

Title: Properties of a non-Abelian chiral spin-1 spin liquid

Date: Jul 28, 2015 11:00 AM

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Abstract: <p>In this talk, we will analyze the properties of the bosonic $\nu = 1$ Moore-Read state when used to build a state
which is strongly believed to be a non-Abelian spin-1 chiral spin liquid state [1]. In this state the bosonic $\nu = 1$
Moore-Read Pfaffian wavefunction is interpreted as a wavefunction for a gas of bosons on a 2D square lattice with one flux quantum per plaquette. We investigate the properties of this wavefunction for the case of planar geometry and for the case of the system living on a torus. For the latter case, there are three distinct states corresponding to the three-fold degeneracy of the $\nu = 1$
bosonic Moore-Read state [2]. Our results show that correlation functions in these states become identical in the limit of large system size. Further issues investigated include the result that the three wavefunctions on the torus are linearly independent and orthogonal in the thermodynamic limit. Additionally, the Renyi entanglement entropy is also calculated for different system partitions in order to extract the topological entanglement entropy γ .

[1] M. Greiter and R. Thomale, Phys. Rev. Lett. 102, 207203 (2009).

[2] S. B. Chung and M. Stone, J. Phys. A: Math. Theor. 40, 4923 (2007).</p>

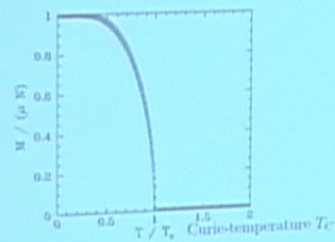
background: topological order I

- classical order

⇒ theoretical framework:
Landau's symmetry breaking

- characterized by:

↳ locally observable order parameter,
e.g. magnetization

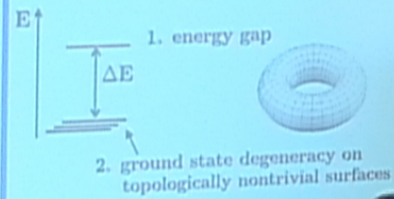


$$F = a|M|^2 + b|\nabla M|^2 + c|M|^4 + \dots$$

- topological order

⇒ no locally
observable symmetry broken

- characterized by:



3. fractionalized quasiparticle excitations with exotic statistics:
Abelian, non-Abelian

background: topological order II

- fractional Quantum Hall effect: incompressible quantum liquid with fractionalized excitations (Laughlin, 1983)
- quantum spin liquids: dimer model on triangular and kagome lattices, \mathbb{Z}_2 quantum spin liquid in kagome (JW and AS), toric code, ...
- specific kind of quantum spin liquids: constructed from quantum Hall-type wave functions
 - ↔ Kalmeyer-Laughlin chiral spin liquid (Abelian)
 - ↔ **Moore-Read chiral spin liquid (non-Abelian)**

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Kalmeyer-Laughlin chiral spin liquid

- precursor: Abelian spin-1/2 chiral spin liquid
Kalmeyer and Laughlin, PRL 59, 2095 (1987)

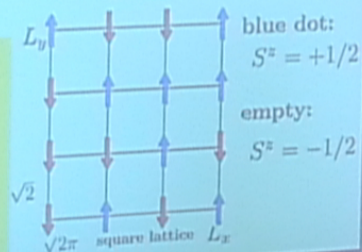
- bosonic FQHE $\nu = 1/2$:

$$\Psi[z_i] = \prod_{i < j} (z_i - z_j)^2 \prod_{i=1}^N e^{-\frac{1}{4}|z_i|^2} \quad \text{Laughlin}$$

- spin- $\frac{1}{2}$ liquid wave function:

$$|\Psi\rangle = \sum_{\{z_1, \dots, z_N\}} \Psi(z_1, \dots, z_N) \prod_{i=1}^N G(z_i) S_{z_i}^+ \dots S_{z_N}^+ |\downarrow \downarrow \dots \downarrow\rangle$$

- lattice system: $z = (n_x + i n_y)\sqrt{2\pi}$
- one flux quantum per plaquette
- singlet in the therm. limit



Kalmeyer-Laughlin chiral spin liquid

- Abelian CSL on the torus: 2-fold ground state degeneracy !

- bosonic FQHE $\nu = 1/2$:

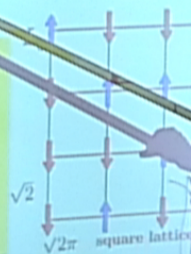
$$\Psi_n[z_i] = \vartheta_1\left[Z - \frac{n}{2}L_x\right]^2 \prod_{i < j}^N \vartheta_1[z_i - z_j]^2 \prod_{i=1}^N e^{-v_i^2/2}$$

$n = 0, 1$

- spin- $\frac{1}{2}$ liquid wave function:

$$|\Psi_n\rangle = \sum_{\{z_1, \dots, z_N\}} \Psi_n(z_1, \dots, z_N) \prod_{i=1}^N G(z_i) S_{z_i}^+ \dots S_{z_N}^+ |\downarrow \downarrow \dots\rangle$$

- lattice system: $z = (n_x + i n_y)\sqrt{2\pi}$
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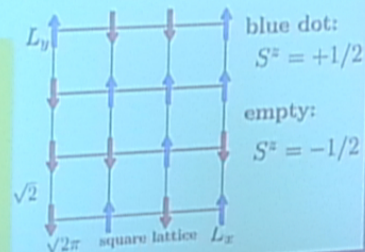
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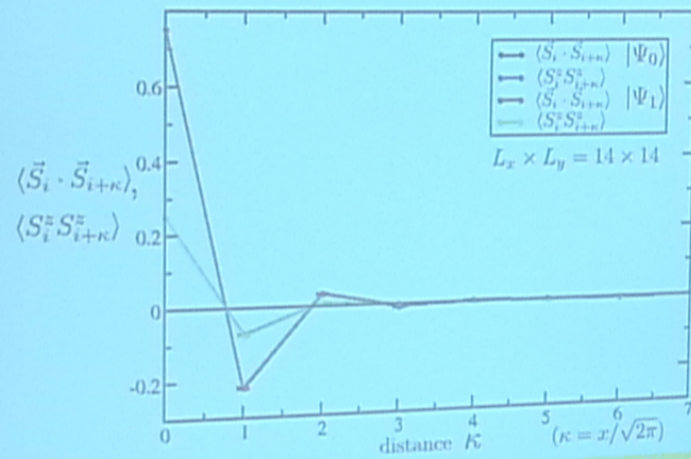
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Kalmeyer-Laughlin CSL: correlations (torus)



- states $|\Psi_0\rangle$ and $|\Psi_1\rangle$ are undistinguishable by their correlations
- $\langle \vec{S}_i \cdot \vec{S}_{i+\alpha} \rangle = 3 \times \langle S_i^\alpha S_{i+\alpha}^\alpha \rangle$, $\alpha = x, y, z$

Non-Abelian chiral spin liquid: droplet

- non-Abelian spin-1 chiral spin liquid for a droplet
 \hookrightarrow Greiter and Thomale, PRL 102, 207203 (2009)

- bosonic FQHE: $\nu = 1$

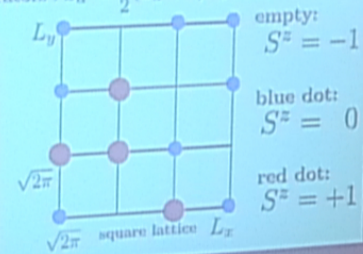
$$\Psi[z_i] = Pf\left(\frac{1}{z_i - z_j}\right) \prod_{i < j}^N (z_i - z_j) \prod_{i=1}^N e^{-\frac{1}{2}|z_i|^2} \quad \text{Moore-Read}$$

- spin-1 liquid wave function:

$$|\Psi\rangle = \sum_{\{z_1, \dots, z_N\}} \Psi(z_1, \dots, z_N) \prod_{i=1}^N G(z_i) \tilde{S}_{z_1}^+ \dots \tilde{S}_{z_N}^+ | -1 \rangle_N$$

'creation operators': $\tilde{S}_n^+ = \frac{1}{2}(S_n^x + 1)S_n^+$

- lattice system: $z = (n_x + i n_y)\sqrt{2\pi}$
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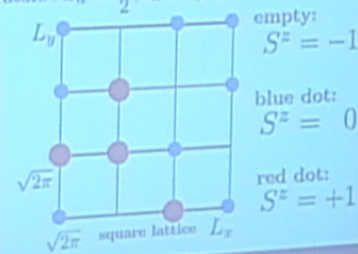
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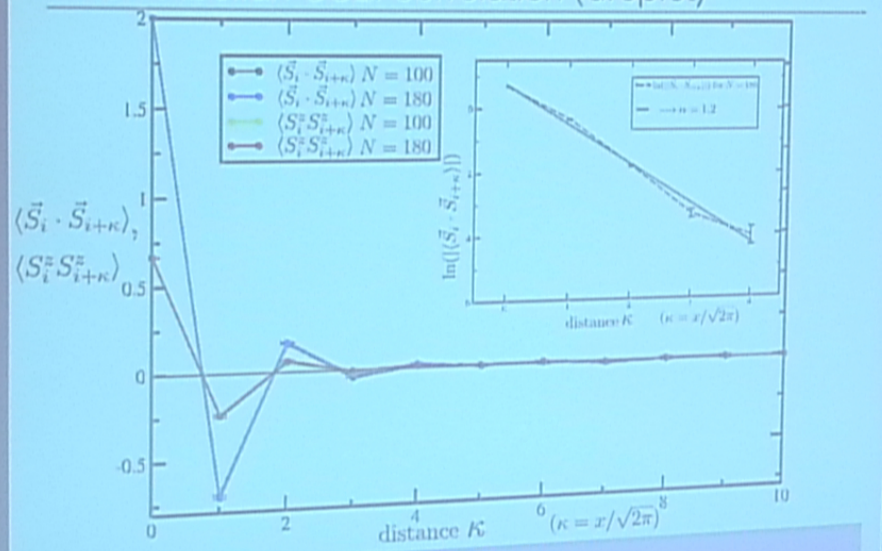
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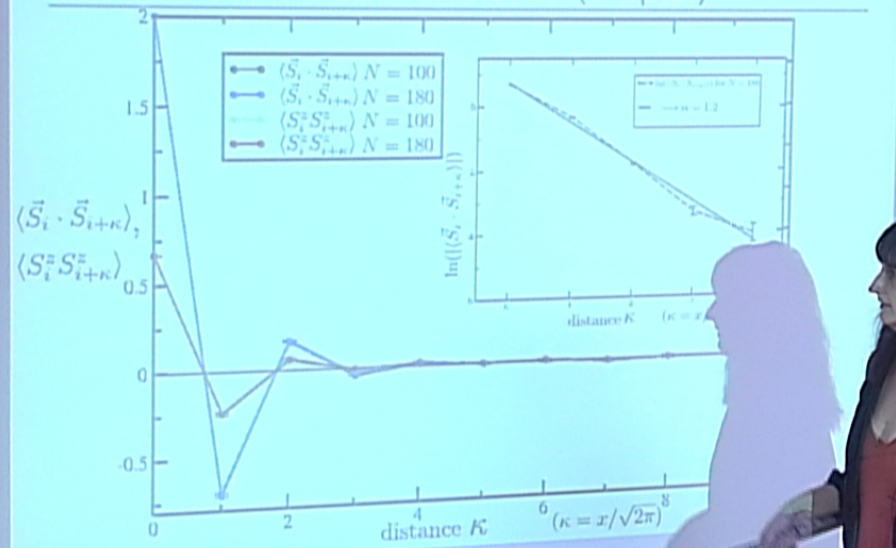
Non-Abelian CSL: correlation (droplet)



→ exponential fall-off

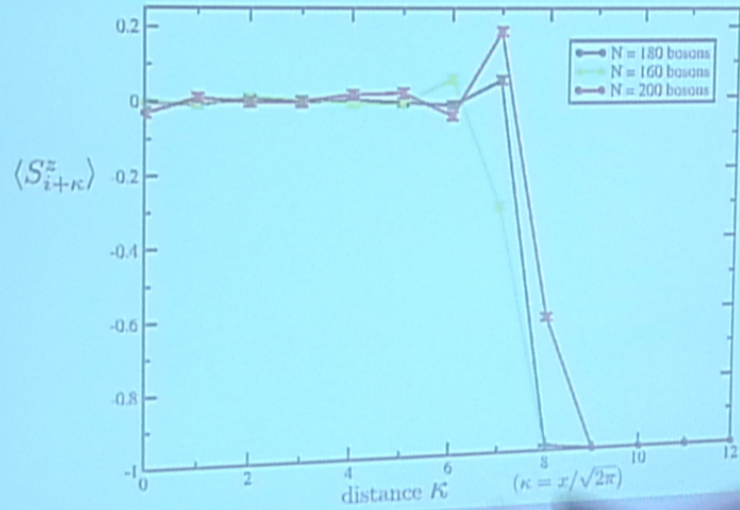


Non-Abelian CSL: correlation (droplet)



→ exponential fall-off

Non-Abelian CSL: correlation (droplet)



→ no singlet: $\langle S_{i+\bar{\kappa}}^z \rangle \neq 0$

Non-Abelian CSL: torus

- N sites on a square lattice for $S = 1$ antiferromagnet
- non-Abelian chiral spin-1 liquid on a torus

- bosonic FQHE $\nu = 1$

$$\Psi_n(z_1, \dots, z_N) = \vartheta_{n+1}(Z) \text{Pf} \left(\frac{\vartheta_{n+1}(z_i - z_j)}{\vartheta_1(z_i - z_j)} \right) \prod_{i < j}^N \vartheta_1(z_i - z_j) \prod_{i=1}^N e^{-y_i^2/2}$$

3-fold degeneracy

Wu, E. Rezayi, Phys. Rev. B 54, 16864 (1996)



$n = 1, 2, 3 \rightarrow \vartheta$ -functions $\vartheta_2(z_i - z_j), \vartheta_3(z_i - z_j), \vartheta_4(z_i - z_j)$

- one flux quantum Φ_0 per plaquette
- exact singlet !

Ground states:

\hookrightarrow 3-fold degeneracy on torus

\Rightarrow spin-1 liquid wave functions: $n = 1, 2, 3$

$$|\Psi_n\rangle = \sum_{\{z_1, \dots, z_N\}} \Psi_n(z_1, \dots, z_N) \prod_{i=1}^N G(z_i) \tilde{S}_{z_1}^+ \dots \tilde{S}_{z_N}^+ | -1 \rangle_N$$

Non-Abelian CSL: overlap

- linear independence of 3 ground states on the square lattice
→ 3x3 overlap matrix corresponding to the 3 ground states

$$\begin{pmatrix} \langle \Psi_1 | \Psi_1 \rangle & \langle \Psi_1 | \Psi_2 \rangle & \langle \Psi_1 | \Psi_3 \rangle \\ \langle \Psi_2 | \Psi_1 \rangle & \langle \Psi_2 | \Psi_2 \rangle & \langle \Psi_2 | \Psi_3 \rangle \\ \langle \Psi_3 | \Psi_1 \rangle & \langle \Psi_3 | \Psi_2 \rangle & \langle \Psi_3 | \Psi_3 \rangle \end{pmatrix} \implies \text{rank of matrix: } r = 3 \\ \text{ground states linearly independent}$$

discussion:

- $L_x = L_y : \langle \Psi_1 | \Psi_2 \rangle = \langle \Psi_2 | \Psi_3 \rangle$ ($\frac{\pi}{2}$ - rotation symmetry)
- $N \rightarrow \infty : \text{off-diagonal elements} \rightarrow 0$
- ground states orthogonal

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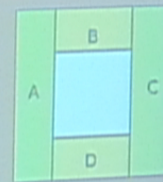
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Entanglement entropy in spin liquids

- reduced density matrix: $\rho_A = \text{Tr}_B |\Psi\rangle\langle\Psi|$
- von Neumann entropy: $S = -\text{Tr}(\rho_A \ln \rho_A)$
- Renyi entropy of order n: $S_n = -\frac{1}{n-1} \ln \text{Tr}(\rho_A^n)$
 $n = 2: S_2 = -\ln \text{Tr}(\rho_A^2)$
 $= -\ln(\langle \text{SWAP} \rangle)$
- area law: $S = \alpha L - \gamma$
- total quantum dimension: $\mathcal{D} = \sqrt{\sum_{\alpha} d_{\alpha}^2}$
- topological entropy: $-\gamma \rightarrow \gamma = \ln(\mathcal{D})$
- Moore-Read non-Abelian spin liquid: $\gamma = \ln(2) \approx 0.69$
- 3-fold degenerate Abelian spin liquid: $\gamma = \frac{1}{2} \ln(3) \approx 0.55$



cylinder

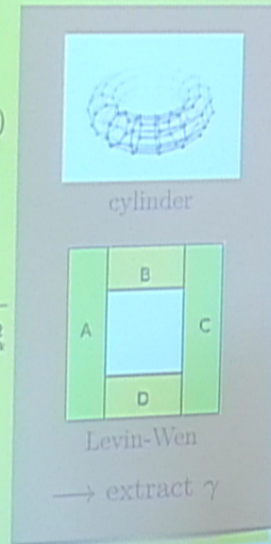


Levin-Wen

→ extract γ

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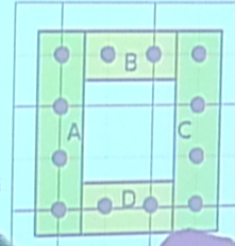


Entanglement entropy: Levin-Wen

$$\rightarrow -2\gamma = (S_{ABCD} - S_{ADC}) - (S_{ABC} - S_{AC})$$

→ convergence issue: re-weighting scheme

$$\hookrightarrow \langle \text{SWAP} \rangle = \langle \text{SWAP}_{\text{sign}} \rangle \times \prod_i \langle \text{SWAP}_{\text{amp}} \rangle_i$$



- theoretical value: $\gamma = -\ln(2) \approx -0.69$
- theoretical value for an Abelian 3-fold degenerate ground state: $\gamma = \ln(3) \approx 1.10$

- 6×6 -system: $\gamma = -1.12 \pm 0.18$
- bigger systems: $8 \times 8, 8 \times 10, 10 \times 10, \dots$

→ $\gamma < 0$: topologically ordered phase !

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- theoretical value: $\gamma = -\ln(2) \approx -0.693$
- theoretical value for an Abelian 3-fold degenerate ground state: $\gamma = -\ln(3) \approx -1.099$
- 6×6 -system: $\gamma = -1.12 \pm 0.18$
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→ $\gamma < 0$: topologically ordered phase !

