

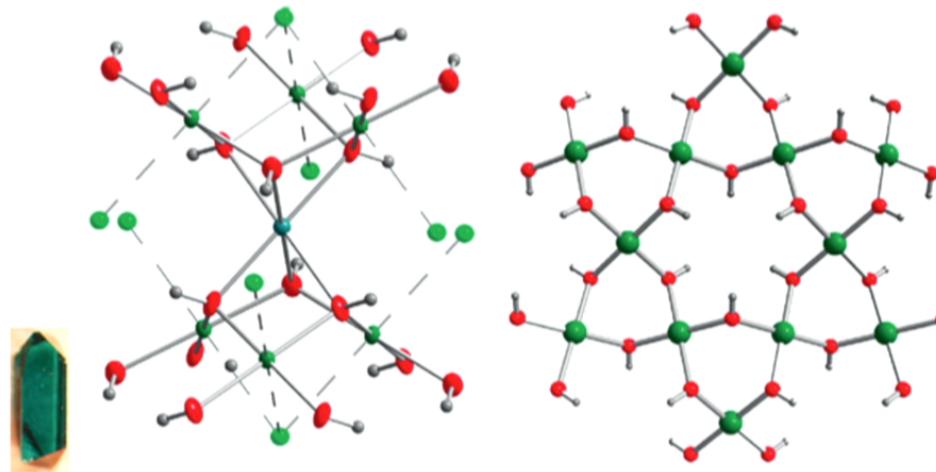
Title: Fractional spin-wave continuum in spin liquid states on the kagome lattice

Date: Jul 10, 2015 11:00 AM

URL: <http://pirsa.org/15070063>

Abstract: Motivated by spin-wave continuum (SWC) observed in recent neutron scattering experiments in Herbertsmithite, we use Gutzwiller-projected wave functions to study dynamic spin structure factor of spin liquid states on the kagome lattice. As their ground state, spin-1 excited states for spin liquids are represented by Gutzwiller-projected two-spinon excited wave functions. We investigate three different spin liquid candidates, spinon Fermi-surface spin liquid (FSL), Dirac spin liquid (DSL) and random-flux spin liquid (RSL). We find that DSL has no explicit contradiction with experiments. Besides a fractionalized spin moment, DSL has a fractionalized crystal momentum which is also detectable directly in the neutron scattering measurements.

Recent progress on high T<sub>c</sub> superconductivity and **related problems**



# Fractional spin-wave continuum in kagome spin liquid states

Jia-Wei Mei

JWM and X.G. Wen, coming soon ....

Z. Liu, X.L. Zou, JWM and F. Liu, arXiv: 1504.00521



## Outline:

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- Motivation  
quantum spin liquid is a key ingredient of high T<sub>c</sub> superconductivity
- Fractional spin-wave continuum in kagome spin liquid states  
Neutron scattering experiments, spin fractionalization, crystal momentum fractionalization
- New kagome material  
Selective doping Barlowite ( $\text{Cu}_4(\text{OH})_6\text{FBr}$ ), LDA,  $\text{Cu}_3\text{Mg}(\text{OH})_6\text{FBr}$ ,  $\text{Cu}_3\text{Zn}(\text{OH})_6\text{FBr}$

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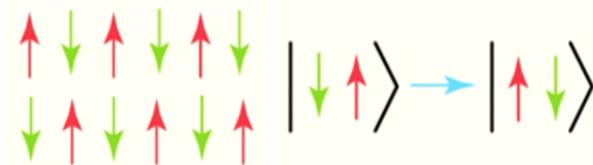


Louis Néel

Quantum AF ( $J > 0$ ) spin model

$$H = J \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j$$

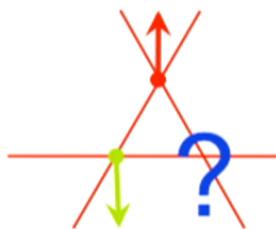
## An End to the Drought of Quantum Spin Liquids -- P.A. Lee



**Ordered spins.** (Left) Néel's picture of antiferromagnet ordering with an alternate spin-up–spin-down pattern across the lattice. (Right) Quantum fluctuations lead to mutual spin flips, which Landau argued would disorder Néel's state.



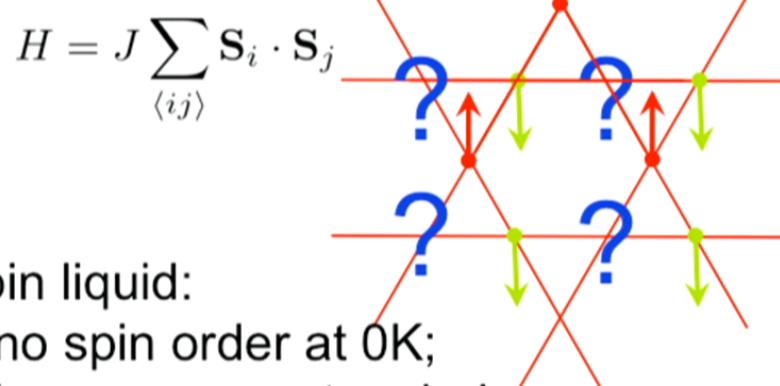
Lev Landau





Louis Néel

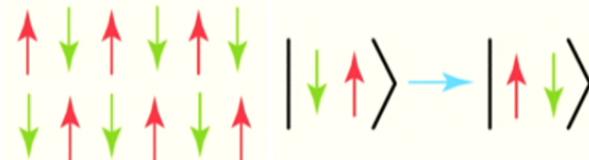
Quantum AF ( $J > 0$ ) spin model



Spin liquid:

- no spin order at 0K;
- Long range entangled;
- Symmetry fractionalization;

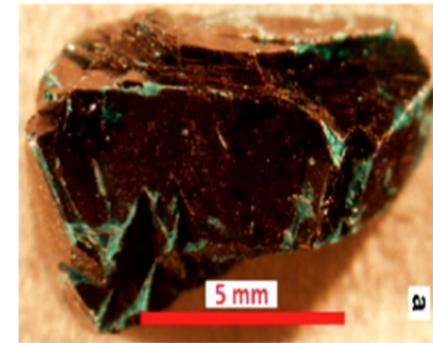
### An End to the Drought of Quantum Spin Liquids -- P.A. Lee, Science (2008)



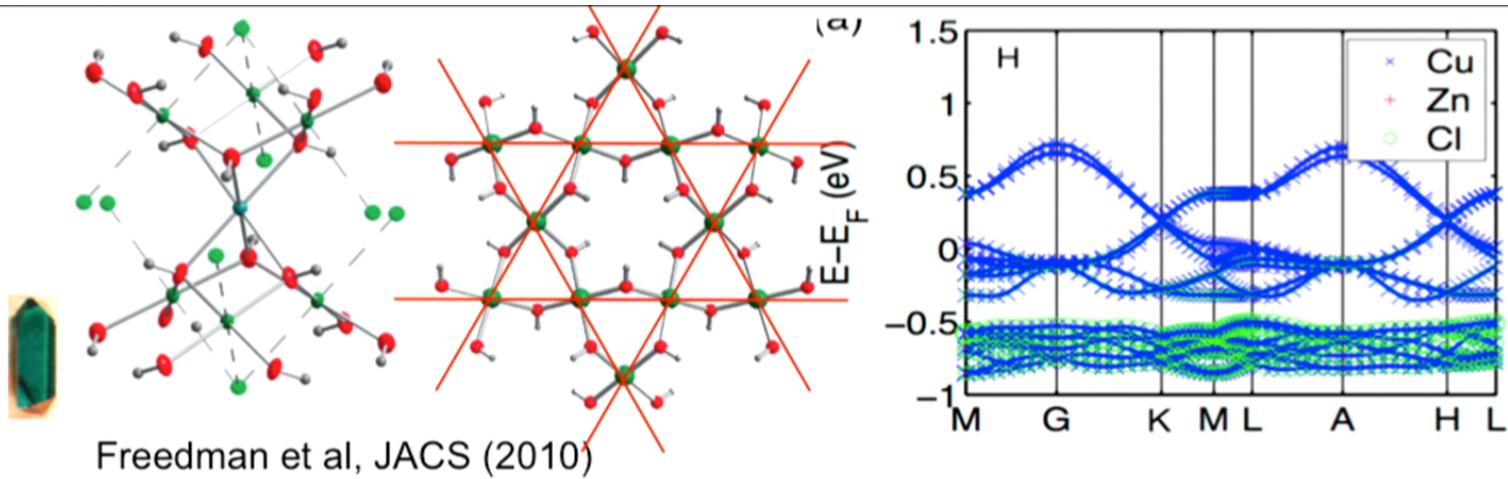
Ordered spins. (Left) Néel's picture of antiferromagnet ordering with an alternate spin-up–spin-down pattern across the lattice. (Right) Quantum fluctuations lead to mutual spin flips, which Landau argued would disorder Néel's state.



Lev Landau



Herbertsmithite (Y. Lee group)



$$H = J \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j$$

Spin wave, Slave particle mean field + VMC, MERA, DMRG, PEPS, ED ...

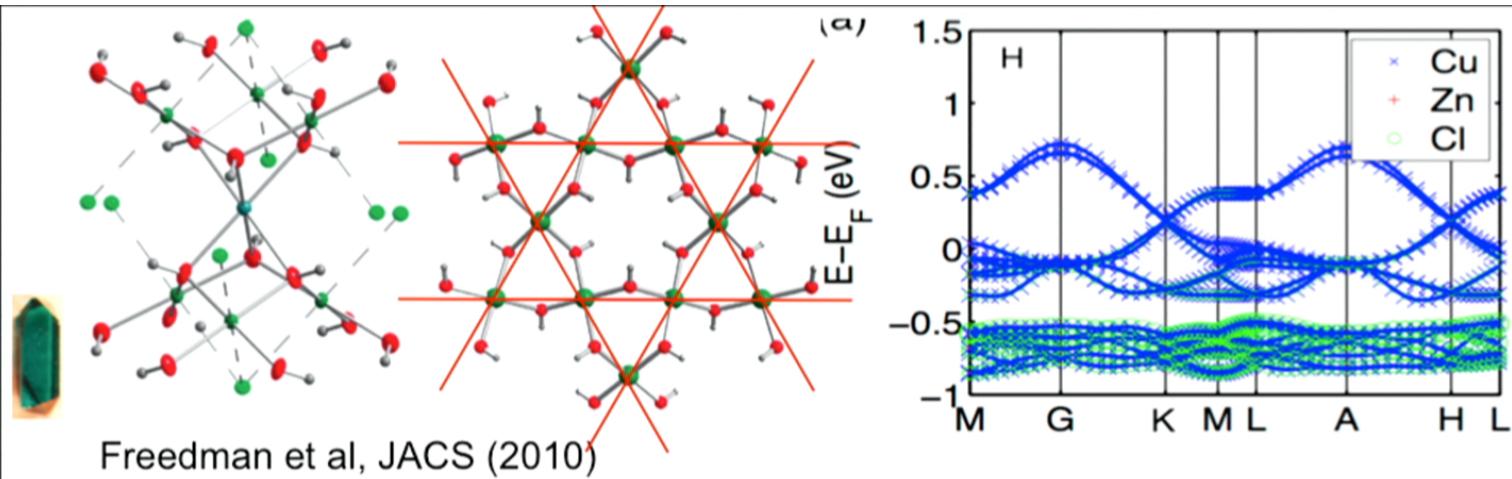
Possible ground state candidates: valence-bond solids, gapless or gapped QSLs ...

DMRG: Z2 spin liquid with small singlet gap  $\Delta_s \approx 0.05J$  and triplet gap  $\Delta_t \approx 0.1J$ .

S. Yan et al, Science (2011)

VMC (Gutzwiller-projected wave functions): U(1) Dirac spin liquid

Y. Ran et al, PRL (2007)



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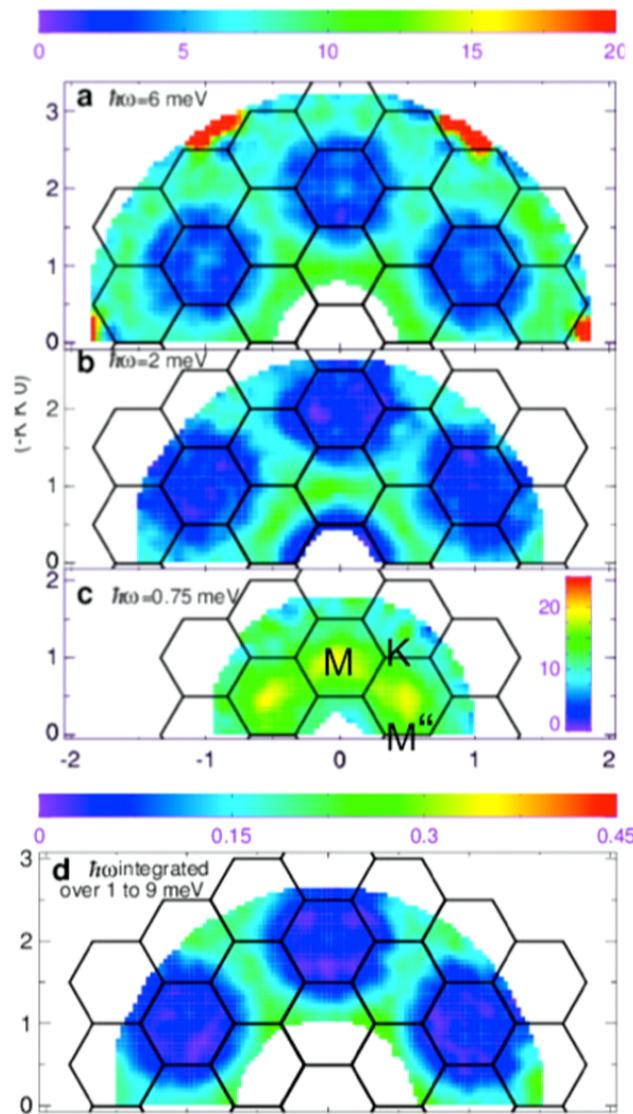
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Properties in neutron scattering:  $J \sim 17$  meV

- It has a spin-wave continuum (SWC) spectrum over a large momentum and energy;
- The magnetic intensity is low in the elementary BZ and high in the extended region of magnetic BZ;
- The integrated intensity up to 11 meV only accounts for 20% of total spin moment and the SWC extends up to  $2 \sim 3J$ ;
- The magnetic intensity is almost energy independent between 1.5 meV and 11 meV;
- Magnetic intensity is quite flat in momentum and has no sharp dispersing boundary edge;
- The low energy pattern of intensity varies in momentum and low energy peaks (at M) are connected through  $M''$ ;

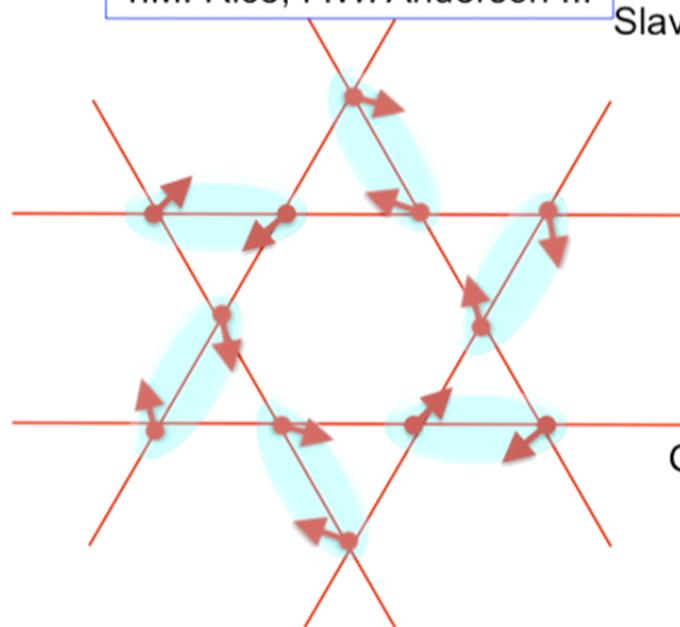


# Gutzwiller-Projected Wave Function (GPWF)

Quantum AF ( $J > 0$ ) spin model:

$$H = J \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j$$

T.M. Rice, P.W. Anderson ...



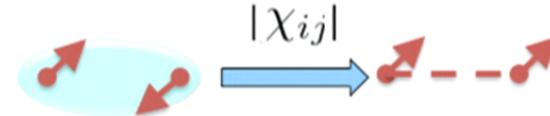
Slave particle decomposition:

$$S_i^a = \frac{1}{2} \sum_{\alpha\beta} f_{i\alpha}^\dagger \sigma^a f_{i\beta}$$

Mean field Hamiltonian:

$$H_{\text{MF}} = - \sum_{\langle ij \rangle} (\chi_{ij} f_{i\sigma}^\dagger f_{j\sigma} + \text{H.C.})$$

Ground state:  $|\Psi\rangle = \mathcal{P}_G |\Psi_{\text{MF}}^{\chi_{ij}}\rangle$



Spin-1 excited states:

$$|\Psi_{ij}^{S=1}\rangle = \mathcal{P}_G f_{e_i\uparrow}^\dagger f_{e_j\downarrow} |\Psi_{\text{MF}}^{\chi_{ij}}\rangle$$

# Gutzwiller-Projected Wave Function (GPWF)

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$$H = J \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j$$

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*Ground-state Assumption*

$$|\Psi_{ij}^{S=1}\rangle = \mathcal{P}_G f_{e_i \uparrow}^\dagger f_{e_j \downarrow} |\Psi_{\text{MF}}^{\chi_{ij}}\rangle$$

*Excited-state Assumption*

The projected Hamiltonian system

$$\{\mathbb{H}, \mathbb{O}\}$$

Tao Li and Fan Yang, PRB(2010)

$$\mathbb{H}(i'j', ij) = \langle i'j'|H|ij\rangle, \quad \mathbb{O}(i'j', ij) = \langle i'j'|ij\rangle$$

MC strategy: single Markov chain for spin configuration samplings.

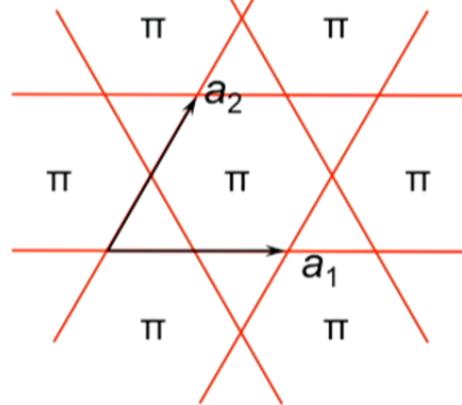
$$S(\mathbf{q}, i\omega_n) = \int_0^\beta d\tau e^{i\omega_n \tau} \frac{1}{N} \sum_{ij} e^{i\mathbf{q} \cdot \mathbf{r}_{ij}} \langle T_\tau S_i^-(\tau) S_j^+(0) \rangle_0$$

Since spin operator commute with Gutzwiller projection operator,

$$S(\mathbf{q}, \omega) = \sum_n \delta(\omega - (\epsilon_n - \epsilon_0)) |\langle \phi_n | S_{\mathbf{q}}^+ \mathcal{P}_G |\Psi_{\text{MF}}^{\chi_{ij}}\rangle|^2 \quad \mathbb{H}|\phi_n\rangle = \epsilon_n \mathbb{O}|\phi_n\rangle$$

## Best trial GPWF: Dirac spin liquid (DSL)

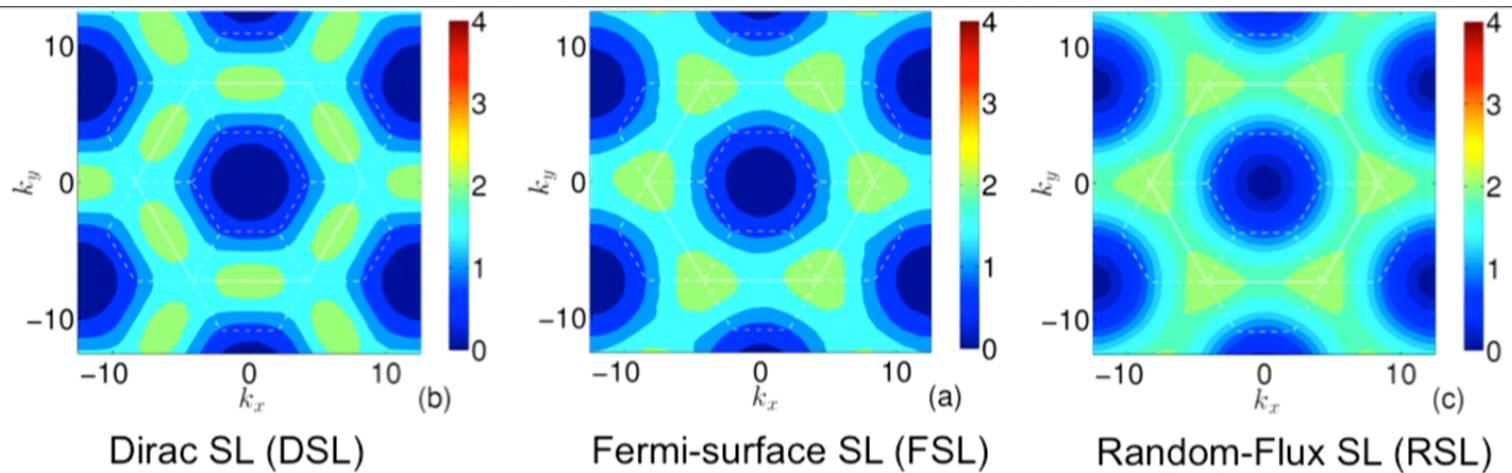
"Q1=Q2 states", Sachdev (1992)



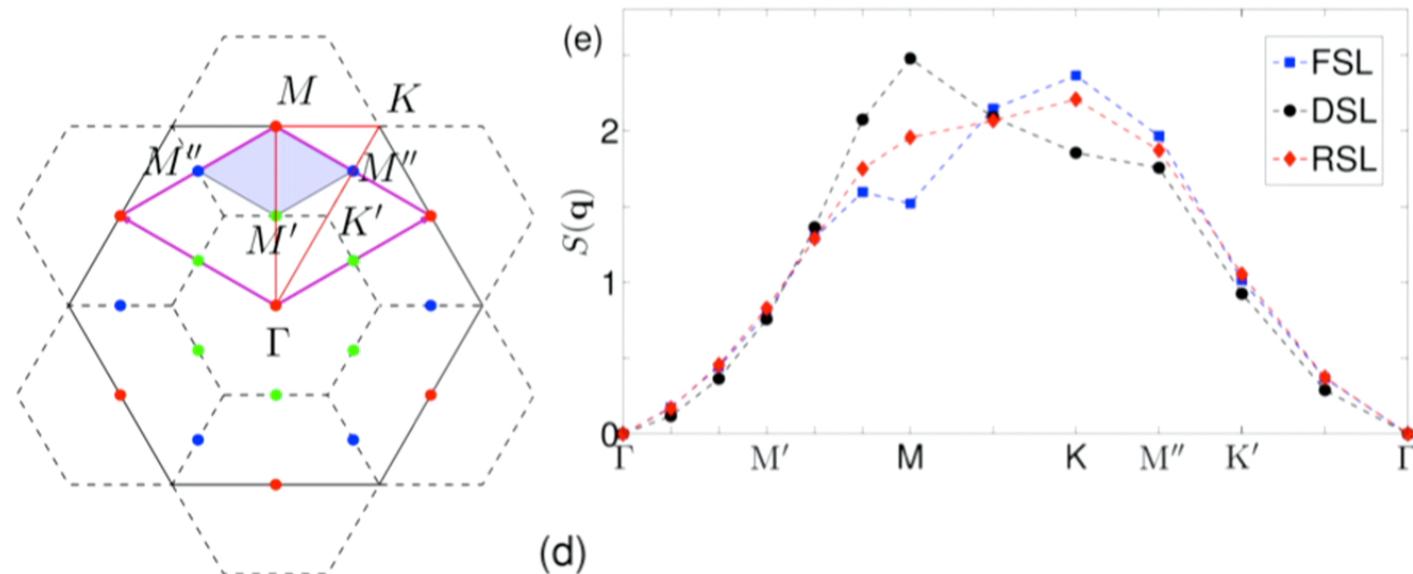
$$|\Psi\rangle = \mathcal{P}_G |\Psi_{\text{MF}}^{\chi_{ij}}\rangle$$

Crystal momentum fractionalization:

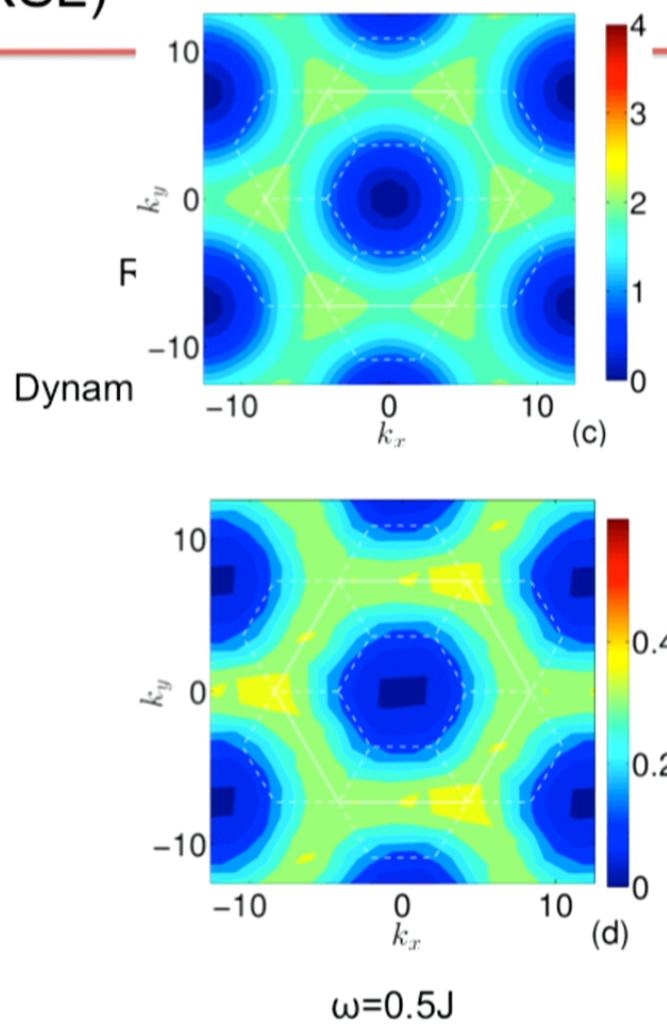
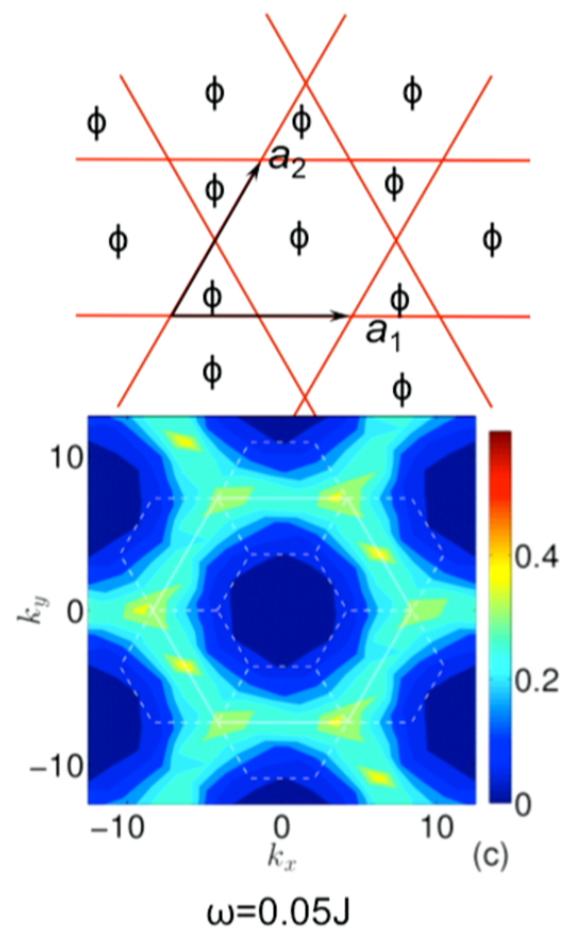
$$T_1 T_2 = -T_2 T_1$$



Equal-time spin structure factor  $S(\mathbf{q})$

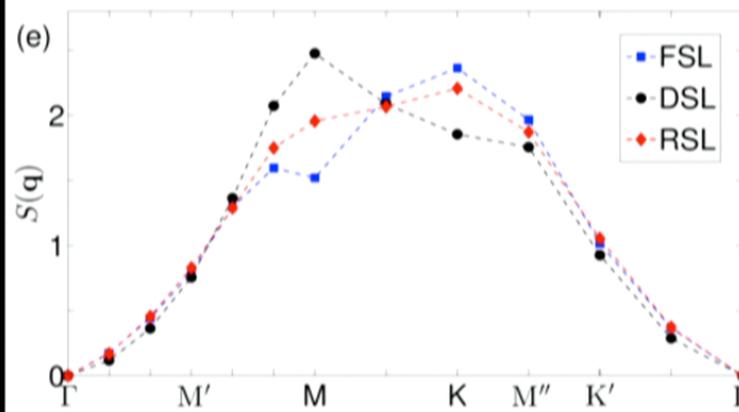


## Random-flux spin liquid (RSL)

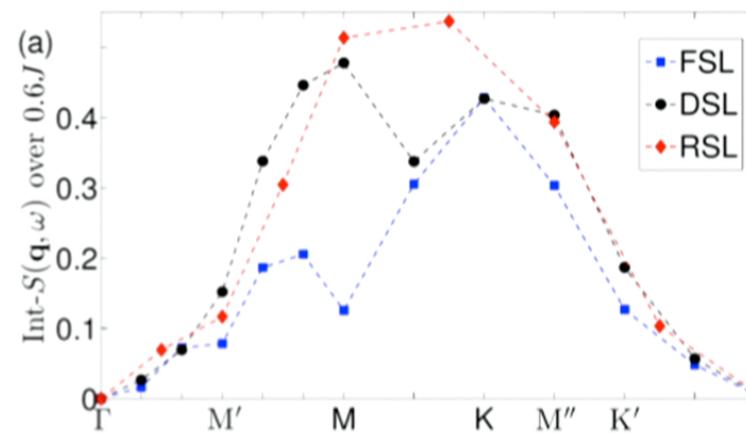


## $S(q, \omega)$ along high symmetry directions

Full  $\omega$ -integrated  $S(q, \omega) = S(q)$

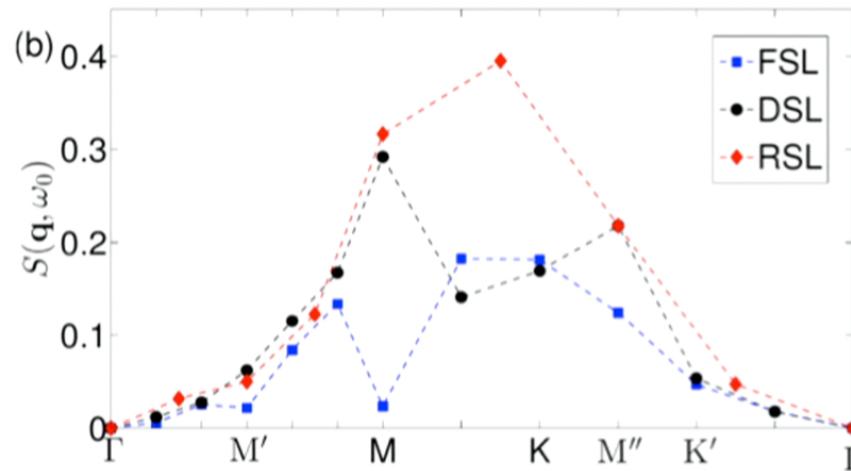


Partial  $\omega$ -integrated  $S(q, \omega)$  from 0 to  $0.6J$



## $S(q,\omega)$ along high symmetry directions

Low energy  $\omega_0$  magnetic intensity



- FSL has a gap at M points;
- FSL and RSL have low energy peaks around K points
- DSL has low energy peaks at M and M“ points.

Properties in neutron scattering:  $J \sim 17$  meV

- It has a spin-wave continuum (SWC) spectrum over a large momentum and energy;
- The magnetic intensity is low in the elementary BZ and high in the extended region of magnetic BZ; short ranged spin liquid (RVB);
- The integrated intensity up to 11 meV only accounts for 20% of total spin moment and the SWC extends up to  $2\sim 3J$ ;
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- Magnetic intensity is quite flat in momentum and has no sharp dispersing boundary edge;
- The low energy pattern of intensity varies in momentum and low energy peaks (at M) are connected through M'';

DSL	FSL	RSL
✓	✓	✓
✓	✓	✓
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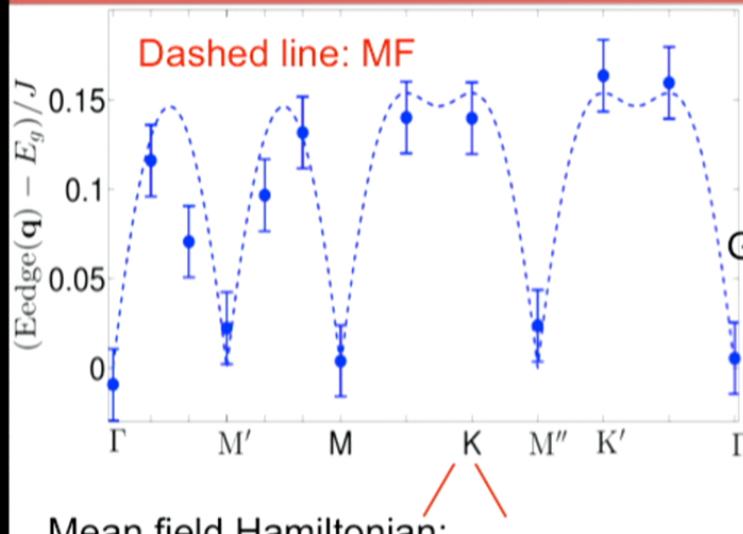
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	?✓	?✓	?✓
	?✓	? ✗	?✓
	✓	✗	✗

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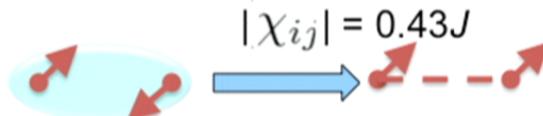
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General eigen equation:

$$\mathbb{H}|\phi_n\rangle = \epsilon_n \mathbb{O}|\phi_n\rangle$$

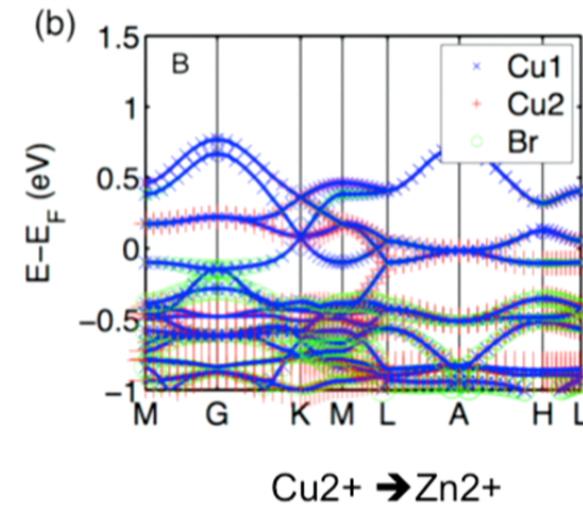
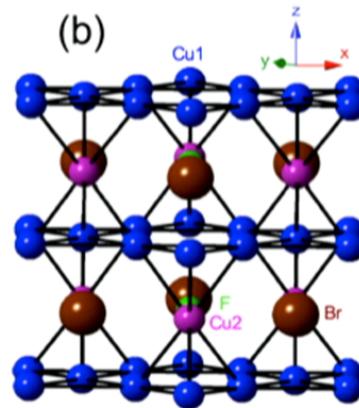
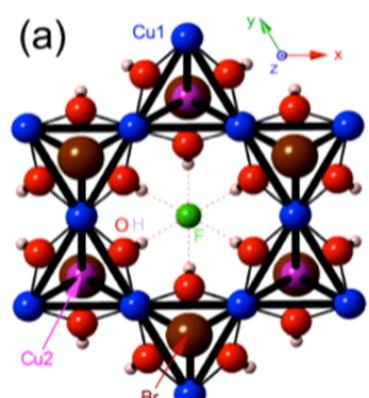


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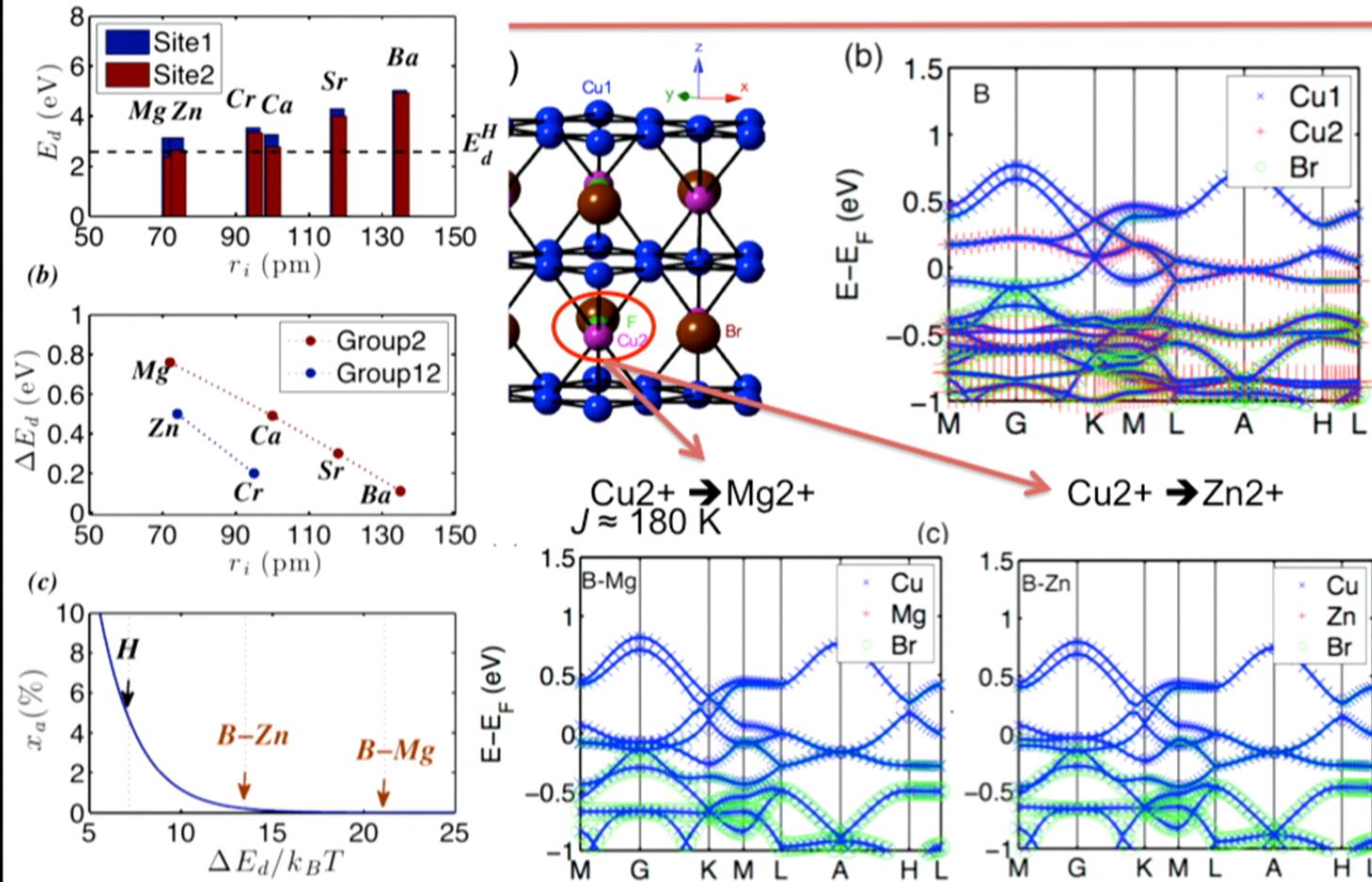
# New kagome antiferromagnet

arXiv: 1504.00521



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## Summary:

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- Dirac spin liquid is very likely in Herbersmithite. A crystal momentum fractionalization is also potential detectable in neutron scattering experiments.
- Selective doping Barlowite varieties ( $\text{Cu}_3\text{Mg}(\text{OH})_6\text{FBr}$  and  $\text{Cu}_3\text{Zn}(\text{OH})_6\text{FBr}$ ) are promising for new kagome materials with much less imperfection.