

Title: TBA

Date: Aug 11, 2015 10:40 AM

URL: <http://pirsa.org/15080014>

Abstract: TBA

Dispersion Distance and Large Scale Structure in Dispersion Space

Kris Sigurdson

The University of British Columbia



Cosmic Flows
(and other Novelties on Large Scales)
Perimeter Institute
August 11, 2015

Dispersion Distance and Large Scale Structure in Dispersion Space
(a potential Novel way to probe
Large Scales)

Kris Sigurdson

The University of British Columbia



Cosmic Flows
(and other Novelties on Large Scales)
Perimeter Institute
August 11, 2015

Dispersion Distance and Large Scale Structure in Dispersion Space

Kris Sigurdson

The University of British Columbia



Cosmic Flows
(and other Novelties on Large Scales)
Perimeter Institute
August 11, 2015

Based On:

“Dispersion Distance and the Matter Distribution of the Universe in Dispersion Space,”

Kiyoshi Wesley Masui and KS.

arXiv:1506.1704

Phys. Rev. Lett. (2015) in Press.

Work with Kiyo Masui @ UBC



(Kiyo made black-background slides)

Dispersion Distance and the Matter Distribution of the Universe in

Kiyoshi Wesley Masui

*Department of Physics and Astronomy, University of British Columbia, Vancouver, BC, C
Canadian Institute for Advanced Research, CIFAR Program in Cosmology and Gravity, Toronto*

Kris Sigurdson

Department of Physics and Astronomy, University of British Columbia, Vancouver, BC,

We propose that ‘standard pings’, brief broadband radio impulses, can be used to study the clustering of matter in the Universe even in the absence of redshift information. The dispersion they travel through the intervening plasma can, like redshift, be used as a cosmological distance measure. In the presence of inhomogeneities in the electron density along the line-of-sight, dispersion is an imperfect distance measure and we show that this leads to calculable dispersion-space distortions in the apparent positions of sources. Fast radio bursts (FRBs) are a new class of radio transients that are the prototypical example of these events due to their high observed dispersion, have been interpreted as originating at cosmological distances. The number of fast radio bursts has been estimated to be several thousand over the whole sky per day, and sources of these events should trace the large-scale structure of the Universe. We calculate the power spectra for a simple model where electrons and FRBs are biased tracers of the large-scale structure of the Universe and show that the clustering signal could be measured using as few as 10000 events.

Based On:

“Dispersion Distance and the Matter Distribution of the Universe in Dispersion Space,”

Kiyoshi Wesley Masui and KS.

arXiv:1506.1704

Phys. Rev. Lett. (2015) in Press.

Work with Kiyo Masui @ UBC



(Kiyo made black-background slides)

Dispersion Distance and the Matter Distribution of the Universe in

Kiyoshi Wesley Masui

*Department of Physics and Astronomy, University of British Columbia, Vancouver, BC, C
Canadian Institute for Advanced Research, CIFAR Program in Cosmology and Gravity, Toronto*

Kris Sigurdson

Department of Physics and Astronomy, University of British Columbia, Vancouver, BC,

We propose that ‘standard pings’, brief broadband radio impulses, can be used to study the clustering of matter in the Universe even in the absence of redshift information. The dispersion they travel through the intervening plasma can, like redshift, be used as a cosmological distance measure. In the presence of inhomogeneities in the electron density along the line-of-sight, dispersion is an imperfect distance measure and we show that this leads to calculable dispersion-space distortions in the apparent positions of sources. Fast radio bursts (FRBs) are a new class of radio transients that are the prototypical example of these events due to their high observed dispersion, have been interpreted as originating at cosmological distances. The number of fast radio bursts has been estimated to be several thousand over the whole sky per day, and sources of these events should trace the large-scale structure of the Universe. We calculate the power spectra for a simple model where electrons and FRBs are biased tracers of the large-scale structure of the Universe and show that the clustering signal could be measured using as few as 10000 events.

Inspired by great paper by Matt McQuinn:

“Locating the Missing Baryons with Extragalactic
Dispersion Measure Estimates,”

Matthew McQuinn

arXiv:1309.4451

ApJ, 780, L33 (2014).

Several other related contributions:

B. Dennison, MNRAS **443**, L11 (2014), arXiv:1403.2263
[astro-ph.HE].

W. Deng and B. Zhang, ApJ **783**, L35 (2014), arXiv:1401.0059
[astro-ph.HE].

B. Zhou, X. Li, T. Wang, Y.-Z. Fan, and D.-M. Wei,
Phys. Rev. D **89**, 107303 (2014), arXiv:1401.2927.

H. Gao, Z. Li, and B. Zhang, ApJ **788**, 189 (2014),
arXiv:1402.2498.

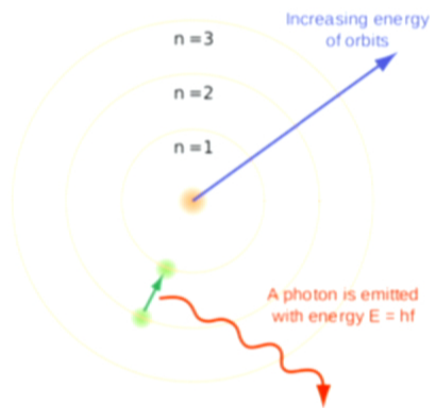
Apologies if I am missing your work!

Distance Measures

- Standard Atoms: Redshift
- Standard Candles: Luminosity Distance
- Standard Rulers: Angular Diameter Distance

Distance Measures

- Standard Atoms: Redshift



$$z \equiv \frac{\nu_e}{\nu_o} - 1 = \frac{\lambda_o}{\lambda_e} - 1$$

Doppler Shift in Special Relativity

$$1 + z = \sqrt{\frac{1 + v/c}{1 - v/c}}$$

Generalized: to geodesics in General Relativity

$$z \approx \frac{v}{c} = \frac{d}{D_H}$$

Formulas: D. Hogg Distance Measures Notes

Distance Measures

- Comoving Radial Distance (Flat Space):

$$\chi = D_C = D_H \int_0^z \frac{dz'}{E(z')}$$

$$H(z) = H_0 E(z)$$

Formulas: D. Hogg Distance Measures Notes

Distance Measures

- Standard Candles: Luminosity Distance



Distance Measures

- Standard Candles: Luminosity Distance

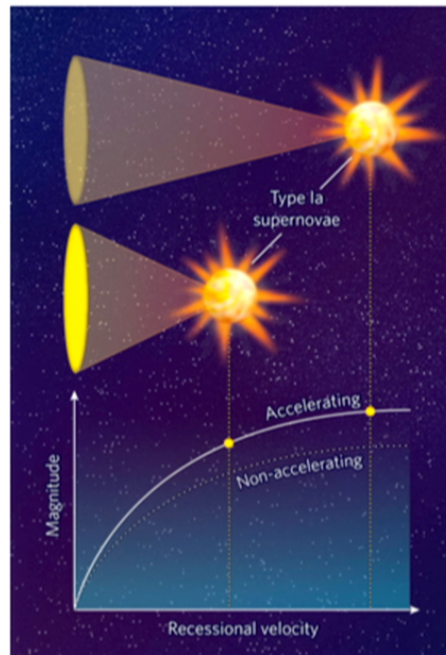


Image: Nature

$$D_L \equiv \sqrt{\frac{L}{4\pi S}}$$

$$D_L = (1 + z) D_C$$

Know Intrinsic Brightness and infer distance

Formulas: D. Hogg Distance Measures Notes

Distance Measures

- Standard Rulers: Angular Diameter Distance

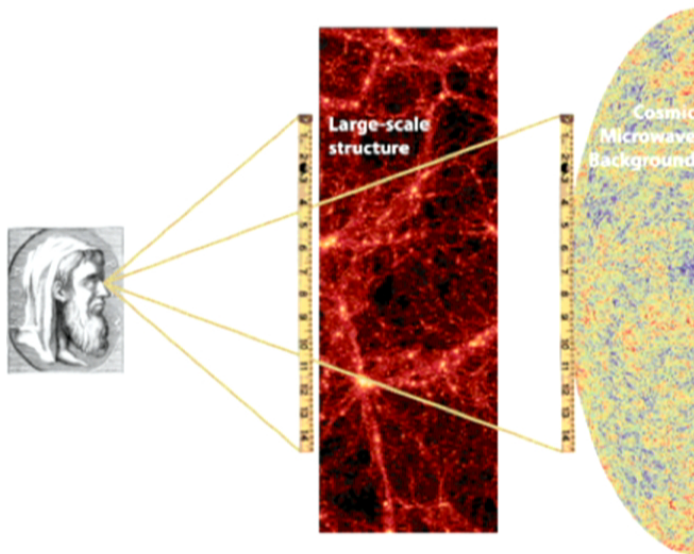


Image: ESA

$$D_A = \frac{D_C}{1 + z}$$

Know Intrinsic Physical Size and infer distance

Formulas: D. Hogg Distance Measures Notes

Dispersion Measure

As radio waves travel through a plasma their propagation speed is reduced below the speed of light due to the column of electrons.

$$DM = \int_0^d n_e dl$$

Frequency-dependent time delay

$$\Delta t = 4.2 \nu_{\text{GHz}}^{-2} \left(\frac{DM}{10^3 \text{ cm}^{-3} \text{ pc}} \right) \text{ s},$$

Formula: Matt McQuinn

Dispersion Measure

As radio waves travel through a plasma their propagation speed is reduced below the speed of light due to the column of electrons.

$$DM = \int_0^d n_e dl$$

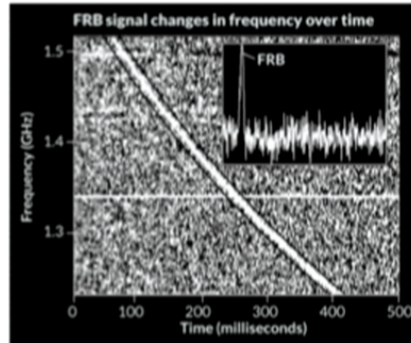
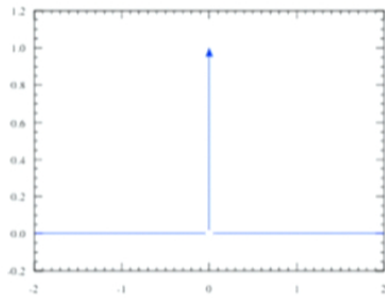
Frequency-dependent time delay

$$\Delta t = 4.2 \nu_{\text{GHz}}^{-2} \left(\frac{DM}{10^3 \text{ cm}^{-3} \text{ pc}} \right) \text{ s},$$

Formula: Matt McQuinn

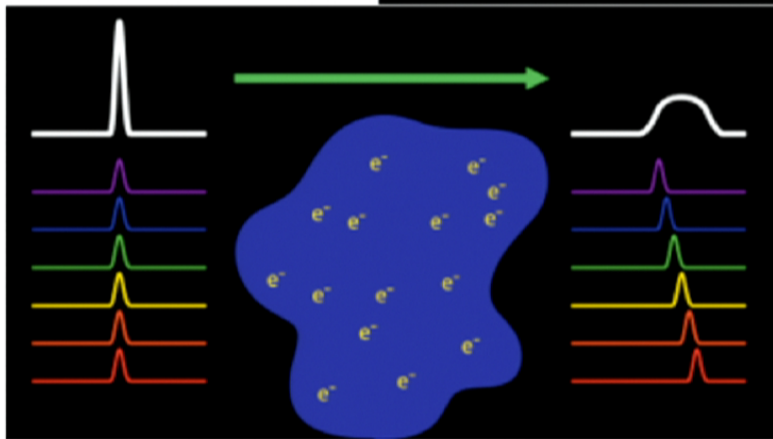
Distance Measures

- Standard Pings: Dispersion Distance



Standard Ping:
Broadband Radio Impulse

Prime Example: Fast Radio Bursts,
~ ms duration radio pulses



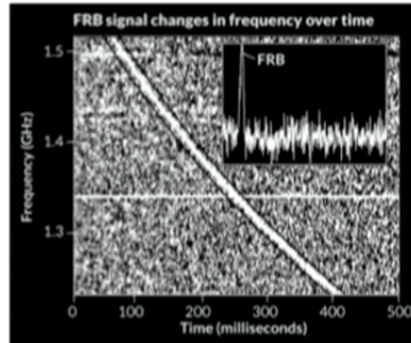
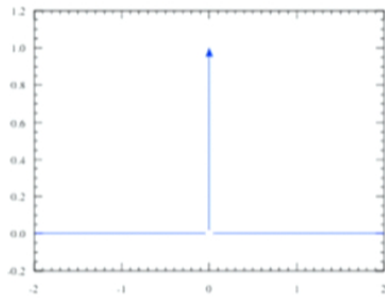
Know pulse is short duration and infer
DM, and Dispersion Distance χ_s

$$DM(\chi_s) = \int_0^{\chi_s} d\chi' a(\chi')^2 \bar{n}_e(\chi').$$

Images: Top Left: Lorimer; Bottom: Erik Madsen

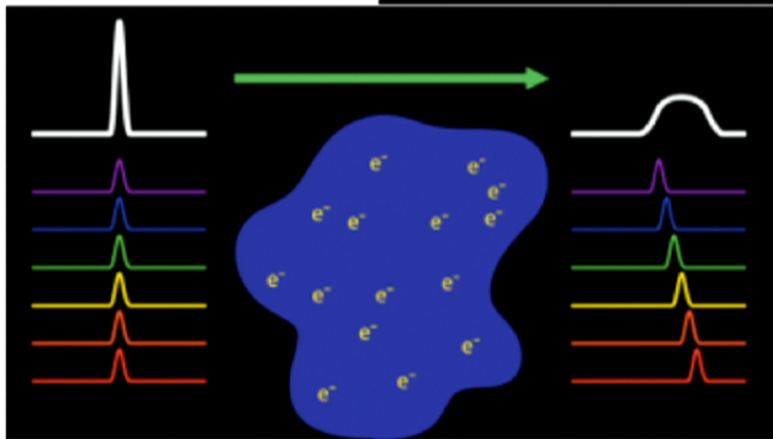
Distance Measures

- Standard Pings: Dispersion Distance



Standard Ping:
Broadband Radio Impulse

Prime Example: Fast Radio Bursts,
~ ms duration radio pulses



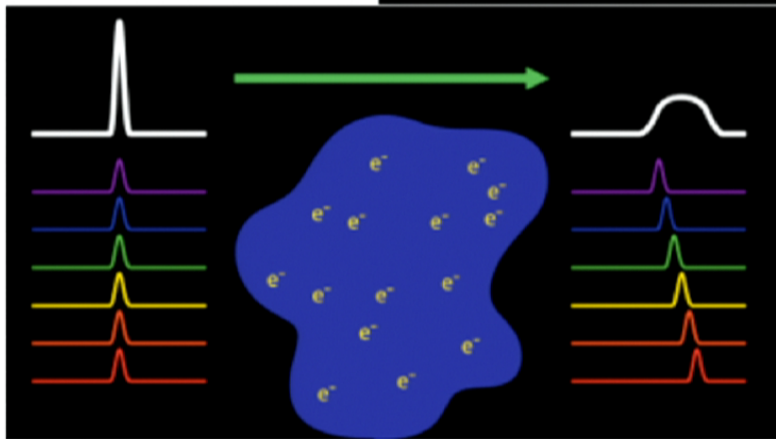
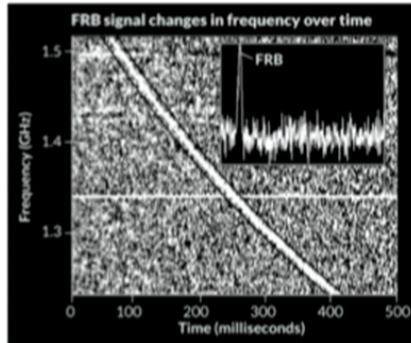
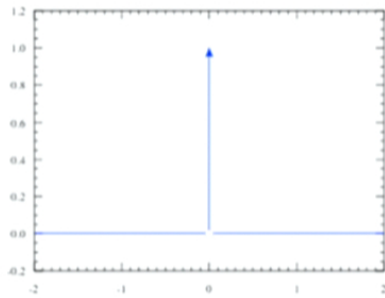
Know pulse is short duration and infer
DM, and Dispersion Distance χ_s

$$DM(\chi_s) = \int_0^{\chi_s} d\chi' a(\chi')^2 \bar{n}_e(\chi').$$

Images: Top Left: Lorimer; Bottom: Erik Madsen

Distance Measures

- Standard Pings: Dispersion Distance



Images: Top Left: Lorimer; Bottom: Erik Madsen

Standard Ping:
Broadband Radio Impulse

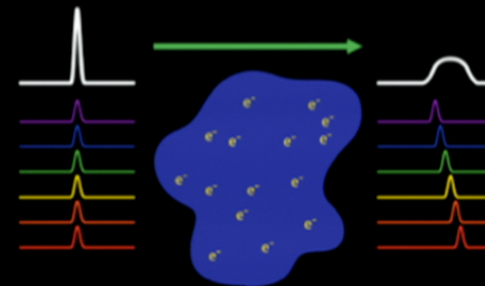
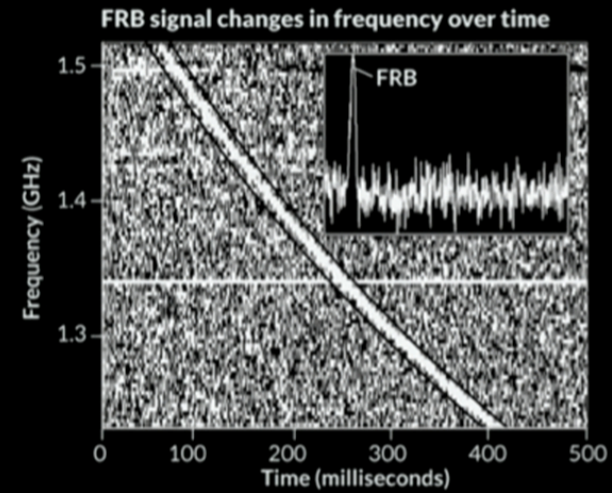
Prime Example: Fast Radio Bursts,
~ ms duration radio pulses

Know pulse is short duration and infer
DM, and Dispersion Distance χ_s

$$DM(\chi_s) = \int_0^{\chi_s} d\chi' a(\chi')^2 \bar{n}_e(\chi').$$

Fast radio bursts

- First detected By Lorimer using the Parkes Multibeam
- Now ~10 published events
- ~1 ms, ~1 Jy bursts
- Estimated ~10 000 similar events per day per sky
- Highly dispersed, interpreted as being extragalactic/cosmological
- If the dispersion is from the IGM, originate from $z \sim 0.5$



images: Lorimer 2007,
Erik Madsen

Fast radio bursts

- If FRBs are of cosmological origin...
- and If we had a survey with many thousands of events...
- then they should be clustered in some way
- How do we describe the clustering?

Location	Model	Galactic scintillation	Faraday rotation	Counterpart	dlnN/dlnS	PA swing	Repeat
Cosmological (>1 hGpc)	Blitzars	X	$\leq 7 \text{ rad/m}^2$	Grav waves	?	X	X
	Merging COs	X	$\leq 7 \text{ rad/m}^2$	Type Ia SNe, X-ray	?	X	X
	Primordial BHs	X	$\leq 7 \text{ rad/m}^2$	~TeV	?	X	X
	NS / Asteroid	X	$\leq 7 \text{ rad/m}^2$	Faint X-ray	?	X	X
Extragalactic, local	Edge-on disk	✓	50-500 rad/m^2	?	-3/2	?	?
	SNR pulsar	✓	20-1000 rad/m^2	Archival CC Sne	-3/2	✓	✓
Galactic	Flare stars	✓	RM_{gal}	Main sequence star	-3/2	X	✓
Terrestrial	RFI	X	RM_{ion}	None	-1/2 if 2D -3/2 if 3D	X	X

Connor et al. 2015 1505.05535

Slide: Liam Connor

What are FRBs? Where do they come from? We don't know yet.

This talk: what if all or some fraction of FRBs are cosmological.

Dispersion Space

- Dispersion measure is a line-of-sight integral of the electron density
- If dominated by the IGM this can be used as a cosmological distance measure
- Can study clustering in 3D!

$$\text{DM}(\hat{n}, \chi) = \int_0^\chi d\chi' a(\chi)^2 n_e(\hat{n}\chi', \chi').$$

Dispersion measure space

- Like redshift, DM is not a perfect proxy for distance due to perturbations

$$\frac{\text{DM}(\hat{n}, \chi)}{\bar{n}_0} = \int_0^\chi \frac{\chi'}{a(\chi')} [1 + \delta_e(\hat{n}\chi')]$$

$$\chi_s - \chi = \int_0^\chi d\chi' \delta_e(\hat{n}\chi')$$

- There will be additional clustering terms: the “DM space distortions”.

Dispersion Space

$$\text{DM}(\hat{n}, \chi) = \int_0^\chi d\chi' a(\chi')^2 \bar{n}_e(\chi') [1 + \delta_e(\hat{n}\chi', \chi')]$$

$$\vec{x}_s = \hat{n}\chi_s$$

$$\text{DM}(\chi_s) = \int_0^{\chi_s} d\chi' a(\chi')^2 \bar{n}_e(\chi').$$

$$\chi_s - \chi = \int_0^\chi d\chi' \delta_e(\hat{n}\chi').$$

Dispersion Space Distortions

We wish to relate the density of a tracer, f , measured in dispersion space to its density in real space. We follow the derivation in Kaiser [16] of the redshift-space distortions. Start by noting that the total number of tracers in a volume element is the same in both spaces:

$$n_{fs}(\vec{x}_s) d^3\vec{x}_s = n_f(\vec{x}) d^3\vec{x}. \quad (6)$$

$$\delta_{fs} = \delta_f - \delta_e - \left(\frac{1}{\bar{n}_f} \frac{d\bar{n}_f}{d\chi} + \frac{2}{\chi} \right) \int_0^\chi d\chi' \delta_e(\hat{n}\chi').$$

$$A(\chi) \equiv \frac{1}{\bar{n}_f} \frac{d\bar{n}_f}{d\chi} + \frac{2}{\chi}.$$

DM Space distortions

$$\delta_{fs} = \delta_f - \delta_e - \frac{2}{\chi}(\chi_s - \chi) - \frac{1}{\bar{n}_f} \frac{d\bar{n}_f}{d\chi}(\chi_s - \chi)$$

- For the most part can just follow Nick Kaiser's (1987) derivation of the (linear) redshift space distortions
- To leading order there are three distinct effects...

Dispersion Space

$$\text{DM}(\hat{n}, \chi) = \int_0^\chi d\chi' a(\chi')^2 \bar{n}_e(\chi') [1 + \delta_e(\hat{n}\chi', \chi')]$$

$$\vec{x}_s = \hat{n}\chi_s$$

$$\text{DM}(\chi_s) = \int_0^{\chi_s} d\chi' a(\chi')^2 \bar{n}_e(\chi').$$

$$\chi_s - \chi = \int_0^\chi d\chi' \delta_e(\hat{n}\chi').$$

Source density gradient

$$\delta_{fs} = \delta_f - \delta_e - \frac{2}{\chi}(\chi_s - \chi) - \frac{1}{\bar{n}_f} \frac{d\bar{n}_f}{d\chi} (\chi_s - \chi)$$

- When we misinterpret the radial distance we compare the density to the wrong background
- Again, in principle present in redshift space but small

Collect terms into a local term and an integral term

$$\delta_{fs} = (\delta_f - \delta_e) - \left(\frac{1}{\bar{n}_f} \frac{d\bar{n}_f}{d\chi} + \frac{2}{\chi} \right) \int_0^{\chi} d\chi' \delta_e(\hat{n}\chi')$$

$$\begin{aligned}
C_{\ell}^{ss}(\chi, \chi') &= \frac{2}{\pi} \int_0^{\infty} dk k^2 j_{\ell}(k\chi) j_{\ell}(k\chi') P_{[ff+ee-2ef]}(k, (\chi + \chi')/2) \\
&+ \frac{2}{\pi} A(\chi) A(\chi') \int_0^{\chi} d\chi'' \int_0^{\chi'} d\chi''' \int_0^{\infty} dk k^2 j_{\ell}(k\chi'') j_{\ell}(k\chi''') P_{ee}(k, (\chi'' + \chi''')/2) \\
&+ \frac{2}{\pi} A(\chi) \int_0^{\chi} d\chi'' \int_0^{\infty} dk k^2 j_{\ell}(k\chi') j_{\ell}(k\chi'') P_{[ee-fe]}(k, (\chi' + \chi'')/2) \\
&+ \frac{2}{\pi} A(\chi') \int_0^{\chi'} d\chi'' \int_0^{\infty} dk k^2 j_{\ell}(k\chi) j_{\ell}(k\chi'') P_{[ee-fe]}(k, (\chi + \chi'')/2).
\end{aligned}$$

Toy Model

- Assume FRBs and electrons are biased tracers of matter
- $b_e = 1$ $b_f = 1.3$
- Most sources at $z \sim 0.5$

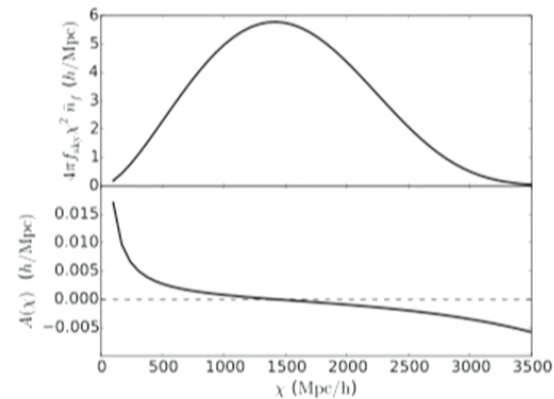
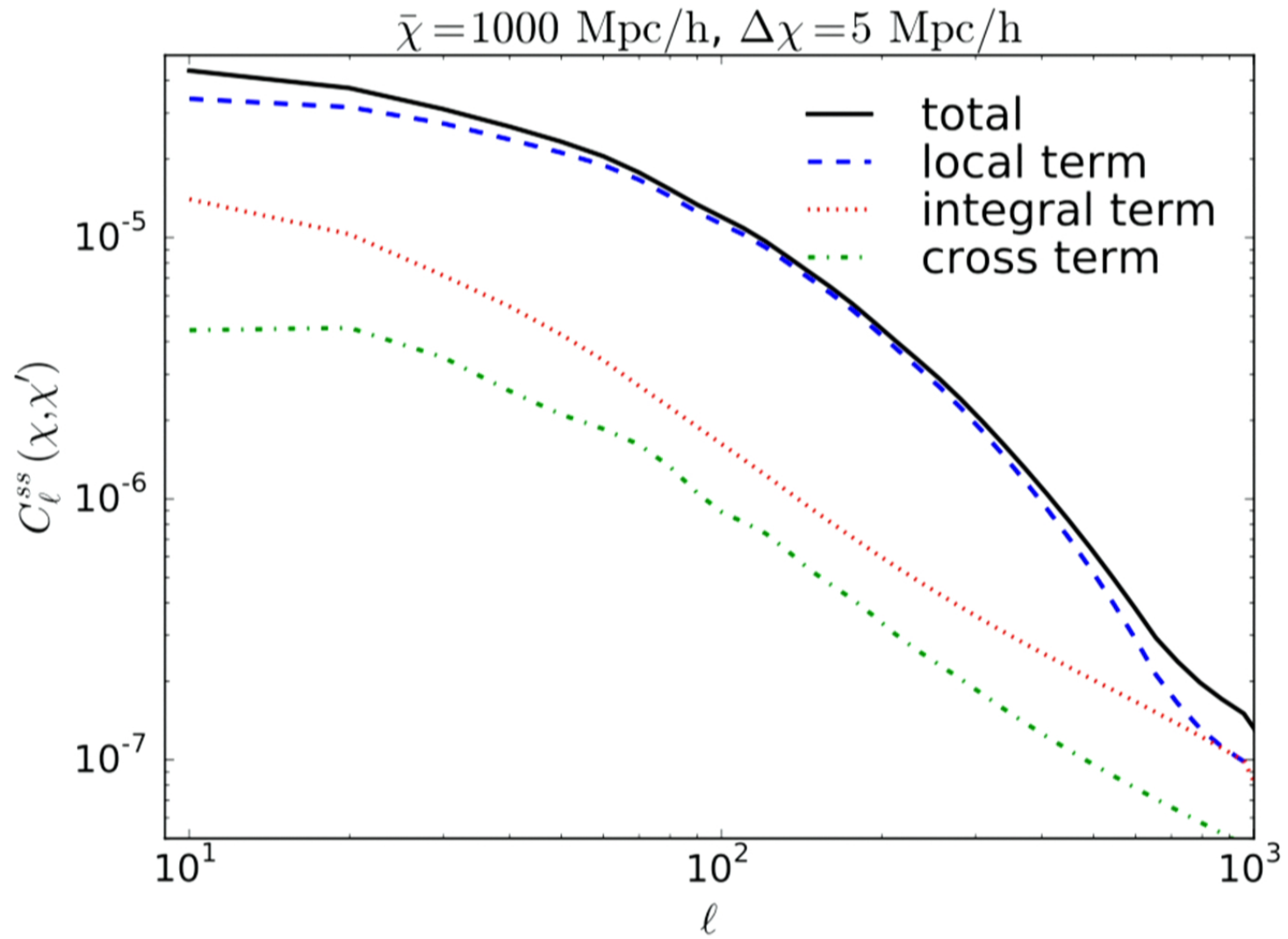
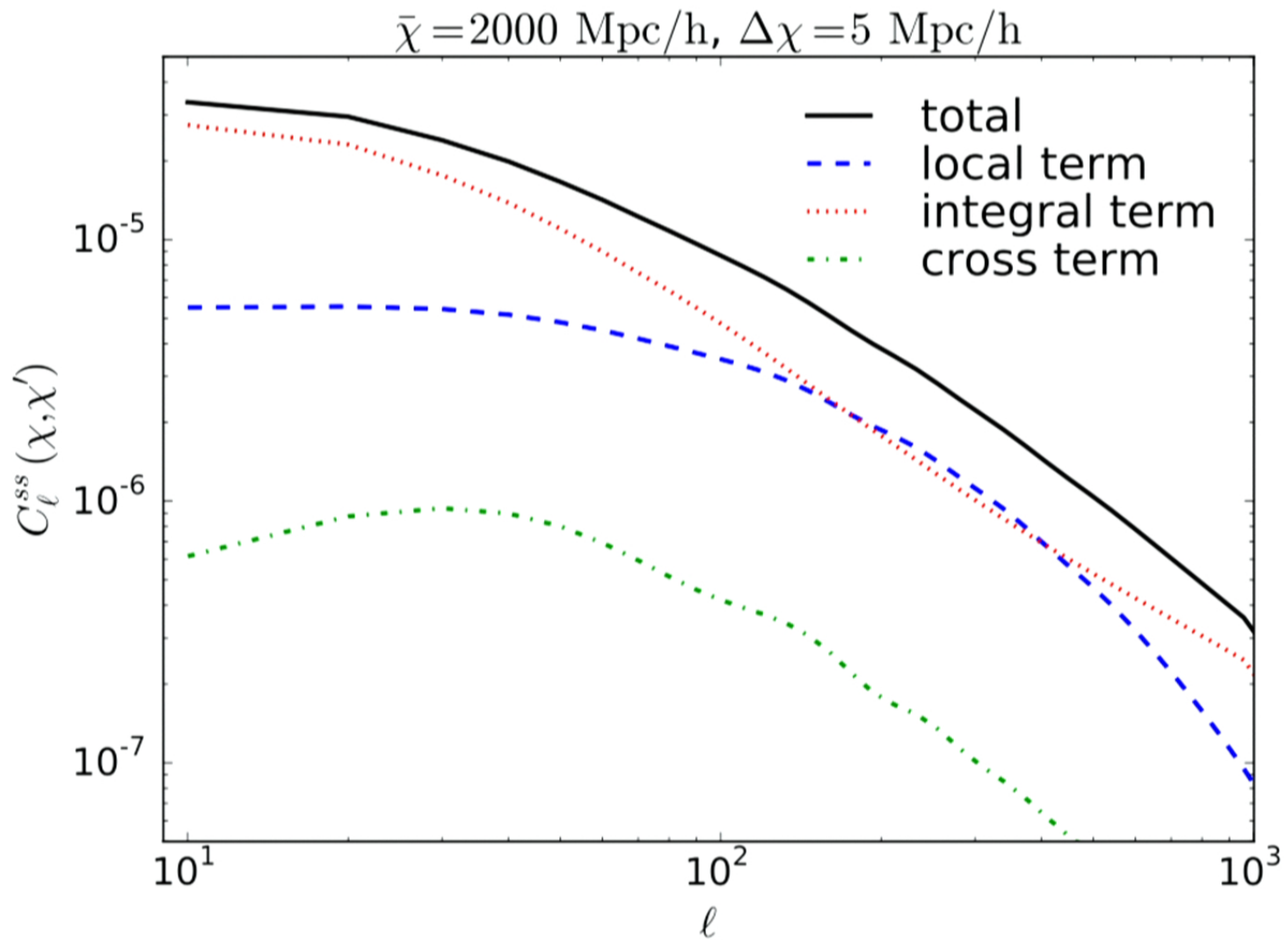


FIG. 1. *Top:* Comoving density of sources in toy model, as described in text. Normalization is such that the survey has a total of 10000 FRB events. *Bottom:* Resulting coefficient $A(\chi)$ as given in Equation 14.





A few comments

- Variations in intrinsic DM will limit slice widths, smooth out some structure
- Unlike lensing, line of sight integral has no weight function: exact tomography

$$\delta_{fs} = (\delta_f - \delta_e) - \left(\frac{1}{\bar{n}_f} \frac{d\bar{n}_f}{d\chi} + \frac{2}{\chi} \right) \int_0^\chi d\chi' \delta_e(\hat{n}\chi')$$

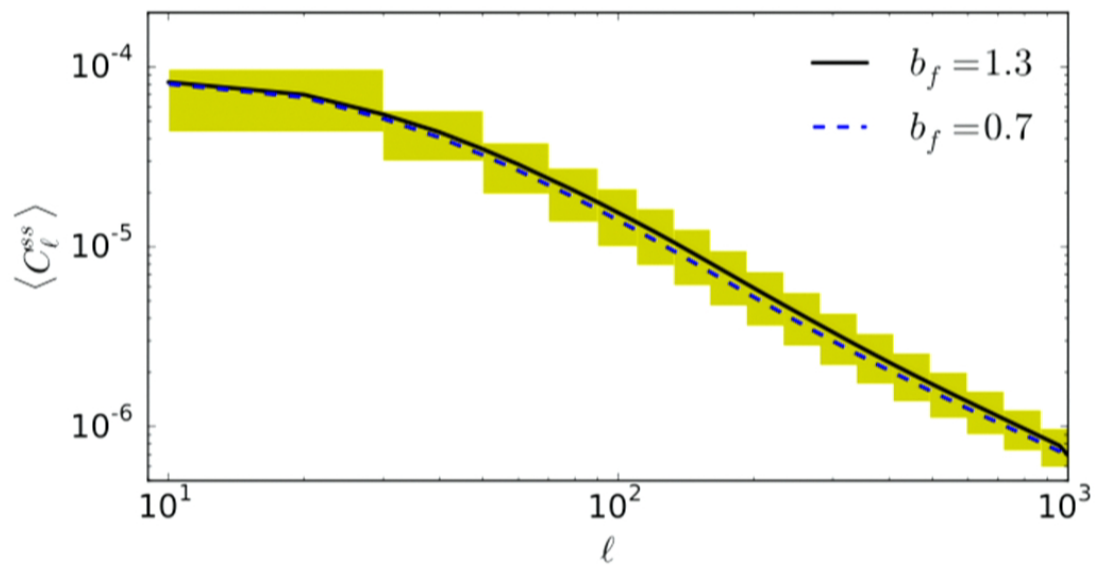


FIG. 3. Sensitivity of a survey with 10000 dispersion measures distributed uniformly over half the sky. Plotted is the cross-correlation power spectrum weighted averaged over all pairs of radial bins. Weights are chosen to maximize signal to noise at $\ell = 100$ for the $b_f = 1.3$ case.

Detectability

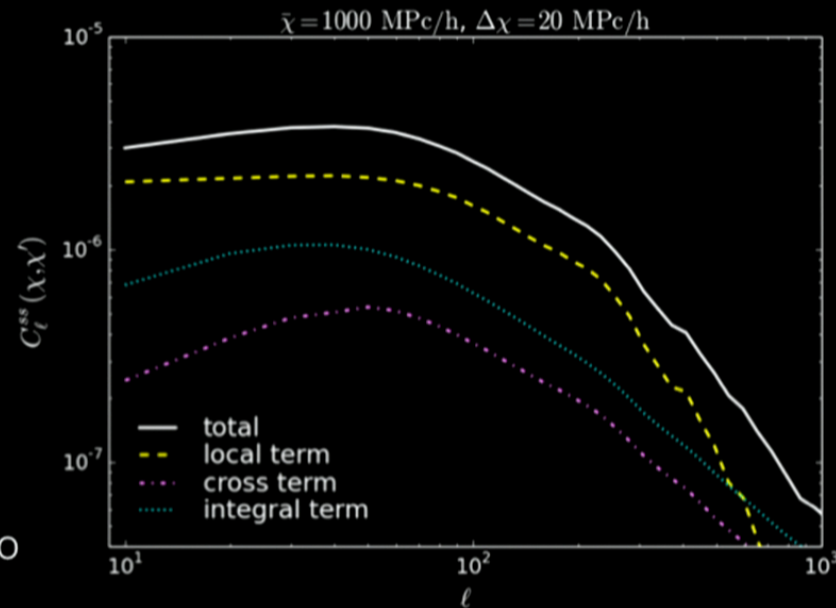
- Full Fisher analysis in progress, but...

$$\ell(\ell + 1)C_\ell^{ss} \sim 10^{-2}$$

- Integral term coherent out to

$$\Delta\chi \sim 1000 \text{ Mpc}/h$$

- Need $\sim 10^4$ events for initial detection
- Local term only coherent out to 20 Mpc/h separations
- Need ~ 50 times more events



Detectability

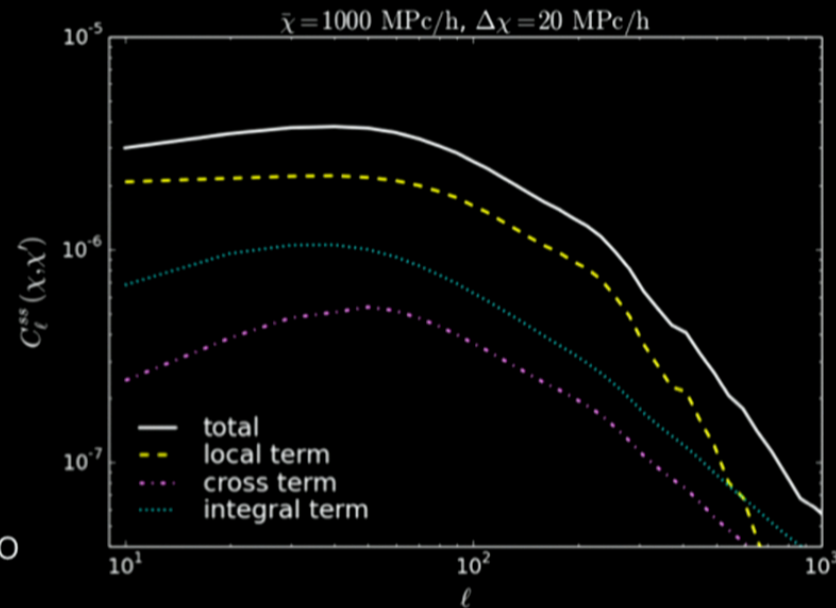
- Full Fisher analysis in progress, but...

$$\ell(\ell + 1)C_\ell^{ss} \sim 10^{-2}$$

- Integral term coherent out to

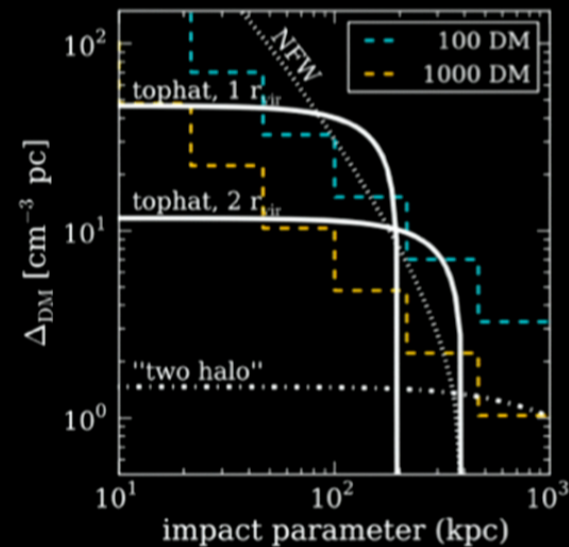
$$\Delta\chi \sim 1000 \text{ Mpc}/h$$

- Need $\sim 10^4$ events for initial detection
- Local term only coherent out to 20 Mpc/h separations
- Need ~ 50 times more events



Applications

- Integral term could be used to find the “missing baryons”, possibly in cross-correlation
- With $\sim 10000/\text{sky}/\text{day}$, could build up LSS survey
- If there are a lot of extra dim events, or spectrum is red...



McQuinn 2013

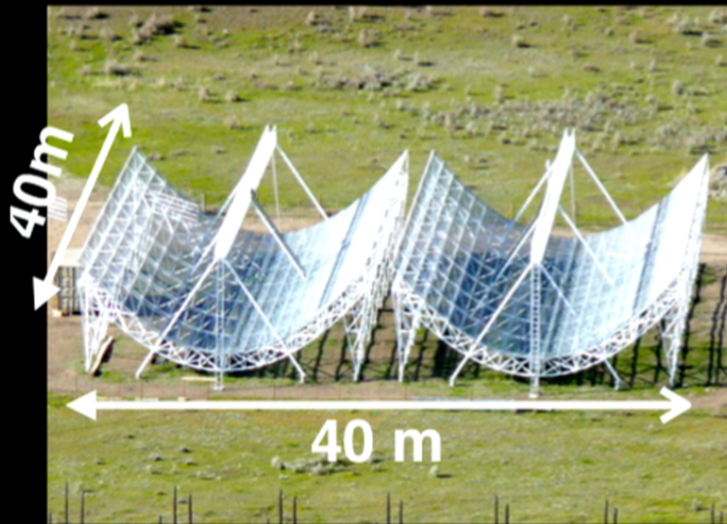
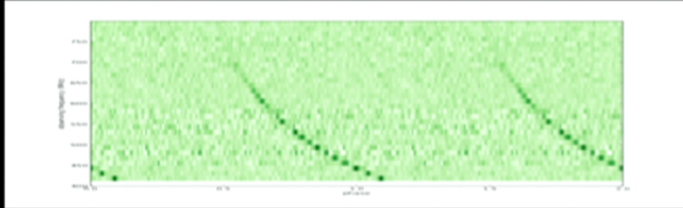
Conclusions

- If FRBs are cosmological then they can use used as 3D tracers of matter
- Have calculated the dispersion measure space distortions
- In toy model, clustering detectable with 10^4 events
- Many opportunities for cross correlation, separation of terms with redshifts, etc.

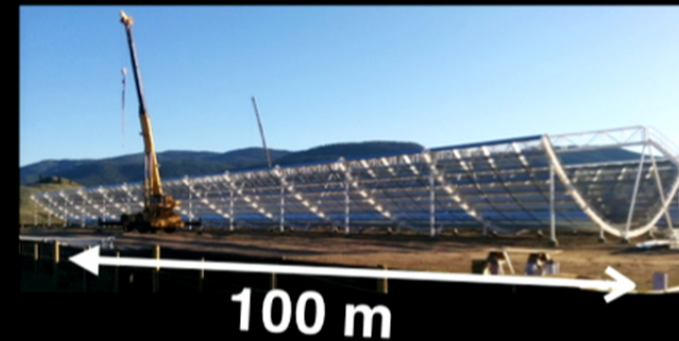
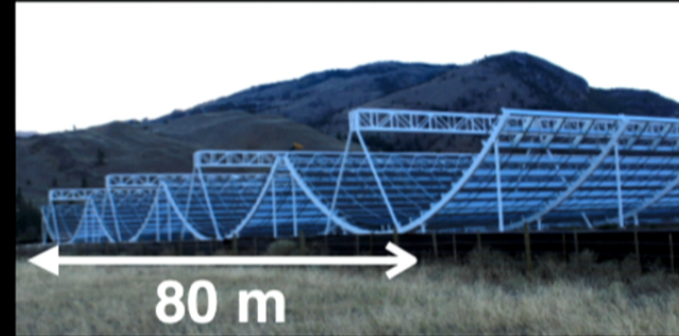
Instrument Status

Pathfinder

Figure: E. Madsen



Full CHIME



FRB sensitivity

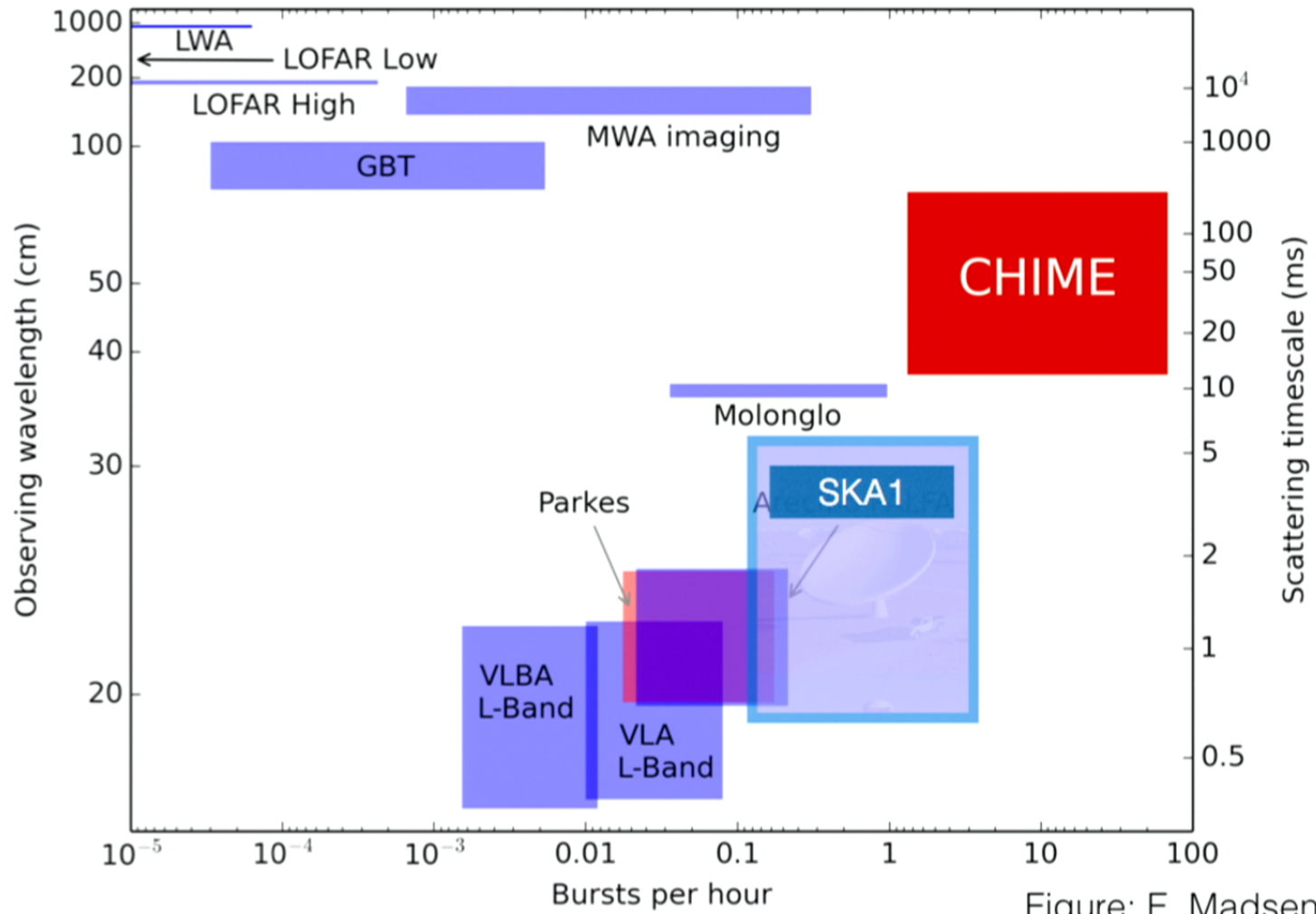


Figure: E. Madsen

FRB sensitivity

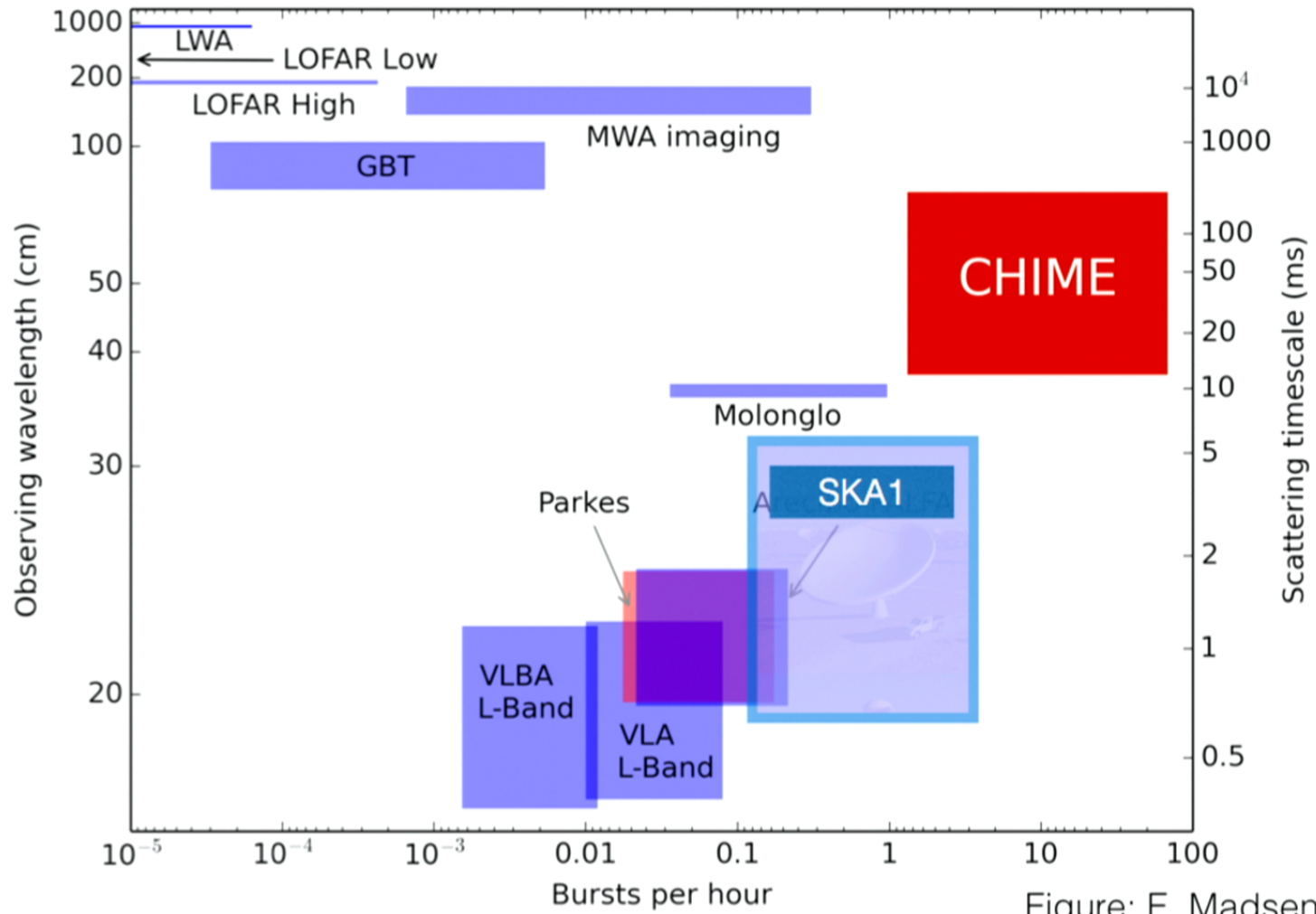


Figure: E. Madsen