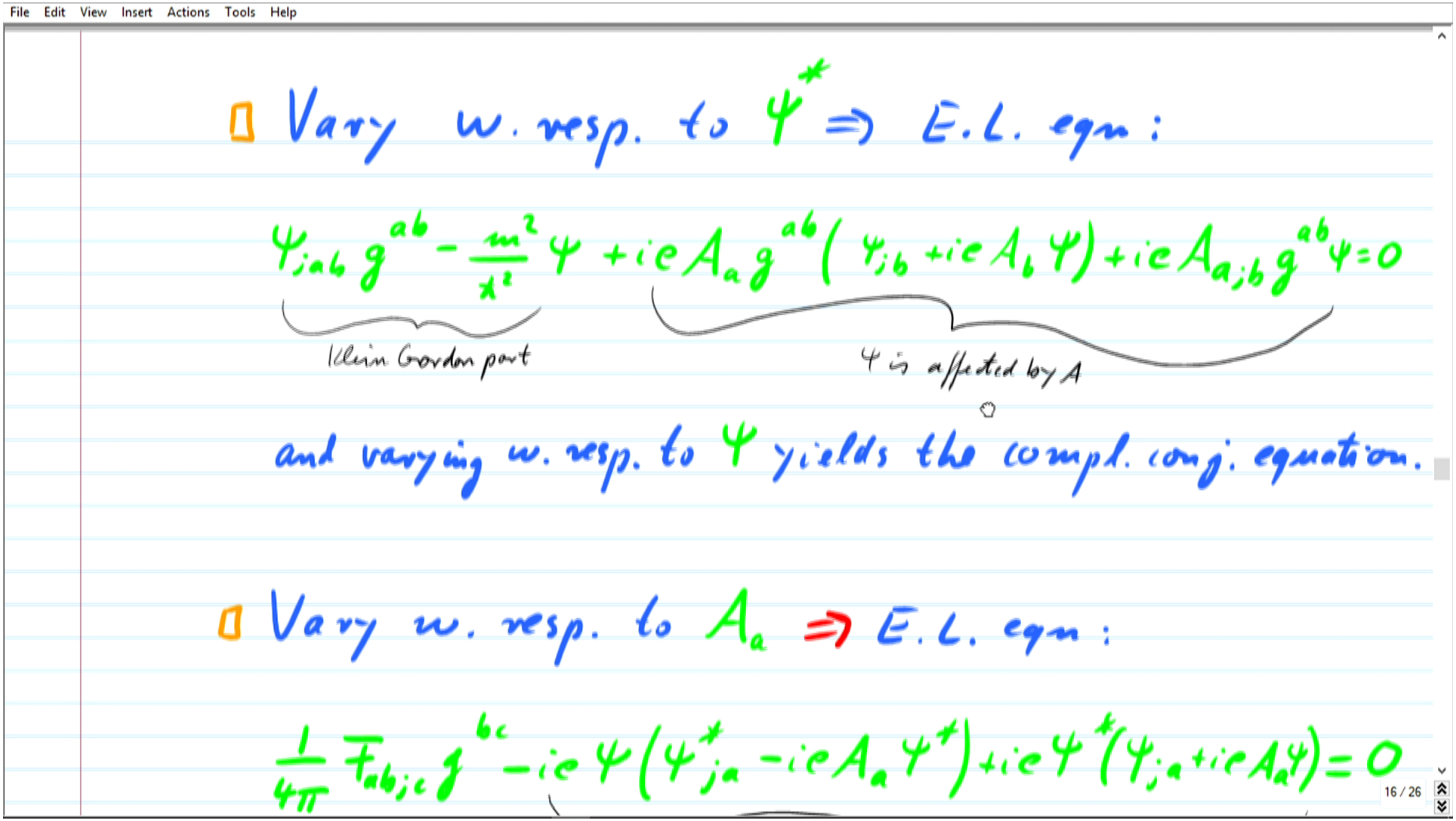
Mixed term is Lorentz force

If ψ was real, it would be absent:

$$-ieA_a\psi^{*}\psi_{;b}g^{ab}$$
$$+ieA_b\psi^{*}_{;a}\psi g^{ab}$$
$$-ieA_a{}^{ab}(\psi^{*}_{,a}\psi-\psi_{,a}\psi^{*})$$

electric charge constant

$$L=-\frac{1}{2}\left(\psi^{*}_{;a}-ieA_a\psi^{*}\right)\left(\psi_{;b}+ieA_b\psi\right)g^{ab}$$
$$-\frac{1}{2}\frac{m^2}{\hbar^2}\psi^{*}\psi-\frac{1}{16\pi}F_{ab}F_{cd}g^{ac}g^{bd}$$



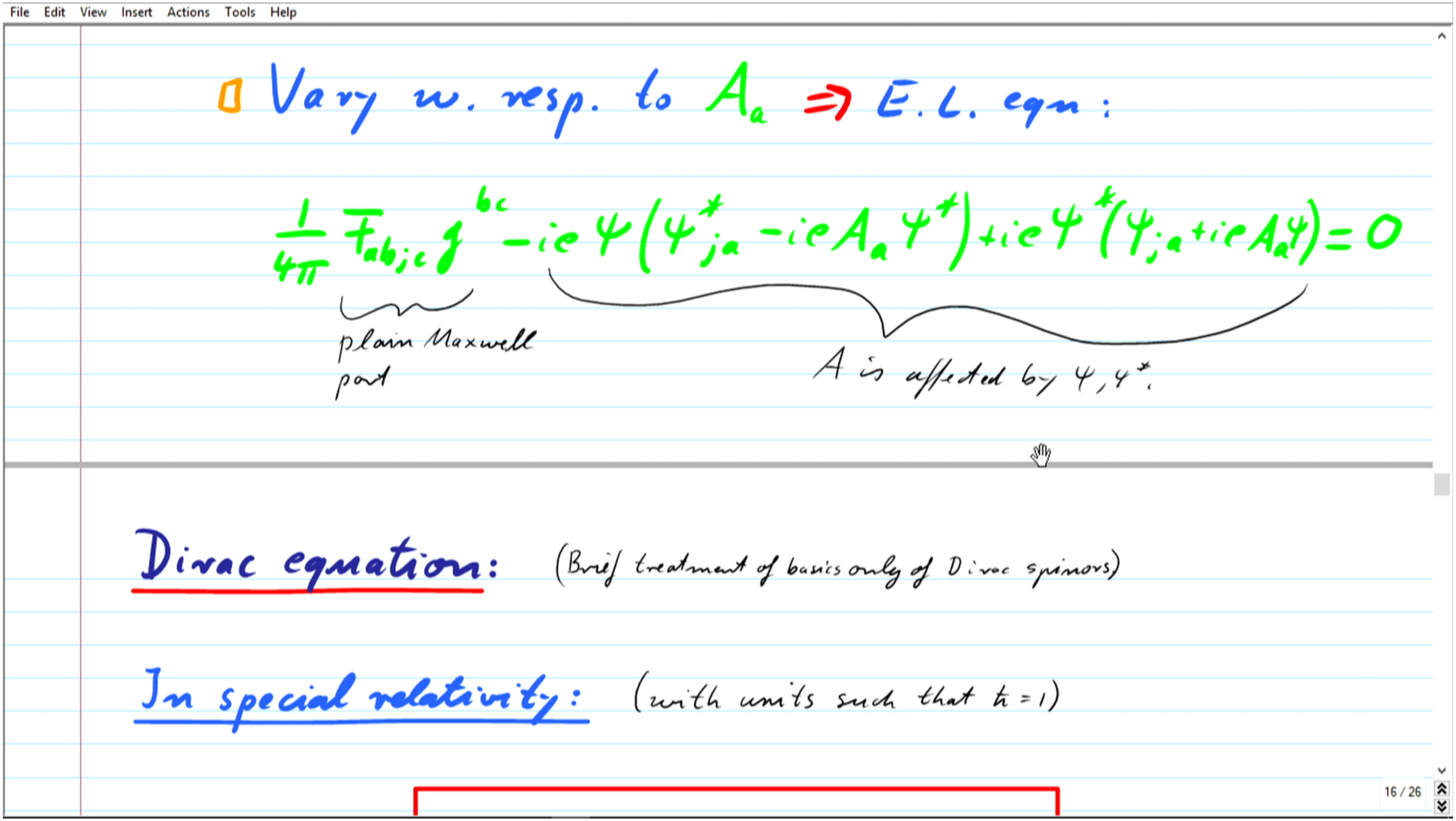
☐ Vary w. resp. to ψ^* ⇒ E.L. eqn:

$$\underbrace{\psi_{;ab}\, g^{ab} - \frac{m^2}{\bar{\lambda}^2}\psi}_{\text{Klein Gordon part}} \underbrace{+ ieA_a g^{ab}(\psi_{;b} + ieA_b\psi) + ieA_{a;b}\, g^{ab}\psi}_{\psi \text{ is affected by } A} = 0$$

and varying w. resp. to ψ yields the compl. conj. equation.

☐ Vary w. resp. to A_a ⇒ E.L. eqn:

$$\frac{1}{4\pi} F_{ab;c}\, g^{bc} - ie\psi(\psi^*_{;a} - ieA_a\psi^*) + ie\psi^*(\psi_{;a} + ieA_a\psi) = 0$$



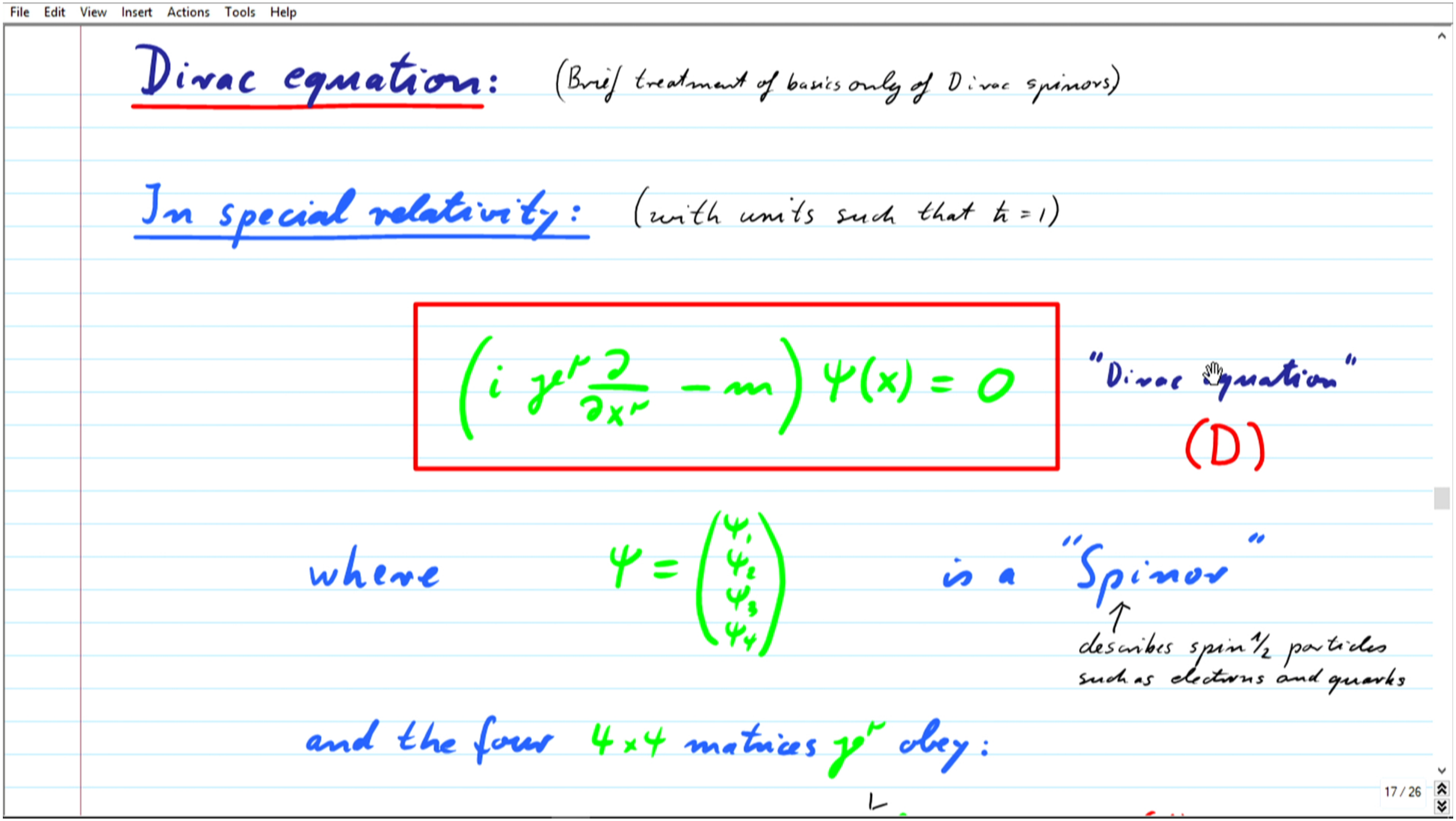
☐ Vary w. resp. to A_a ⇒ E.L. eqn:

$$\underbrace{\frac{1}{4\pi}F_{ab;c}\,g^{bc}}_{\text{plain Maxwell part}}\underbrace{-ie\psi\left(\psi^*_{;a}-ieA_a\psi^*\right)+ie\psi^*\left(\psi_{;a}+ieA_a\psi\right)}_{A\text{ is affected by }\psi,\psi^*.}=0$$

Dirac equation:

(Brief treatment of basics only of Dirac spinors)

In special relativity: (with units such that $\hbar=1$)



Dirac equation: (Brief treatment of basics only of Dirac spinors)

In special relativity: (with units such that $\hbar = 1$)

$$\boxed{\left(i\gamma^\mu \frac{\partial}{\partial x^\mu} - m\right)\psi(x) = 0}$$

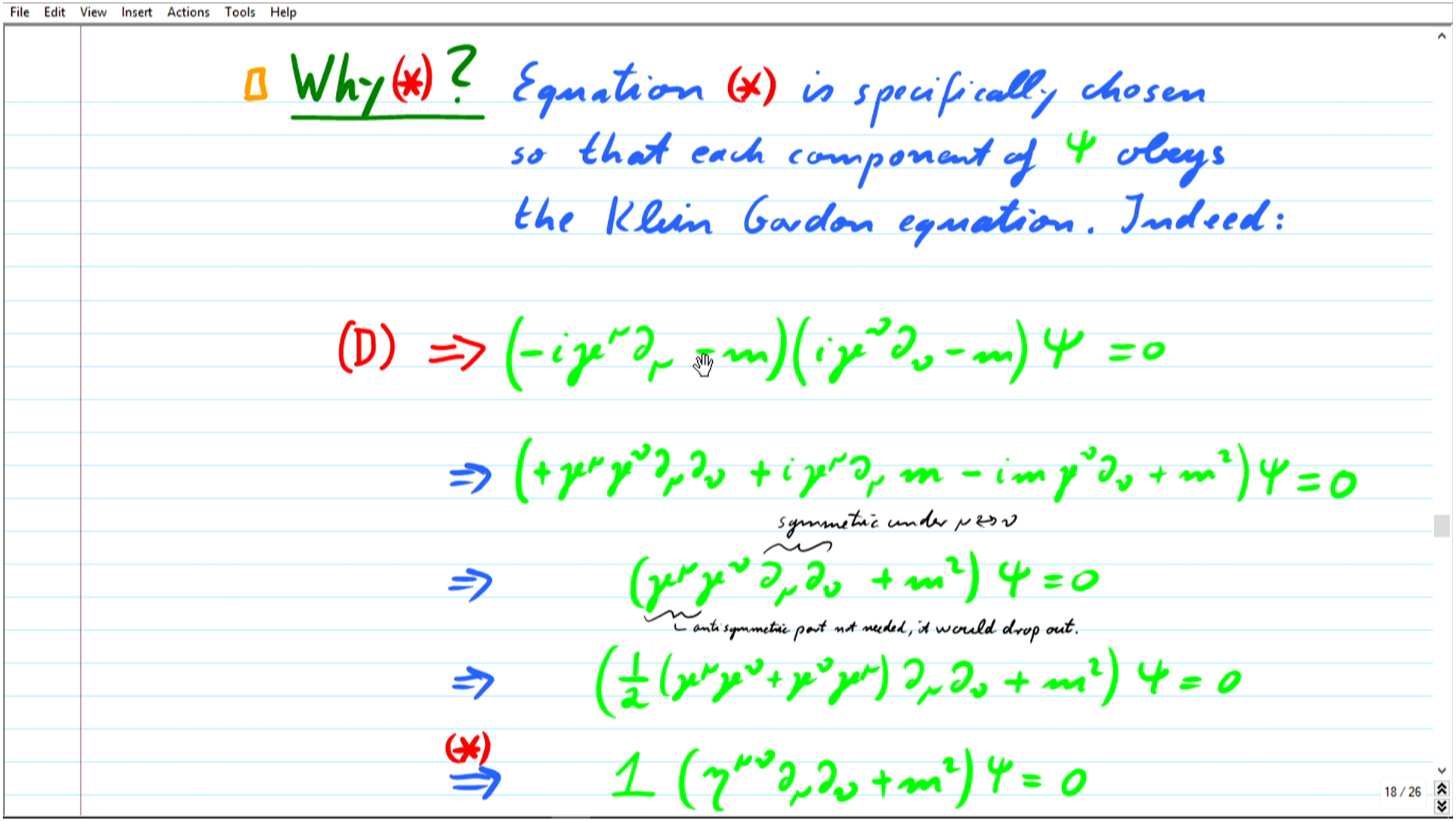
"Dirac equation"

(D)

where $\psi = \begin{pmatrix}\psi_1\\ \psi_2\\ \psi_3\\ \psi_4\end{pmatrix}$ is a "Spinor"

↑ describes spin ½ particles such as electrons and quarks

and the four 4×4 matrices γ^μ obey:



☐ **Why (*)?** Equation (*) is specifically chosen so that each component of ψ obeys the Klein Gordon equation. Indeed:

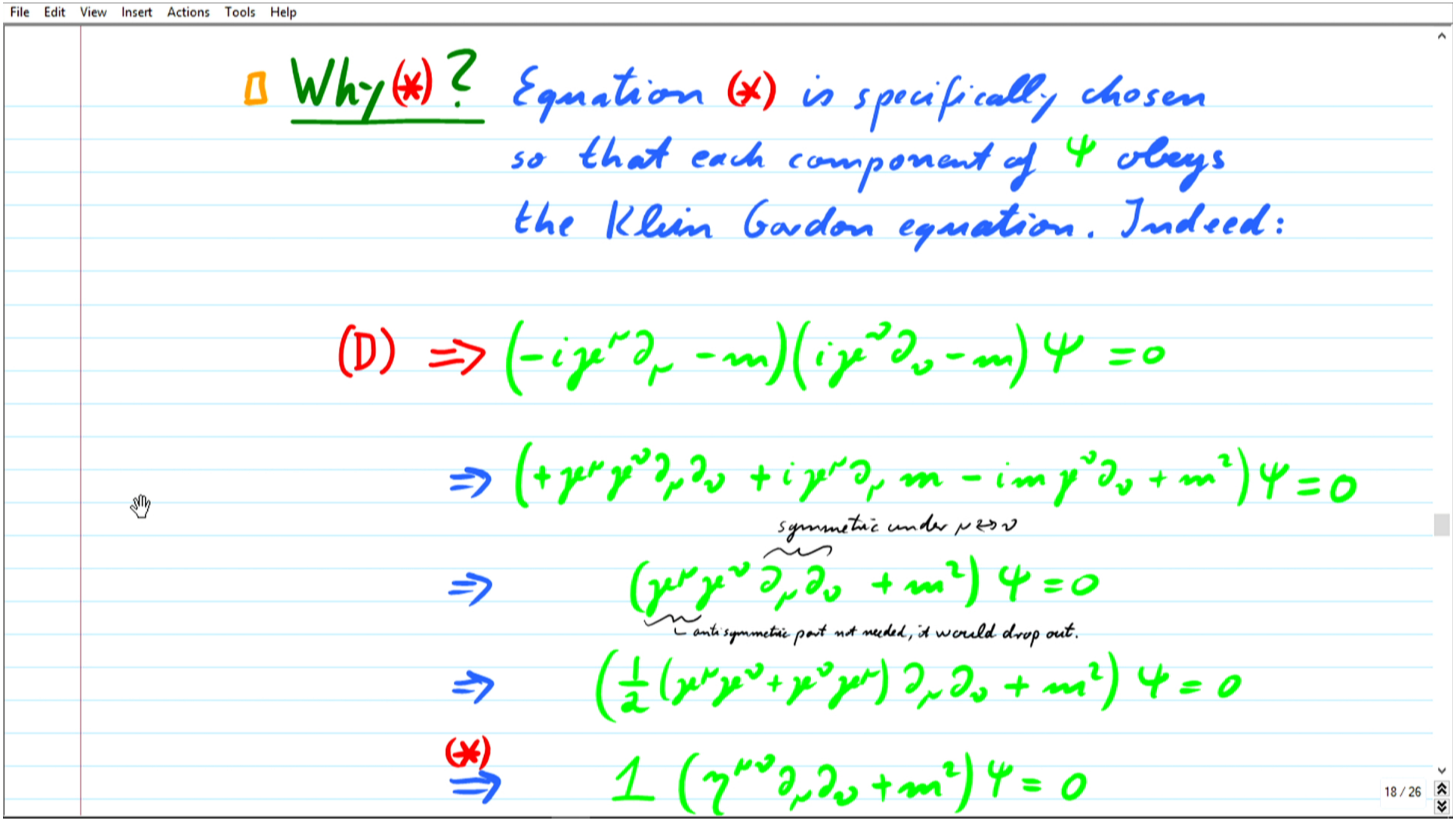
$$(D) \Rightarrow \left(-i\gamma^\mu\partial_\mu - m\right)\left(i\gamma^\nu\partial_\nu - m\right)\psi = 0$$

$$\Rightarrow \left(+\gamma^\mu\gamma^\nu\partial_\mu\partial_\nu + i\gamma^\mu\partial_\mu m - im\gamma^\nu\partial_\nu + m^2\right)\psi = 0$$

$$\Rightarrow \left(\underbrace{\gamma^\mu\gamma^\nu}_{\text{antisymmetric part not needed, it would drop out.}}\overbrace{\partial_\mu\partial_\nu}^{\text{symmetric under } \mu\leftrightarrow\nu} + m^2\right)\psi = 0$$

$$\Rightarrow \left(\frac{1}{2}\left(\gamma^\mu\gamma^\nu + \gamma^\nu\gamma^\mu\right)\partial_\mu\partial_\nu + m^2\right)\psi = 0$$

$$\overset{(*)}{\Rightarrow} 1\left(\eta^{\mu\nu}\partial_\mu\partial_\nu + m^2\right)\psi = 0$$



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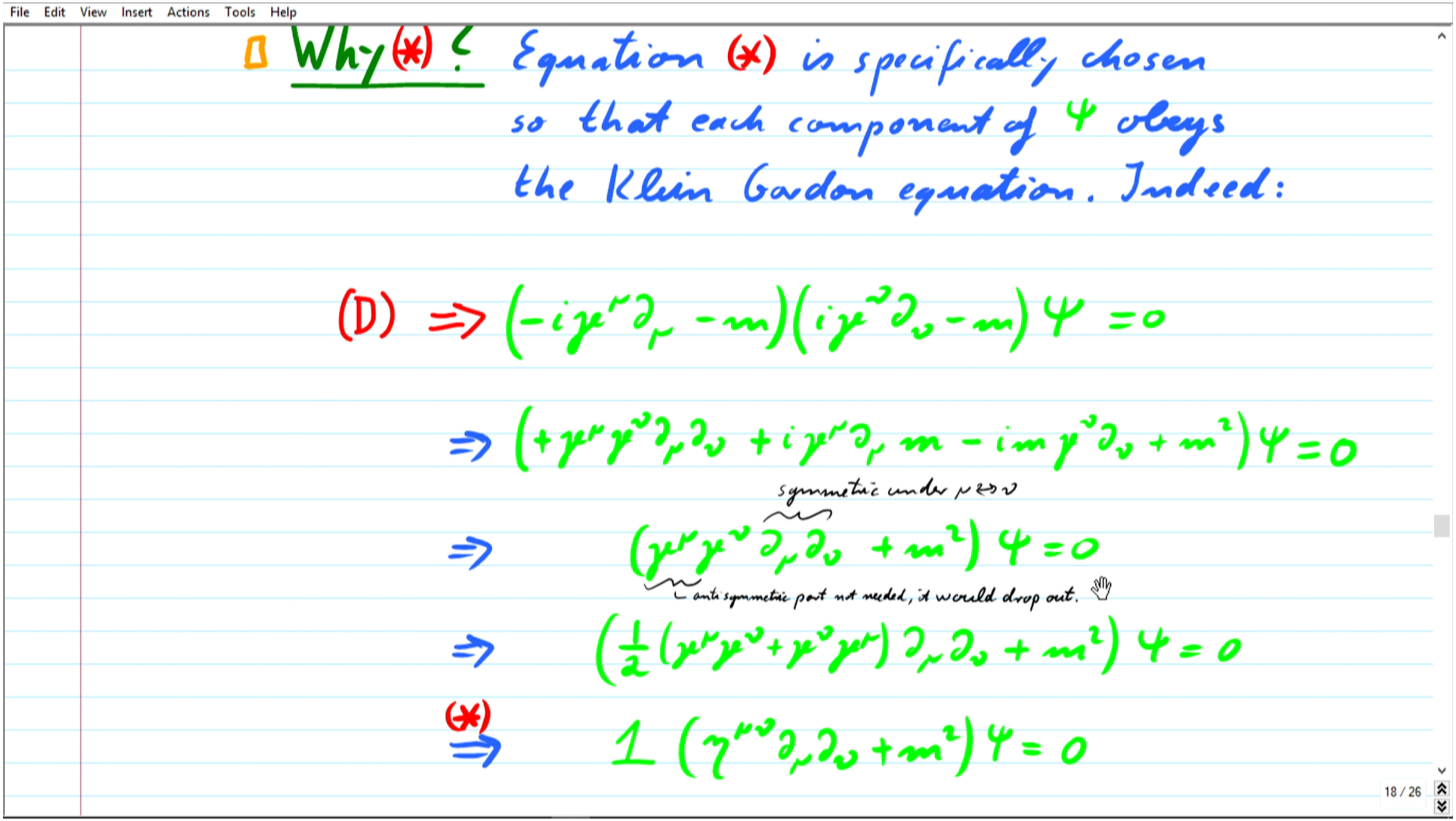

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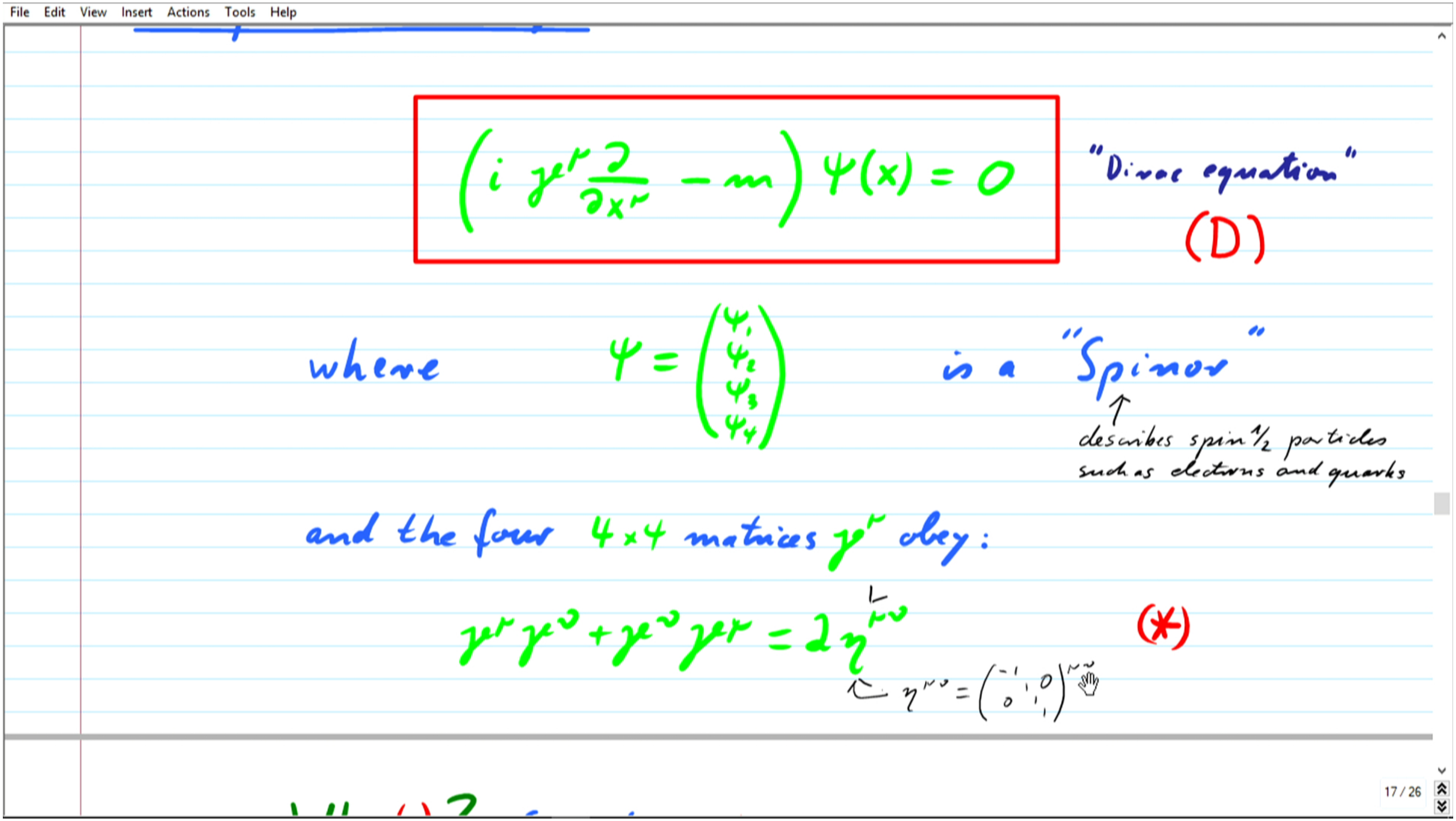

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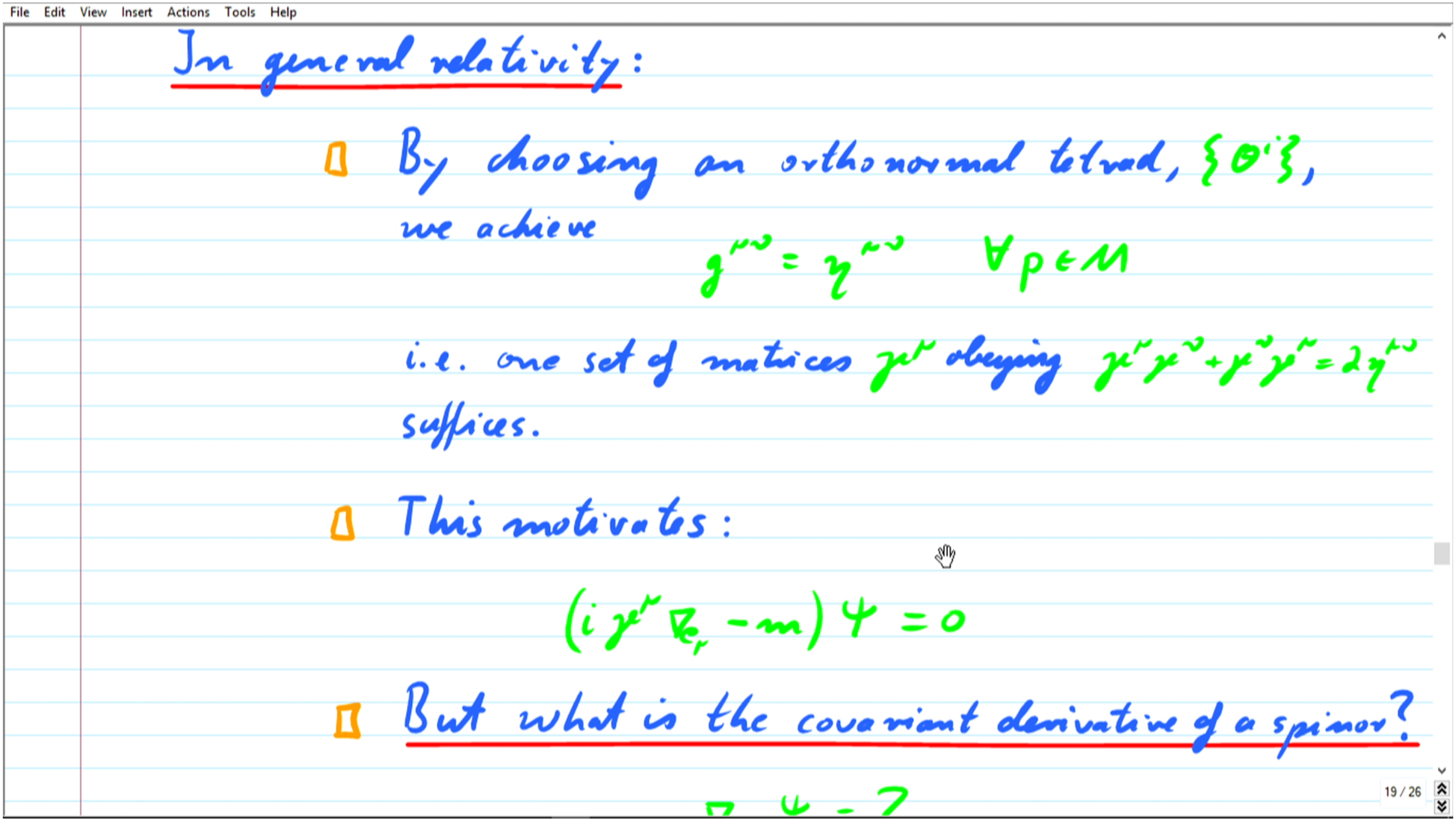
$$\gamma^\mu\gamma^\nu + \gamma^\nu\gamma^\mu = 2\eta^{\mu\nu} \qquad (*)$$

$$\eta^{\mu\nu} = \begin{pmatrix} -1 & & & 0 \\ & 1 & & \\ & & 1 & \\ 0 & & & 1 \end{pmatrix}^{\mu\nu}$$



In general relativity:

- By choosing an orthonormal tetrad, $\{\theta^i\}$, we achieve

$$g^{\mu\nu} = \eta^{\mu\nu} \quad \forall p \in M$$

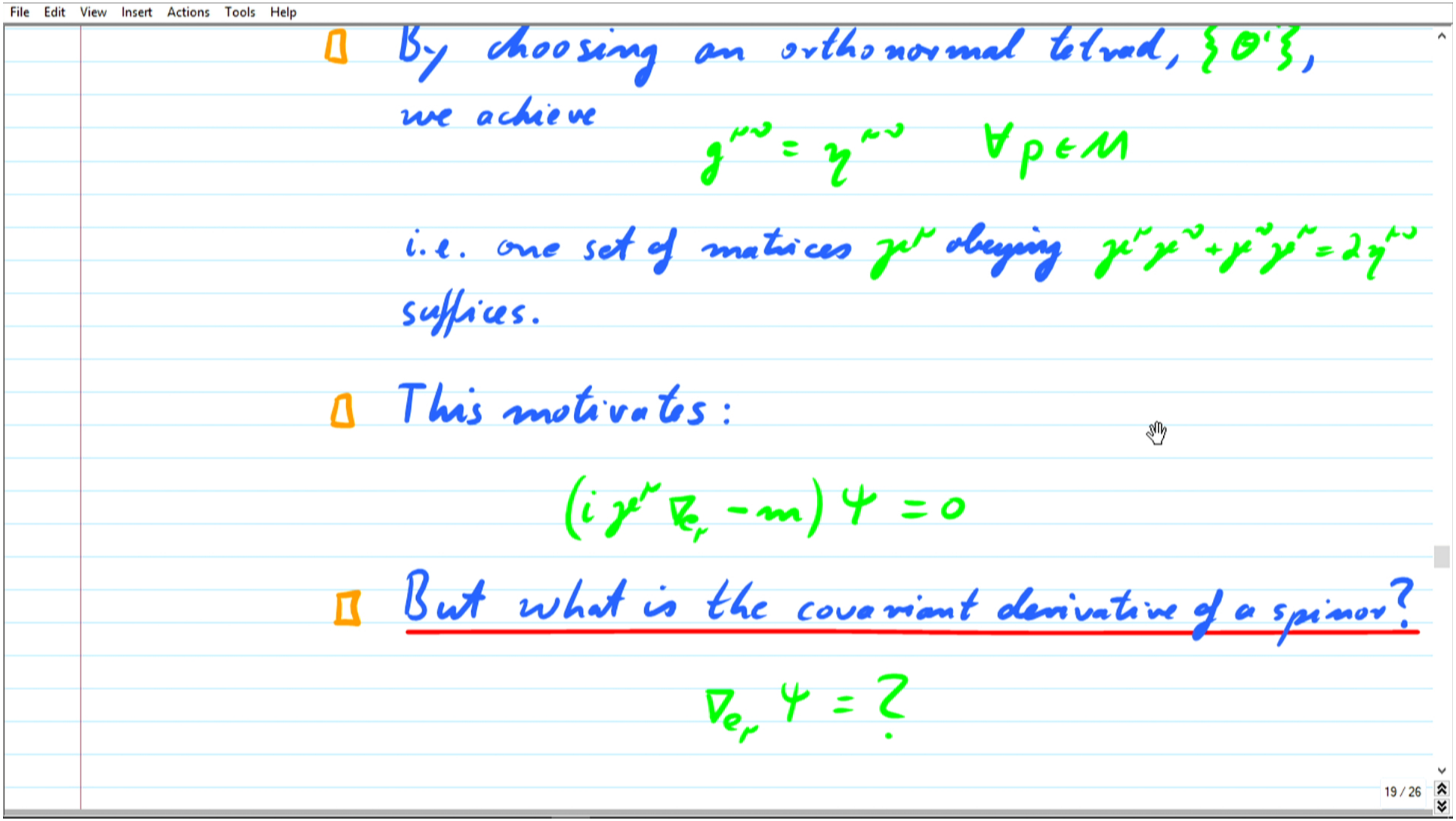
i.e. one set of matrices γ^μ obeying $\gamma^\mu\gamma^\nu + \gamma^\nu\gamma^\mu = 2\eta^{\mu\nu}$ suffices.

- This motivates:

$$(i\gamma^\mu \nabla_{e_\mu} - m)\psi = 0$$

- **But what is the covariant derivative of a spinor?**

$$\nabla \psi = ?$$



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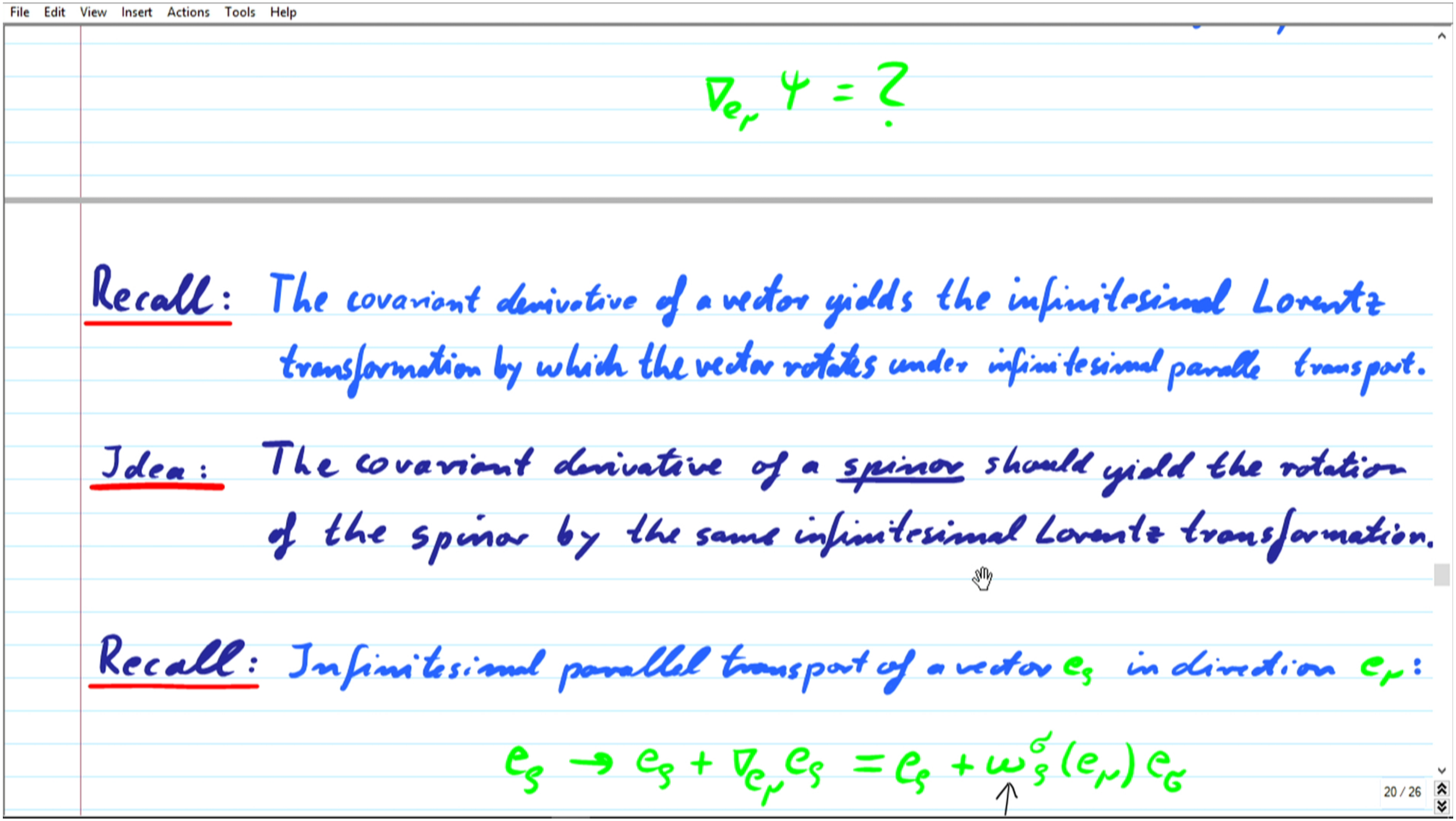


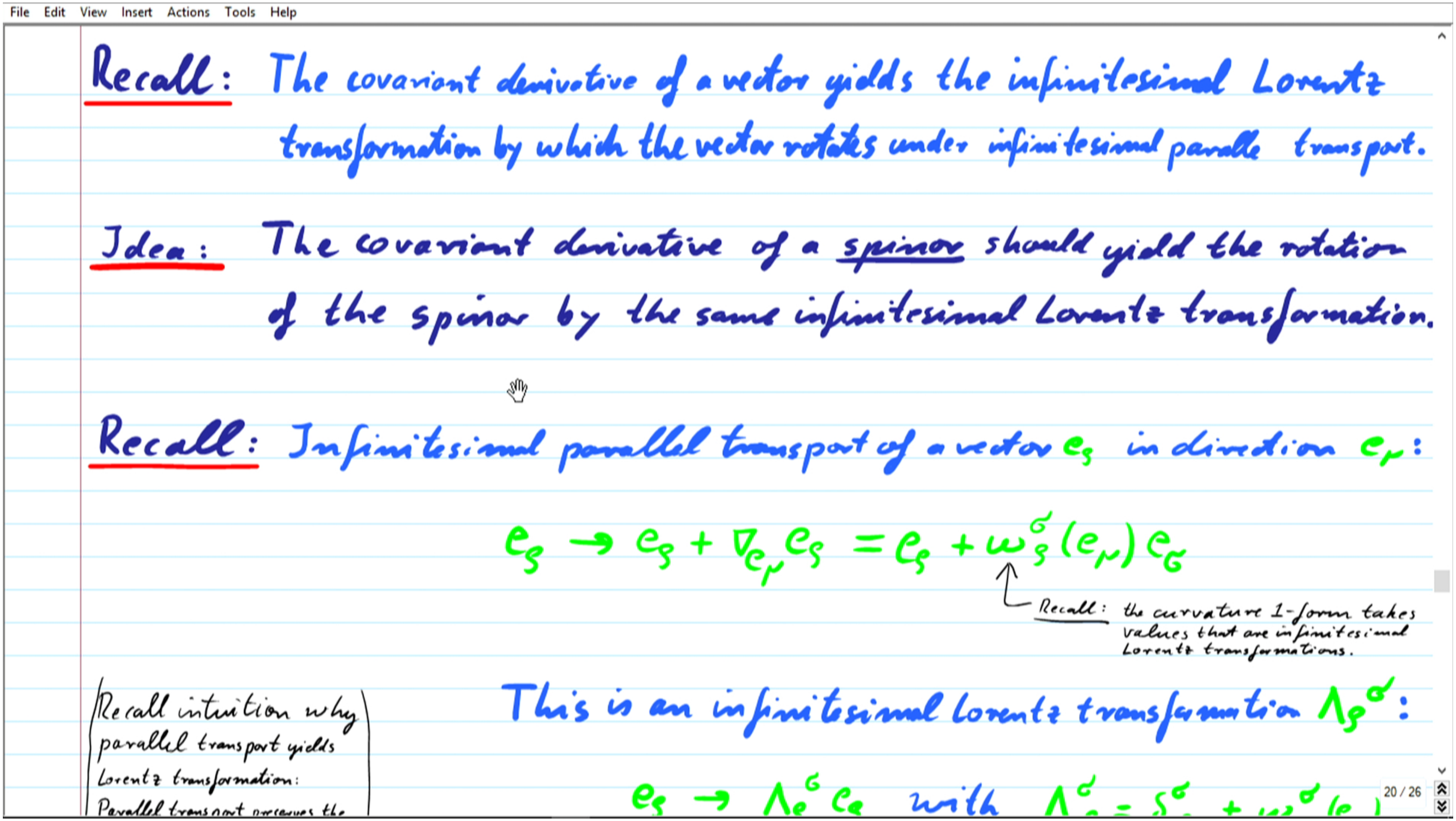
$$\nabla_{e_\mu} \psi = ?$$

Recall: The covariant derivative of a vector yields the infinitesimal Lorentz transformation by which the vector rotates under infinitesimal parallel transport.

Idea: The covariant derivative of a spinor should yield the rotation of the spinor by the same infinitesimal Lorentz transformation.

Recall: Infinitesimal parallel transport of a vector e_ς in direction e_μ:

$$e_\varsigma \to e_\varsigma + \nabla_{e_\mu} e_\varsigma = e_\varsigma + \underset{\uparrow}{\omega}{}^{\sigma}{}_{\varsigma}(e_\mu)\, e_\sigma$$



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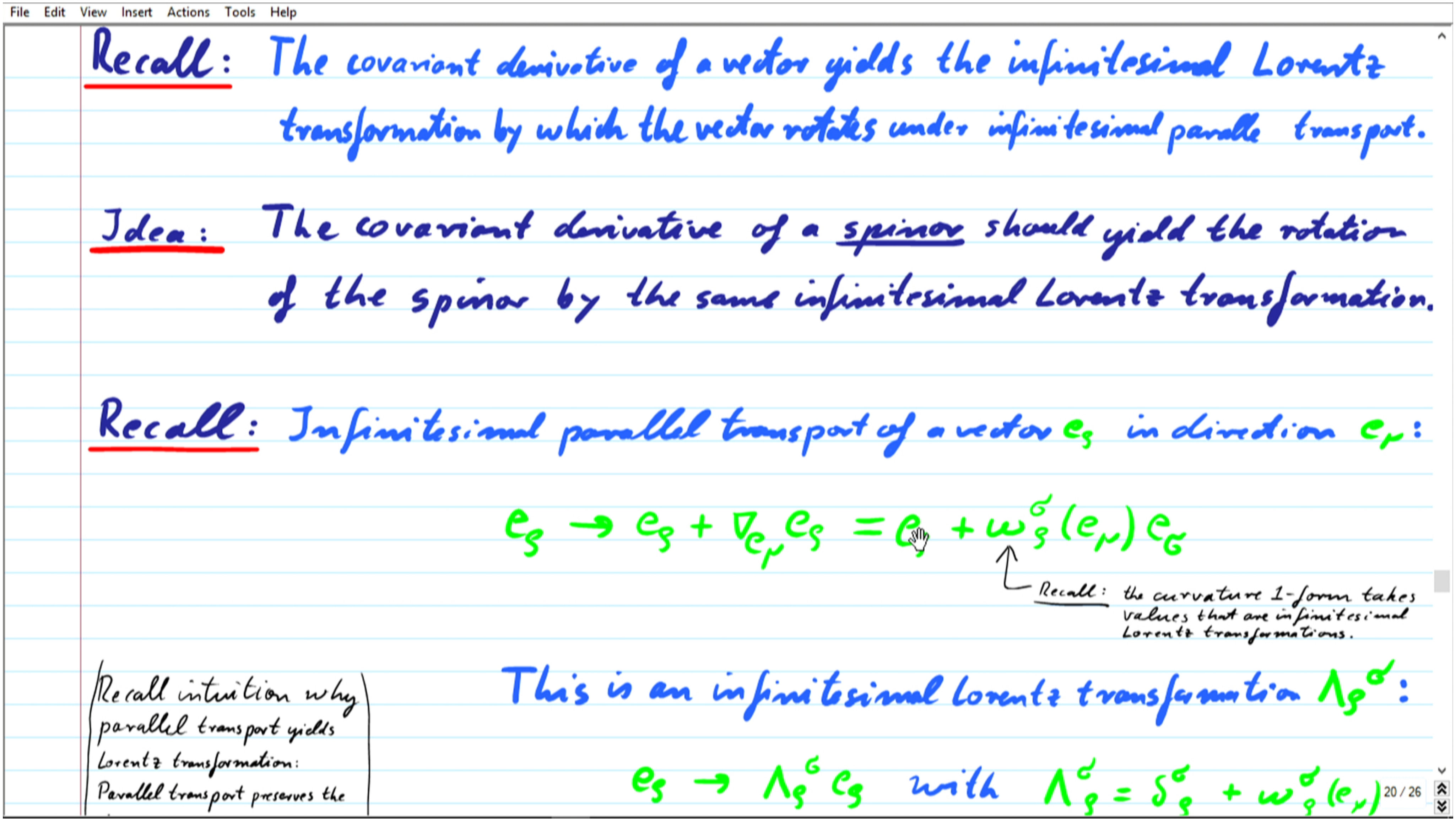
Recall: the curvature 1-form takes values that are infinitesimal Lorentz transformations.

Recall intuition why parallel transport yields Lorentz transformation:
Parallel transport preserves the

This is an infinitesimal Lorentz transformation Λ_ρ^σ:

$$e_\rho \to \Lambda_\rho^\sigma e_\sigma \quad \text{with} \quad \Lambda^\sigma = \delta^\sigma + \omega^\sigma (e$$



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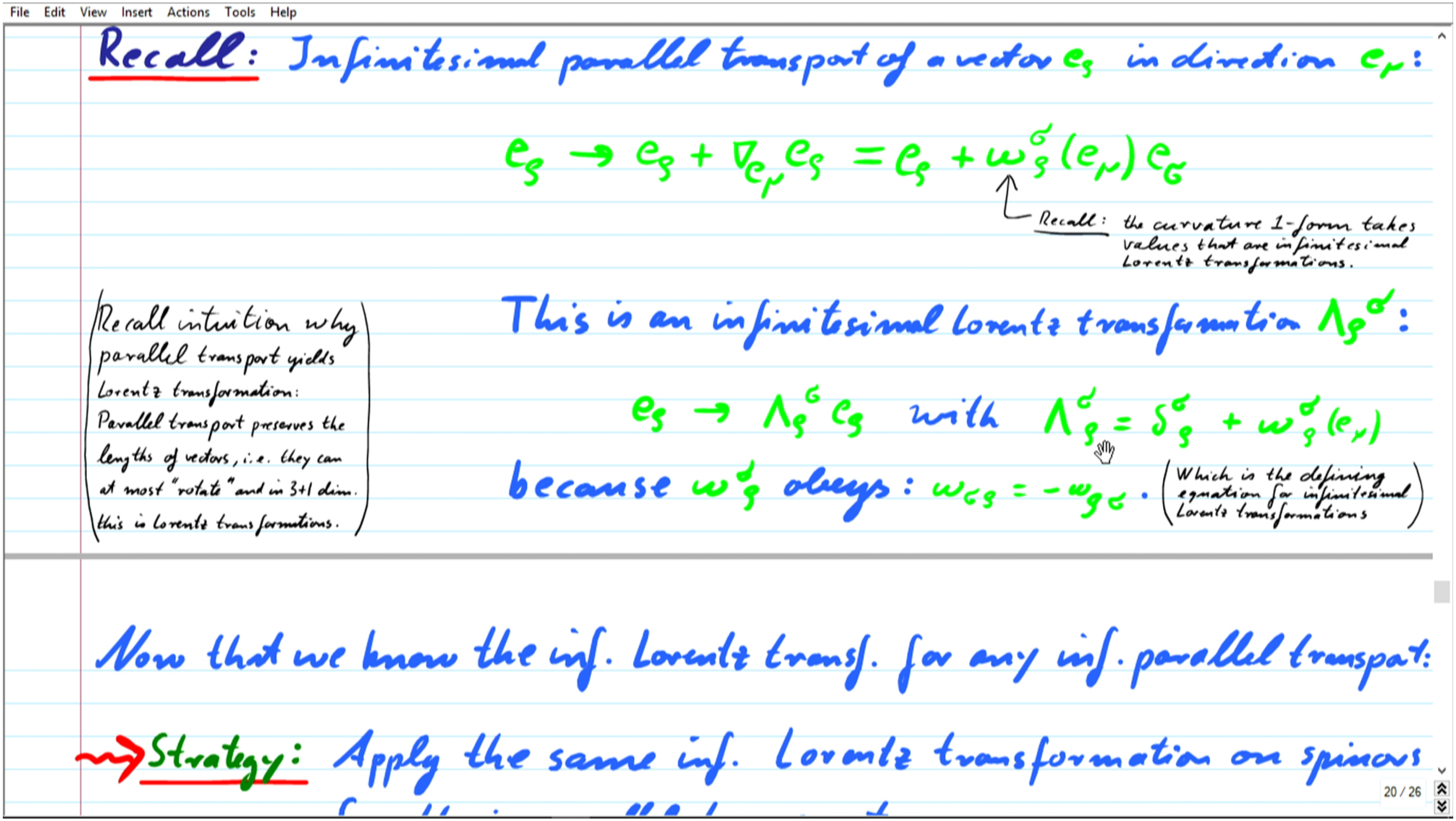

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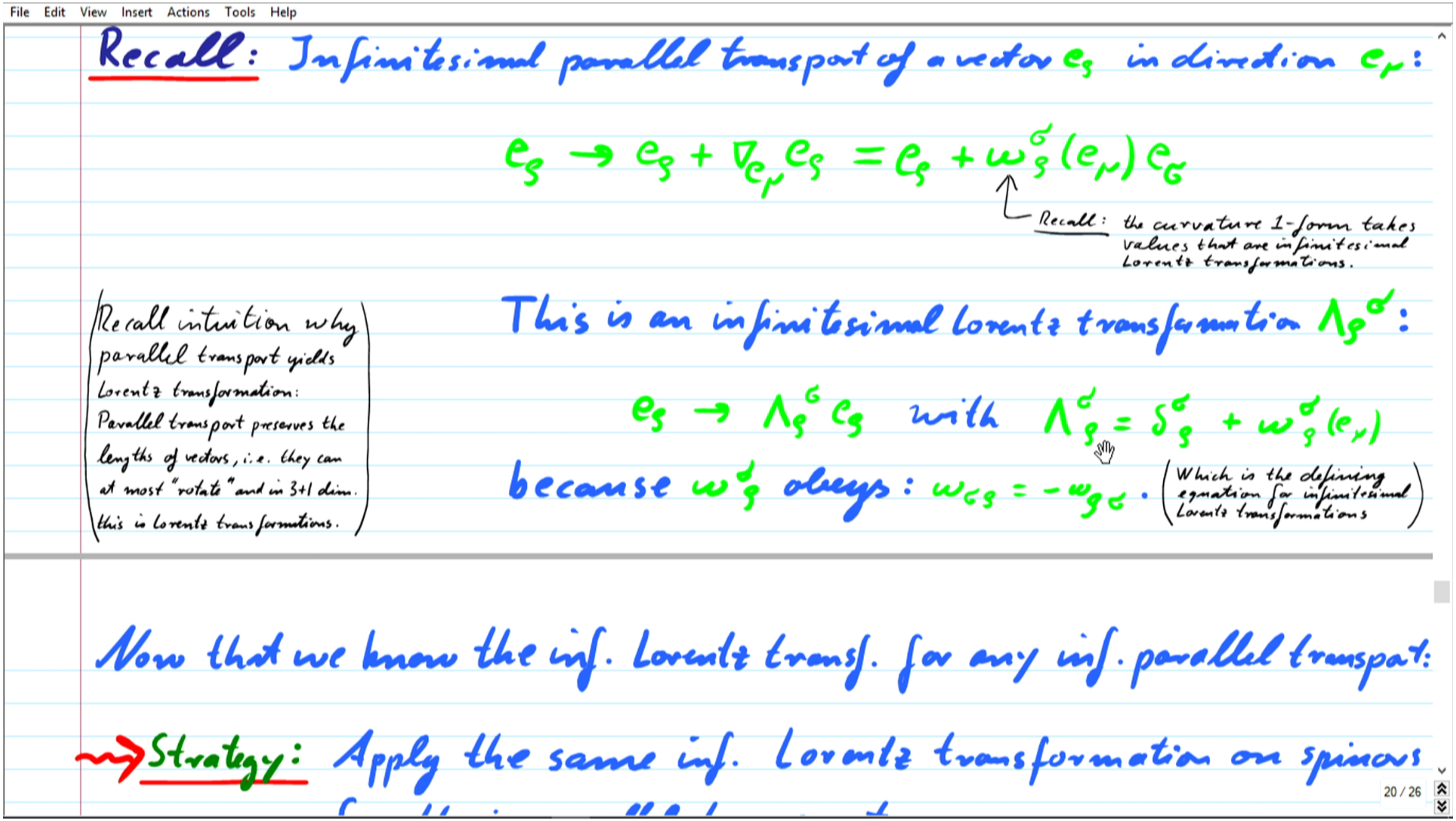
→ **Strategy:** Apply the same inf. Lorentz transformation on spinors



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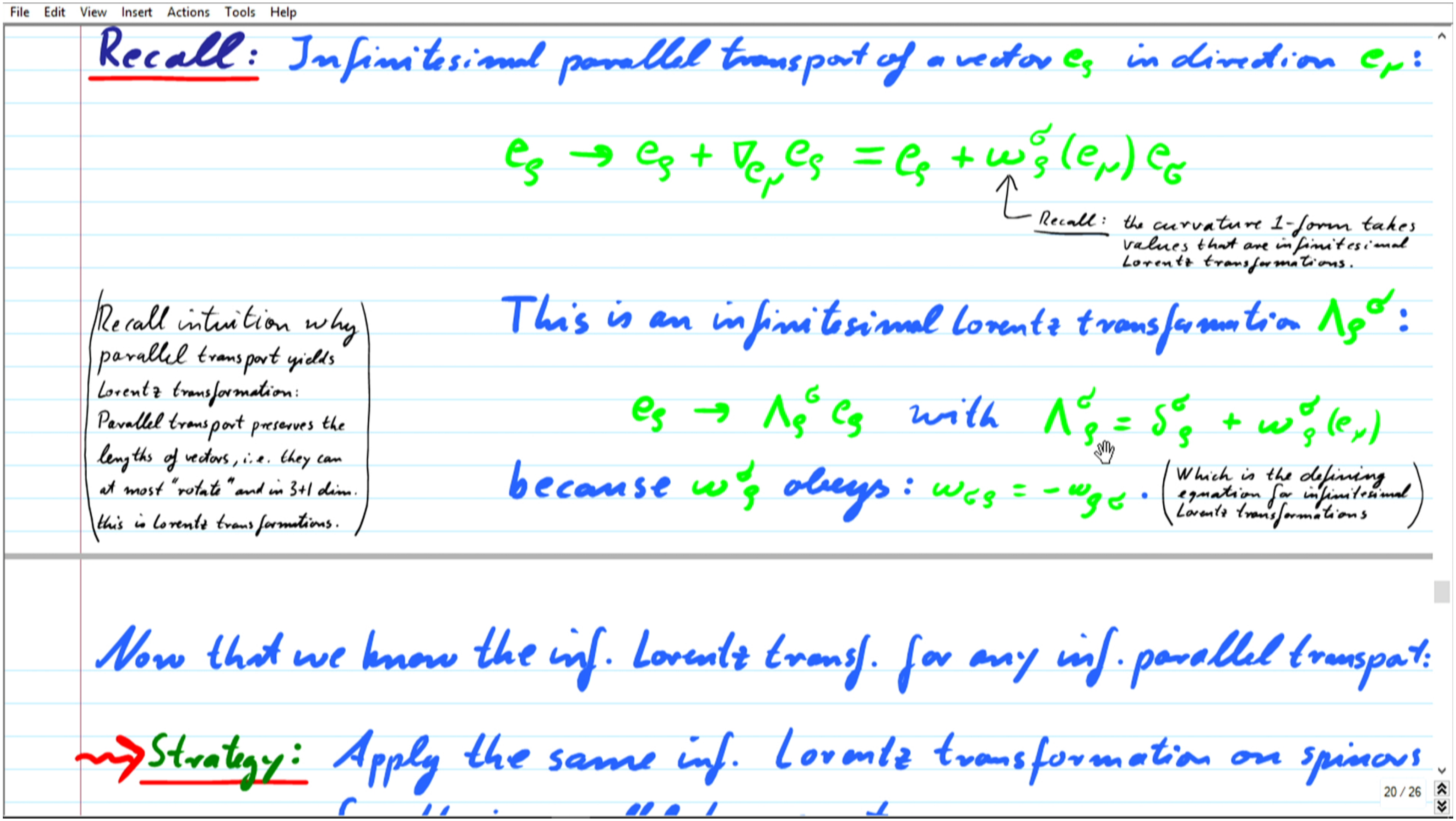

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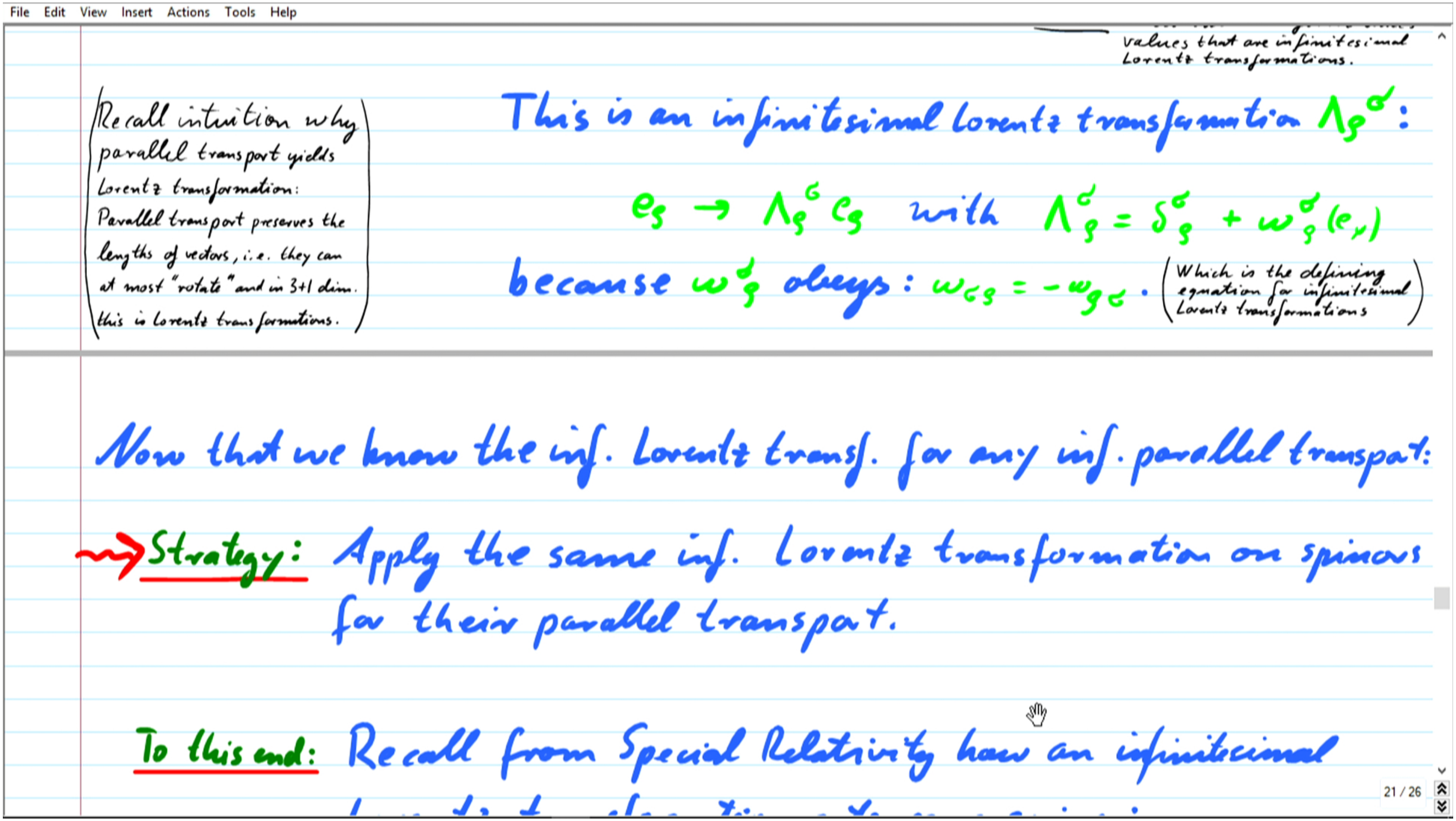

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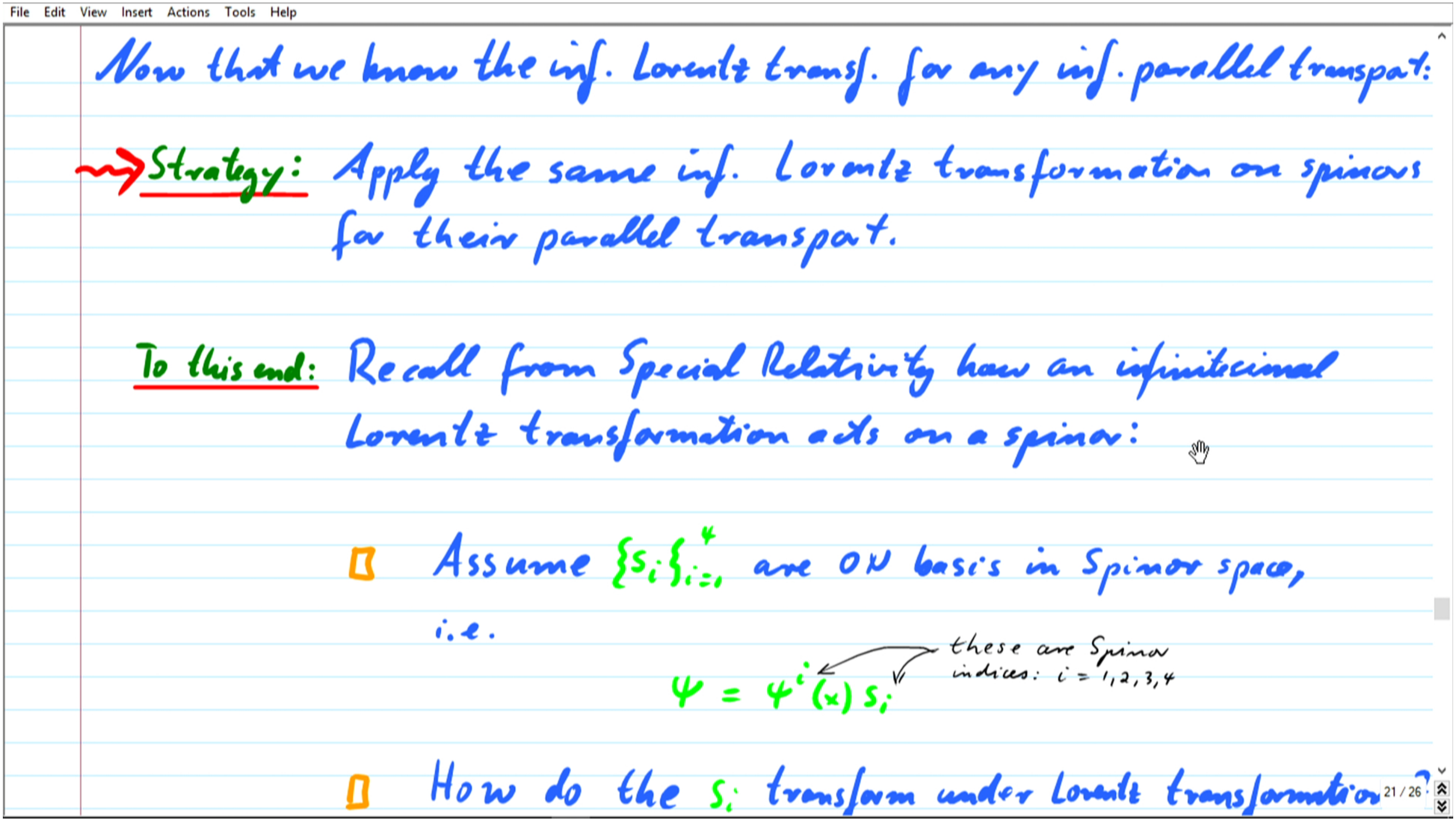

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Strategy: Apply the same inf. Lorentz transformation on spinors for their parallel transport.

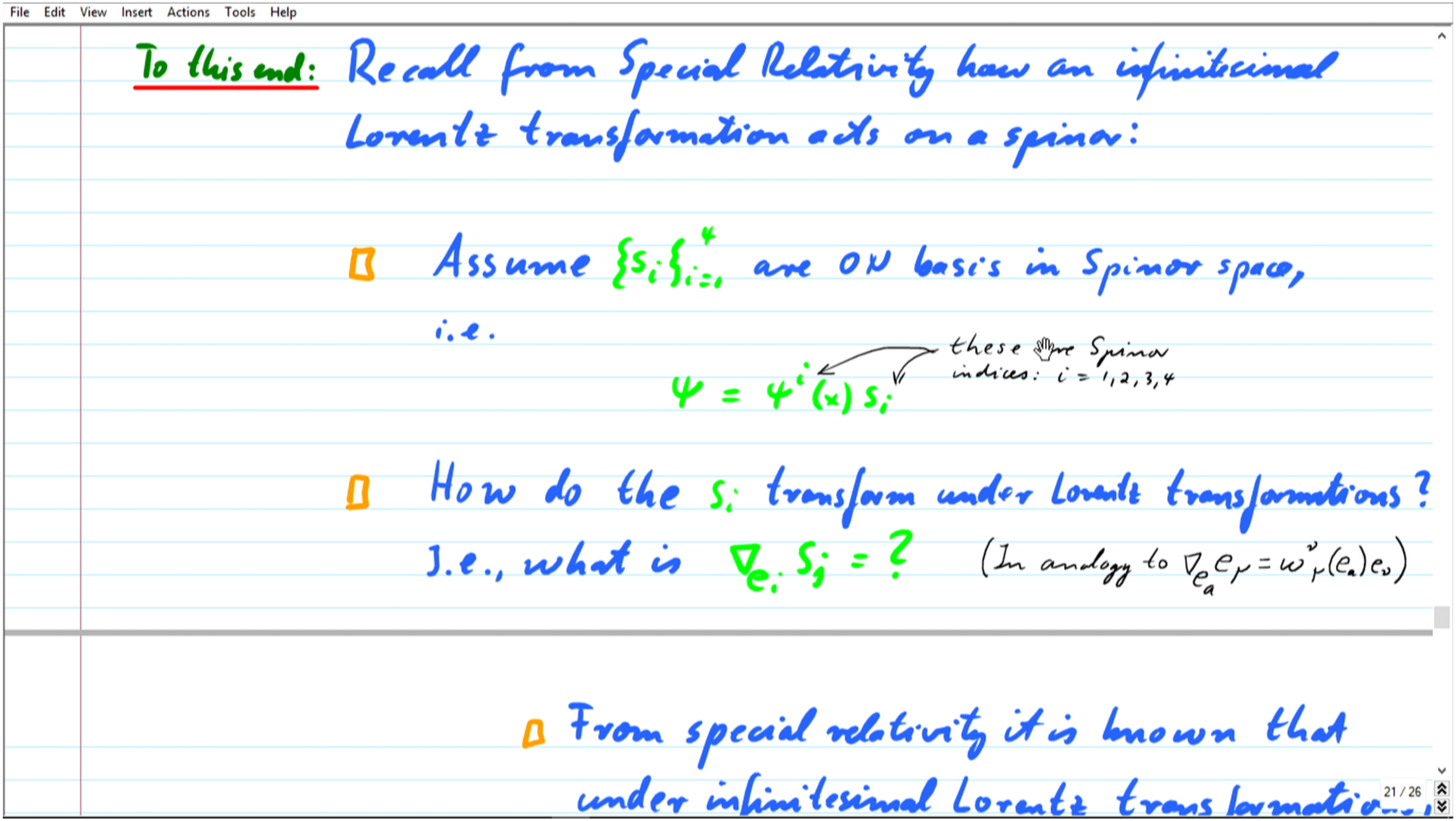
To this end: Recall from Special Relativity how an infinitesimal Lorentz transformation acts on a spinor:

☐ Assume $\{s_i\}_{i=1}^{4}$ are ON basis in Spinor space, i.e.

$$\psi = \psi^i(x)\, s_i$$

these are Spinor indices: $i = 1,2,3,4$

☐ How do the s_i transform under Lorentz transformation?



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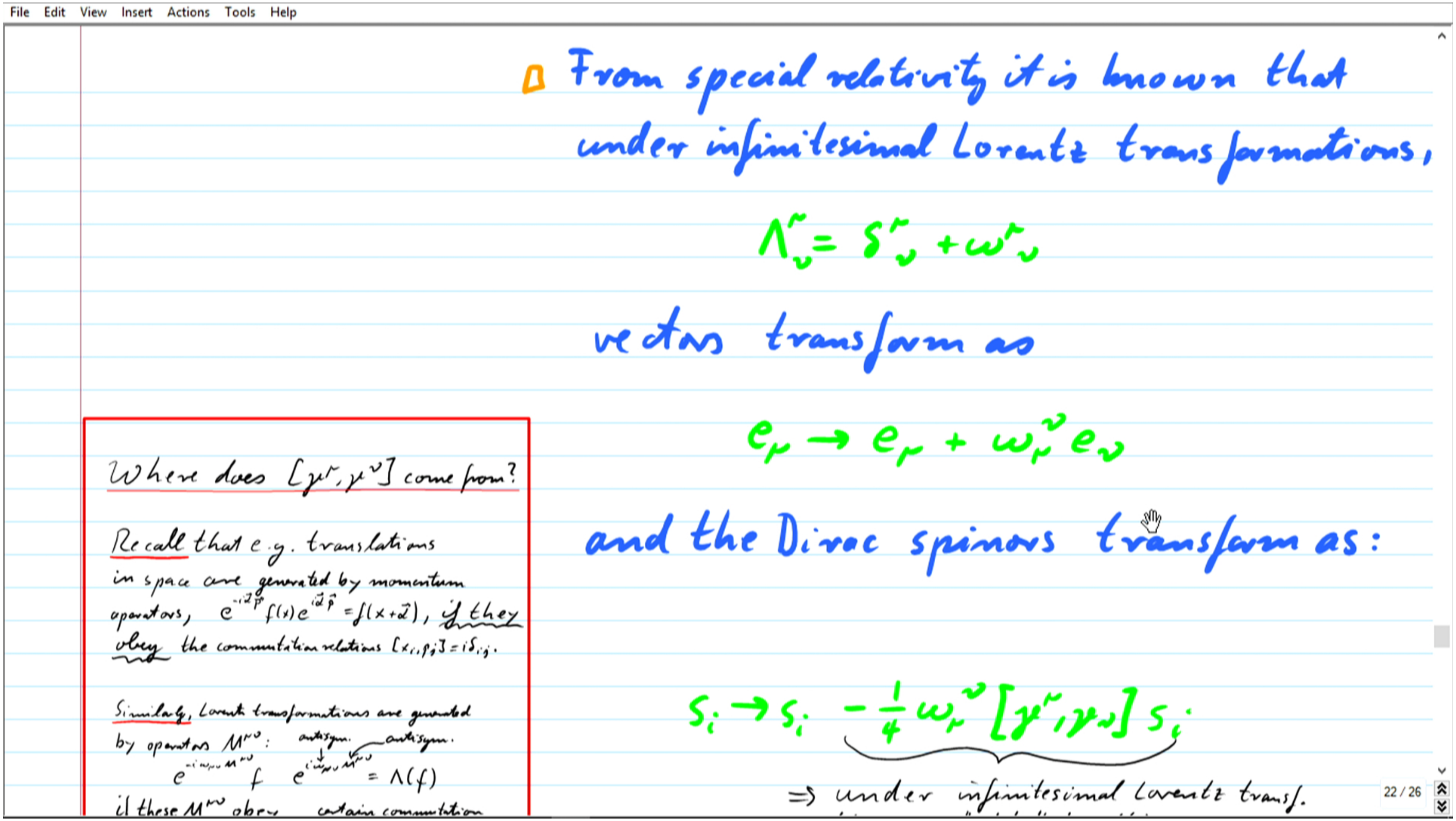
☐ How do the s_i transform under Lorentz transformations? I.e., what is $\nabla_{e_i} S_j = ?$ (In analogy to $\nabla_{e_a} e_\mu = \omega^\nu{}_\mu(e_a)\, e_\nu$)

☐ From special relativity it is known that under infinitesimal Lorentz transformatio...



From special relativity it is known that under infinitesimal Lorentz transformations,

$$\Lambda^\mu{}_\nu = \delta^\mu{}_\nu + \omega^\mu{}_\nu$$

vectors transform as

$$e_\mu \rightarrow e_\mu + \omega_\mu{}^\nu e_\nu$$

and the Dirac spinors transform as:

$$s_i \rightarrow s_i \underbrace{- \frac{1}{4} \omega_\mu{}^\nu [\gamma^\mu, \gamma_\nu] s_i}$$

$\Rightarrow$ under infinitesimal Lorentz transf.

Where does $[\gamma^\mu, \gamma^\nu]$ come from?

Recall that e.g. translations in space are generated by momentum operators, $e^{-i\vec{a}\vec{p}} f(x) e^{i\vec{a}\vec{p}} = f(x+\vec{a})$, if they obey the commutation relations $[x_i, p_j] = i\delta_{ij}$.

Similarly, Lorentz transformations are generated by operators $M^{\mu\nu}$:

$$e^{-i\omega_{\mu\nu}M^{\mu\nu}} f \; e^{i\overset{\overset{\text{antisym.}}{\downarrow}}{\omega_{\mu\nu}} \overset{\text{antisym.}}{M^{\mu\nu}}} = \Lambda(f)$$

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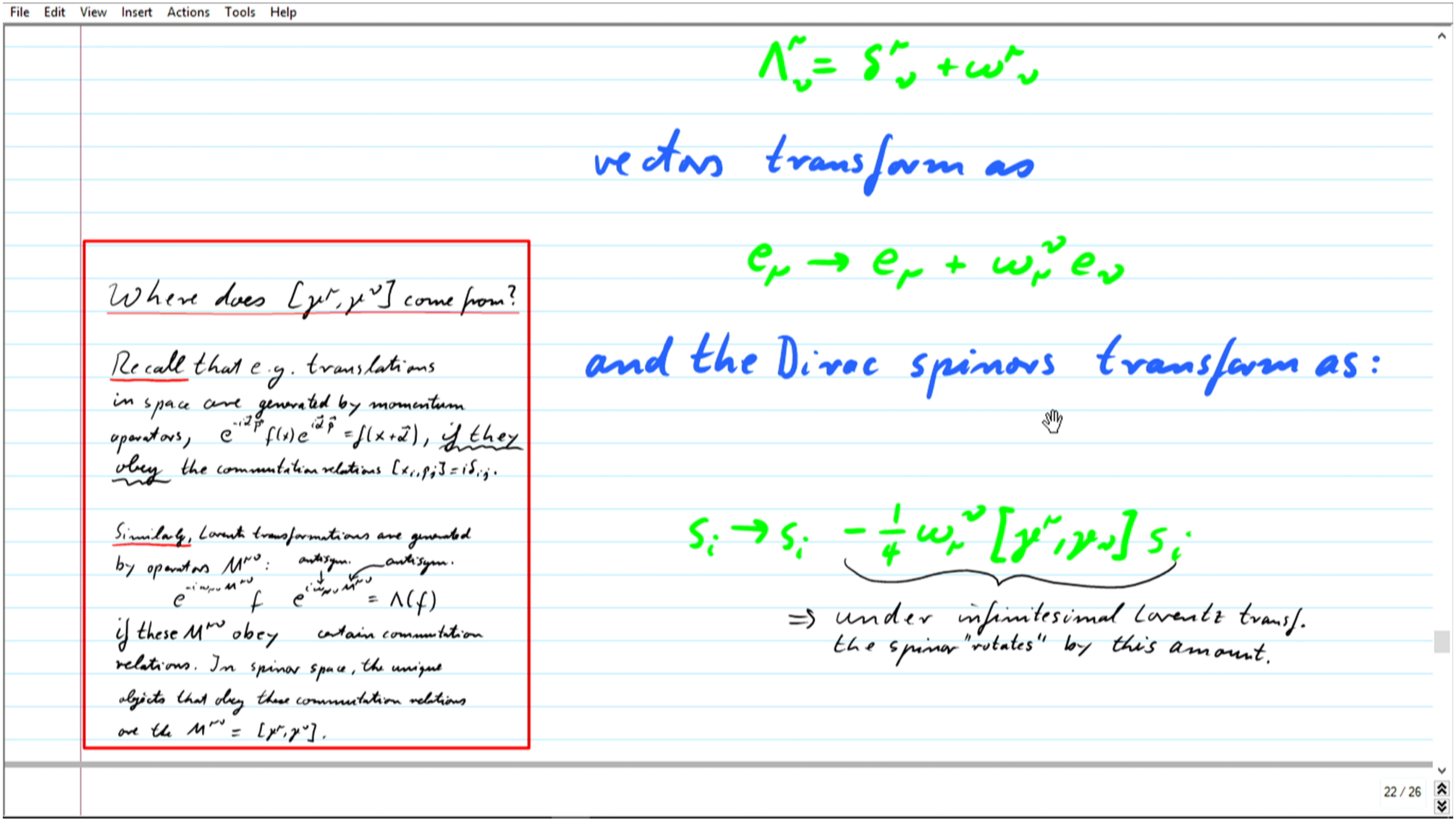

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⇒ under infinitesimal Lorentz transf. the spinor "rotates" by this amount.

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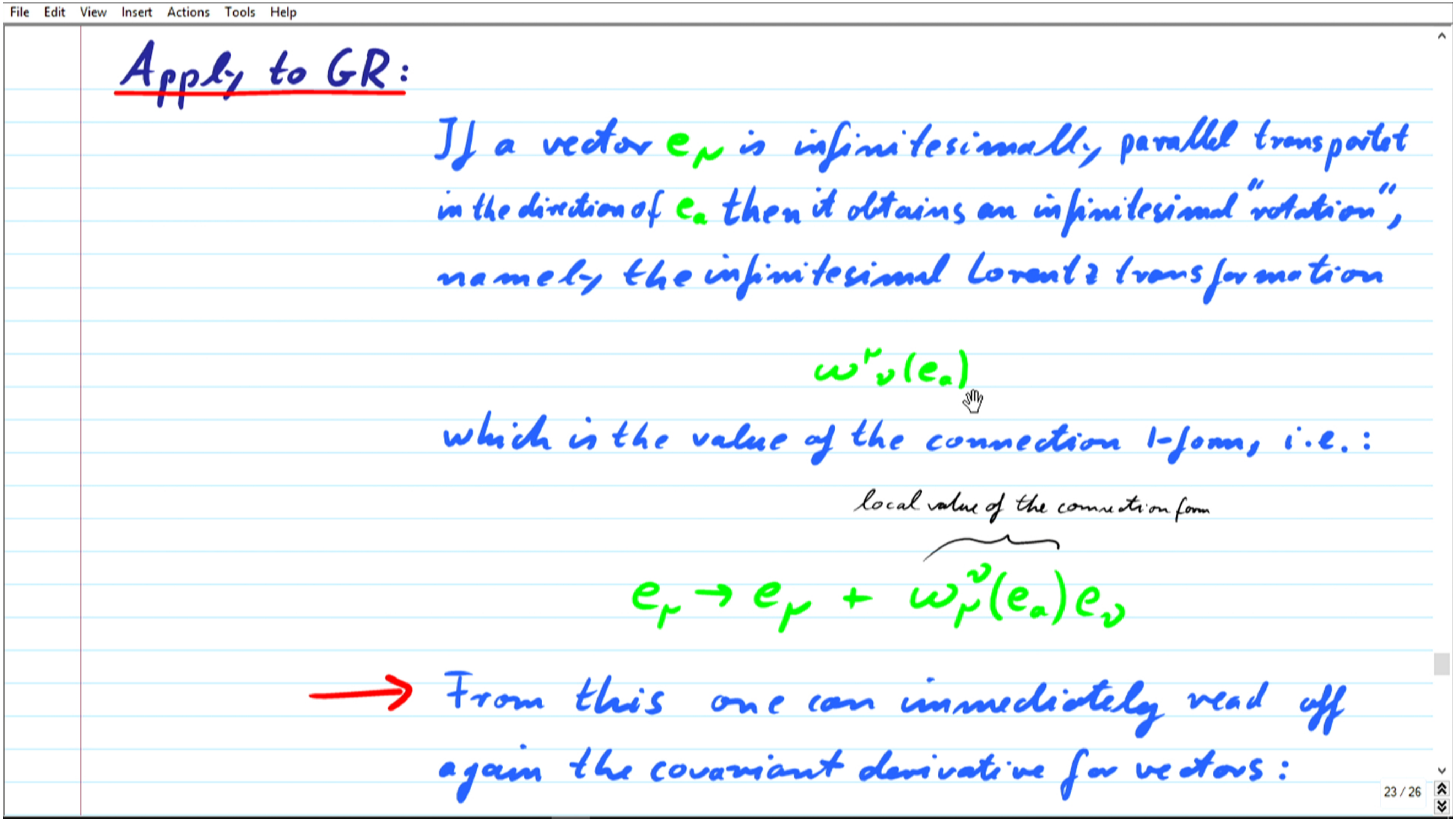
$$e^{-i\omega_{\mu\nu}M^{\mu\nu}} f \quad e^{i\overset{\text{antisym.}}{\omega_{\mu\nu}}\overset{\text{antisym.}}{M^{\mu\nu}}} = \Lambda(f)$$

if these $M^{\mu\nu}$ obey certain commutation relations. In spinor space, the unique objects that obey these commutation relations are the $M^{\mu\nu} = [\gamma^\mu, \gamma^\nu]$.



Apply to GR:

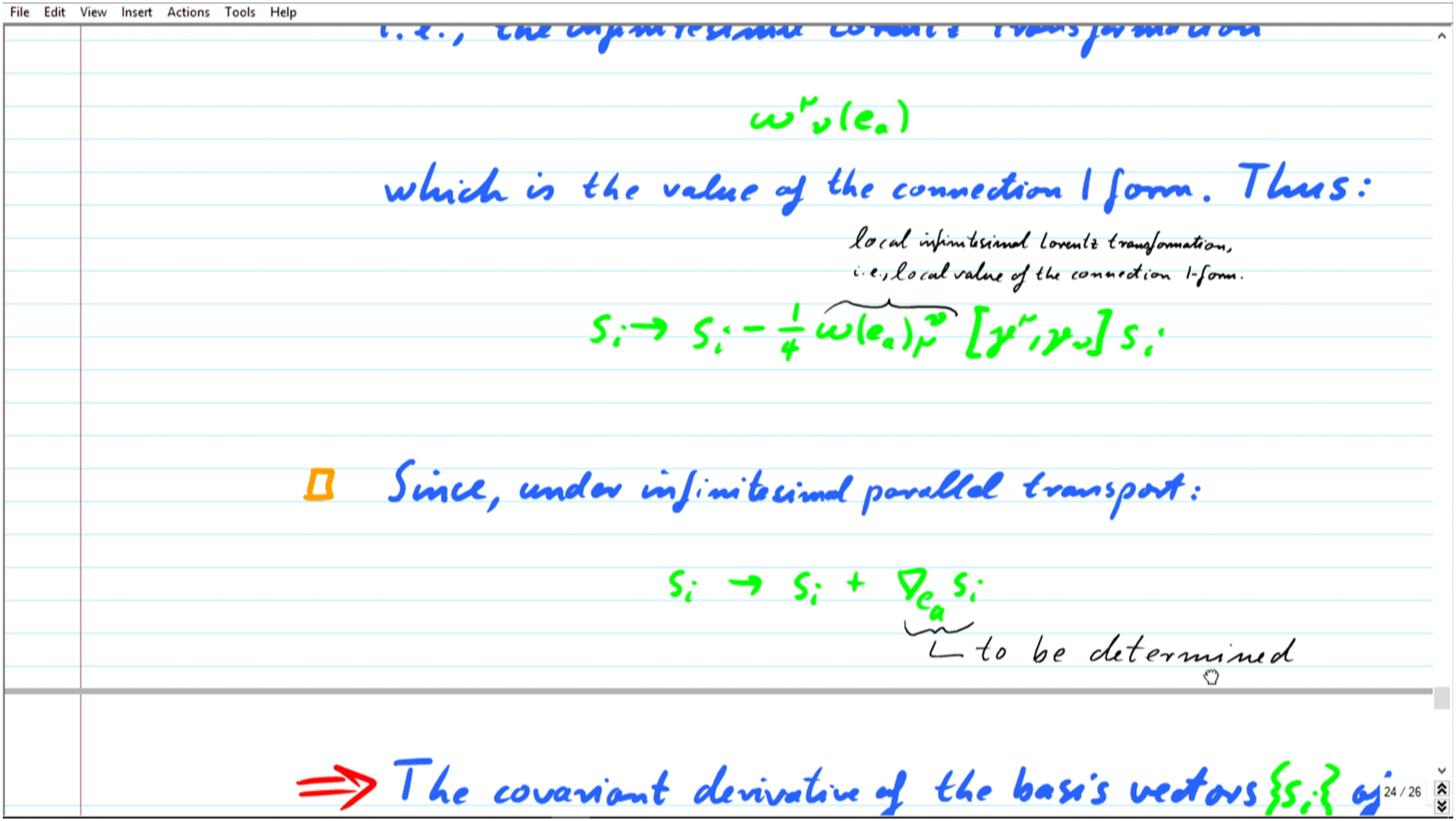
If a vector e_μ is infinitesimally parallel transportet in the direction of e_a then it obtains an infinitesimal "rotation", namely the infinitesimal Lorentz transformation

$$\omega^\mu{}_\nu(e_a)$$

which is the value of the connection 1-form, i.e.:

$$e_\mu \to e_\mu + \overbrace{\omega_\mu^\nu(e_a)}^{\text{local value of the connection form}} e_\nu$$

→ From this one can immediately read off again the covariant derivative for vectors:



i.e., the infinitesimal Lorentz transformation

$$\omega^\mu{}_\nu(e_a)$$

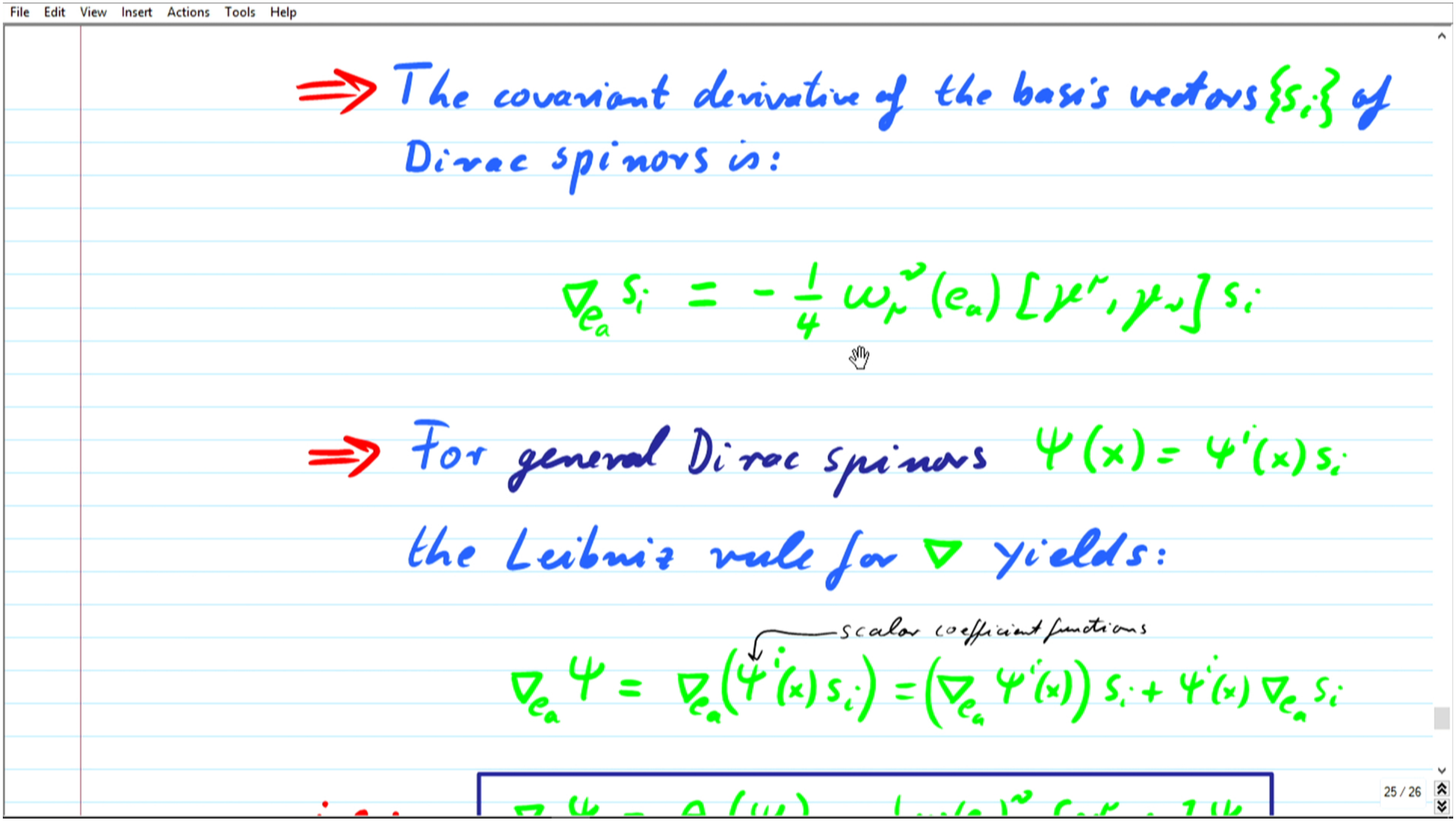
which is the value of the connection 1 form. Thus:

$$s_i \to s_i - \frac{1}{4}\overbrace{\omega(e_a)_\mu{}^\nu}^{\substack{\text{local infinitesimal Lorentz transformation,}\\ \text{i.e., local value of the connection 1-form.}}}[\gamma^\mu, \gamma_\nu]\, s_i$$

☐ Since, under infinitesimal parallel transport:

$$s_i \to s_i + \underbrace{\nabla_{e_a}}_{\text{to be determined}} s_i$$

$\Rightarrow$ The covariant derivative of the basis vectors $\{s_i\}$ of



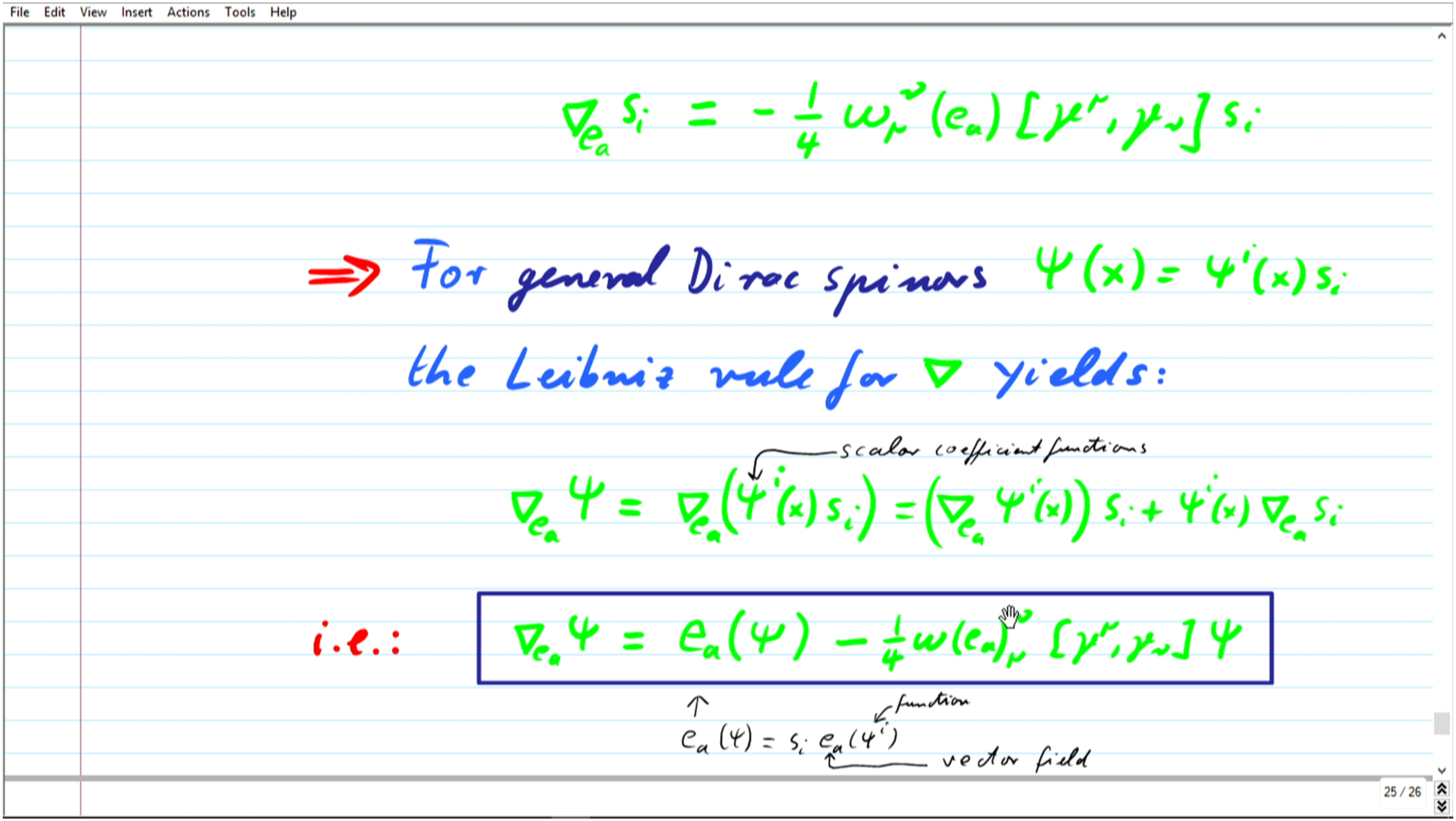
⇒ The covariant derivative of the basis vectors $\{s_i\}$ of Dirac spinors is:

$$\nabla_{e_a} s_i = -\frac{1}{4}\,\omega_\mu{}^\nu(e_a)\,[\gamma^\mu,\gamma_\nu]\,s_i$$

⇒ For general Dirac spinors $\psi(x) = \psi^i(x)\,s_i$ the Leibniz rule for ∇ yields:

scalar coefficient functions

$$\nabla_{e_a}\psi = \nabla_{e_a}\left(\psi^i(x)\,s_i\right) = \left(\nabla_{e_a}\psi^i(x)\right)s_i + \psi^i(x)\,\nabla_{e_a}s_i$$



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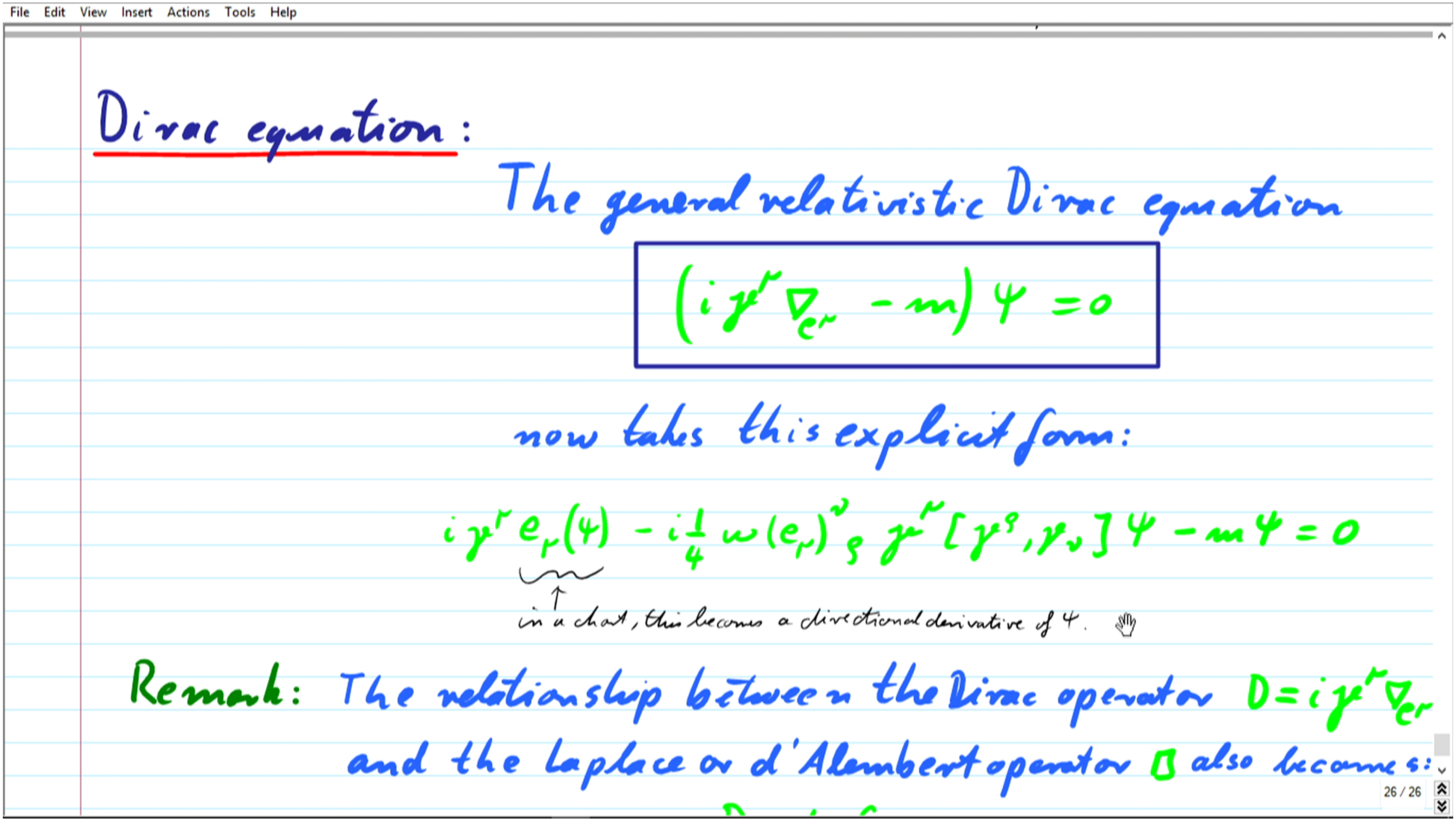
$$\nabla_{e_a}\psi = \nabla_{e_a}\big(\overset{\text{scalar coefficient functions}}{\overset{\downarrow}{\psi^i(x)}}\,s_i\big) = \big(\nabla_{e_a}\psi^i(x)\big)\,s_i + \psi^i(x)\,\nabla_{e_a}s_i$$

i.e.:

$$\boxed{\nabla_{e_a}\psi = e_a(\psi) - \frac{1}{4}\,\omega(e_a)_\mu{}^\nu\,[\gamma^\mu,\gamma_\nu]\,\psi}$$

$$\underset{\substack{\uparrow\\ e_a(\psi)\,=\,s_i\,\underset{\substack{\uparrow\\ \text{vector field}}}{e_a}\,(\overset{\substack{\text{function}\\ \swarrow}}{\psi^i})}}{}$$



Dirac equation:

The general relativistic Dirac equation

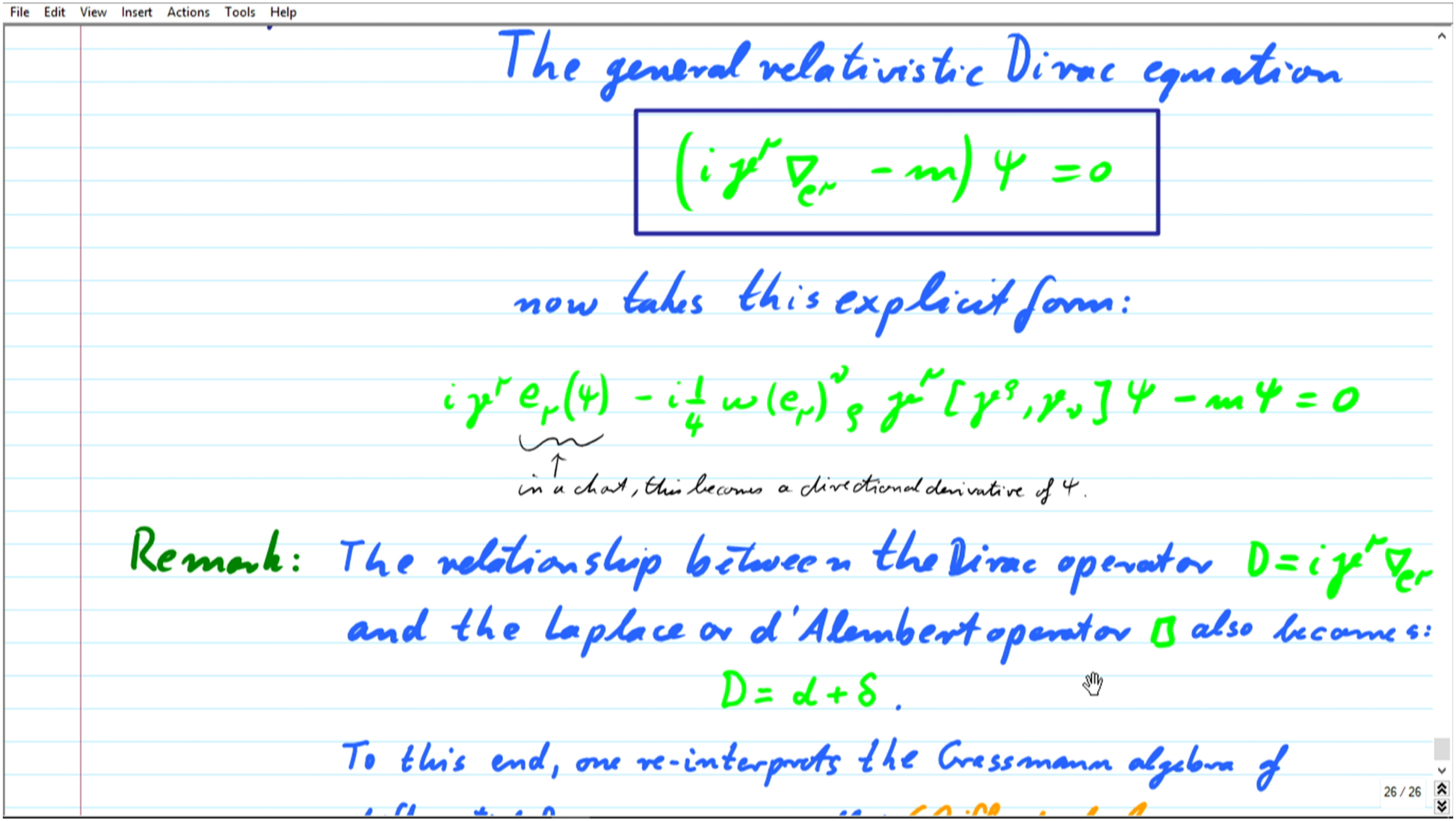
$$\boxed{\left(i\gamma^\mu \nabla_{e_\mu} - m\right)\psi = 0}$$

now takes this explicit form:

$$i\gamma^\mu \underbrace{e_\mu(\psi)}_{\uparrow} - i\frac{1}{4}\,\omega(e_\mu)^\nu{}_\rho\, \gamma^\mu[\gamma^\rho,\gamma_\nu]\psi - m\psi = 0$$

in a chart, this becomes a directional derivative of ψ.

Remark: The relationship between the Dirac operator $D = i\gamma^\mu\nabla_{e_\mu}$ and the Laplace or d'Alembert operator $\square$ also becomes:



The general relativistic Dirac equation

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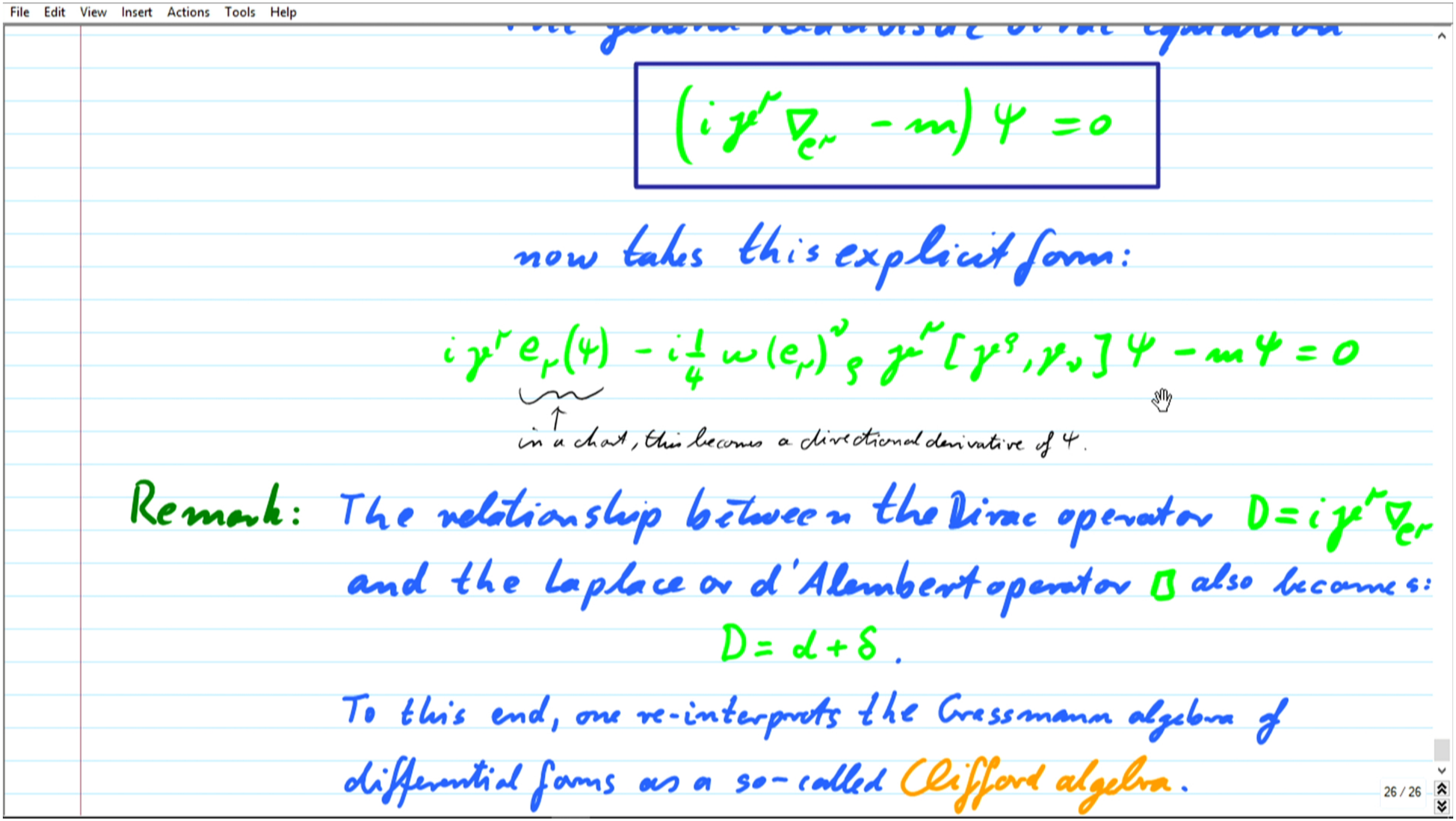
Remark: The relationship between the Dirac operator $D = i\gamma^\mu \nabla_{e_\mu}$ and the Laplace or d'Alembert operator $\Box$ also becomes:

$$D = d + \delta.$$

To this end, one re-interprets the Grassmann algebra of



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Remark: The relationship between the Dirac operator $D = i\gamma^\mu \nabla_{e_\mu}$ and the Laplace or d'Alembert operator $\Box$ also becomes:

$$D = d + \delta .$$

To this end, one re-interprets the Grassmann algebra of differential forms as a so-called Clifford algebra.