

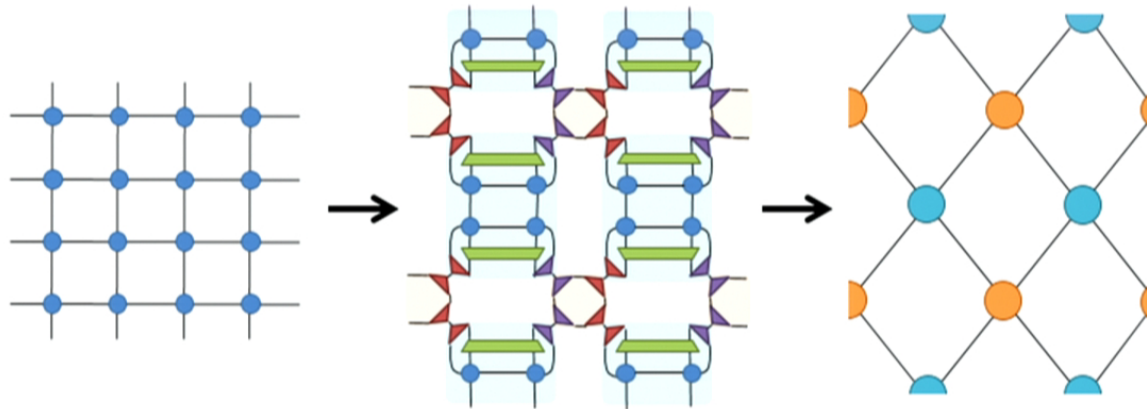
Title: Tensor Network Renormalization and the MERA

Date: Aug 18, 2015 09:00 AM

URL: <http://pirsa.org/15080065>

Abstract: I describe a class of non-perturbative renormalization group (RG) transformations which, when applied to the (discrete time) Euclidean path integrals of a quantum systems on the lattice, can give results consistent with conformal transformations of quantum field theories. In particular, this class of transformation, which we call Tensor Network Renormalization (TNR), is shown to generate a scale-invariant RG flow for quantum systems at a critical point. Applications of TNR towards study of quantum critical systems, and its relationship to the multi-scale-entanglement renormalization ansatz (MERA) for ground and thermal states of quantum systems, will be discussed.

# Tensor Network Renormalization and the MERA



Glen Evenbly  
Guifre Vidal

Tensor Network Renormalization, arXiv:1412.0732

Tensor Network Renormalization yields the MERA, arXiv:1502.05385



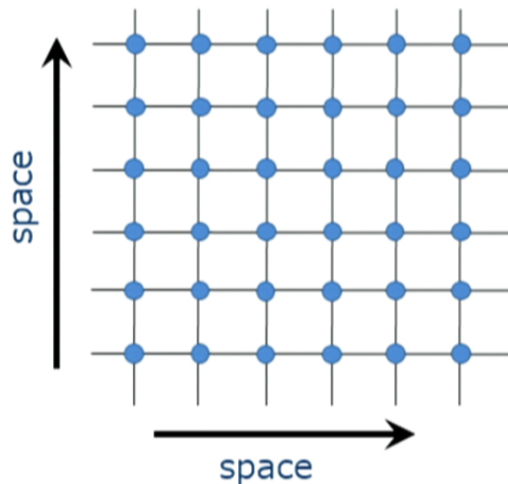


# Overview

express many-body system  
as a tensor network:

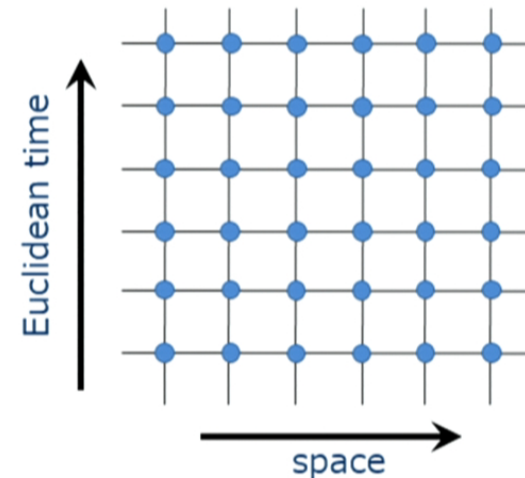
$$A_{ijkl} \longleftrightarrow \begin{array}{c} i \\ | \\ l - \textcircled{A} - j \\ | \\ k \end{array}$$

partition function of 2D  
classical statistical model



- tensors encode Boltzmann weights
- contraction of tensor network equals weighted sum over all microstates

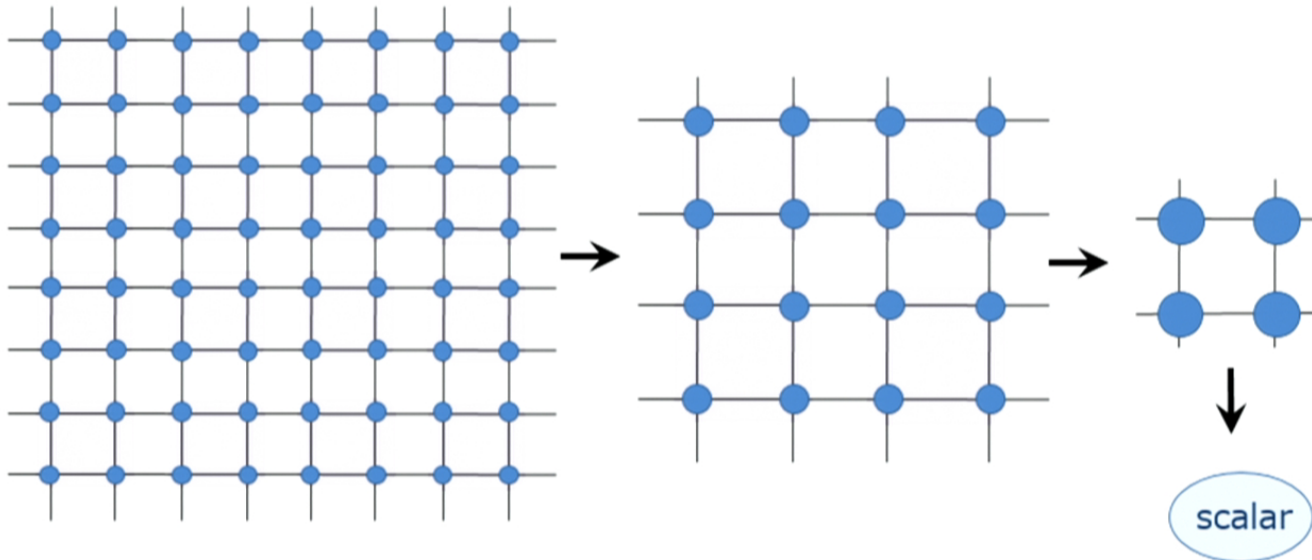
Euclidean path integral of  
1D quantum model



- row of tensors encodes small evolution in imaginary time
- contraction of tensor network equals weighted sum over all trajectories

# Overview

local scale transformations on the lattice:



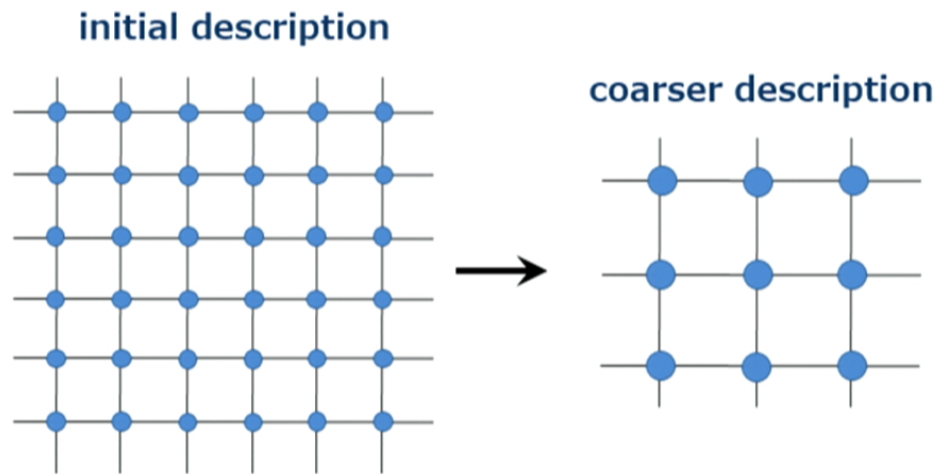
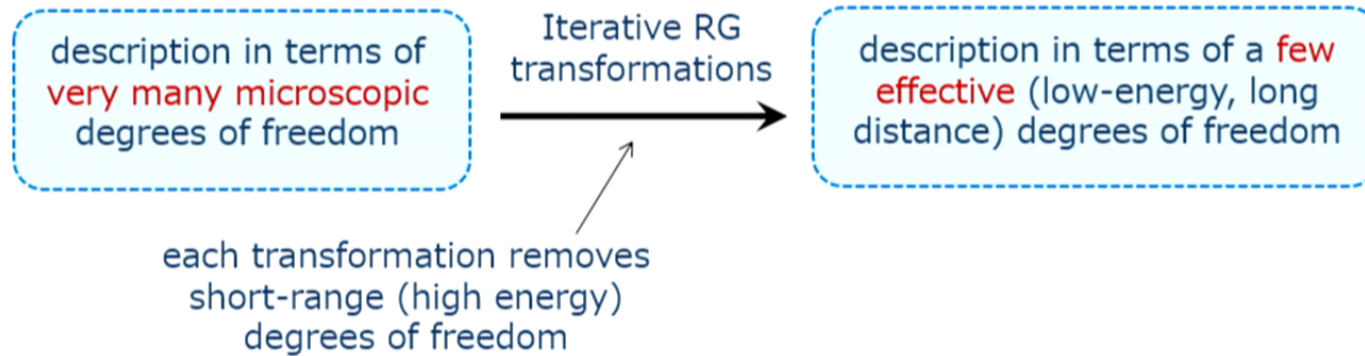
**Practical goal:** efficient and accurate contraction of a tensor network to a scalar

**Conceptual goal:** achieve a proper RG flow

i.e. could represent an expectation value in the quantum system  $\langle \psi | o | \psi \rangle$

# Overview: real-space RG

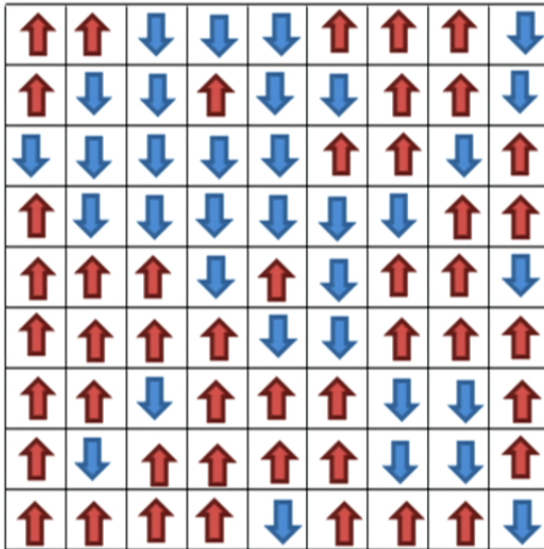
Basic idea of RG:



# Overview: real-space RG

Early real-space RG: Kadanoff's "spin blocking" (1966)

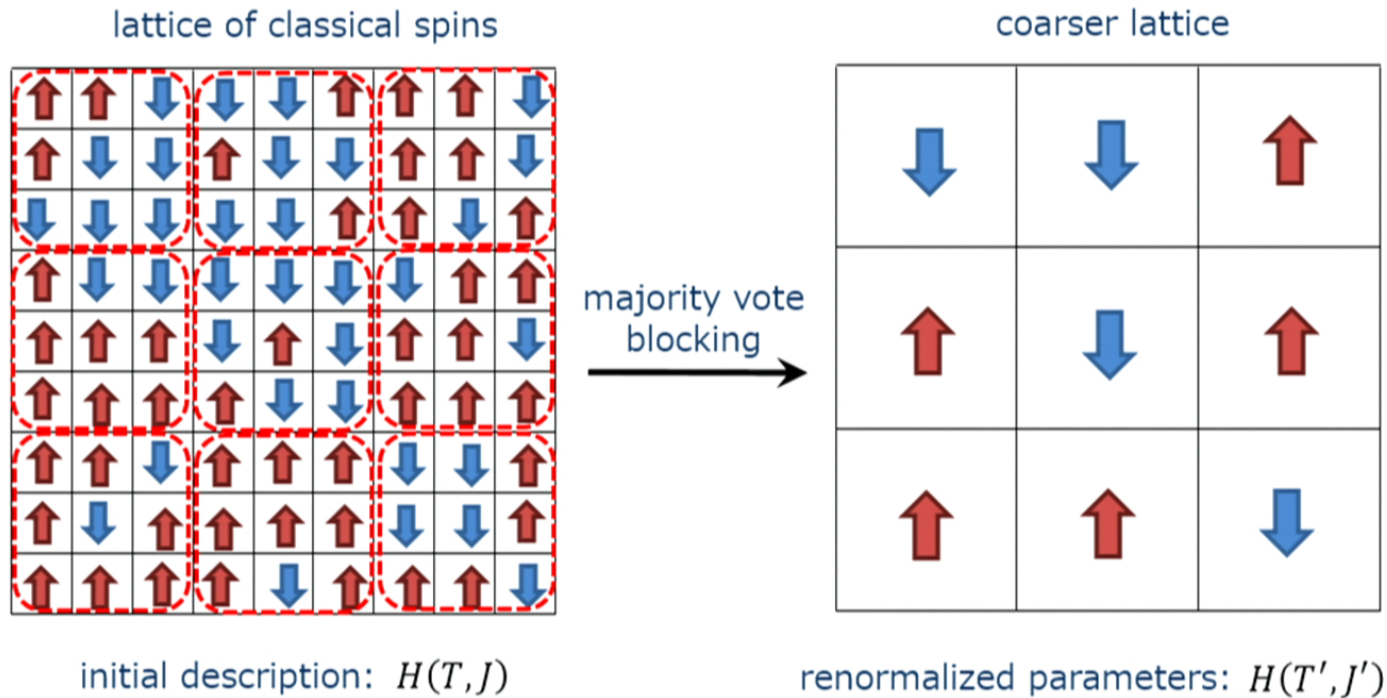
lattice of classical spins



initial description:  $H(T, J)$

# Overview: real-space RG

Early real-space RG: Kadanoff's "spin blocking" (1966)



...successful only for certain systems

# Overview: real-space RG

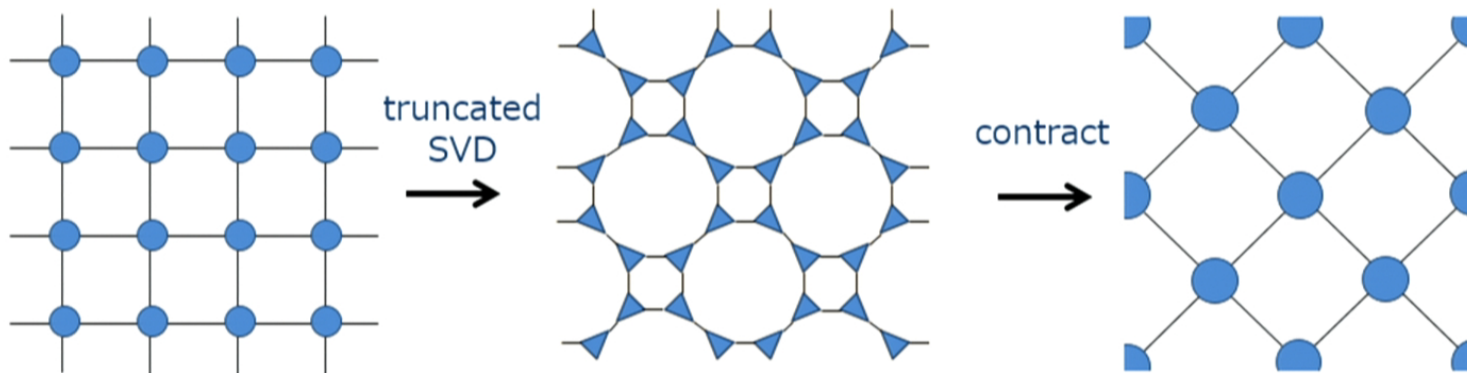
L.P. Kadanoff (1966): “Spin blocking”

spiritual  
successor



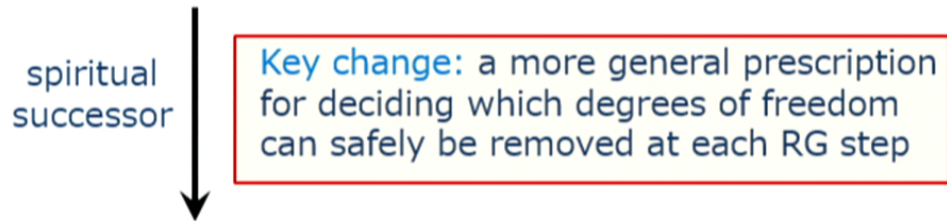
**Key change:** a more general prescription  
for deciding which degrees of freedom  
can safely be removed at each RG step

Levin, Nave (2006) : “Tensor renormalization group (LN-TRG)”



# Overview: real-space RG

L.P. Kadanoff (1966): “Spin blocking”



Levin, Nave (2006) : “Tensor renormalization group (LN-TRG)”

+ many improvements and generalizations:

Xie, Jiang, Weng, Xiang (2008): “Second Renormalization Group (SRG)”

Gu, Levin, Wen (2008): “Tensor Entanglement Renormalization Group (TERG)”

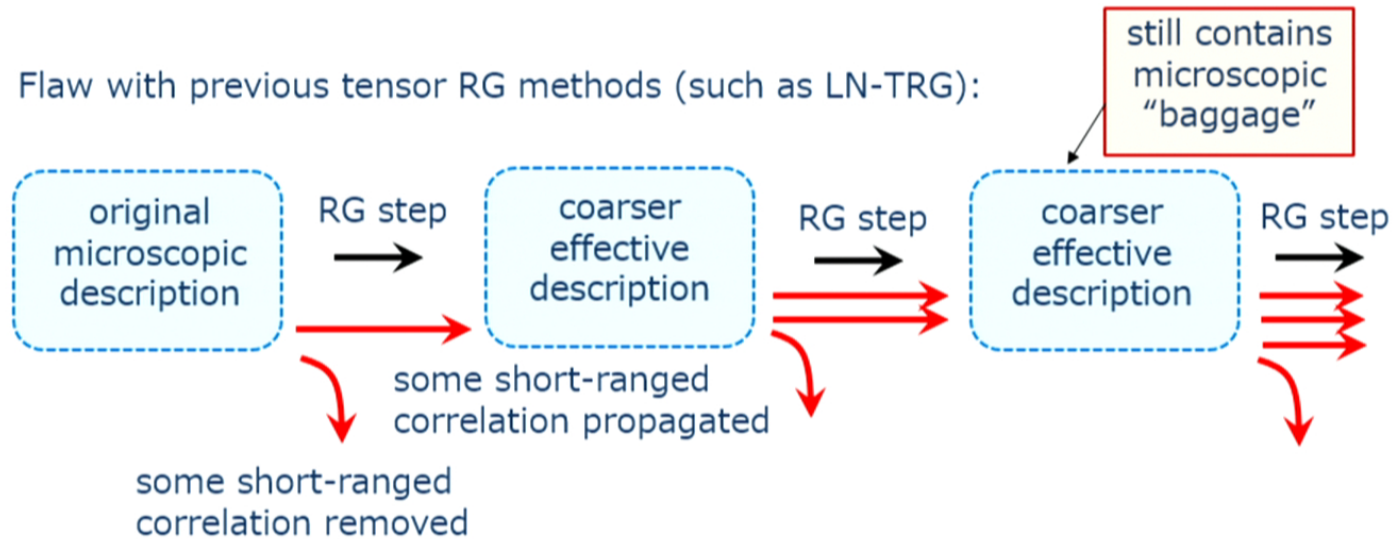
Gu, Wen (2009): “Tensor Entanglement Filtering Renormalization (TEFR)”

Xie, Chen, Qin, Zhu, Yang, Xiang (2012): “Higher Order Tensor Renormalization Group (HOTRG)”



# Overview

Flaw with previous tensor RG methods (such as LN-TRG):



**Flaw:** each RG step removes some (but not all) of the short-ranged degrees freedom

## Consequences:

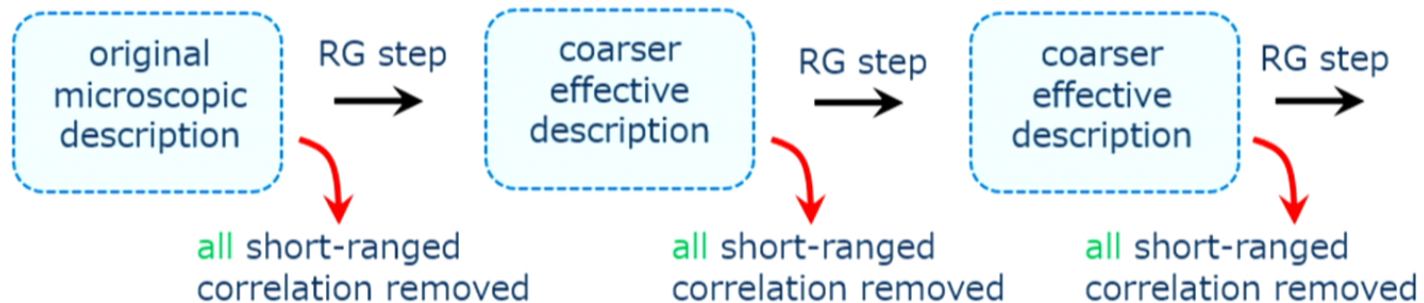
- Computational breakdown near (or at) a **critical point** **×**
- Effective theory still contains **unwanted microscopic detail**; one does not recover proper structure of RG fixed points **×**



# Overview

New approach: “Tensor Network Renormalization (TNR)” arXiv:1412.0732

A way of implementing real-space RG that addresses all short-ranged degrees of freedom at each RG step



## Advantages:

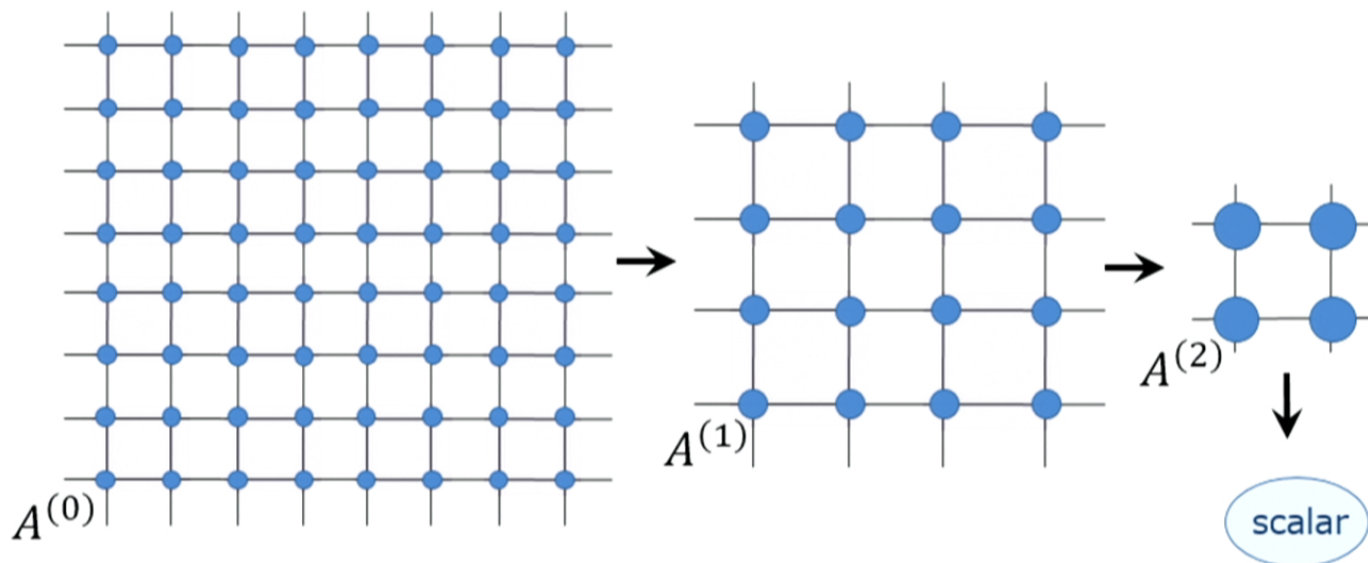
- **Proper RG flow is achieved**, TNR reproduces the correct structure of RG fixed points (including critical fixed points)
- Prevents harmful accumulation of short-ranged detail, allowing for a computationally sustainable RG flow



# Overview

local scale transformations on the lattice:

RG flow in the space of tensors:  $A^{(0)} \rightarrow A^{(1)} \rightarrow A^{(2)} \rightarrow \dots \rightarrow A^{(s)} \rightarrow \dots$



**Practical goal:** efficient and accurate contraction of a tensor network to a scalar

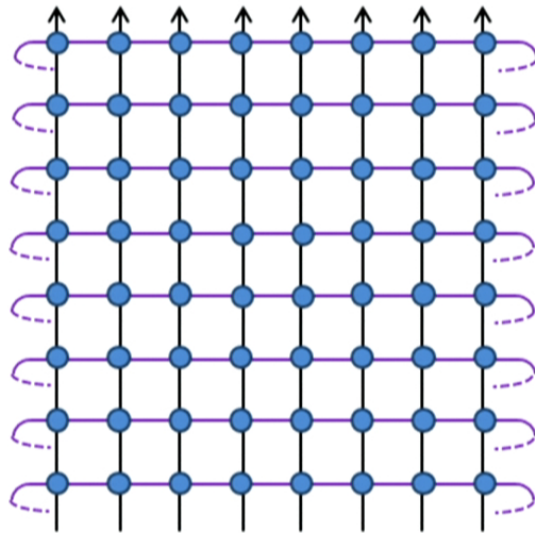
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i.e. could represent an expectation value in the quantum system  $\langle \psi | o | \psi \rangle$

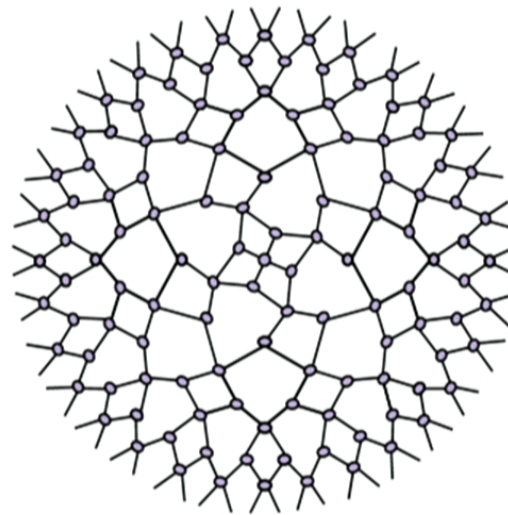
# Overview

Local scale transformation can be applied inhomogeneously to do interesting things...

**Euclidean path integral  
(1D quantum system, PBC)**



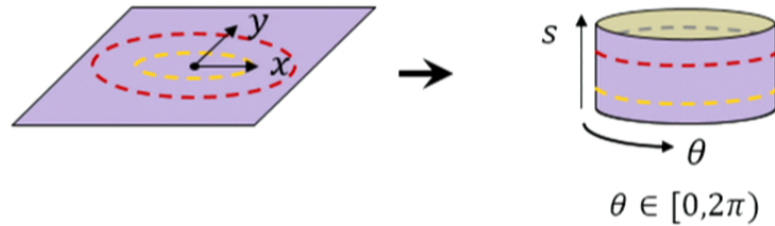
**Multi-scale entanglement  
renormalization ansatz (MERA)**



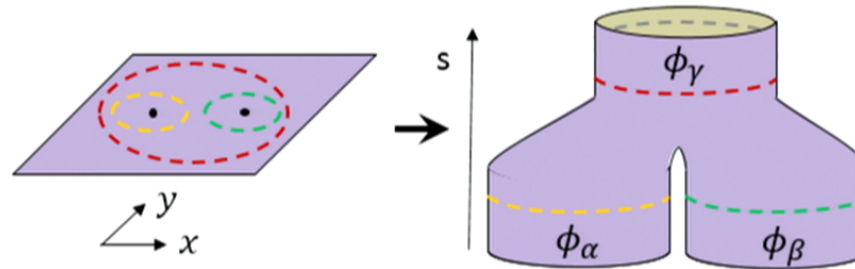
# Local scale transformations on the lattice with TNR

(G. E., D. Gaiotto, R. Myers, G. Vidal, *in prep.* )

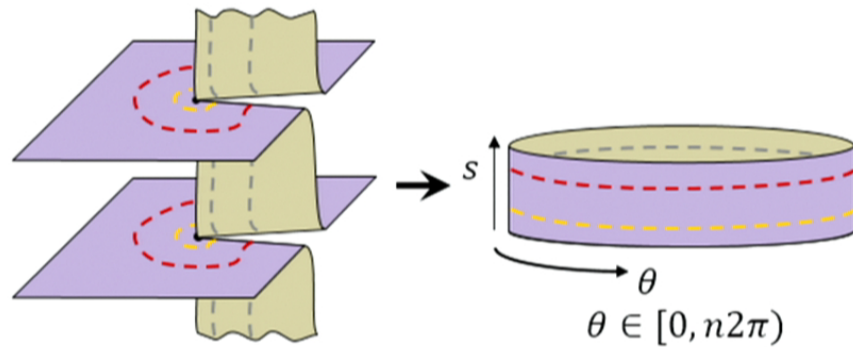
punctured plane to  
cylinder



plane with two  
punctures into  
pair of pants



$n$  - sheeted  
Riemann surface  
to cylinder

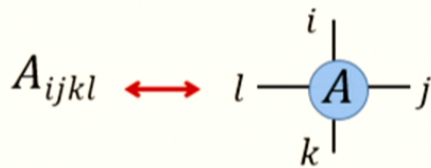


# Overview: Tensor Networks

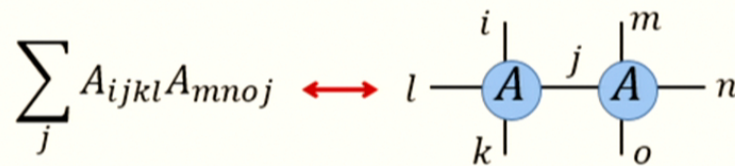
bond  
dimension  
↓

Let  $A_{ijkl}$  be a four index tensor with  $i, j, k, l \in \{1, 2, 3, \dots, \chi\}$   
i.e. such that the tensor is a  $\chi \times \chi \times \chi \times \chi$  array of numbers

Diagrammatic notation:



Contraction of two tensors:

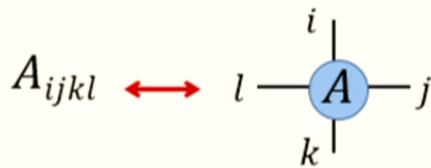


# Overview: Tensor Networks

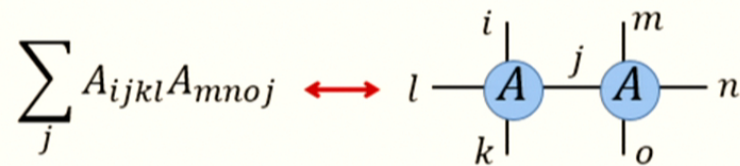
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bond  
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Diagrammatic notation:



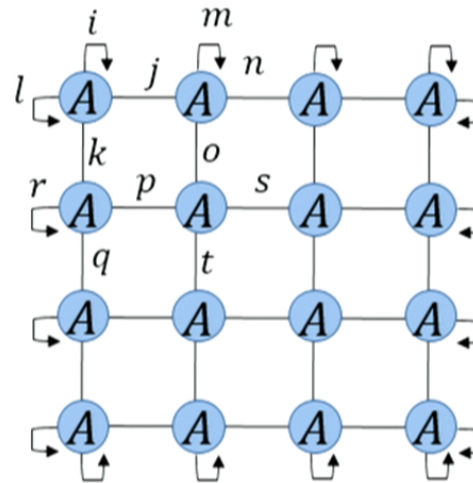
Contraction of two tensors:



Square lattice network (PBC):

$$\sum_{ijklmn\dots} A_{ijkl} A_{mno j} A_{kpqr} A_{ostp} \dots$$

$$\equiv \text{tTr} \left( \bigotimes_{x=1}^N A \right) = Z$$



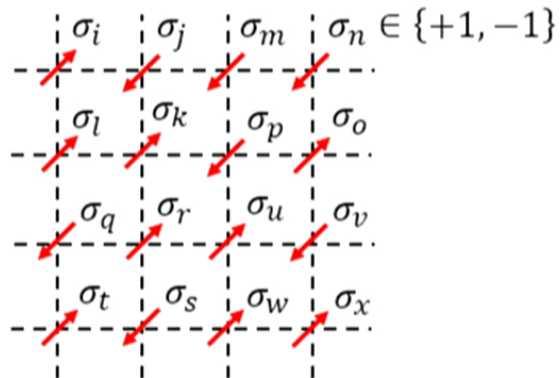
Contracts to a  
scalar:





# Partition functions as Tensor Networks

Square lattice of Ising spins:



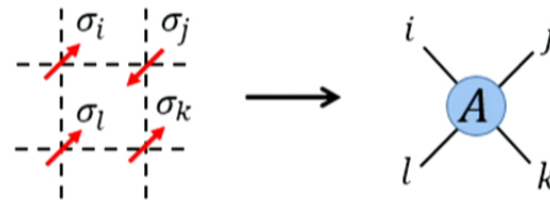
Hamiltonian functional for Ising ferromagnet:

$$H(\{\sigma\}) = - \sum_{\langle i,j \rangle} \sigma_i \sigma_j$$

Partition function:

$$Z = \sum_{\{\sigma\}} e^{-H(\{\sigma\})/T}$$

Encode the Boltzmann weights of a plaquette of spins in a four-index tensor

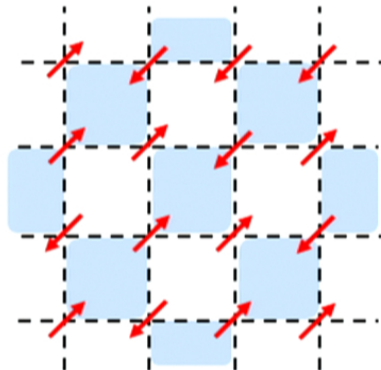


where:

$$A_{ijkl} = e^{(\sigma_i \sigma_j + \sigma_j \sigma_k + \sigma_k \sigma_l + \sigma_l \sigma_i)/T}$$

# Partition functions as Tensor Networks

Square lattice of Ising spins:

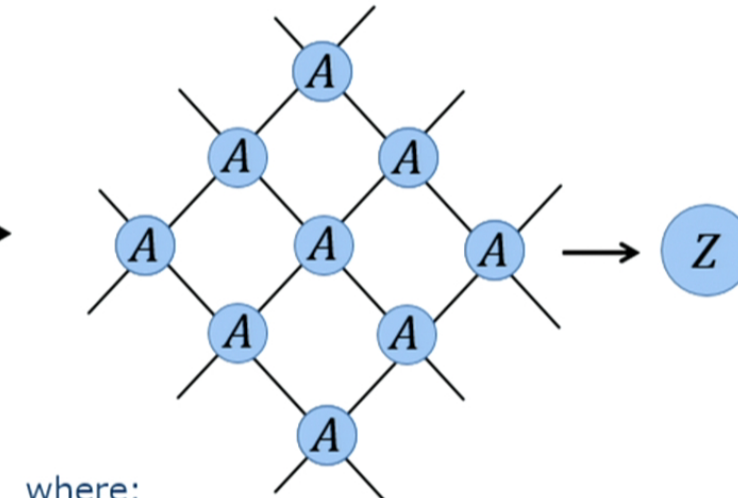


Hamiltonian functional for Ising ferromagnet:

$$H(\{\sigma\}) = - \sum_{\langle i,j \rangle} \sigma_i \sigma_j$$

Partition function:

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where:

$$A_{ijkl} = e^{(\sigma_i \sigma_j + \sigma_j \sigma_k + \sigma_k \sigma_l + \sigma_l \sigma_i)/T}$$

← Partition function given by contraction of tensor network



# Path Integrals as Tensor Networks

Nearest neighbour Hamiltonian for a 1D quantum system:

$$\begin{aligned} H = \sum_r h(r, r+1) &= \sum_{r \text{ even}} h(r, r+1) + \sum_{r \text{ odd}} h(r, r+1) \\ &= H_{\text{even}} + H_{\text{odd}} \end{aligned}$$

We want to express  $e^{-\beta H}$  as a tensor network:

$$\text{recall: } |\psi_{\text{GS}}\rangle\langle\psi_{\text{GS}}| = \lim_{\beta \rightarrow \infty} [e^{-\beta H}]$$



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$$\text{recall: } |\psi_{\text{GS}}\rangle\langle\psi_{\text{GS}}| = \lim_{\beta \rightarrow \infty} [e^{-\beta H}]$$

Expand in small time steps:

$$e^{-\beta H} = e^{-\tau H} e^{-\tau H} e^{-\tau H} e^{-\tau H} \dots$$

Suzuki-Trotter expansion:

$$e^{-\tau H} = e^{-\tau H_{\text{even}}} e^{-\tau H_{\text{odd}}} + o(\tau^2)$$

# Path Integrals as Tensor Networks

Separate Hamiltonian into even and odd terms:

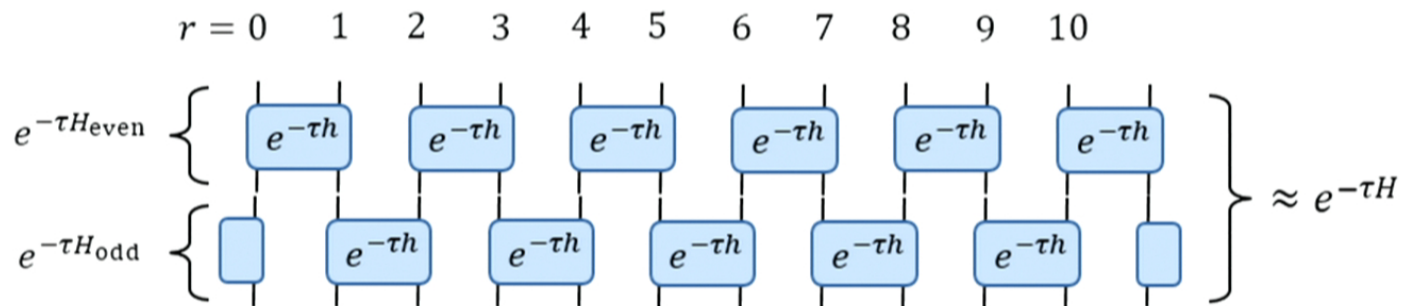
$$H = \sum_{r \text{ even}} h(r, r+1) + \sum_{r \text{ odd}} h(r, r+1) = H_{\text{even}} + H_{\text{odd}}$$

Expand path integral in small discrete time steps:

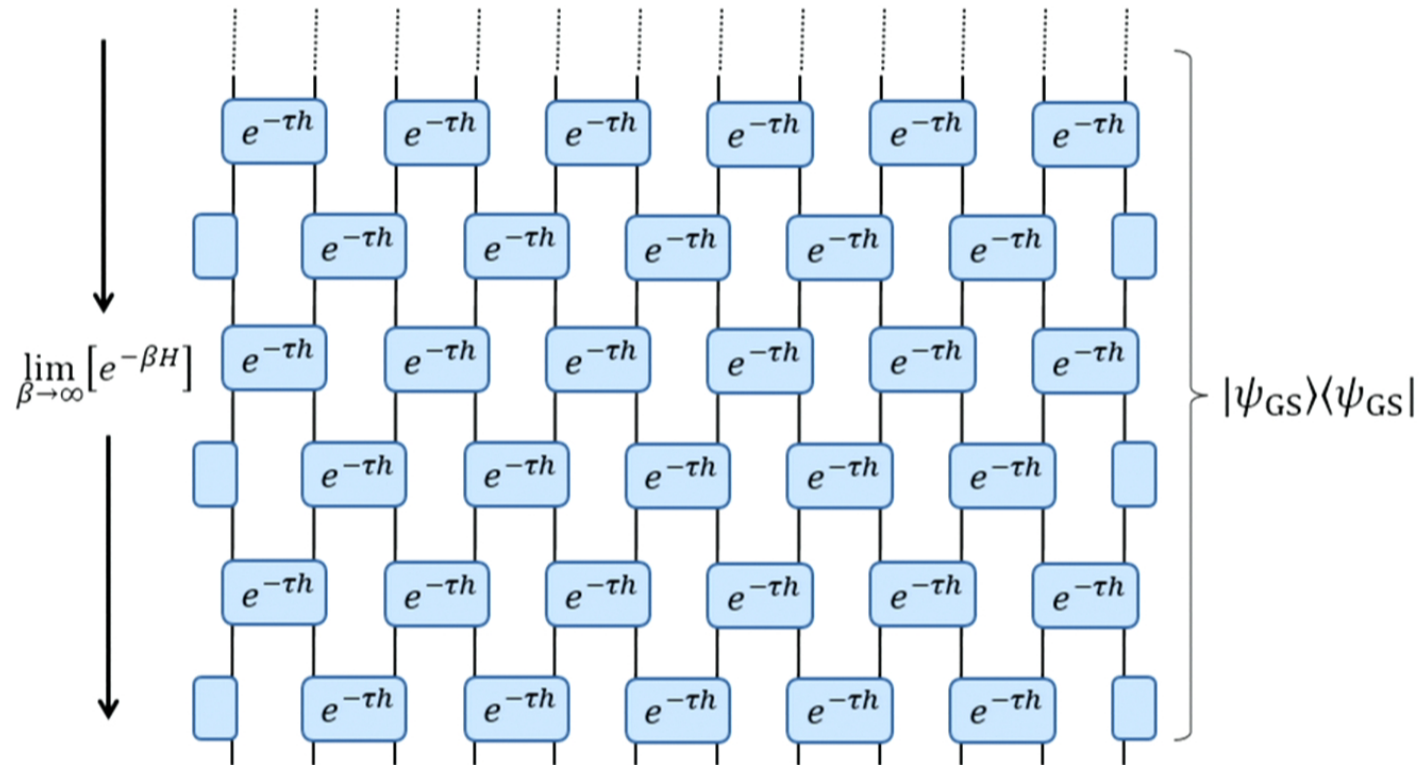
$$e^{-\beta H} = e^{-\tau H} e^{-\tau H} e^{-\tau H} e^{-\tau H} \dots$$

$$e^{-\tau H} = e^{-\tau H_{\text{even}}} e^{-\tau H_{\text{odd}}} + o(\tau^2)$$

Exponentiate even and odd separately :

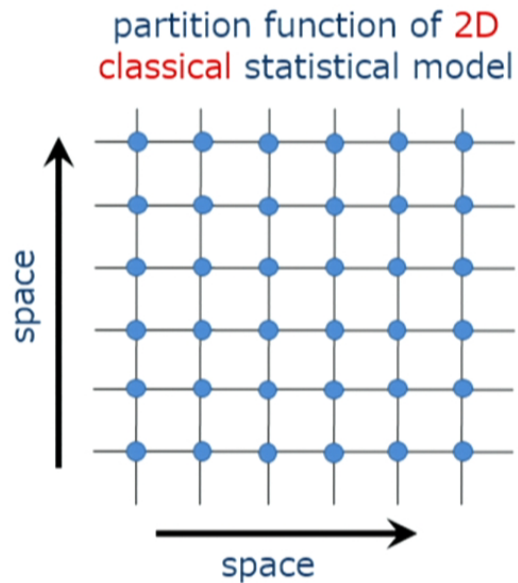


# Path Integrals as Tensor Networks

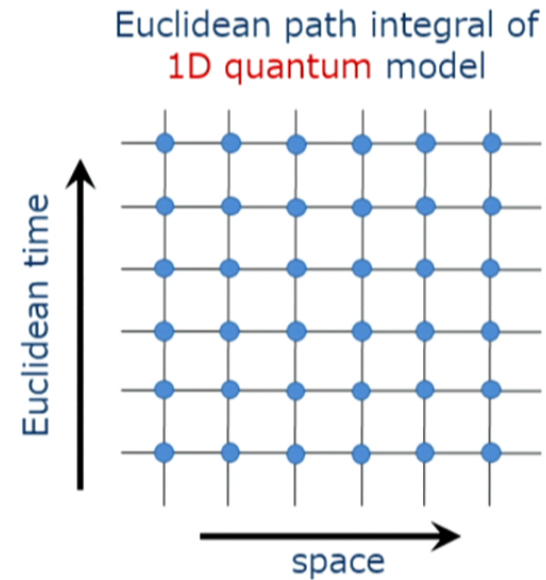


# Overview

encode many-body systems as a **tensor network**:



- tensors encode Boltzmann weights
- contraction of tensor network equals weighted sum over all microstates

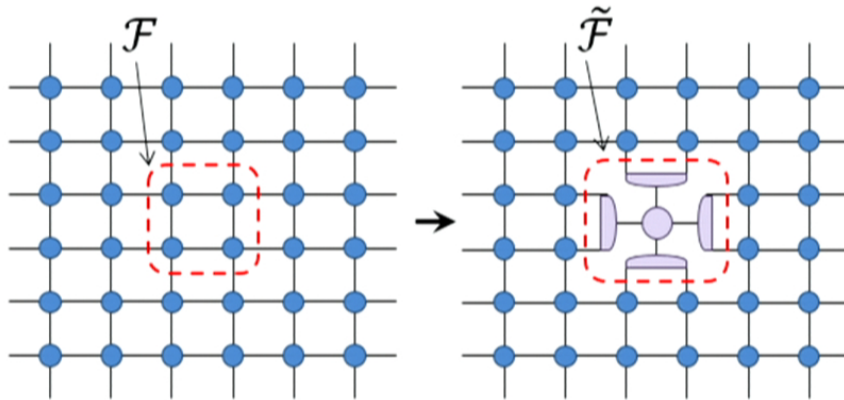


- row of tensors encodes small evolution in imaginary time
- contraction of tensor network equals weighted sum over all trajectories

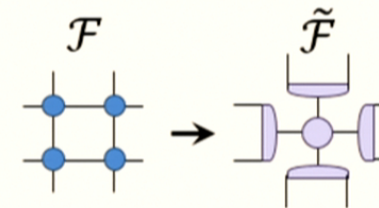


# Local Approximations:

want to replace a local subnetwork of tensors  
(e.g. a 2 x 2 block) with something else...



Rule: any local replacement  
is potentially valid...



...so long as the difference  
is small:

$$\delta = \|\mathcal{F} - \tilde{\mathcal{F}}\|$$

Replacement error:

$$\delta^2 = \|\mathcal{F} - \tilde{\mathcal{F}}\|^2 = \|\mathcal{F}\|^2 - \text{tTr}(\mathcal{F} \otimes \tilde{\mathcal{F}}^\dagger) - \text{tTr}(\tilde{\mathcal{F}} \otimes \mathcal{F}^\dagger) + \|\tilde{\mathcal{F}}\|^2$$

$$\left\| \begin{array}{c} \text{2x2 grid of blue dots} \\ - \\ \text{2x2 grid of purple and blue shapes} \end{array} \right\|^2 = \left\{ \begin{array}{c} \text{2x2 grid of blue dots} \\ - \\ \text{2x2 grid of blue dots with purple shapes} \\ - \\ \text{2x2 grid of blue dots with purple shapes} \\ + \\ \text{2x2 grid of purple and blue shapes} \end{array} \right\}$$

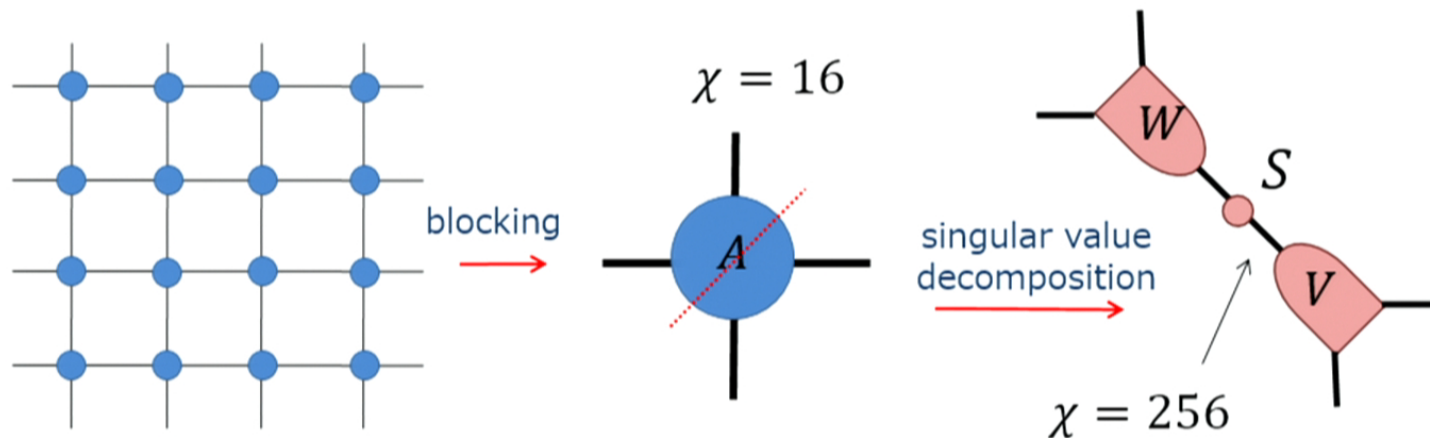
# Tensor Renormalization Group (LN-TRG)

Levin, Nave (2006)

**Tensor renormalization group (LN-TRG)** is a method for coarse-graining tensor networks based upon **blocking** and **truncation steps**

**Example of blocking + truncation:** 2D classical Ising (critical temp)

- take a  $(4 \times 4)$  block of tensors from the partition function
- contract to a single tensor; each (16-dim) index describes the state of four classical spins
- can the block tensor be truncated?



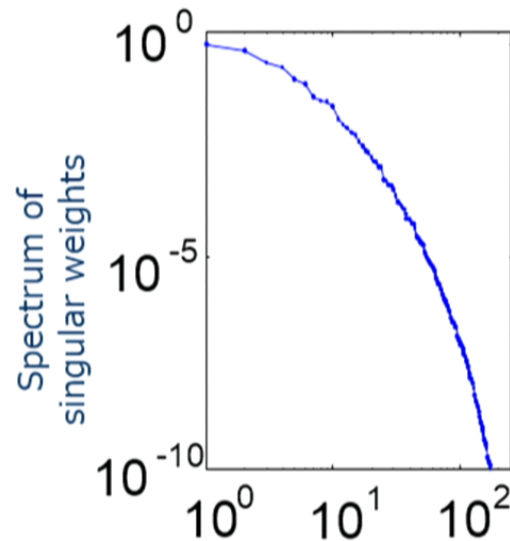
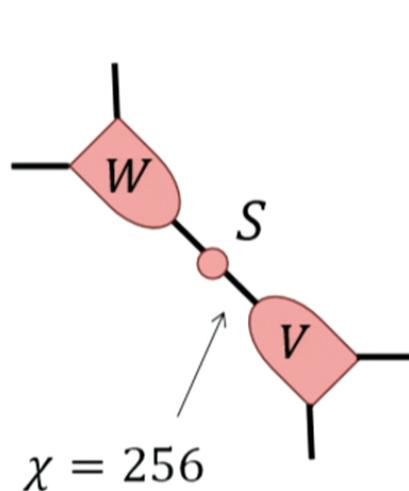
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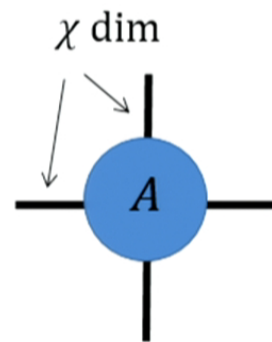
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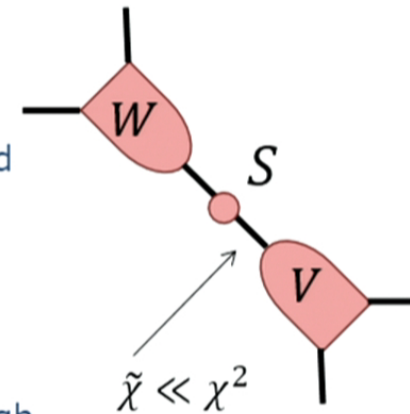


# Tensor Renormalization Group (LN-TRG)

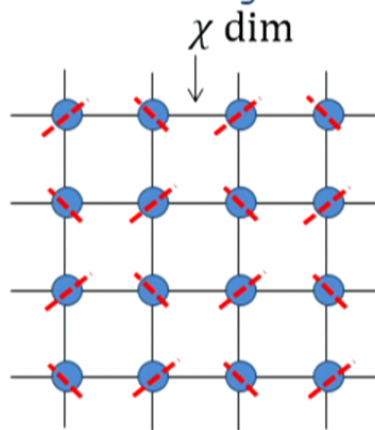
discard singular values  
smaller than desired  
truncation error  $\delta$



truncated  
SVD

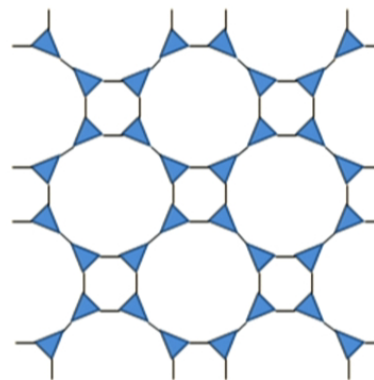


**Tensor Renormalization Group (LN-TRG)** works through  
alternating truncated SVD and contraction steps:

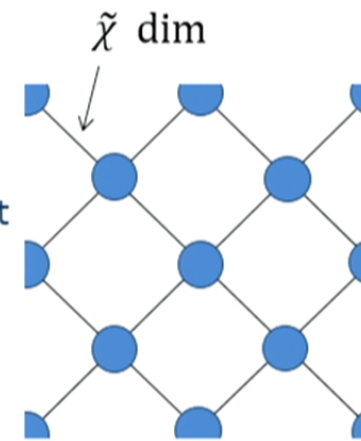


initial network

truncated  
SVD



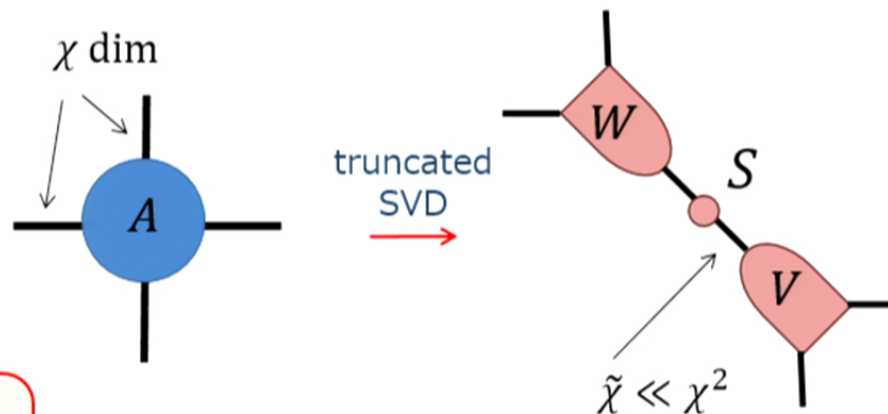
contract



coarser network

# Tensor Renormalization Group (LN-TRG)

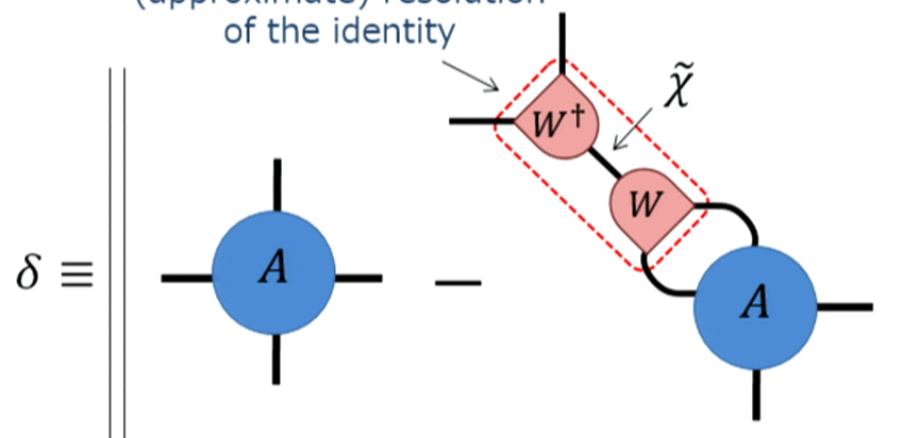
discard singular values  
smaller than desired  
truncation error  $\delta$



**equivalent approach:**  
implement truncation through  
projector of the form  $W^\dagger W$  for  
isometric  $W$

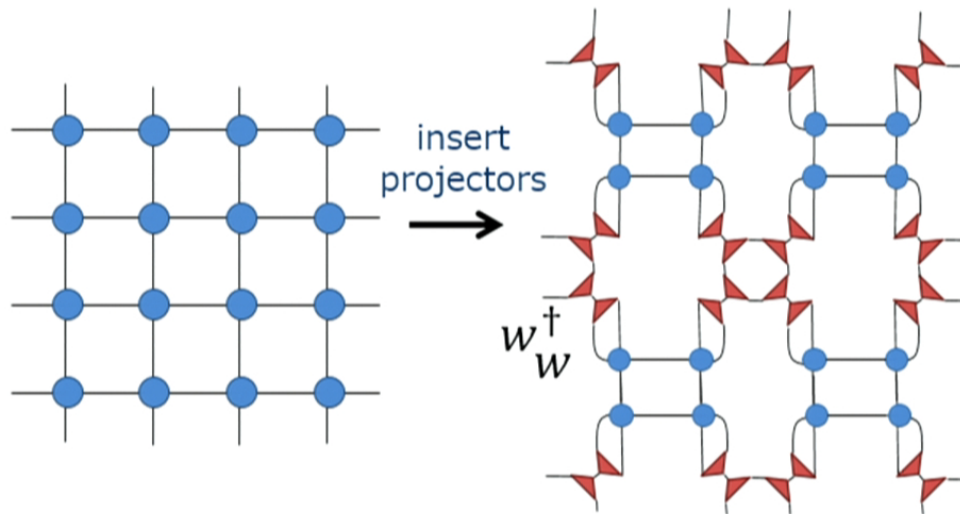
i.e. choose isometry  $W$   
to minimise truncation  
error  $\delta$

projector  $P$  acts as a  
(approximate) resolution  
of the identity



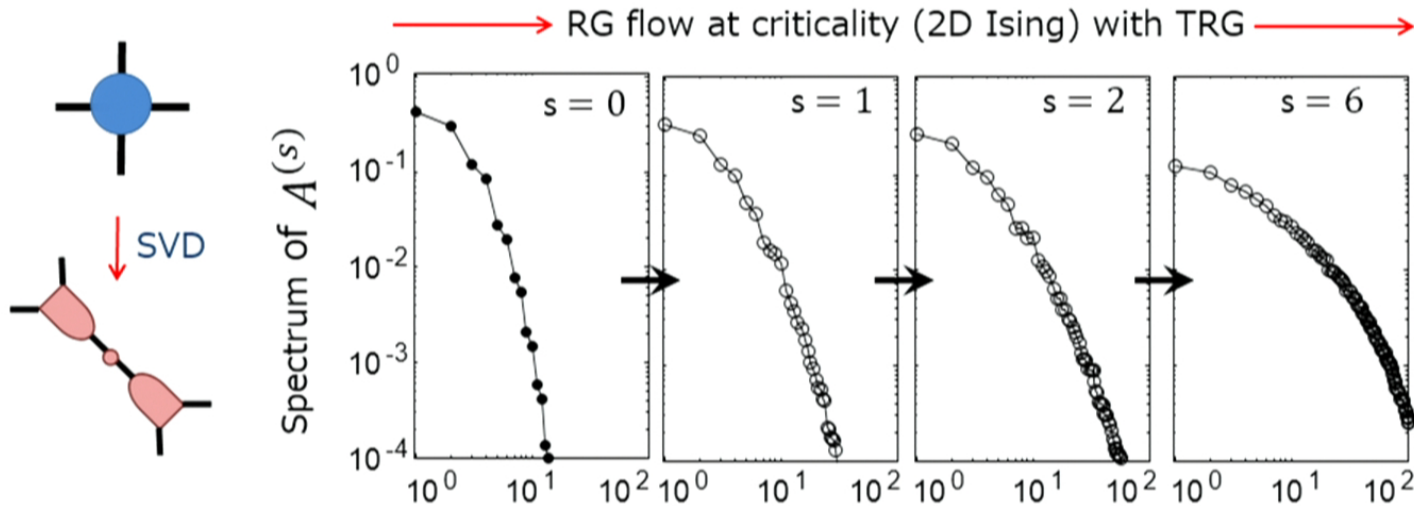
# Tensor Renormalization Group (LN-TRG)

Levin, Nave (2006)



# Tensor Renormalization Group (LN-TRG)

RG flow in the space of tensors:  $A^{(0)} \rightarrow A^{(1)} \rightarrow A^{(2)} \rightarrow \dots \rightarrow A^{(s)} \rightarrow \dots$

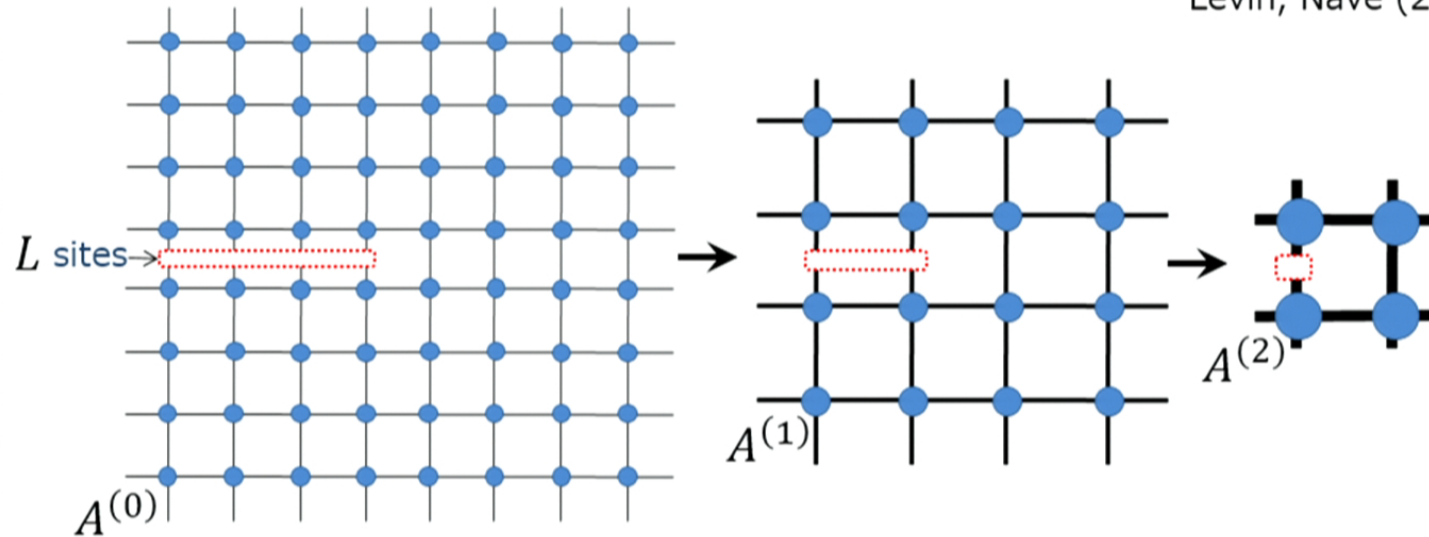


Bond dimension  $\chi$  required for truncation error  $< 10^{-3}$ :  $\sim 10 \rightarrow \sim 20 \rightarrow \sim 40 \rightarrow > 100$

Spectra become increasingly flat with each RG step! ❌

# Tensor Renormalization Group (LN-TRG)

Levin, Nave (2006)



Is LN-TRG a good RG transformation for critical systems?

- need to grow bond dimension  $\chi$  with each RG step
- does not converge to scale-invariant fixed point

Can be understood logarithmic scaling of entanglement entropy in 1D quantum systems:  $S_L = \frac{c}{3} \log(L) + k$

# Tensor Renormalization Group (LN-TRG)

Levin, Nave (2006)

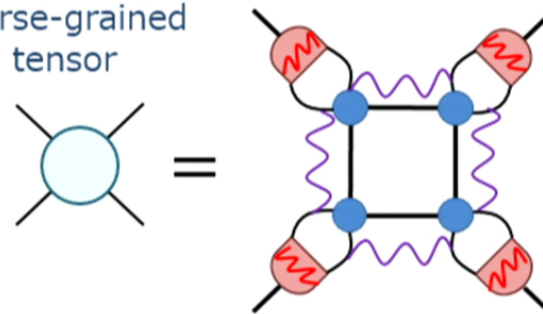
basic step of LN-TRG  
(projective truncation):

$$\delta = \left\| \begin{array}{c} \text{2x2 grid of blue tensors} \end{array} - \begin{array}{c} \text{2x2 grid of blue tensors with red corner tensors } W \text{ and } W^\dagger \end{array} \right\|$$

- isometries remove some (but not all!) short-ranged correlated degrees of freedom
- **LN-TRG fails to remove** some short-ranged correlations, which propagate to next length scale

**Example: corner-double line (CDL) tensors**

coarse-grained tensor



# Tensor Renormalization Group (LN-TRG)

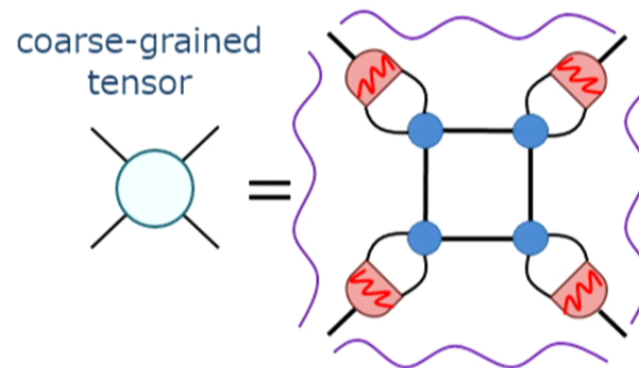
Levin, Nave (2006)

basic step of LN-TRG  
(projective truncation):

$$\delta = \left\| \begin{array}{c} \text{2x2 grid of blue tensors} \end{array} - \begin{array}{c} \text{2x2 grid of blue tensors with red tensors } W \text{ and } W^\dagger \text{ on the edges} \end{array} \right\|$$

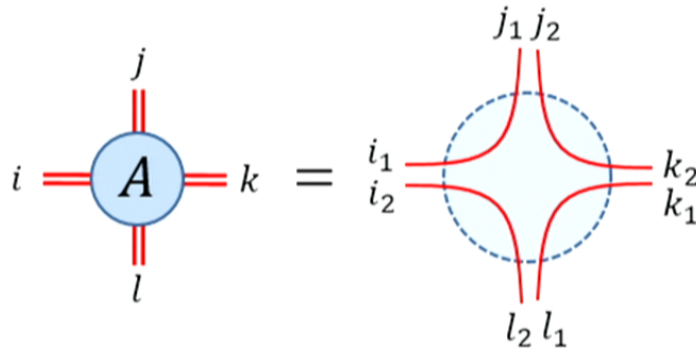
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- **LN-TRG fails to remove** some short-ranged correlations, which propagate to next length scale

**Example: corner-double line (CDL) tensors**





# Fixed points of LN-TRG



Imagine “A” is a special tensor such that each index can be decomposed as a product of smaller indices,

$$A_{ijkl} = A_{(i_1 i_2)(j_1 j_2)(k_1 k_2)(l_1 l_2)}$$

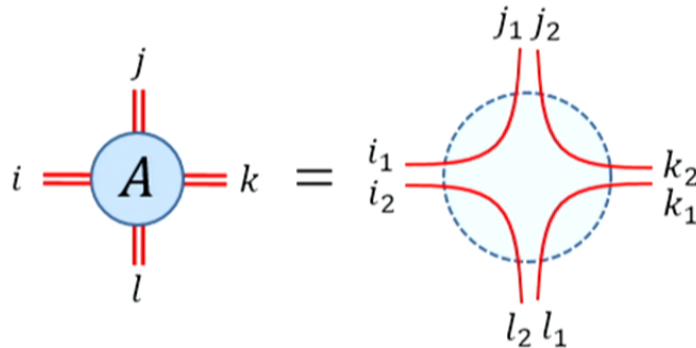
such that certain pairs of indices are perfectly correlated:

$$A_{(i_1 i_2)(j_1 j_2)(k_1 k_2)(l_1 l_2)} \equiv \delta_{i_1 j_1} \delta_{j_2 k_2} \delta_{k_1 l_1} \delta_{l_2 i_2}$$

These are called **corner double line** (CDL) tensors. CDL tensors are fixed points of TRG.



# Fixed points of LN-TRG



Imagine “A” is a special tensor such that each index can be decomposed as a product of smaller indices,

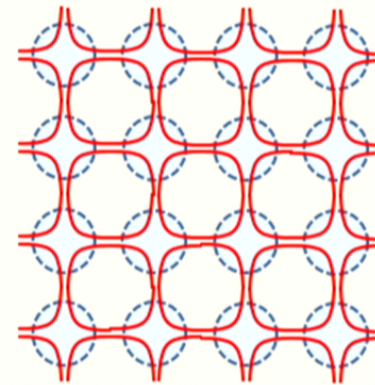
$$A_{ijkl} = A_{(i_1 i_2)(j_1 j_2)(k_1 k_2)(l_1 l_2)}$$

such that certain pairs of indices are perfectly correlated:

$$A_{(i_1 i_2)(j_1 j_2)(k_1 k_2)(l_1 l_2)} \equiv \delta_{i_1 j_1} \delta_{j_2 k_2} \delta_{k_1 l_1} \delta_{l_2 i_2}$$

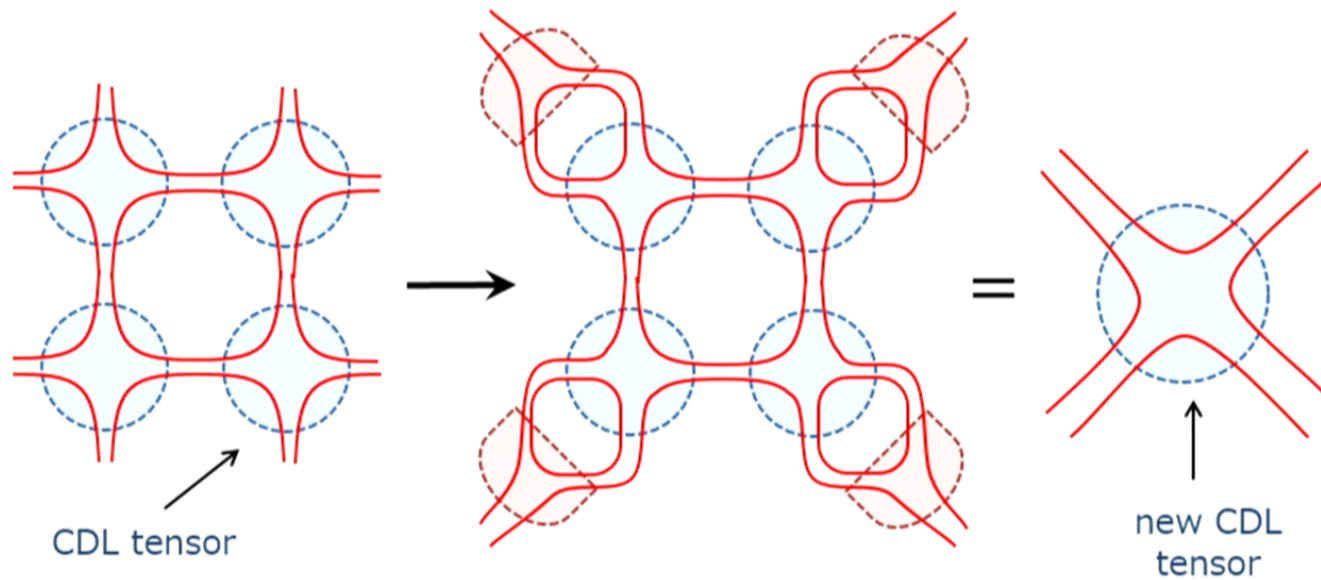
These are called **corner double line** (CDL) tensors. CDL tensors are fixed points of TRG.

Partition function built from CDL tensors represents a state with short-ranged correlations



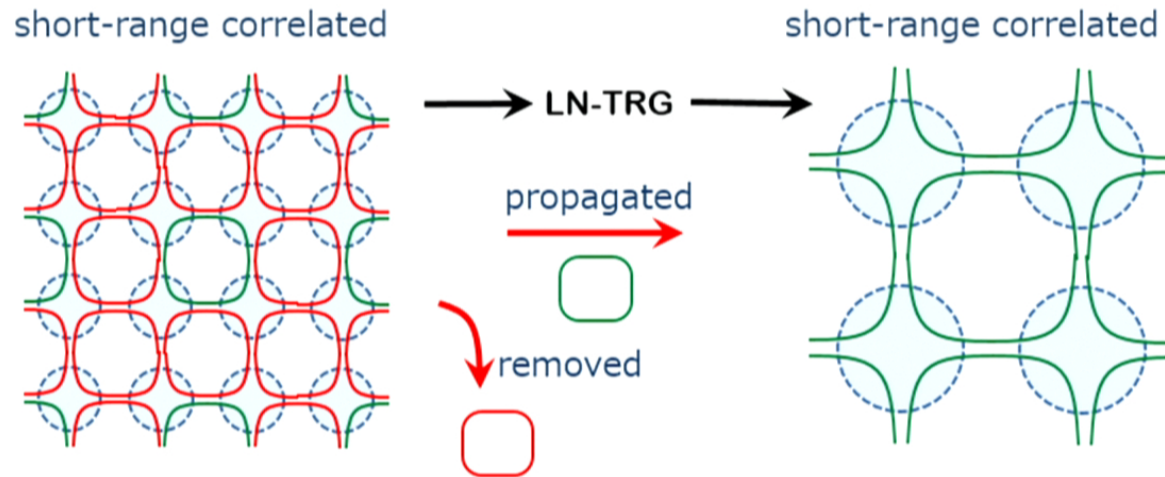
# Fixed points of LN-TRG

single iteration of LN-TRG:



Some short-ranged always  
correlations remain under LN-TRG!

# Fixed points of LN-TRG



TRG removes some short ranged correlations, but...  
others are **artificially promoted** to the next length scale

**Accumulation** of short-ranged details causes computational breakdown when near (or at) criticality

Is there some way to 'fix' tensor renormalization such that **all short-ranged** correlations are addressed?

# Outline: Tensor Network Renormalization

**The set-up:** Representation of partition functions and path integrals as tensor networks

**Previous approaches:** Levin and Nave's Tensor Renormalization Group (LN-TRG), conceptual and computation problems.

**New approach:** Tensor network renormalization (TNR): proper removal of all short-ranged degrees of freedom via disentanglers

**Benchmark results**

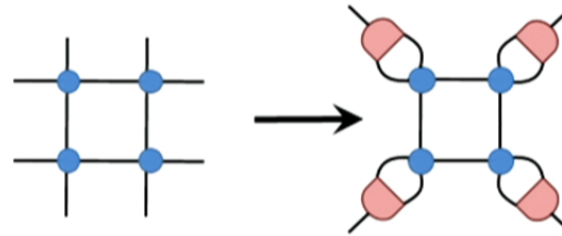
**Extensions**

# Tensor Network Renormalization

arXiv:1412.0732

previous RG schemes for tensor networks based upon **blocking**:

i.e. isometries responsible for combining and truncating indices



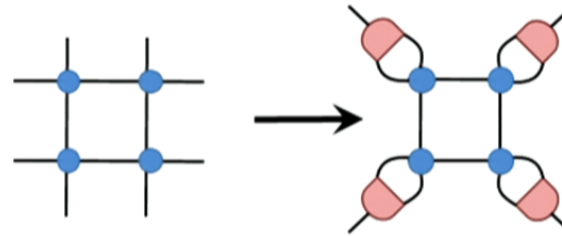
but blocking alone fails to remove short-ranged degrees of freedom...  
...can one incorporate some form of **unitary disentangling** into a tensor RG scheme?

# Tensor Network Renormalization

arXiv:1412.0732

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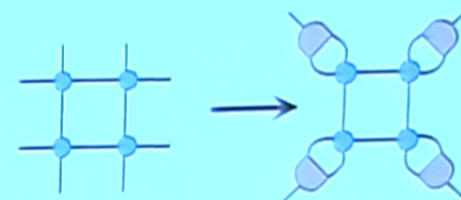


# Tensor Network Renormalization

arXiv:1412.0732

previous RG schemes for tensor networks based upon **blocking**:

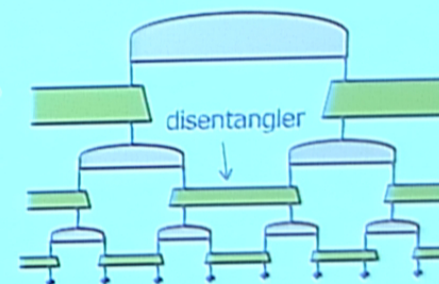
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but blocking alone fails to remove short-ranged degrees of freedom...  
...can one incorporate some form of unitary disentangling into a tensor RG scheme?

Tree tensor network (TTN)

Multi-scale entanglement renormalization ansatz (MERA)

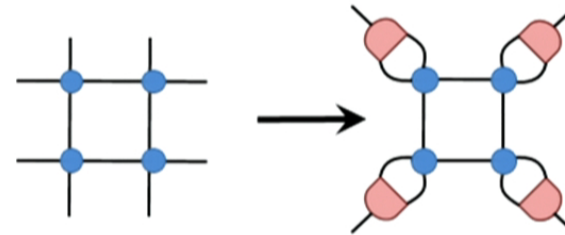


# Tensor Network Renormalization

arXiv:1412.0732

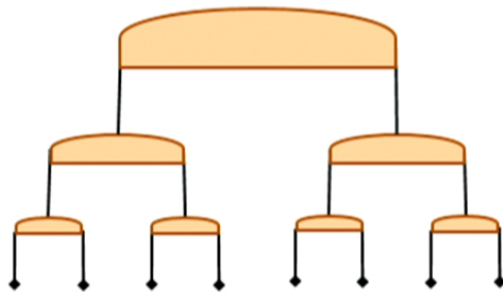
previous RG schemes for tensor networks based upon **blocking**:

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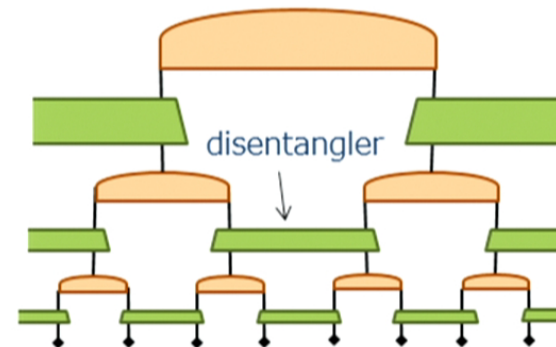


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Tree tensor network (TTN)



Multi-scale entanglement renormalization ansatz (MERA)



# Tensor Network Renormalization

previous approach

basic step of LN-TRG  
(projective truncation):

$$\delta = \left\| \begin{array}{c} \text{2x2 grid of blue nodes} \end{array} - \begin{array}{c} \text{2x2 grid of blue nodes with red nodes } W, W^\dagger \text{ on edges} \end{array} \right\|$$

new approach

basic step of Tensor  
Network Renormalization  
(projective truncation):

$$WW^\dagger = I \quad UU^\dagger = I^{\otimes 2}$$

$$\begin{array}{c} \text{red node} \end{array} = \begin{array}{c} | \end{array} \quad \begin{array}{c} \text{green node} \end{array} = \begin{array}{c} | \end{array}$$

$$\delta = \left\| \begin{array}{c} \text{2x2 grid of blue nodes} \end{array} - \begin{array}{c} \text{2x2 grid of blue nodes with green nodes } U, U^\dagger \text{ on edges} \end{array} \right\|$$

disentangler  $U^\dagger$

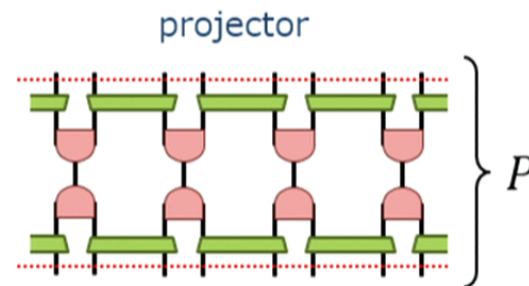
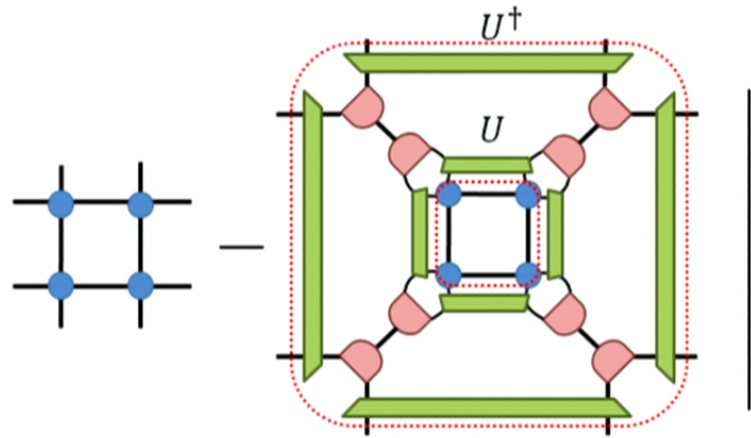
# Tensor Network Renormalization

new approach

basic step of Tensor Network Renormalization  
(projective truncation):

use iterative methods for  
optimizing isometry  $W$   
and disentangler  $U$

$$\delta =$$





# Tensor Network Renormalization

previous approach

basic step of LN-TRG  
(projective truncation):

$$\delta = \left\| \begin{array}{c} \text{2x2 grid of blue nodes} \end{array} - \begin{array}{c} \text{2x2 grid of blue nodes with red nodes } W, W^\dagger \text{ on edges} \end{array} \right\|$$

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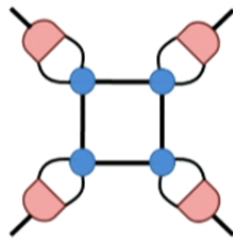
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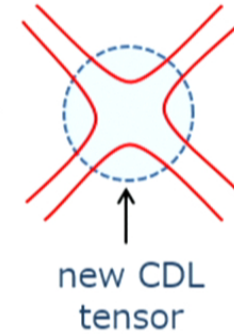
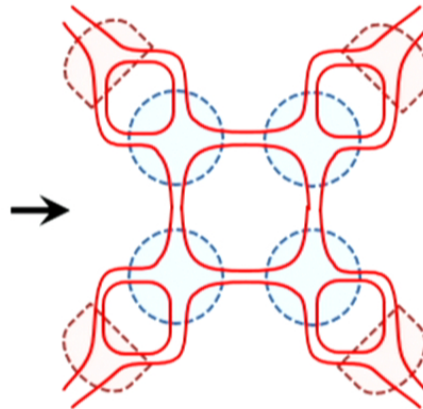
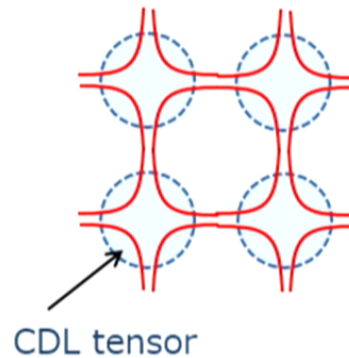
disentangler  
 $\downarrow U^\dagger$

# Corner double line tensors revisited

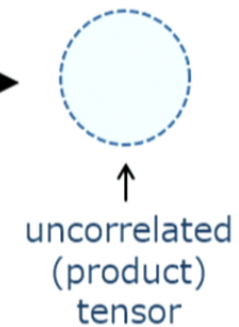
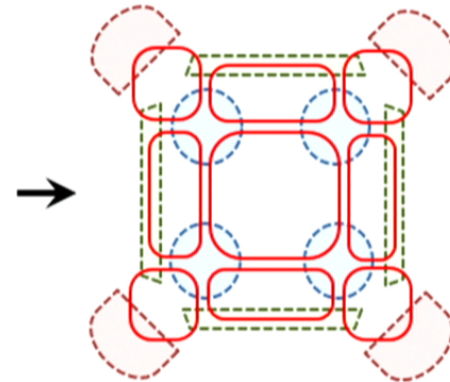
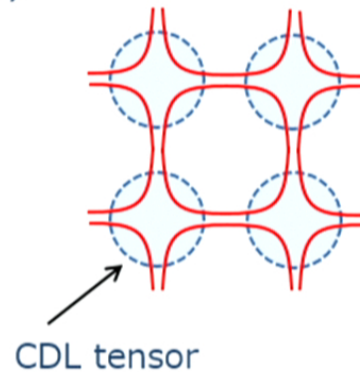
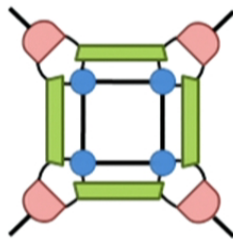
Isometries only  
(LN-TRG)



example: corner double  
line (CDL) tensors

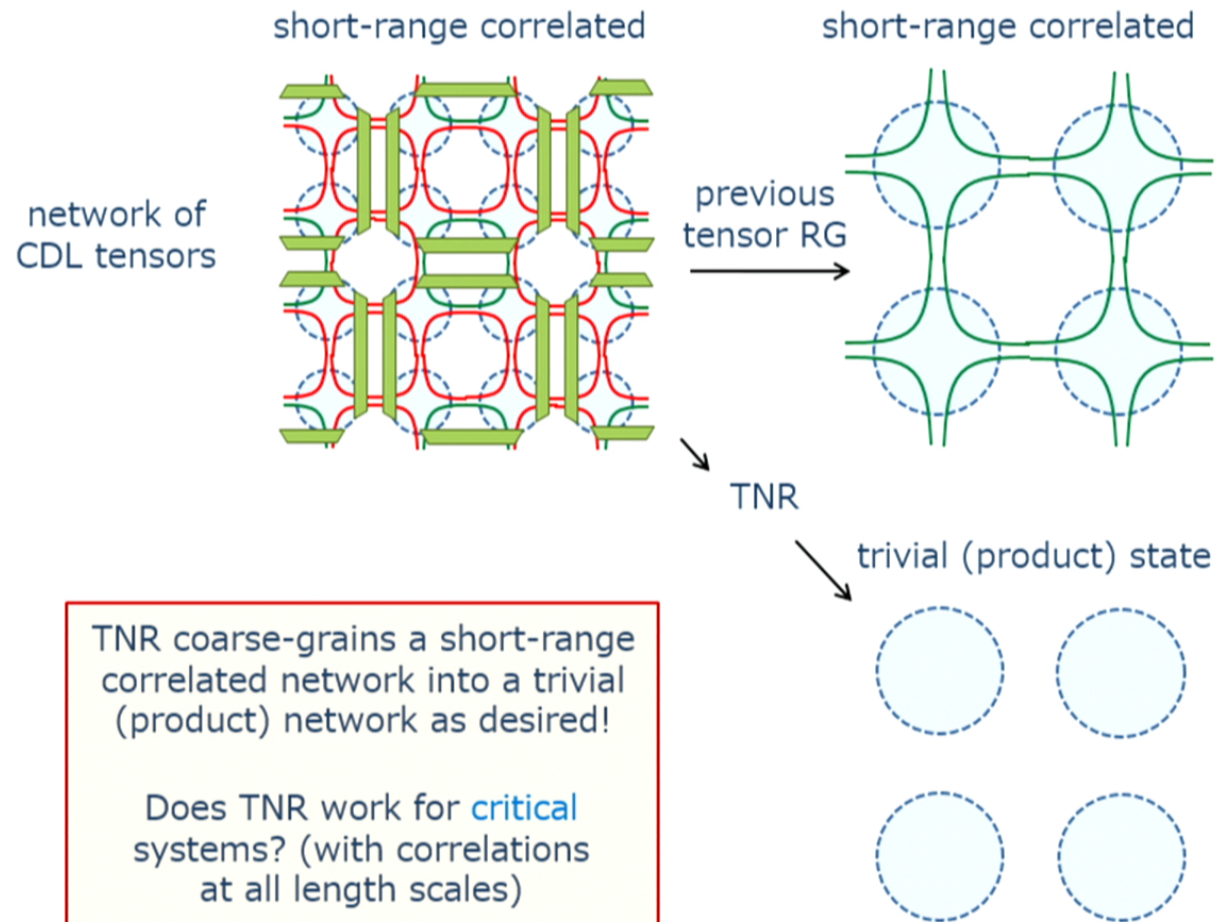


Isometries and  
disentangler (TNR)



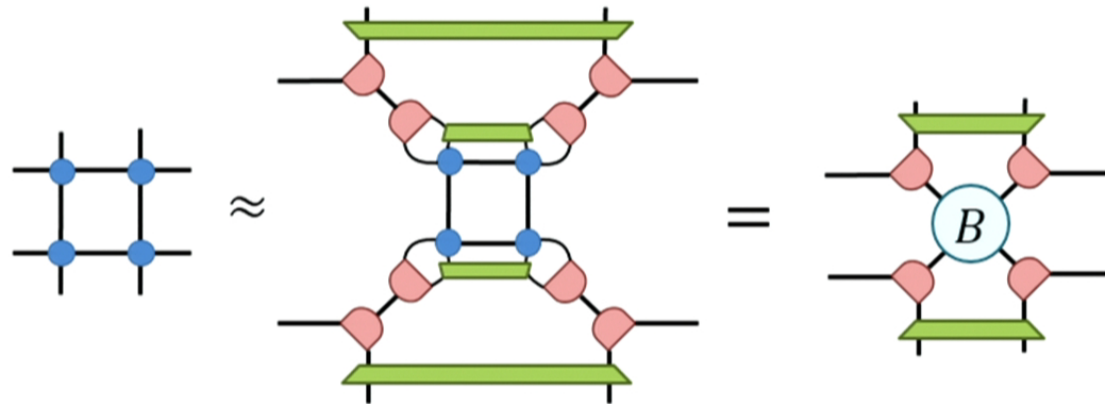


# Corner double line tensors revisited



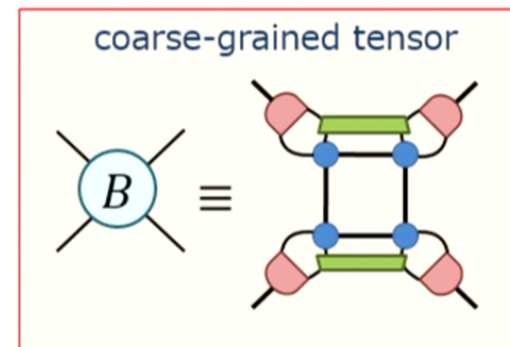
# Tensor Network Renormalization

key coarse-graining step of TNR:

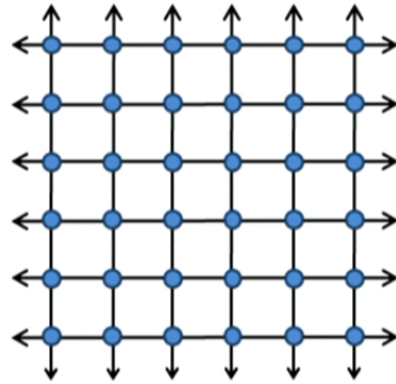


**Simplification:** disentangling is expensive... are all disentanglers necessary? No!

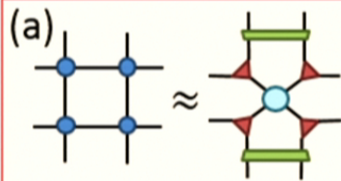
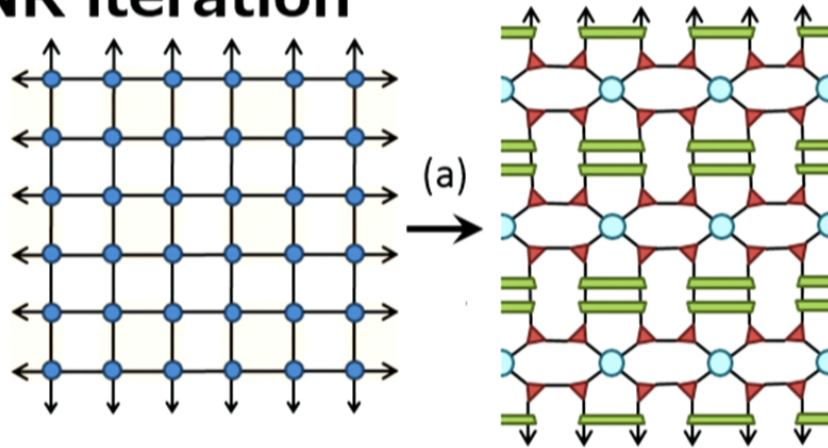
Can still address all short-range degrees of freedom...



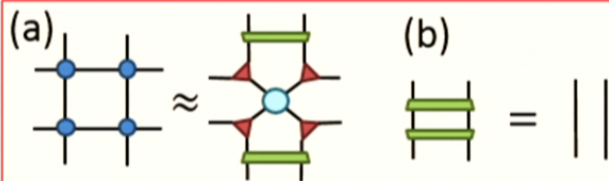
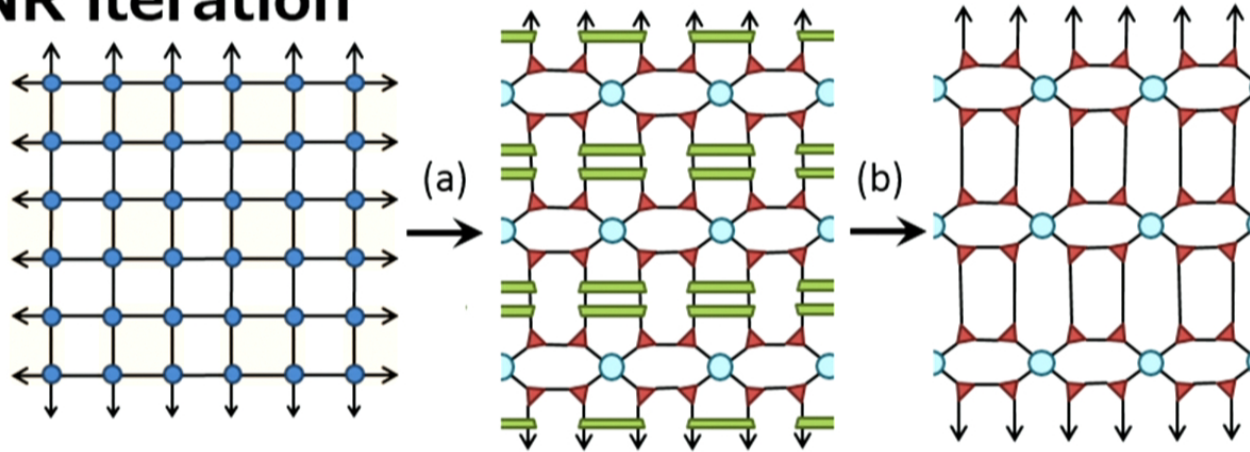
## TNR iteration



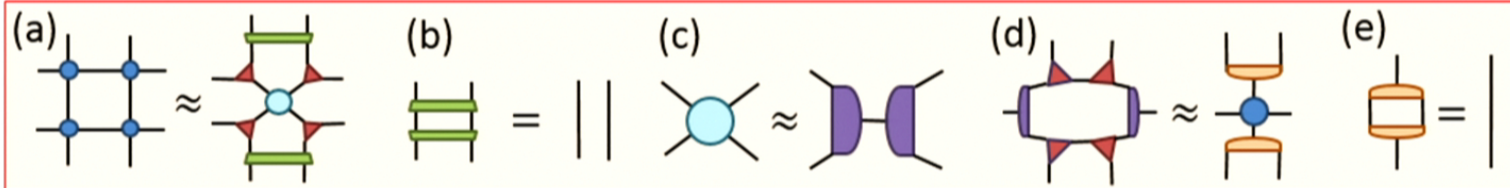
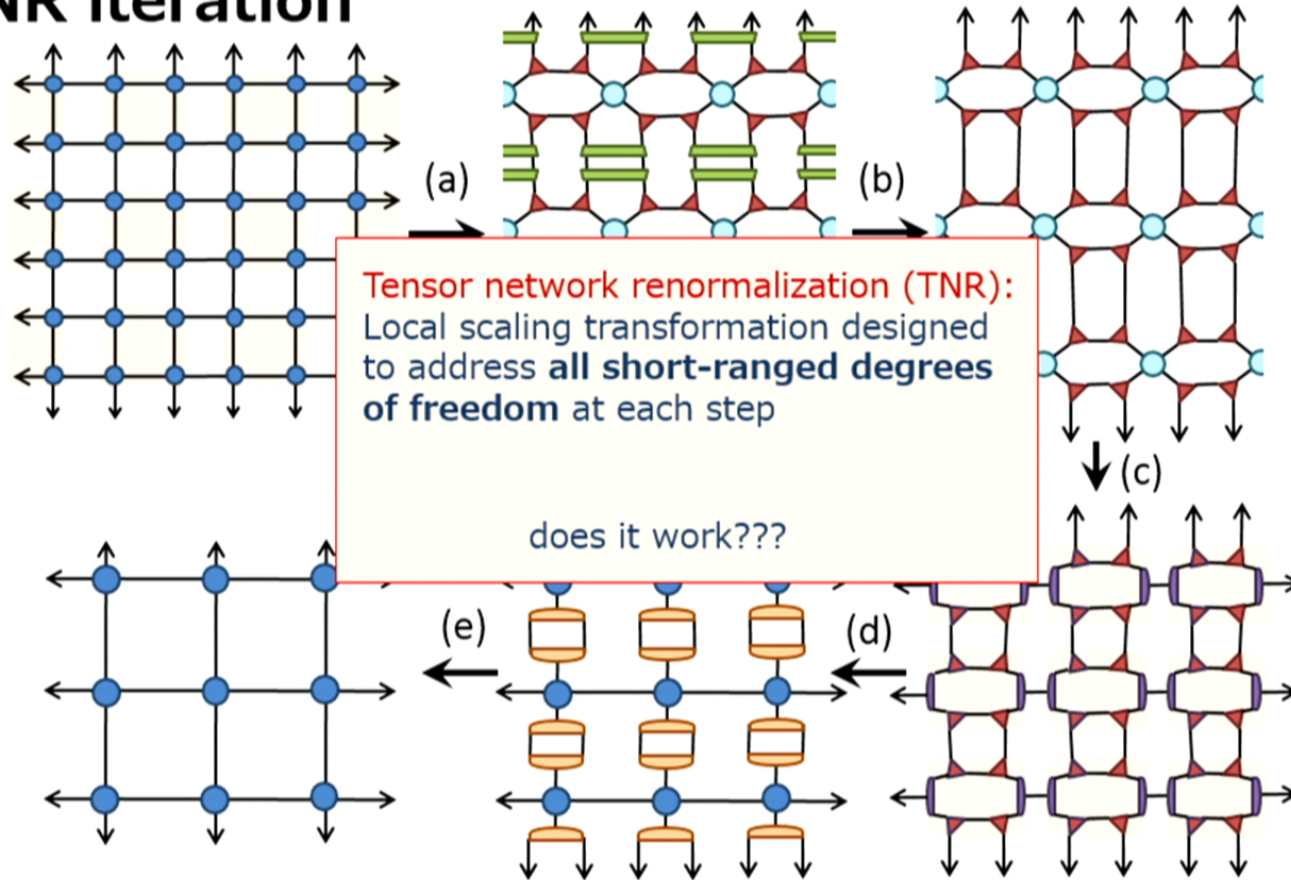
# TNR iteration



# TNR iteration



# TNR iteration

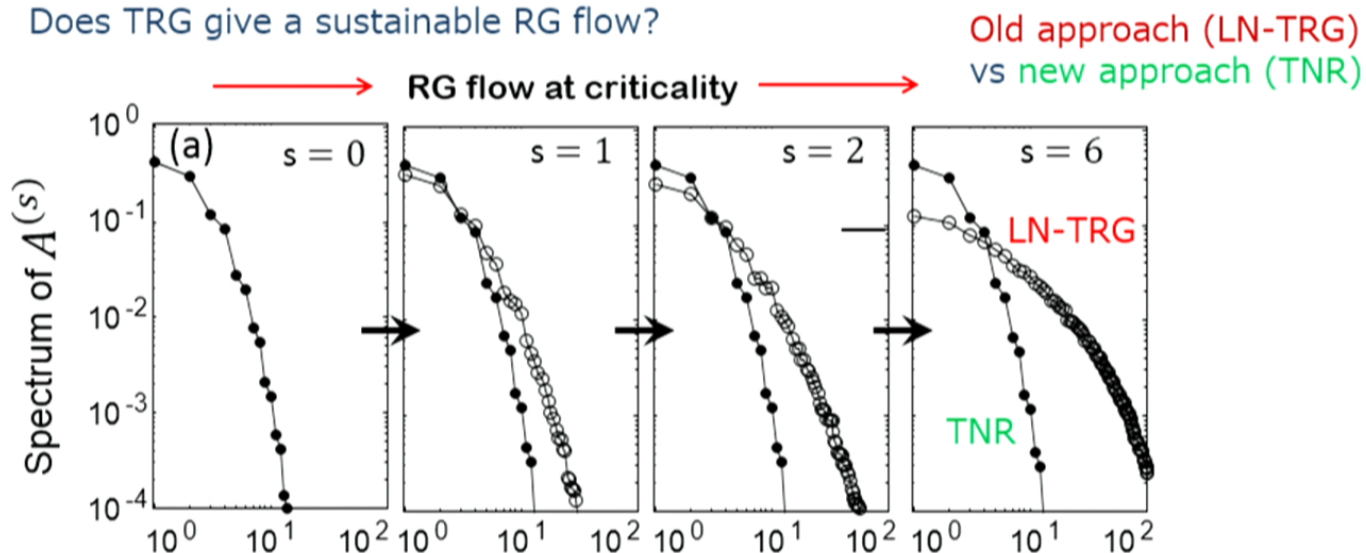




# Sustainable RG flow

RG flow in the space of tensors:  $A^{(0)} \rightarrow A^{(1)} \rightarrow A^{(2)} \rightarrow \dots \rightarrow A^{(s)} \rightarrow \dots$

Does TRG give a sustainable RG flow?

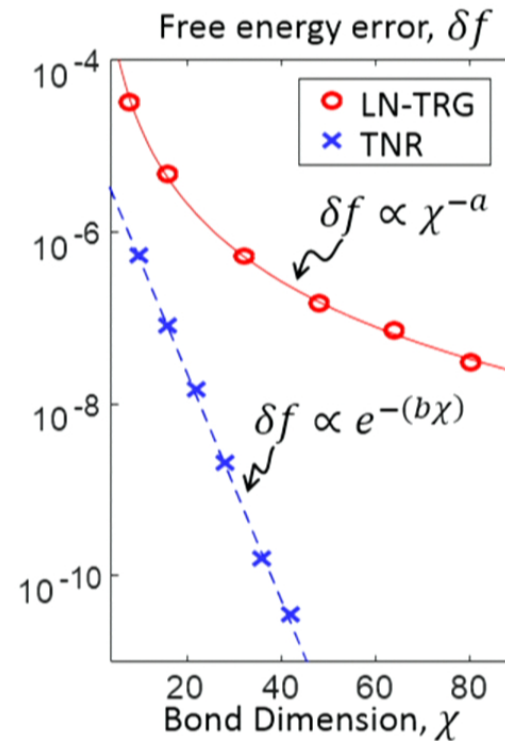
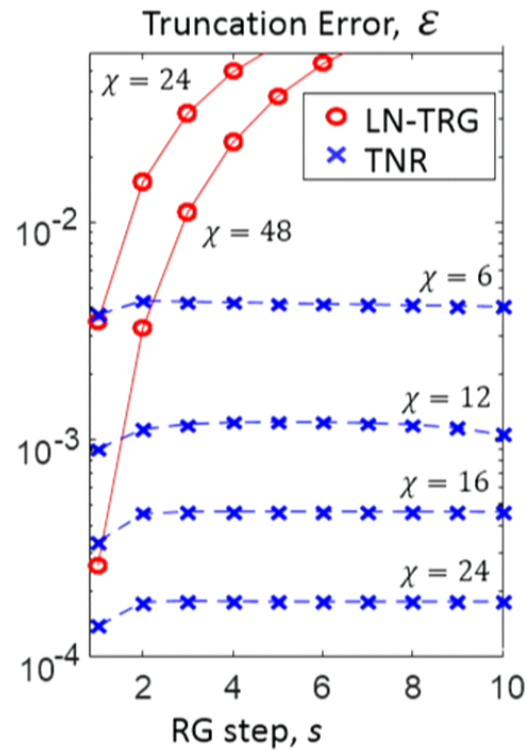


LN-TRG: Spectra become increasingly flat with each RG step! ❌

TNR: Spectra unchanged over many RG steps

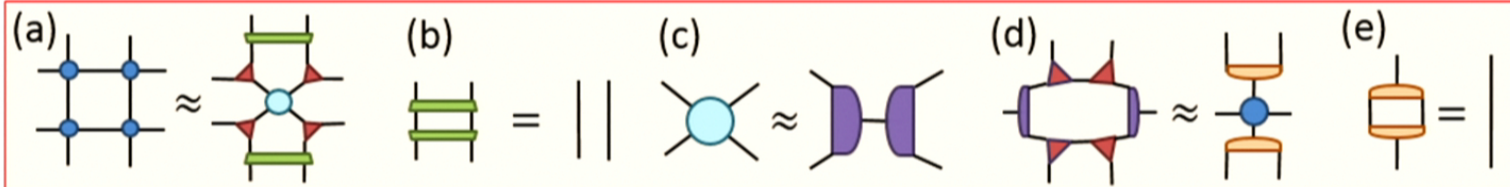
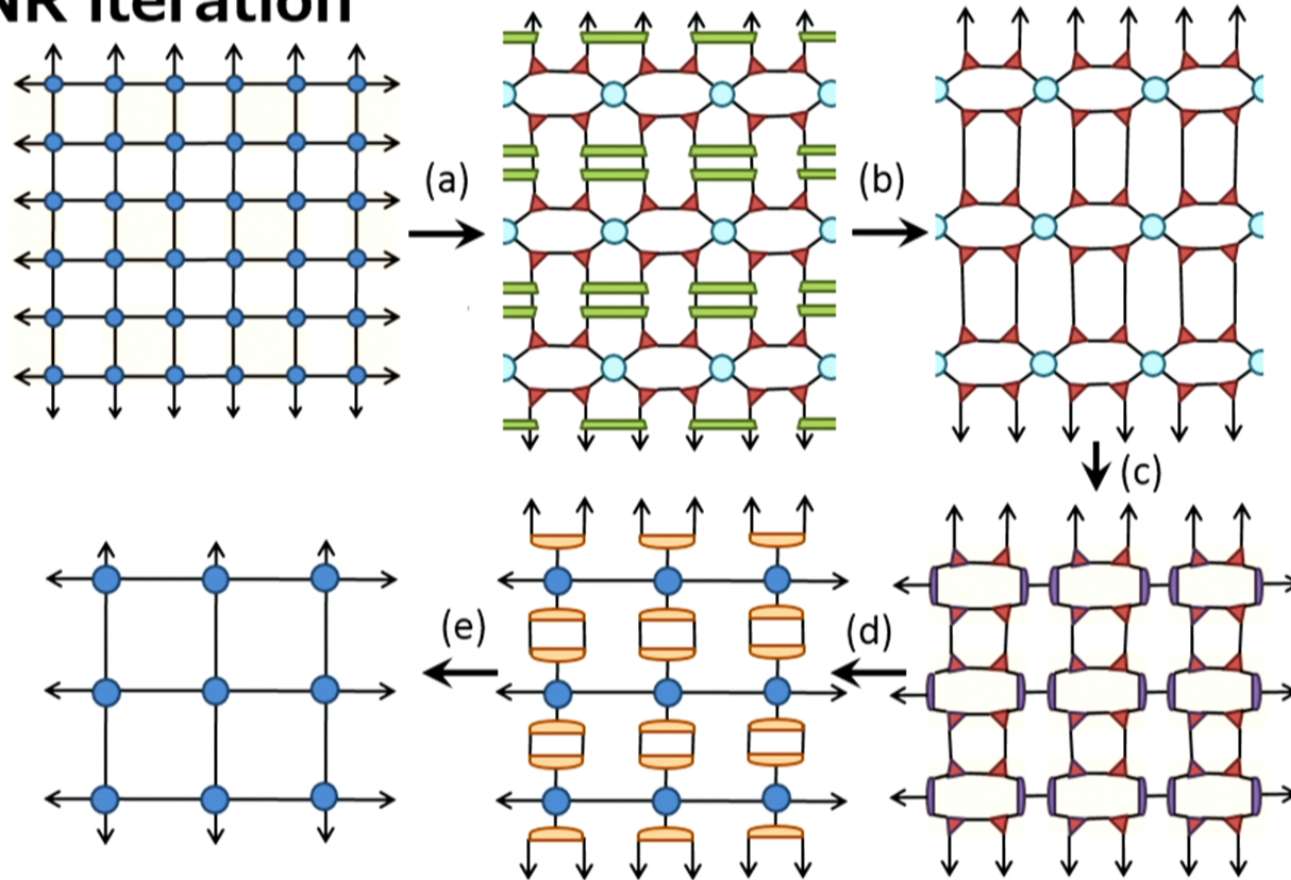


# Benchmark results, 2D classical Ising at criticality



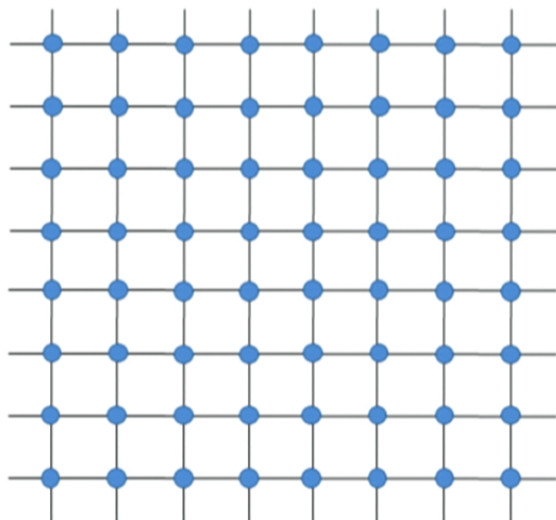
Old approach (LN-TRG):	truncation error <b>grows</b> with RG step (fixed $\chi$ )	free energy converges <b>polynomially</b> in $\chi$
New approach (TNR):	truncation error <b>remains constant</b>	free energy converges <b>exponentially</b> in $\chi$

# TNR iteration



# Overview

RG flow in the space of tensors:  $A^{(0)} \rightarrow A^{(1)} \rightarrow A^{(2)} \rightarrow \dots \rightarrow A^{(s)} \rightarrow \dots$



contraction  $\longrightarrow$

partition  
function

$$Z$$

free energy  
(per site)

$$f = \frac{-kT \log(Z)}{N}$$

**Practical goal:** efficient and accurate contraction of a tensor network to a scalar

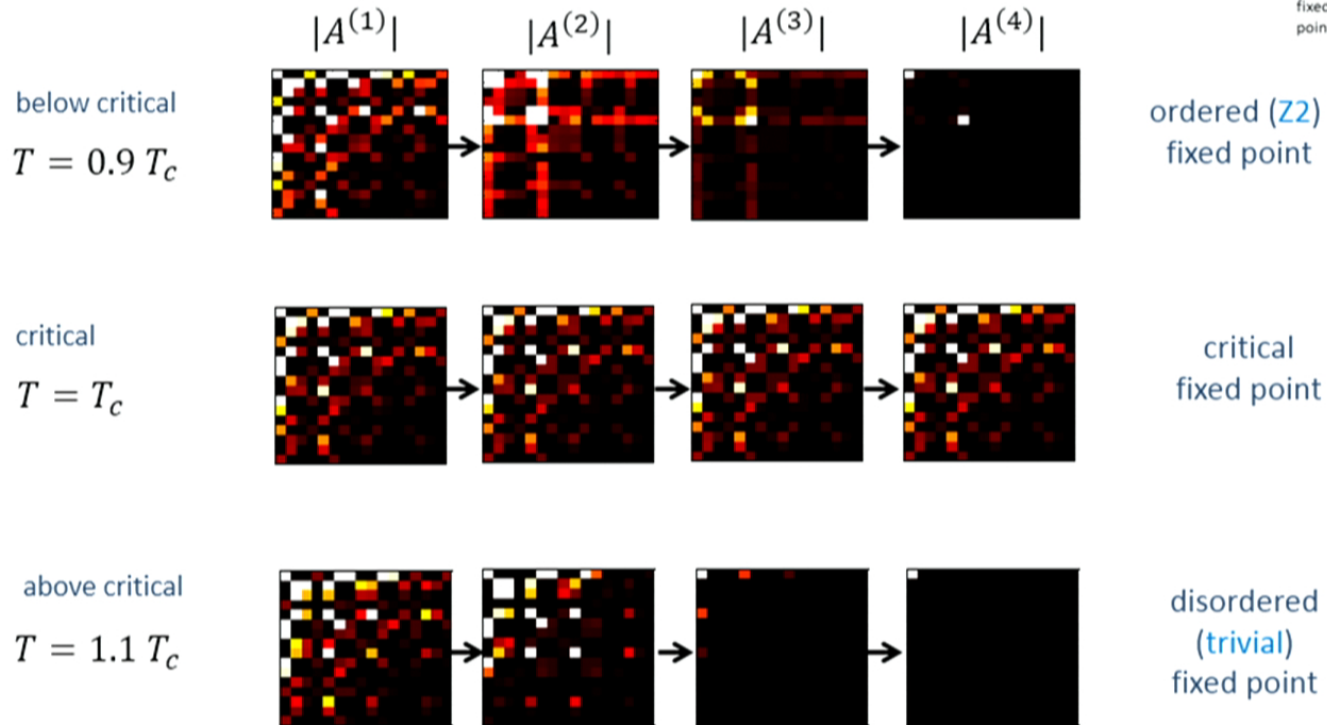
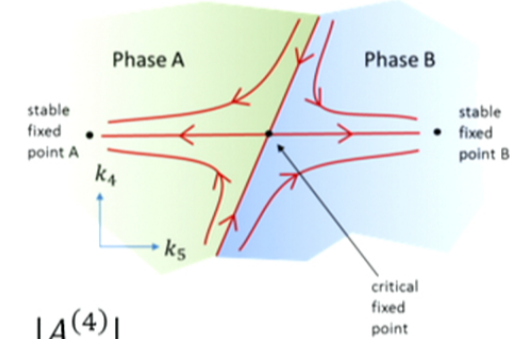


**Conceptual goal:** achieve a proper RG flow

# TNR: proper RG flow

2D classical Ising,  $\chi = 6$  TNR

RG flow in the space of tensors:  $A^{(0)} \rightarrow A^{(1)} \rightarrow A^{(2)} \rightarrow \dots \rightarrow A^{\text{fp}}$

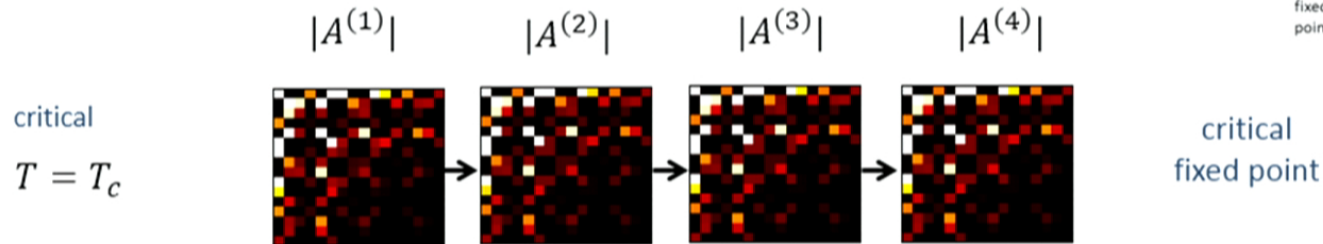
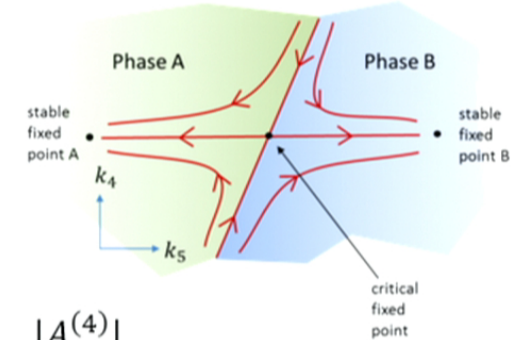




# TNR: proper RG flow

2D classical Ising

RG flow in the space of tensors:  $A^{(0)} \rightarrow A^{(1)} \rightarrow A^{(2)} \rightarrow \dots \rightarrow A^{\text{fp}}$



Previous approach (LN-TRG): RG flow does not converge at criticality

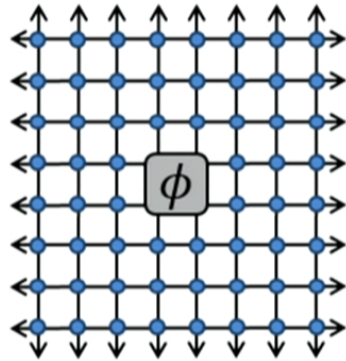
New approach (TNR): RG flow converges at criticality (to arbitrarily high precision)

Is the critical fixed point RG map given by TNR a good approximation to the Ising CFT???

...it should be if the truncation error has been kept small, but let's check anyway

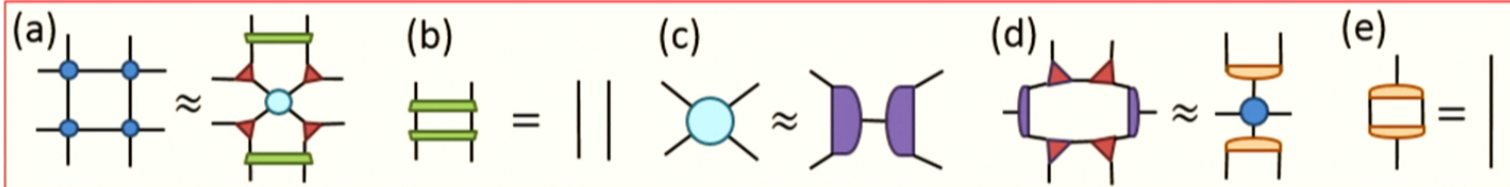
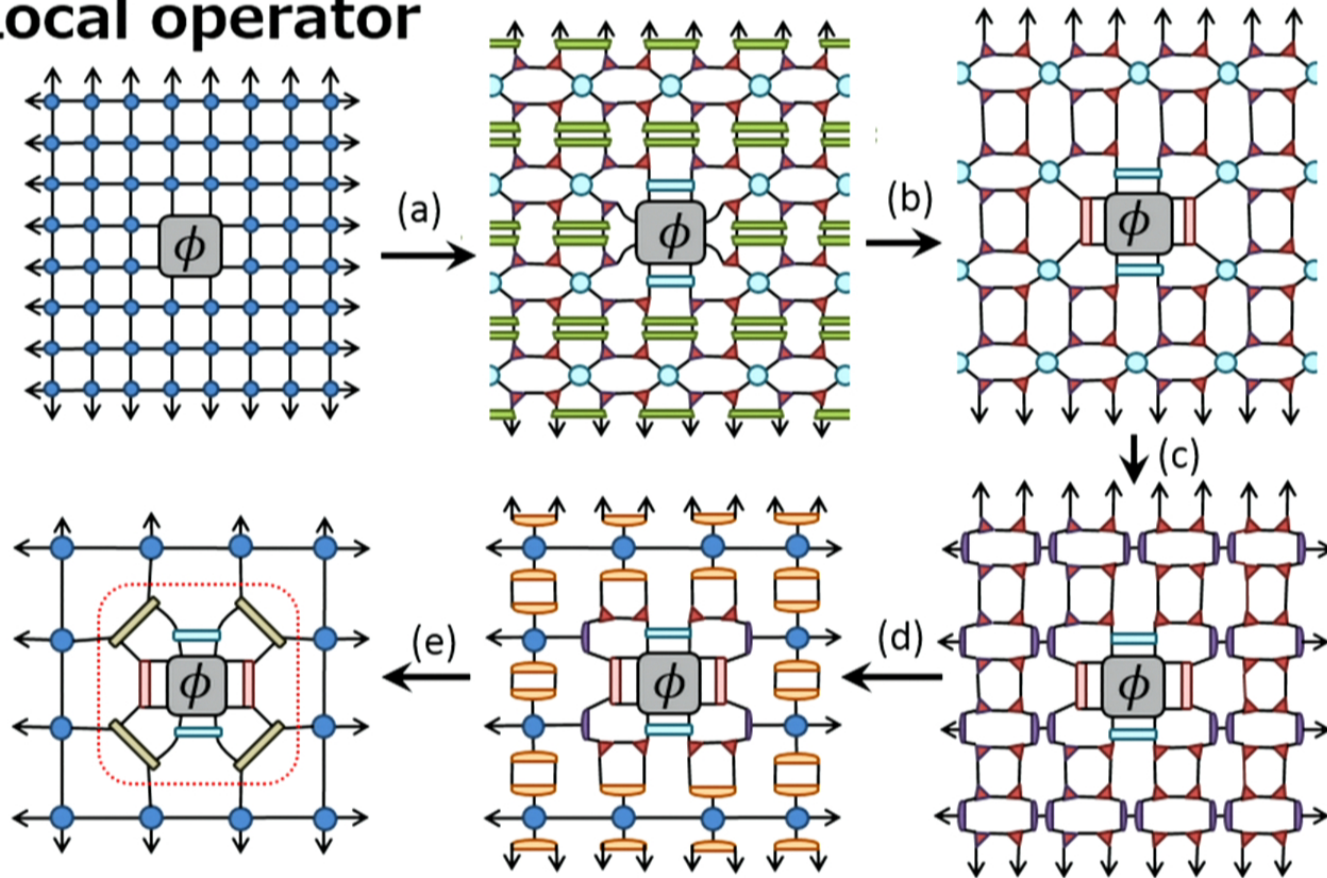


# Local operator

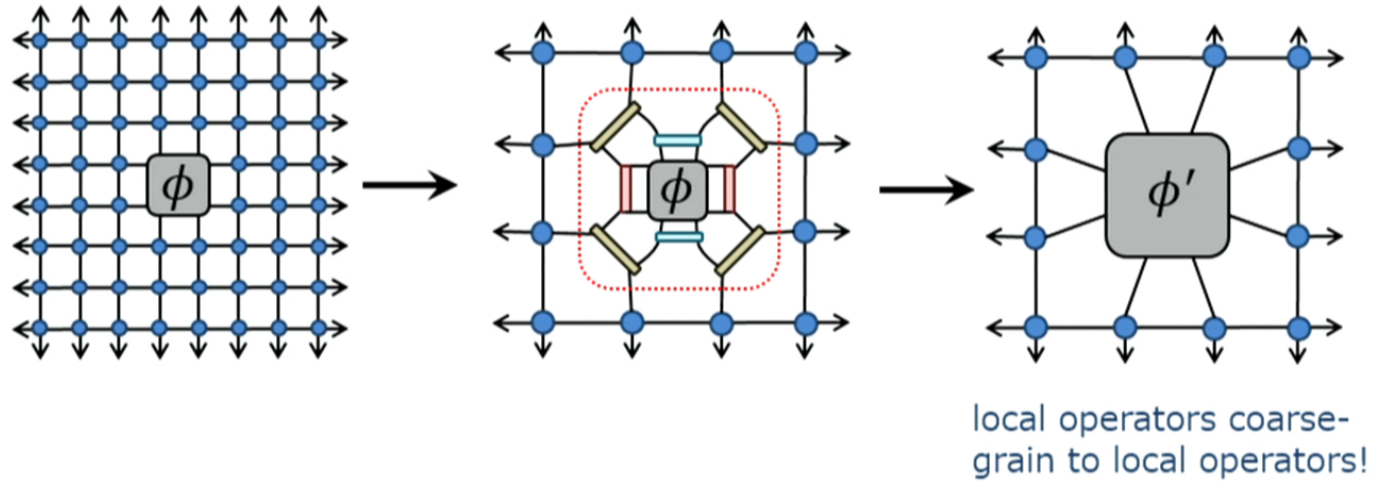


how does a local operator (inserted as in impurity) transform under the fixed point RG map?

# Local operator

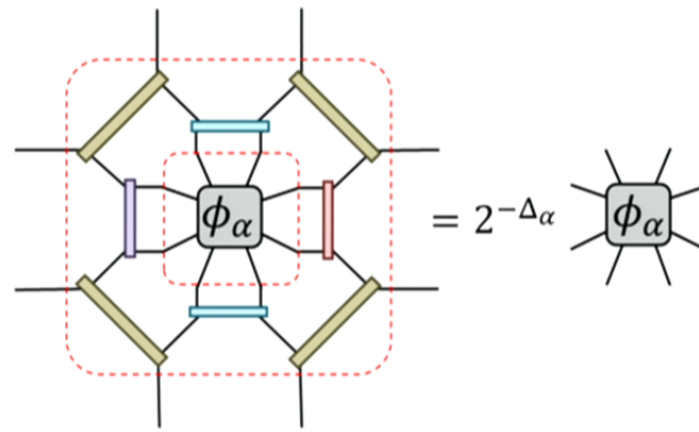


## Local operator



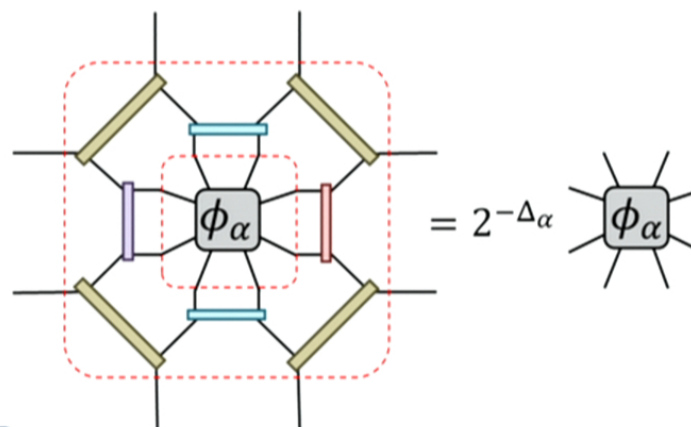
# Local impurity

diagonalize transfer operator  
to find scaling operators

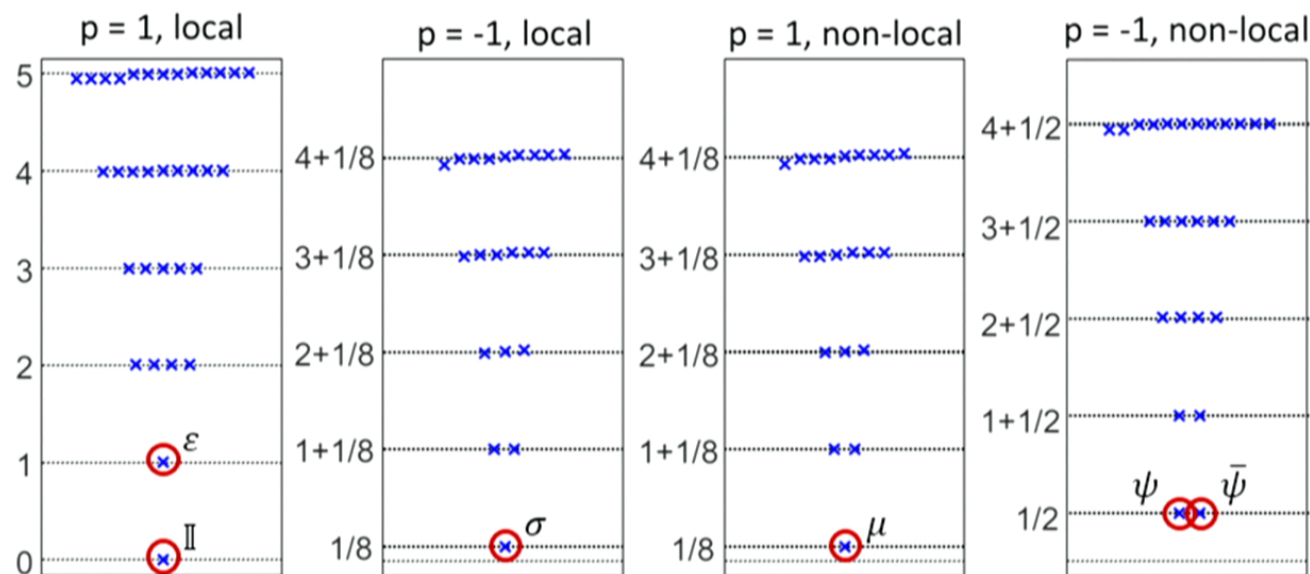


# Local impurity

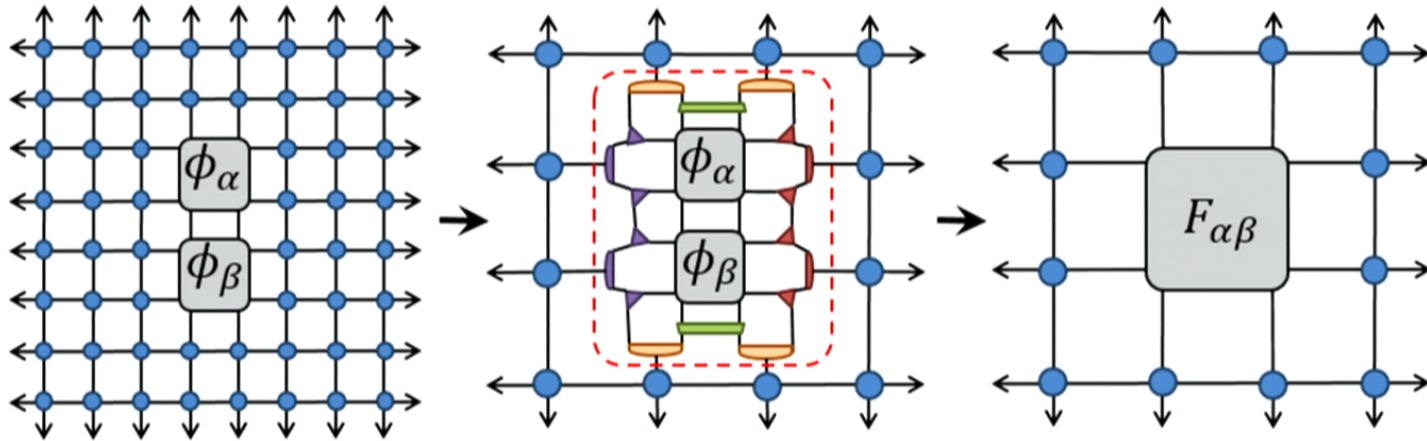
diagonalize transfer operator  
to find scaling operators



2D classical Ising at criticality,  $\chi = 6$  TNR



# Fusion of operators



expand fused operator in  
basis of scaling operators:

$$F_{\alpha\beta} = \sum_{\gamma} C_{\alpha\beta\gamma} \underbrace{\phi_{\gamma}}_{\text{operator product expansion (OPE) coefficients}}$$

exact

$\chi = 6$  TNR

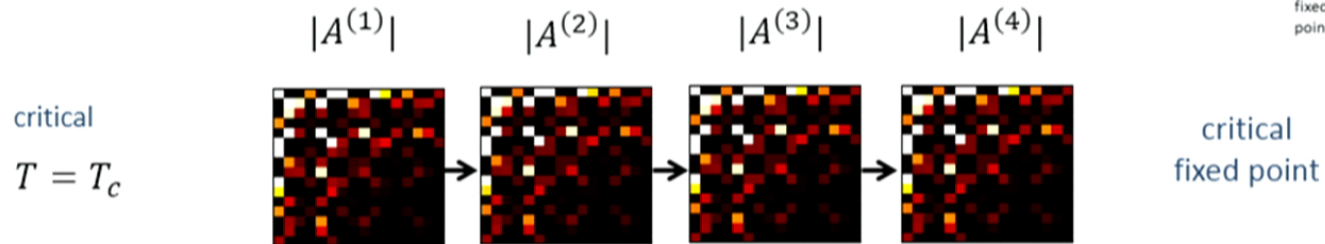
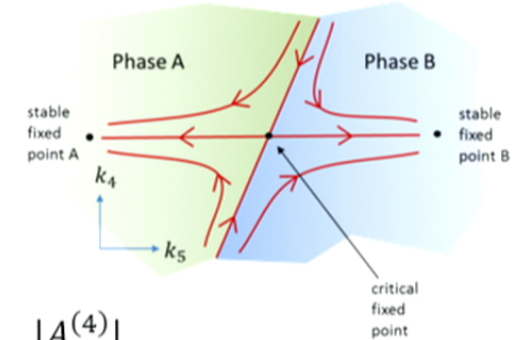
$C_{\varepsilon,\sigma,\sigma}$	$1/2$	0.50105
$C_{\varepsilon,\mu,\mu}$	$-1/2$	-0.49941
$C_{\psi,\mu,\sigma}$	$(1-i)/2$	$0.49957 - 0.49957i$
$C_{\bar{\psi},\mu,\sigma}$	$(1+i)/2$	$0.49957 + 0.49957i$
$C_{\varepsilon,\psi,\bar{\psi}}$	$i$	$1.00009i$
$C_{\varepsilon,\bar{\psi},\psi}$	$-i$	$-1.00009i$



# TNR: proper RG flow

2D classical Ising

RG flow in the space of tensors:  $A^{(0)} \rightarrow A^{(1)} \rightarrow A^{(2)} \rightarrow \dots \rightarrow A^{\text{fp}}$



Previous approach (LN-TRG): RG flow does not converge at criticality

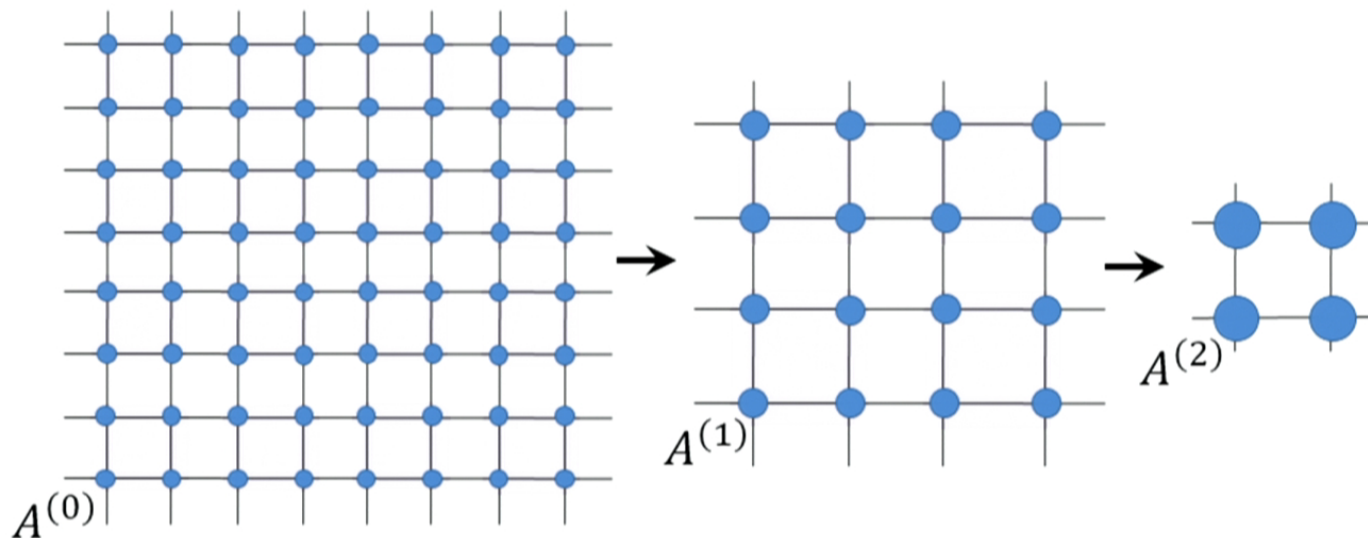
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# Overview

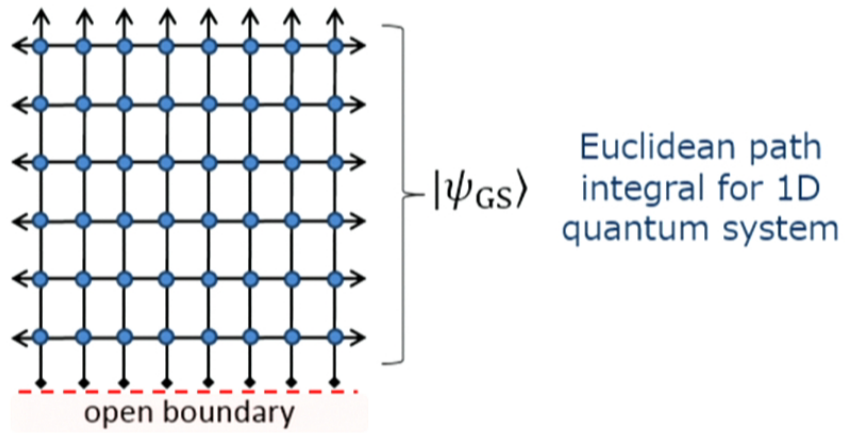
local scale transformations on the lattice:

RG flow in the space of tensors:  $A^{(0)} \rightarrow A^{(1)} \rightarrow A^{(2)} \rightarrow \dots \rightarrow A^{(s)} \rightarrow \dots$

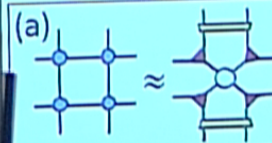
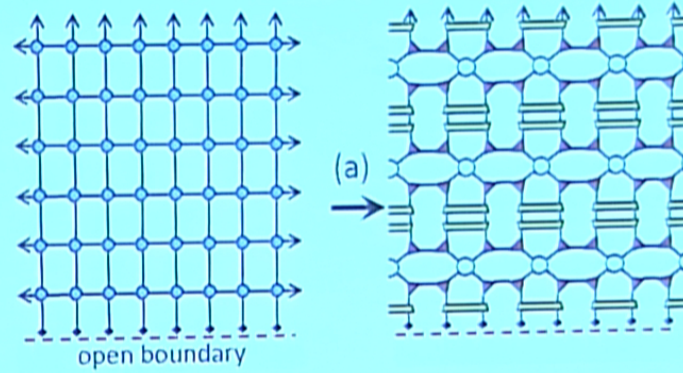


**Practical goal:** efficient and accurate contraction of a tensor network to a scalar

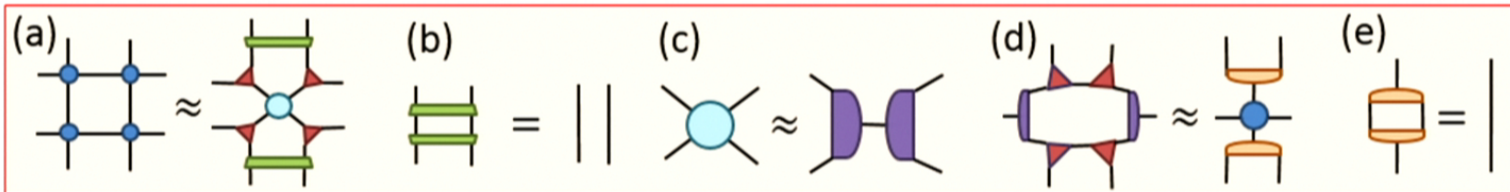
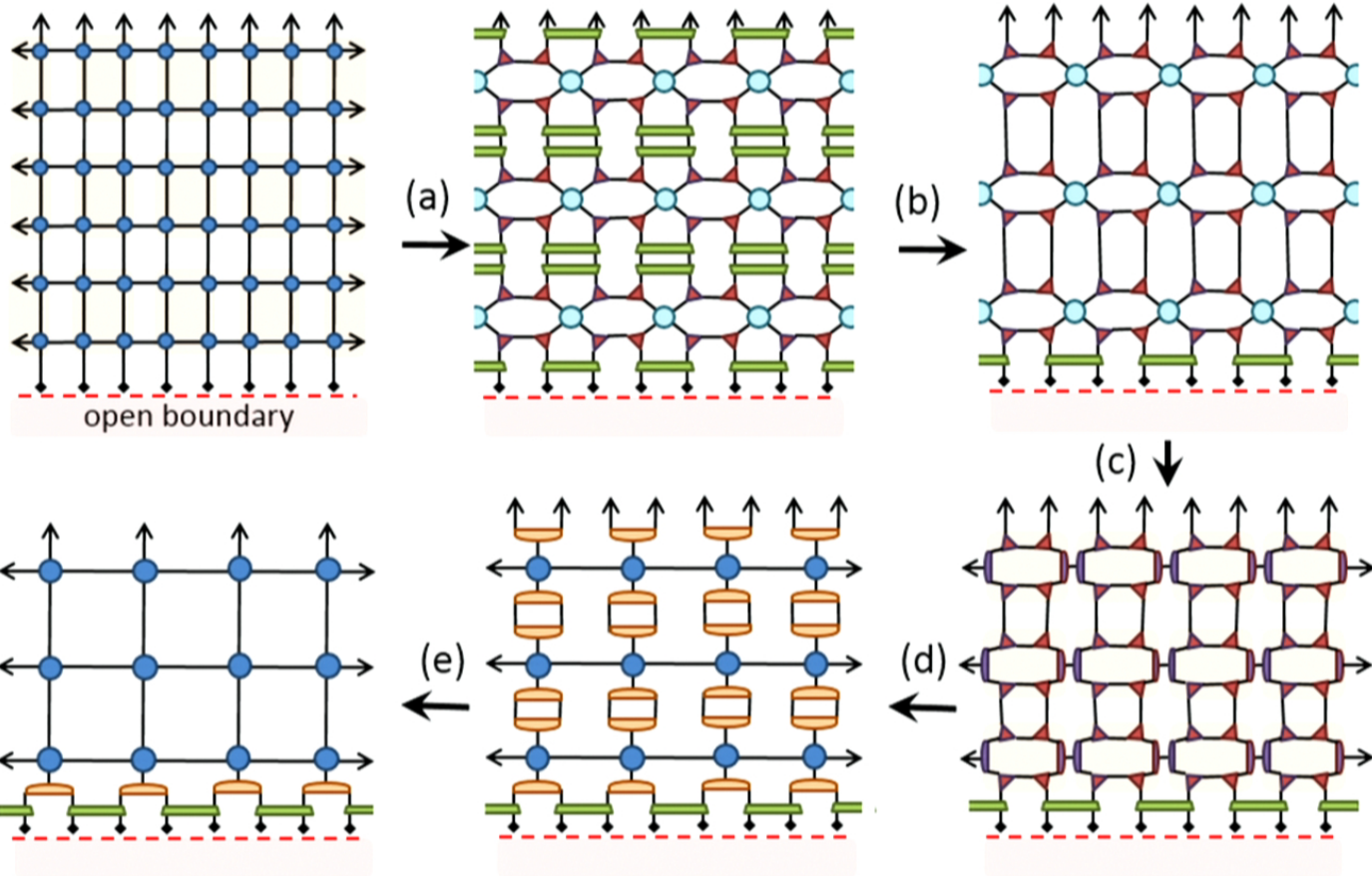
**Conceptual goal:** achieve a proper RG flow



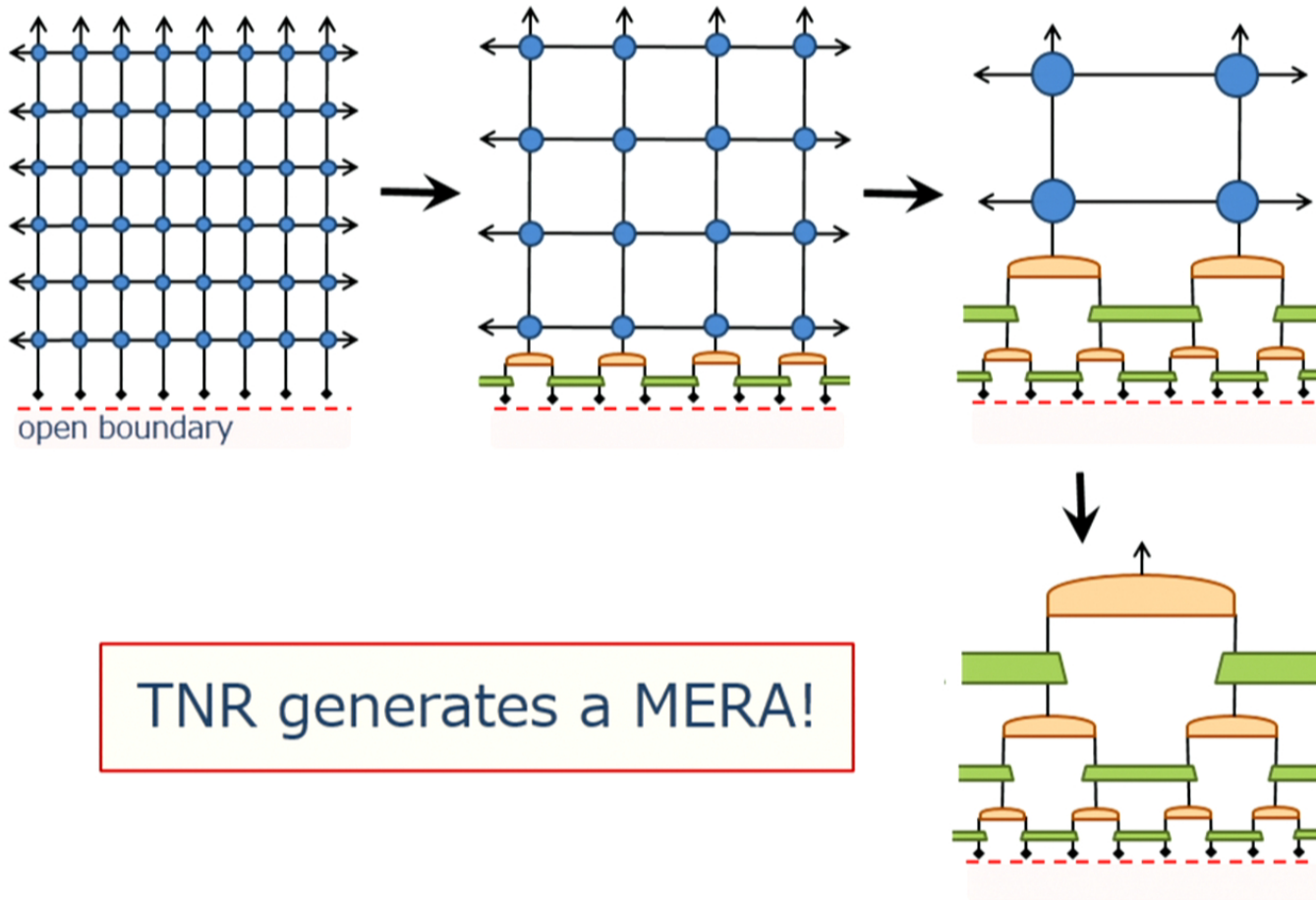
Apply TNR to the  
half plane???







# TNR yields the MERA (Evenbly, Vidal, arXiv:1502.05385)

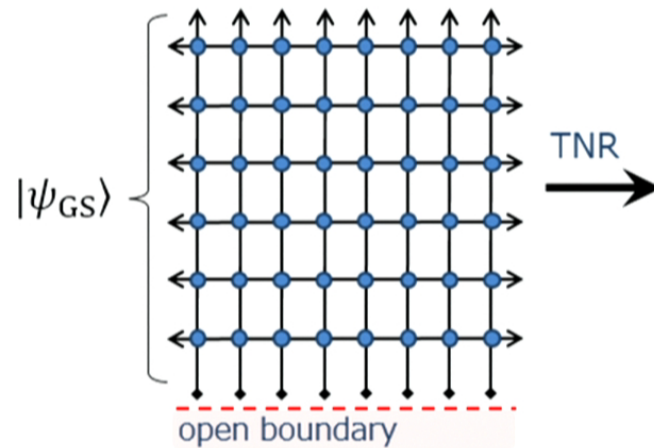




# TNR yields the MERA

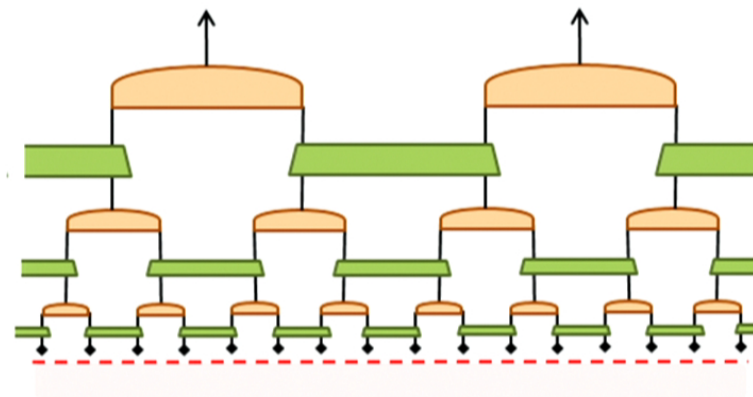
(Evenbly, Vidal, arXiv:1502.05385)

exact representation of ground state as a path integral



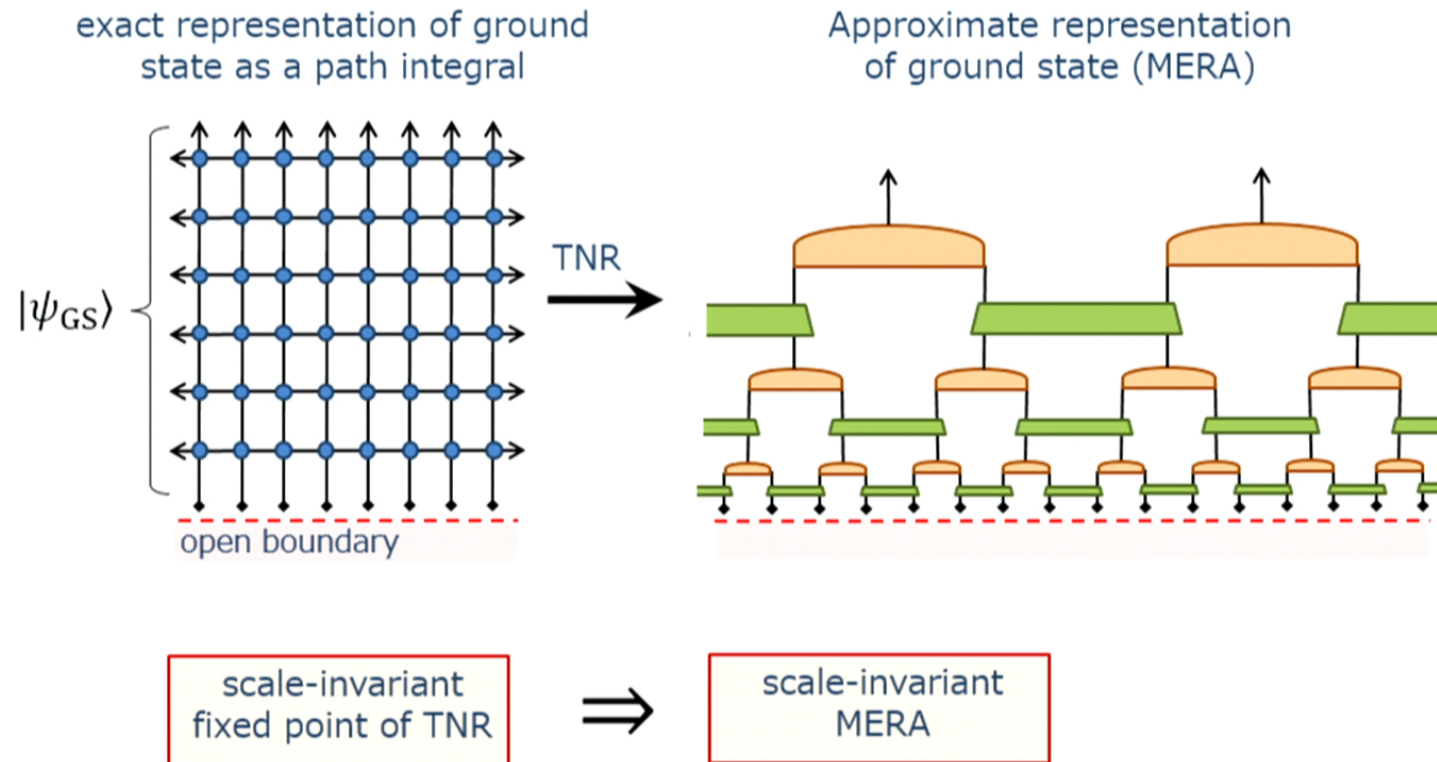
TNR

Approximate representation of ground state (MERA)

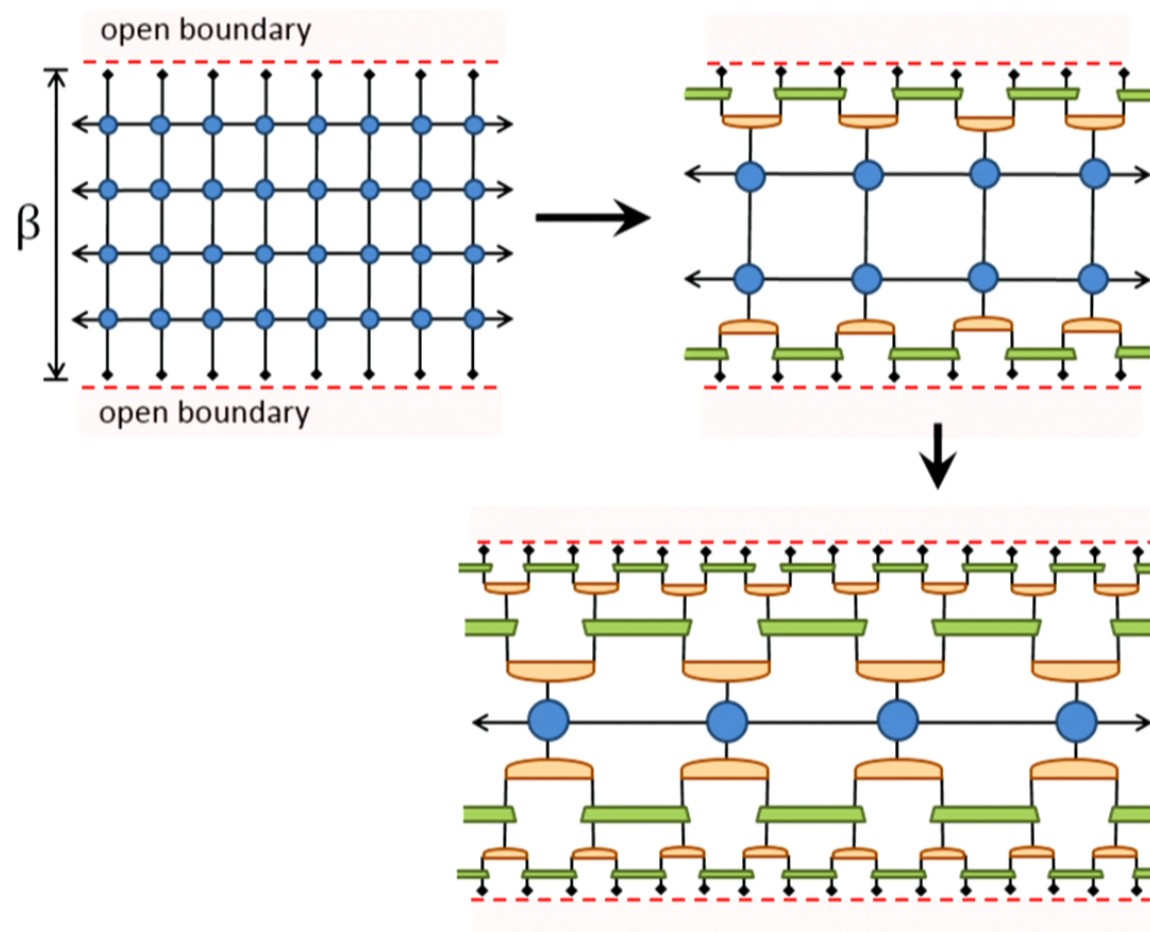


# TNR yields the MERA

(Evenbly, Vidal, arXiv:1502.05385)



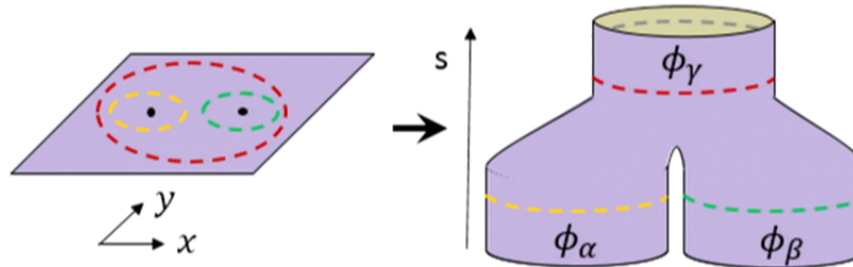
# MERA for a thermal state (or black hole in holography)



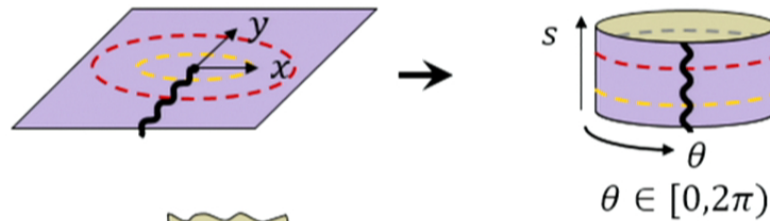
# Local scale transformations on the lattice with TNR

(G. E., D. Gaiotto, R. Myers, G. Vidal, *in prep.* )

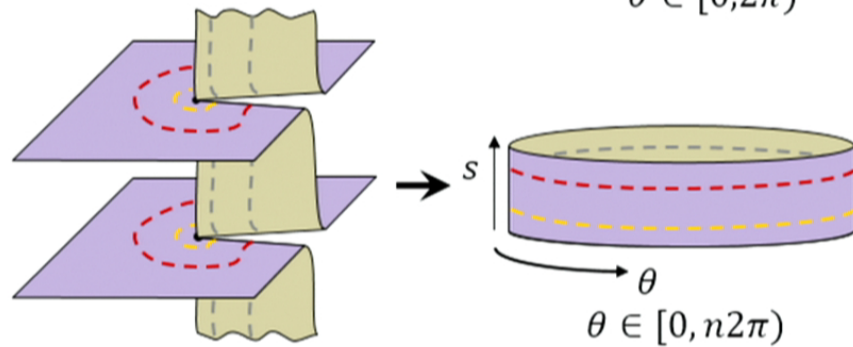
plane with two  
punctures into  
pair of pants



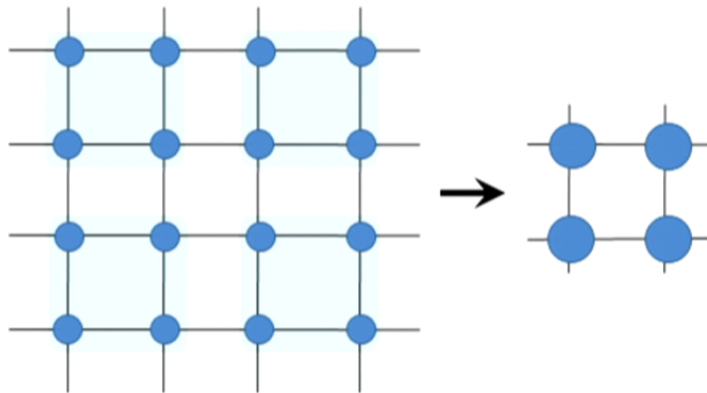
plane with semi-  
infinite defect line  
to cylinder



$n$  - sheeted  
Riemann surface  
to cylinder



# Summary



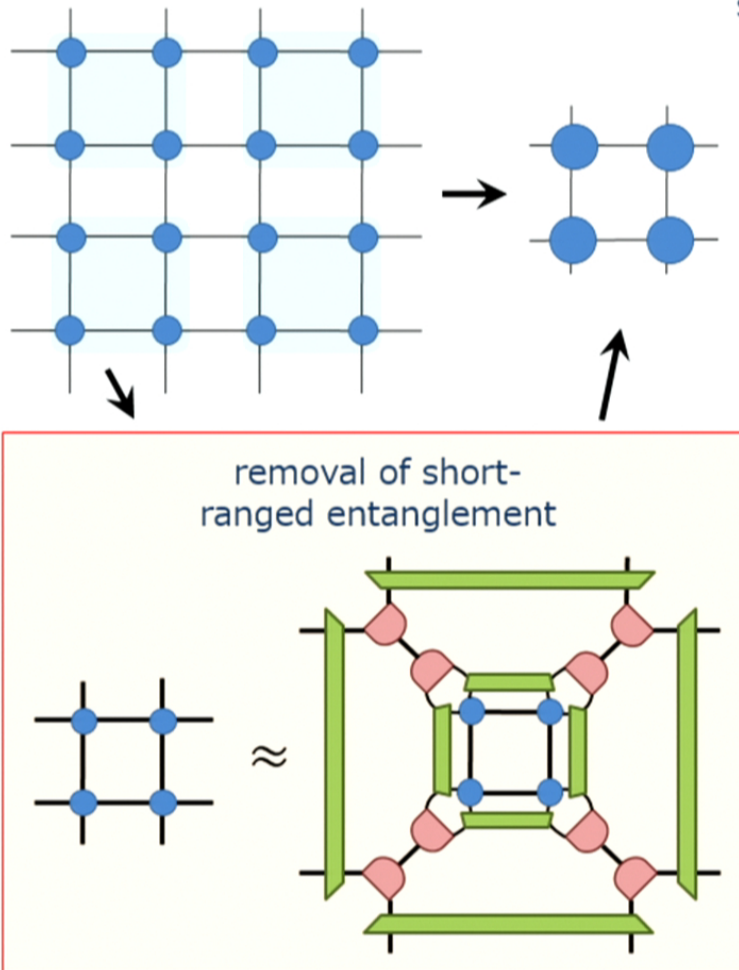
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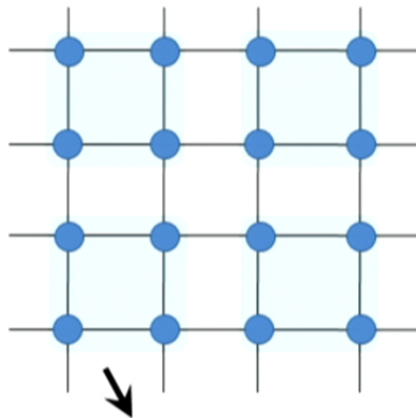


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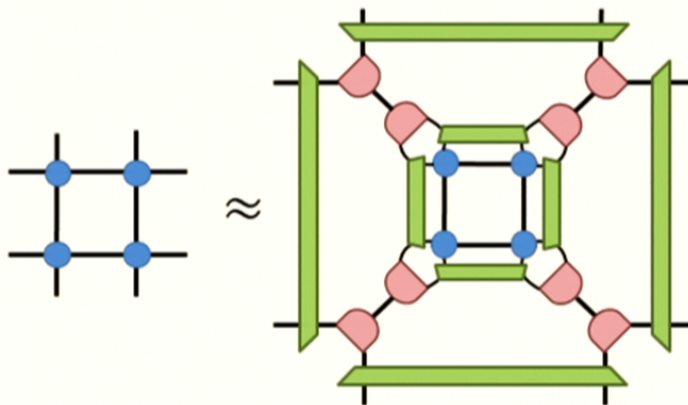
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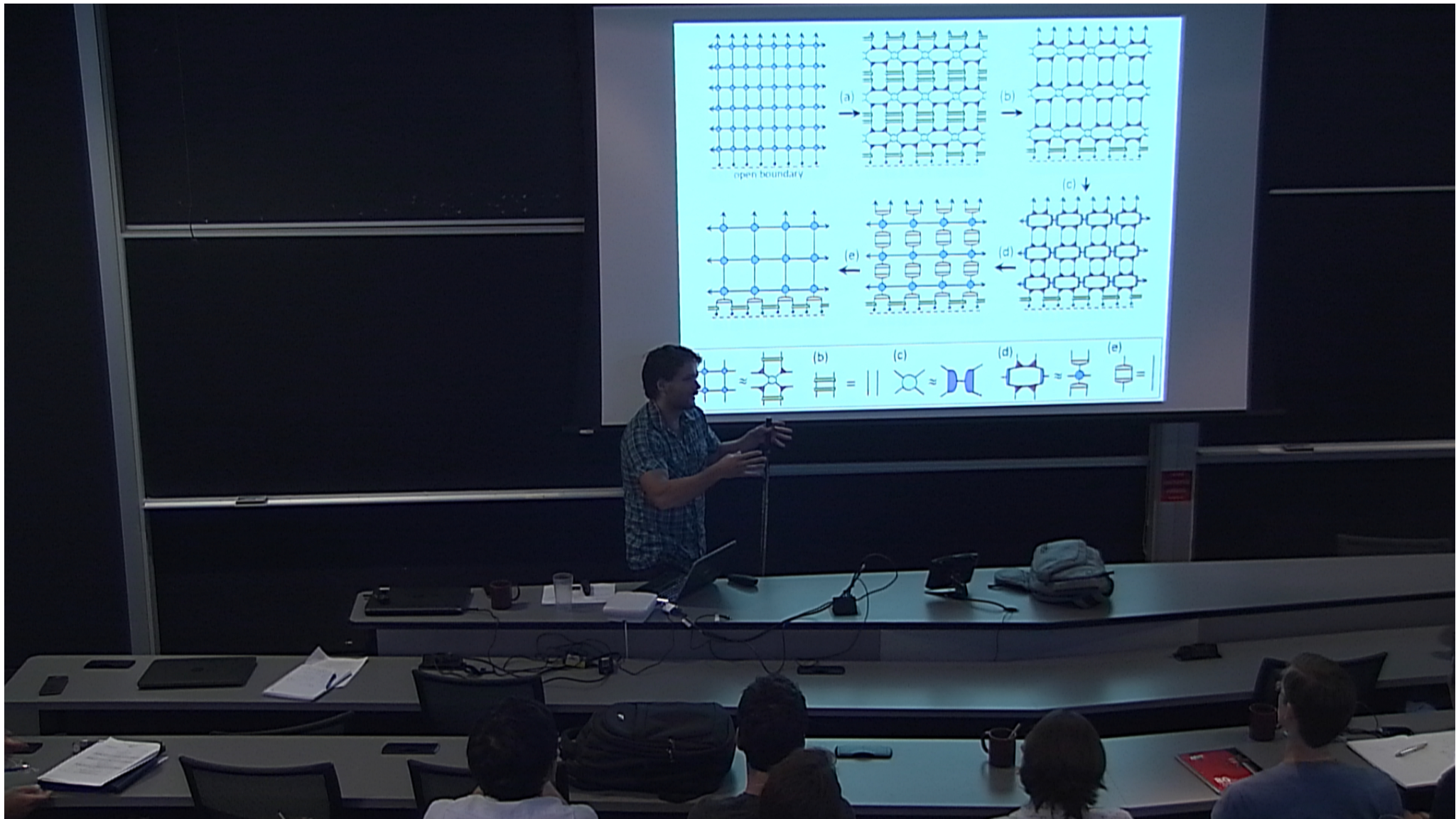
**Thanks!**

removal of short-ranged entanglement



- Base for powerful numeric algorithms
- Nice conceptual connections

Future: higher dimensions, other applications, blah blah blah...



# TNR yields the MERA

(Evenbly, Vidal, arXiv:1502.05385)

