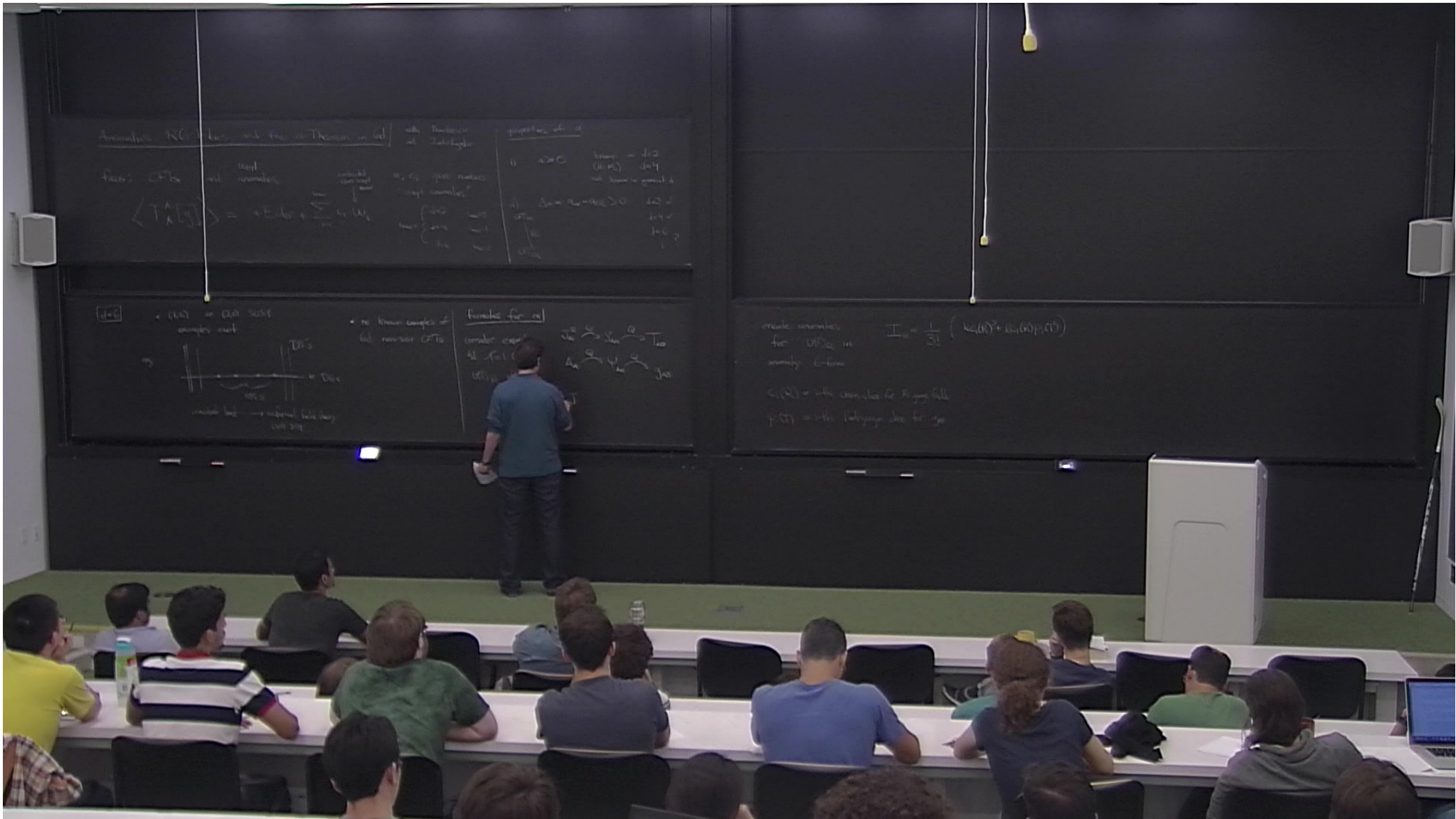


Title: Anomalies, RG Flows, and the a-Theorem in 6d

Date: Sep 29, 2015 02:00 PM

URL: <http://pirsa.org/15090018>

Abstract: <p>I will overview recent progress in understanding Weyl anomalies and the a-theorem in supersymmetric six-dimensional field theories.</p>



encode anomalies  
for  $U(1)_R$  in  
anomaly 6-form

$$I_6 = \frac{1}{3!} \left( k c_1(R)^3 + l c_1(R) p_1(T) \right)$$

$c_i(R)$  =  $i$ -th Chern class for  $R$  gauge field

$p_i(T)$  =  $i$ -th Pontryagin class for  $g_{AB}$



• no known examples of  
6d non-susy CFTs

### formulas for a

consider example

4d  $\mathcal{N}=1$  CFTs

$U(1)_R$  global sym

$$J_A^R \xrightarrow{Q} S_{AA} \xrightarrow{Q} T_{AB}$$

$$A_A \xrightarrow{Q} \psi_{Ax} \xrightarrow{Q} g_{AB}$$

$$J_A^A \xrightarrow{Q} T_A^A$$

$$\underline{I}_G = \frac{1}{3!} \left( k c_1(R)^3 + l c_1(R) p_1(T) \right)$$

$$a = \frac{9}{8} k - \frac{3}{8} l$$

units  $a \left( \frac{N=4}{\text{vol}} \right) = 1$ .

$k, l$  are computable away from conformal limit

$$I_6 = \frac{1}{3!} \left( k_G(R)^3 + l_G(R)P_1(T) \right)$$

$$a = \frac{9}{8}k - \frac{3}{8}l$$

units  $a_{(D=4)} = 1$ .

class for R gauge field  
origin class for g<sub>AB</sub>

k, l are computable away from conformal limit  
explicit evidence for a-thm



6d (1,0)

$SU(2)_R$

$$\mathbb{I}_8 = \frac{1}{4!} \left( \alpha C_2(R)^2 + \beta C_2(R) P_1(T) + \gamma P_1(T)^2 + \delta P_2(T) \right)$$

$$a = \frac{16}{7} (\alpha - \beta + \gamma) + \frac{6}{7} \delta$$

main result  
(restricted action)

conformal limit  $\rightarrow$  conformal field theory  
(1,0) susy

Examples

• (2,0) SCFTs labeled by ADE Lie algebra  $\mathfrak{g}$

$$a = \frac{16}{7} \left[ h_{\mathfrak{g}} d_{\mathfrak{g}} + r_{\mathfrak{g}} \right] \stackrel{\text{SUSY}}{=} \frac{16}{7} N^3 - \frac{9}{7} N - 1$$

$\swarrow$  dimension  
 $\uparrow$  Dual Coxeter #  
 $\uparrow$  rank



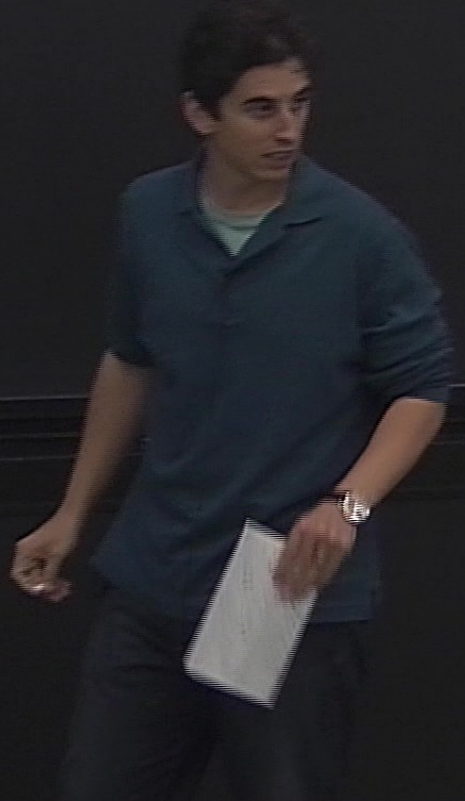
field theory  
on

labeled by ADE Lie algebra  $\mathfrak{g}$

on

$$\stackrel{\text{SUSY}}{=} \frac{16}{7} N^3 - \frac{9}{7} N - 1 \quad \left( \frac{1}{N} \text{ corrections to } N^3 \text{ observed via holography} \right)$$

ink

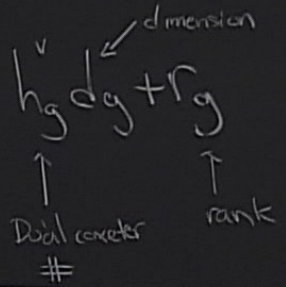


(1,0) SUSY

# Examples

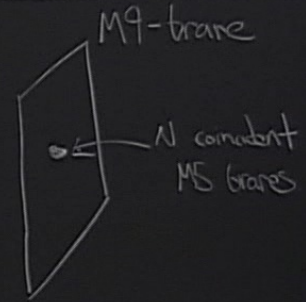
• (2,0) SCFTs labeled by ADE Lie algebra  $g$

$$a = \frac{16}{7} \left( h_g d_g + r_g \right) \stackrel{\text{SUSY}}{=} \frac{16}{7} N^3 - \frac{9}{7} N - 1$$



(1,1) corrections

(1,0)  $E_N$   $N$  small  $E_8$  instantons



(1/3) SUSY

labelled by ADE Lie algebra  $\mathfrak{g}$

$$S(U(N)) \equiv \frac{16}{7} N^3 - \frac{9}{7} N - 1$$

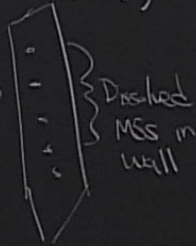
( $1/N$  corrections to  $N^3$  observed via holography)

dimension  
 $\mathfrak{g} + \mathfrak{r}_\mathfrak{g}$   
rank  
N small instanton

M9-brane

N coincident MS branes

RG  $\rightarrow$



$$a(E_N) = \frac{64}{7} N^3 + \frac{144}{7} N^2 + \frac{99}{7} N$$

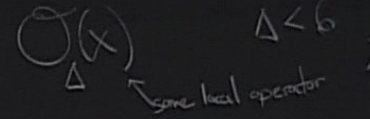
$\mathcal{M}_{Higgs}$  = moduli space of N  $E_8$  instantons

$$\Delta a = \frac{64}{7} N^3 + \frac{144}{7} N^2 + \frac{60}{7} N > 0$$



- Explicitly break conformal symmetry

$$\delta\mathcal{L} = \text{Tr}(\mathcal{O}(x)) \quad \Delta < 6$$


  
 ← some local operator

impossible if SUSY is preserved

- Spontaneously break conformal symmetry

$\Leftrightarrow$  move onto moduli space of vacua |  $\exists$  Goldstone boson  $\varphi$  dilaton  
 $\langle \mathcal{O} \rangle \neq 0$

• Higgs branch  $SU(2)_R$

$\varphi \in$  hypermultiplet

4 scalars = 3 + 1  
↑ ↑  
goldstons dilatons  
for  $SU(2)_R$

• Tensor branch  $SU(2)_R$

$\varphi \in$  tensor multiplet

$\varphi, \psi_i$  ← doublet of  $SU(2)_R$

$\mathcal{L}_3 = H_{ABC} = *H$

$$\mathcal{L}_\varphi = \frac{1}{2} (\partial\varphi)^2 + b (\partial\varphi)^4 + \underbrace{\Delta a (\partial\varphi)^6}_{\text{why?}} + \dots$$

$$\Delta a = a_{UV} - a_{IR}$$

WZ term makes  $\varphi$  a mismatch

not Weyl invariant  
↓

$$\int \supseteq \Delta a \int \log(\varphi) \text{Euler}(g) + \dots + (\partial\varphi)^6 \leftarrow$$

or even in flat space  
 $g = \eta$

$$\iint_g = \Delta a \int \sigma \text{Euler}(g)$$



$$\Delta a (\partial\phi)^6 + \dots$$

$$\Delta a = a_{UV} - a_{IR}$$

why?

n even in  
-lat space  
g=1

$$+ (\partial\phi)^6$$

Need to SUSY complete

$$\mathcal{L} = \mathcal{L}_{\text{Free}}^{(1,0) \text{ tensor multiplet}} + Q^8(\sigma)$$

8 is total # of Qs

eg/  
6 term

$$Q^8(\phi^4) \supset (\partial\phi)^4 + \dots$$

# Anomalies, RG Flows, and the a-Theorem in $6d$

with Donatello  
and Intriligator

focus: CFTs and Weyl anomalies.

$$\langle T_A^A[g] \rangle = a \text{ Euler} + \sum_{i=1}^{i_{\max}} c_i W_i$$

constructed  
from Weyl  
tensor

$a, c_i$  pure numbers  
"Weyl anomalies"

$$i_{\max} = \begin{cases} d=2 & i_{\max}=0 \\ d=4 & i_{\max}=1 \\ d=6 & i_{\max}=3 \end{cases}$$

problem in  $\mathcal{G}el$

with Dumitrescu  
and Intriligator

constructed  
from Weyl  
tensor.

$c_i, W_i$

$a, c_i, p$  users

Weyl anomaly

$l_{max} =$

$l_{max} = 0$

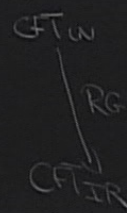
$l_{max} = 1$

$l_{max} = 3$

### properties of $a$

1)  $a \geq 0$  known in  $d=2$   
(H. M.)  $d=4$   
not known in general  $d$

2)  $\Delta a = a_{UV} - a_{IR} > 0$   $d=2 \checkmark$   
 $d=4 \checkmark$   
 $d=6 ?$   
 $\vdots$

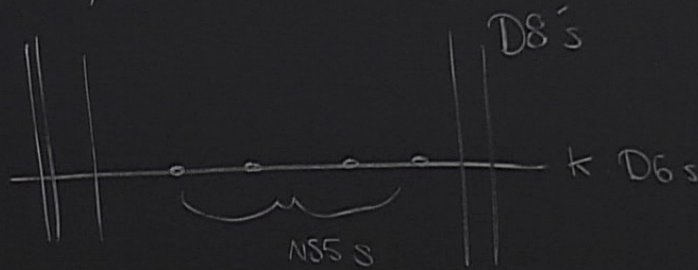




$d=6$

- $(1,0)$  or  $(2,0)$  SUSY examples exist

eg.



coincident limit  $\rightarrow$  conformal field theory  
 $(1,0)$  SUSY

- no known examples of 6d non-susy CFTs

formulas for

- no known examples of 6d non-susy CFTs

D8's

$k$  D6's

al field theory  
susy

### formulas for a

consider example

4d  $\mathcal{N}=1$  CFTs

$U(1)_R$  global sym

$$J_A^R \xrightarrow{Q} S_{AAx} \xrightarrow{Q} T_{AB}$$

$$A_A \xrightarrow{Q} \psi_{Ax} \xrightarrow{Q} g_{AB}$$

$(\partial\varphi)^6$  does not arise in  $Q^S(\sigma)$  for any  $\sigma$ !

$Q^S (\partial_A \varphi \partial^A \varphi \varphi^4) \stackrel{\text{EOM}}{\sim} Q^S \partial_A (\partial^A \varphi \varphi^4)$  no change in action

$\sigma = \varphi^2 \begin{pmatrix} \psi^i_1 \\ \alpha_1 \end{pmatrix} \begin{pmatrix} \psi^i_2 \\ \alpha_2 \end{pmatrix} \begin{pmatrix} \psi^i_3 \\ \alpha_3 \end{pmatrix} \begin{pmatrix} \psi^i_4 \\ \alpha_4 \end{pmatrix} \in \text{dir } \sigma_4$  no R inv component!

↑  
Lorentz scalar  
 $SU(2)_R$  inv



$Q^S(\sigma)$  for any  $\sigma$  ?!

EQM  $\sim Q^S \int dA (d^4\varphi \ q^4)$

change  
ctm

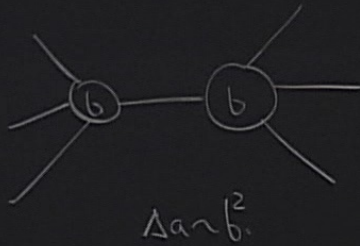
$\cup_{\alpha_3} \cup_{\alpha_4} \in \mathbb{R}^4$  no R

wrong to view  $\Delta a$  as infinitesimal def  
instead it arises via SUSY completion  $(\partial\varphi)^4$

universal  
relation!

$$\Delta a = \frac{9304\pi^2}{7} b^2$$

$S(2)_R$  inv



6pt  $b$  derivative  
amplitude

$$\Delta \Gamma_3 =$$

6pt 6 derivative  
amplitude

$$\Delta I_g = (x c_2(R) + y p_1(T))^2$$

$x, y$  same coeffs

$$\Delta \alpha = x^2; \Delta B = 2xy; \Delta \gamma = y^2; \Delta S = 0$$

GS interactions  $J_{GS} = \int B_1 (x c_2(R) + y p_1(T))$

implies  $\Delta I_g = ( )^2$

$$b \sim (x, y)$$



SUSY  
complete  
conformal  
 $R^2$  SUGRA

(Salam 1987)

$$\mathcal{L}_{\text{SUGRA}} = B_1 (x c_2(R) + y p_1(T)) + (y - \frac{x}{4}) R^{ABCD} R_{ABCD}$$

$$I_{\text{SUGRA}} = B \wedge (x c_2(R) + y p_1(T)) + (y - \frac{x}{4}) R^{ABCD} R_{ABCD} + \frac{3}{2} x R_{[AB} R_{CD]}$$

restrict to conformally flat

$$b = (x - y) \Rightarrow \boxed{\Delta a = \frac{16}{7} (\Delta x - \Delta y)} \implies a = \frac{16}{7} (\alpha - \beta + \gamma) + \frac{6}{7} \delta$$

Lorentz scalar  
 $SU(2)_R$  inv

### Open problems

- $a \geq 0$  Why? (true in examples)
- $\Delta a > 0$  for non SSB flows  
Higgs branch

$$\Delta I_8 = (x C_2(R) + y P_1(T))^2$$

$$\Delta \alpha = x^2, \Delta \beta = 2xy, \Delta \gamma = y^2, \Delta \delta = 0$$

$$\text{GS interactions } \int_{\text{dss}} = \int B_1 (x C_2(R) + y P_1(T))$$

$$b \sim (x, y)$$

$x, y$  same cells

implies