

Title: Cosmology & Observations: Observables in an expanding universe

Date: Jul 13, 2015 11:45 AM

URL: <http://pirsa.org/15070038>

Abstract:

Plan

- Getting comfortable in the expanding universe
- Cosmic microwave background fluctuations
- Clustering of galaxies and large scale structure
- Gravitational lensing

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$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot \vec{v} = 0$$

$$\frac{\partial \rho}{\partial t} + (\vec{v} \cdot \vec{\nabla}) \rho = -\frac{1}{P} \vec{v} \cdot \vec{\nabla} P$$

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$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot \vec{v} = 0$$

$$\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \vec{\nabla}) \vec{v} = -\frac{1}{\rho} \vec{\nabla} \rho - \nabla \phi$$

$$\nabla^2 \phi = 4\pi G\rho$$

* $\rho = \text{constant in space}$



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$$\frac{d\vec{v}}{dt} + (\vec{v} \cdot \vec{\nabla})\vec{v} = -\frac{1}{\rho} \vec{\nabla} p - \nabla \phi$$

$$\nabla^2 \phi$$

$G\rho$

$(v \cdot \nabla)v$

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot \vec{v} = 0$$

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* $\rho = \text{constant in space}$
 $p = 0$ (for simplicity)



Sph. symmetry

$$\frac{\partial \ln \rho}{\partial t} = -\frac{1}{r} \frac{\partial}{\partial r} (r^2 v)$$

$$-v(r) = -\frac{1}{3} \left(\frac{\partial \ln \rho}{\partial t} \right) r$$

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$$\frac{\partial \rho}{\partial t} = -\frac{1}{r} \frac{\partial}{\partial r} (\rho v)$$

$$-\dot{\rho} = -\frac{1}{3} \left(\frac{\partial \rho}{\partial t} \right) r$$

$$\rho = \frac{1}{6\pi G} t^{-2} \Rightarrow \frac{\dot{\rho}}{\rho} = -\frac{2}{t}$$

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$$\frac{d \ln \rho}{dt} = -\frac{1}{r} \frac{d}{dr} (r^2 v)$$

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$$\rho = \frac{1}{6\pi G} t^{-2} \Rightarrow \frac{\dot{\rho}}{\rho} = -\frac{2}{t}$$



$$\rho = \frac{m}{\frac{4}{3}\pi a^3}$$

$$\frac{d \ln \rho}{dt} = \frac{\dot{\rho}}{\rho} = -\frac{3\dot{a}}{a}$$

$$\hookrightarrow \left(\frac{\dot{a}}{a} \right)^2 = \frac{8}{3} \pi G \rho$$

Friedmann equation

(pressure) need GR

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Friedmann equation

(Pressure)

$$\left(\frac{\dot{a}}{a} \right)^2 = \frac{8}{3} \pi G (\rho + 3p)$$

$$\rho_m + \rho_m + p$$

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Friedmann equation

Pressure) need GR

curvature

$$\left(\frac{\dot{a}}{a} \right)^2 = \frac{8}{3}\pi G \rho_{\text{total}} + \left(\frac{k}{a^2} \right)$$

$\rho_{\text{total}} = \rho_m + \rho_{\text{rad}} + \rho_{\Lambda}$

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lan

in the expanding

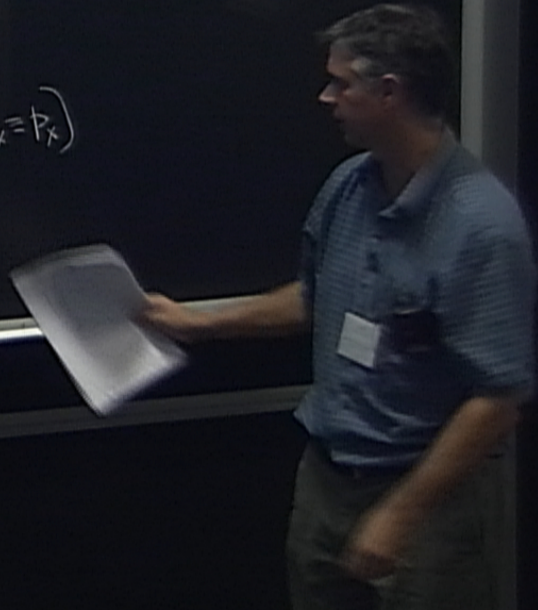
background fluctuations
s and large scale

g

$$\ddot{\frac{a}{a}} = -\frac{4\pi G}{3} \left[\sum_i (\rho_i + 3p_i) \right]$$

$i = \text{matter, radiation, ??}$

$$\rho_m = \rho_{m0} \left(\frac{a}{a_0} \right)^{-3}$$
$$\rho_r = \rho_{r0} \left(\frac{a}{a_0} \right)^{-4}$$
$$\rho_x = \rho_{x0} \left(\frac{a}{a_0} \right)^{3(1+w_x)} \quad \left[\rho_x w_x = P_x \right]$$



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$$\rho_r = \rho_{r0} \left(\frac{a}{a_0} \right)^{-4} \quad w=\frac{1}{3}$$
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CONSERVATION
OF TIME + SPACE

→

[Handwritten notes on a separate sheet of paper]

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CONSERVATION
OF TIME + SPACE
 $\Rightarrow w = -1$

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UNUSUAL
w
time + space
 $\Rightarrow w = -1$
Drives expansion!
 $\frac{GM}{R}$

lan

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$$\ddot{\frac{a}{a}} = -\frac{4\pi G}{3} \left[\sum_i (\rho_i + 3p_i) \right]$$

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WIDEALL
LA
time + space

$$\Rightarrow w = -1$$

Drives expansion!

$$\frac{GM}{R^2} \sim \frac{W_G}{R^2}$$

lan

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$$w=0$$


$$w=\frac{1}{3}$$

$$\left[\rho_x w_x = P_x \right]$$

WIDE
LO
TIME + SPACE

$$\Rightarrow w = -1$$

Drives expansion!


$$\frac{-GM}{R} \sim w_c$$

$$\frac{d \rho_p}{dt} = \frac{\dot{\rho}}{\rho} = -3 \frac{\dot{a}}{a}$$

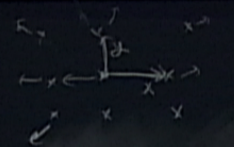
$$\hookrightarrow \left(\frac{\dot{a}}{a}\right)^2 = \frac{8}{3} \pi G \rho$$

Friedmann equation

(pressure) need GR

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8}{3} \pi G \rho_{\text{curvature}} + \left(\frac{k}{a^2}\right)$$

$\rho_{\text{curvature}} = \rho_m + \rho_{\text{rad}} + \rho_p$



$$\frac{a^2(t) dr^2}{c^2} - \frac{c^2 dt^2}{a^2(t)}$$

com
dr =

$$\frac{a}{\delta} = (1+z)^{-1}$$

redshift

$$r = \int \frac{dz}{\left(\frac{\dot{a}}{a}\right) c}$$

$$H^2 \equiv \left(\frac{\dot{a}}{a}\right)^2$$

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$$\frac{d(\rho a^3)}{dt} = \dot{\rho} a^3 = -3\frac{\dot{a}}{a} \rho a^3$$

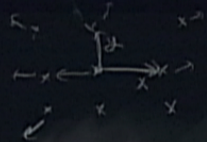
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$\rho_{\text{curvature}} = \rho_m + \rho_{\text{rad}} + \rho_{\Lambda}$



$$a^2(t) dr^2 = c^2 dt^2$$

'comoving'

$$dr = \frac{c dt}{a}$$

$$r = \int \frac{c dt}{a}$$

$$\frac{a}{a_0} = (1+z)^{-1}$$

↑
redshift

$$r = \int \frac{dz}{\left(\frac{\dot{a}}{a}\right)}$$

$$H^2 \equiv \left(\frac{\dot{a}}{a}\right)^2 = H_0^2 \left[\frac{\Omega_m}{2(1+z)^3} + \frac{\Omega_{\text{rad}}}{4(1+z)^4} + \frac{\Omega_{\Lambda}}{1+z} \right]$$

↑
Hubble Parameter

$H_0 = \text{Hubble constant}$

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ding

uctuations

scale

$\gamma x \quad \gamma x_0 (a_0) \quad L^2$

$$H^2 = H_0^2 \left[\Omega_m (1+z)^3 + \Omega_{\text{rad}} (1+z)^4 + \underbrace{\Omega_{\text{de}} (1+z)^{-3(w)}}_{(\Omega_{\Lambda} \quad w=-1)} \right]$$

inding

uctuations

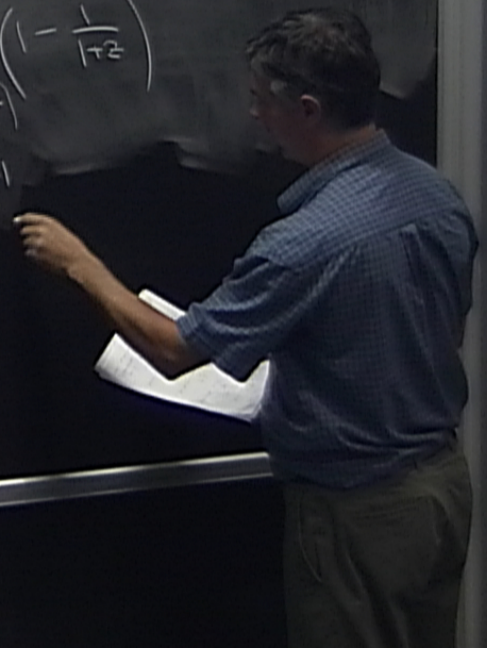
scale

$\gamma x \quad \gamma x_0 (a_0) \quad L^2$

$$H^2 = H_0^2 \left[\Omega_m (1+z)^3 + \Omega_{\text{rad}} (1+z)^4 + \underbrace{\Omega_{\Lambda} (1+z)^{-3}}_{\Omega_{\Lambda} w = -1} \right]$$

$$r = \int \frac{c dz'}{H(z')} \stackrel{\text{RAD DOM}}{=} \underbrace{c H_0^{-1}}_{(3000 \text{ Mpc})} \left(1 - \frac{1}{1+z} \right)$$

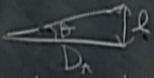
$$t = \int \frac{dz}{H(z) (1+z)} \stackrel{\text{RAD DOM}}{=} H_0^{-1}$$



$$\frac{\partial \ln \rho}{\partial t} = -\frac{1}{r} \frac{\partial \rho}{\partial r} (r \dot{r})$$

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Angular size distance
physical length l

$$D = \frac{\text{comoving length}}{D_A}$$

=

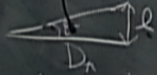
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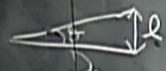
Angular size distance
physical length l

$$D = \frac{\text{comoving length}}{r} = \frac{l(1+z)}{\text{distance}}$$

$$D_A = \frac{r}{1+z}$$

Plan

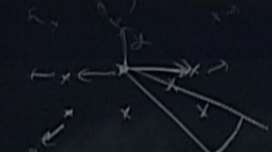
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Angular size distance
Physical length l

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$$\frac{a}{a_0} = (1+z)^{-1}$$

↑
redshift

$$r = \int \frac{dz}{\left(\frac{\dot{a}}{a}\right) c}$$

$$H^2 \equiv \left(\frac{\dot{a}}{a}\right)^2 = H_0^2 \left[\frac{\Omega_m}{2(1+z)} + \frac{\Omega_{vac}}{\Omega_0} \right]$$

↑
Hubble parameter

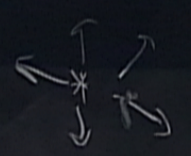
H_0 = Hubble = velocity constant / today

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$$D_A = \frac{r}{1+z}$$

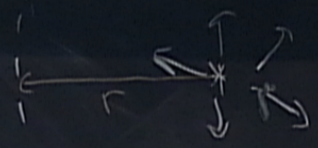
distance



$$\text{flux} = \frac{\text{luminosity}}{4\pi R^2} = \frac{L}{4\pi D_L^2}$$

$$D_A = \frac{r}{1+z}$$

distance



$$\text{flux} = \frac{\text{Luminosity}}{4\pi R^2} = \frac{L}{4\pi D_L^2} = \frac{L}{4\pi r^2 (1+z)^2}$$

Luminosity = energy emitted per second at time of emission

flux = energy received per s per cm² at time of receiving

$$\Rightarrow D_L = r(1+z)$$

$$D = \frac{\text{comoving length}}{r} = \frac{l(1+z)}{\text{distance}}$$

$$D_A = \frac{r}{1+z}$$

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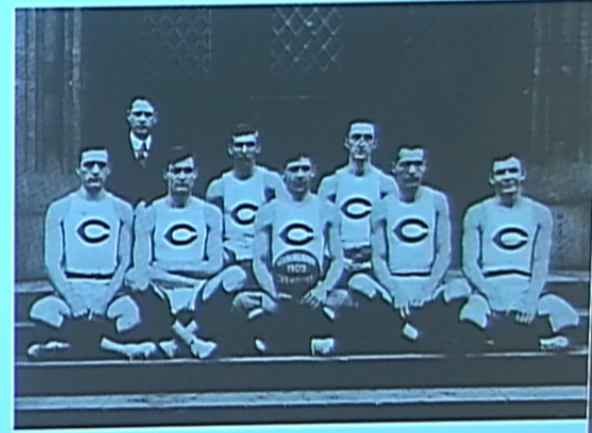
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$$\Rightarrow D_L = r(1+z)$$

March Madness: Edwin Hubble

- University of Chicago graduate
- Rhodes scholar
- Varsity basketball player
 - U Chicago won Big 10 title in 1909



“a man of magnificent physique, admirable scholarship, and worthy and lovable character.” – Robert Millikan

<http://chicagomaroon.com/2009/04/10/before-revolutionizing-astronomy-hubble-helped-rewrite-record-books/>

$$D_A = \frac{r}{1+z}$$

$$D = \frac{\text{comoving length}}{r} = \frac{l(1+z)}{\text{distance}}$$

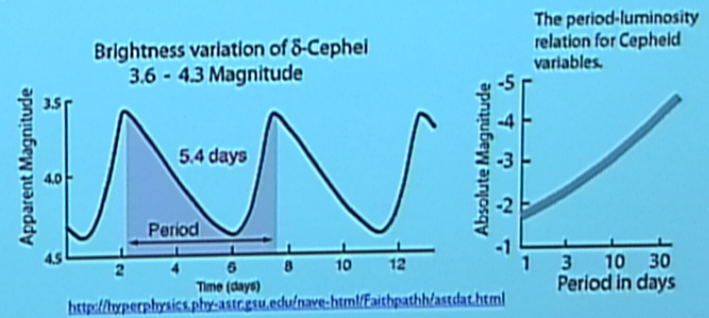
$$\text{flux} = \frac{\text{Luminosity}}{4\pi R^2} = \frac{L}{4\pi D_L^2}$$

$$\text{Luminosity} = \text{energy emitted per second at top of emission} \rightarrow D_L$$

$$\text{flux} = \text{energy received per s per cm}^2 \text{ at top of receiving}$$

Cepheid Variables

- Large variations over several days
- Empirical tight correlation between timescale of variation and luminosity
- So, measure the period, you know its brightness!



$$D_A = \frac{r}{1+z}$$

$$D = \frac{\text{comoving length}}{r} = \frac{l(1+z)}{\text{distance}}$$

$$\text{flux} = \frac{\text{luminosity}}{4\pi R^2} = \frac{L}{4\pi D_L^2}$$

$$\text{luminosity} = \text{energy emitted per second at time of emission}$$

$$\text{flux} = \text{energy received per s per cm}^2 \text{ at time of receiving} \Rightarrow D_L =$$

Hubble's Law

Things that are further away are moving away from us faster!

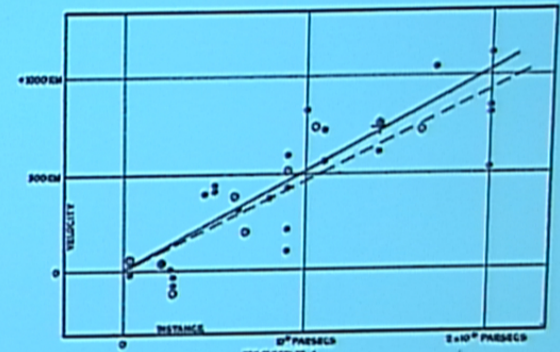


FIGURE 1
Velocity-Distance Relation among Extra-Galactic Nebulae.
 Radial velocities, corrected for solar motion, are plotted against distances estimated from involved stars and mean luminosities of nebulae in a cluster. The black discs and full line represent the solution for solar motion using the nebulae individually; the circles and broken line represent the solution combining the nebulae into groups; the cross represents the mean velocity corresponding to the mean distance of 22 nebulae whose distances could not be estimated individually.

Hubble (1929)

$$D = \frac{\text{comoving length}}{r} = \frac{l(1+z)}{\text{distance}}$$

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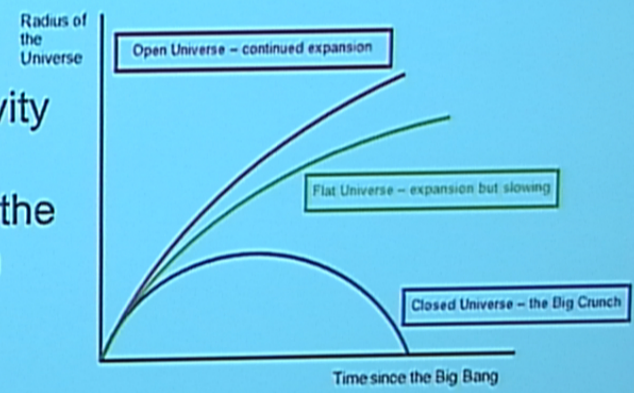
Luminosity = energy emitted per second at time of emission

flux = energy received (per s) per cm² at time of receiving

$$D_L = r(1+z)$$

The Expanding Universe

- Since gravity is always pulling in, the expansion should be always slowing, right?



<http://www.schoolphysics.co.uk/age16-19/Astronomy/text/>

$$D = \frac{\text{comoving length}}{r} = \frac{l(1+z)}{\text{distance}}$$

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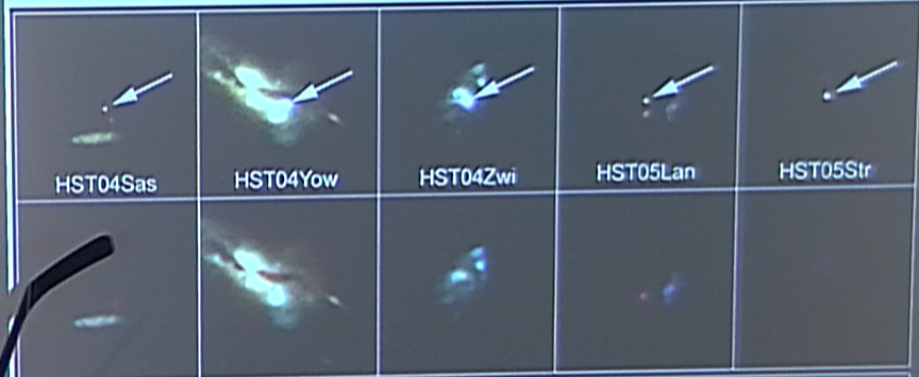
$$\text{luminosity} = \text{energy emitted per second at time of emission}$$

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$$\rightarrow D_L = r(1+z)$$

Measuring cosmic expansion

- Type Ia supernovae can be seen across the universe, with known intrinsic brightnesses



Host Galaxies of Distant Supernovae
Hubble Space Telescope • Advanced Camera for Surveys

NASA, ESA, and A. Riess (STScI)

STScI-PRC06-52

h. sym.

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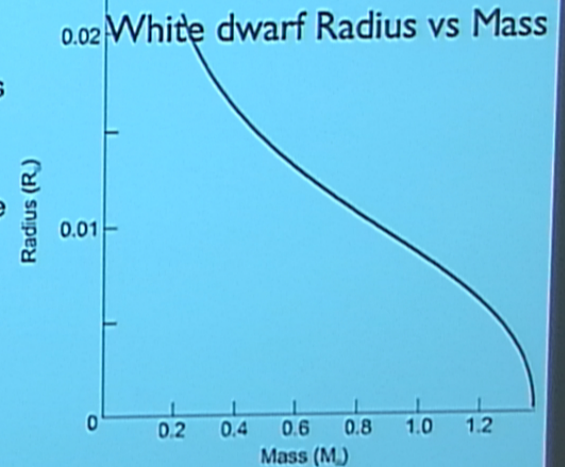
angular size distance
physical length l

$D = \text{comoving length}$

$$D_A = \frac{r}{1+z}$$

Degenerate Matter Mass-Radius relation

- As mass gets larger, equilibrium size gets smaller as gravity starts to dominate
- Smaller radius means higher temperature as mass gets piled on core
- Something new will happen when size approaches 0
- "degenerate" in this case means lowest-energy states all filled



<http://hendrix2.uoregon.edu/~imamura/122/lecture-8/white-dwarf.html>

h. sym.

$$\frac{\partial \ln \rho}{\partial E} = -\frac{1}{r} \frac{\partial}{\partial r} (r^2 \sigma)$$

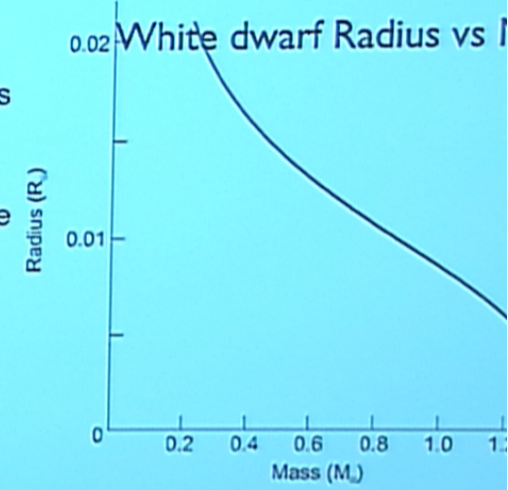
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$$\rho = \frac{1}{6\pi G} t^{-2} \Rightarrow \frac{\dot{\rho}}{\rho} = -\frac{2}{t}$$



Degenerate Matter Mass-Radius relation

- As mass gets larger, equilibrium size gets smaller as gravity starts to dominate
- Smaller radius means higher temperature as mass gets piled on core
- Something new will happen when size approaches 0
- "degenerate" in this case means lowest-energy states all filled

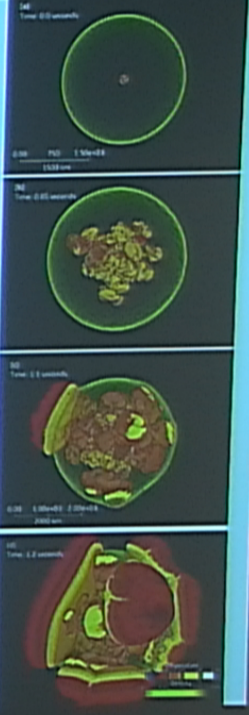


<http://hendrix2.uoregon.edu/~imamura/122/lecture-8/white-dwarf.htm>

White dwarf explodes

- Details are not at all certain
 - Detonation starts at center or off-center?
 - Thermonuclear propagating “flame front” has features on scales of a few cm, most likely, propagating billions of cm in a second or so

https://en.wikipedia.org/wiki/Type Ia_supernova



h. sym.

$$\frac{\partial \ln \rho}{\partial t} = -\frac{1}{r} \frac{\partial}{\partial r} (r^2 v)$$

$$-v(r) = -\frac{1}{3} \left(\frac{\partial \ln \rho}{\partial t} \right) r$$

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Quantum

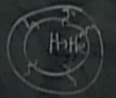
$E_g \sim \frac{GM^2}{R}$
 $E_f \sim N E_f \sim N \left(\frac{N^{3/4} h^2}{R} \right)$

h. sym.

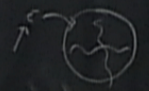
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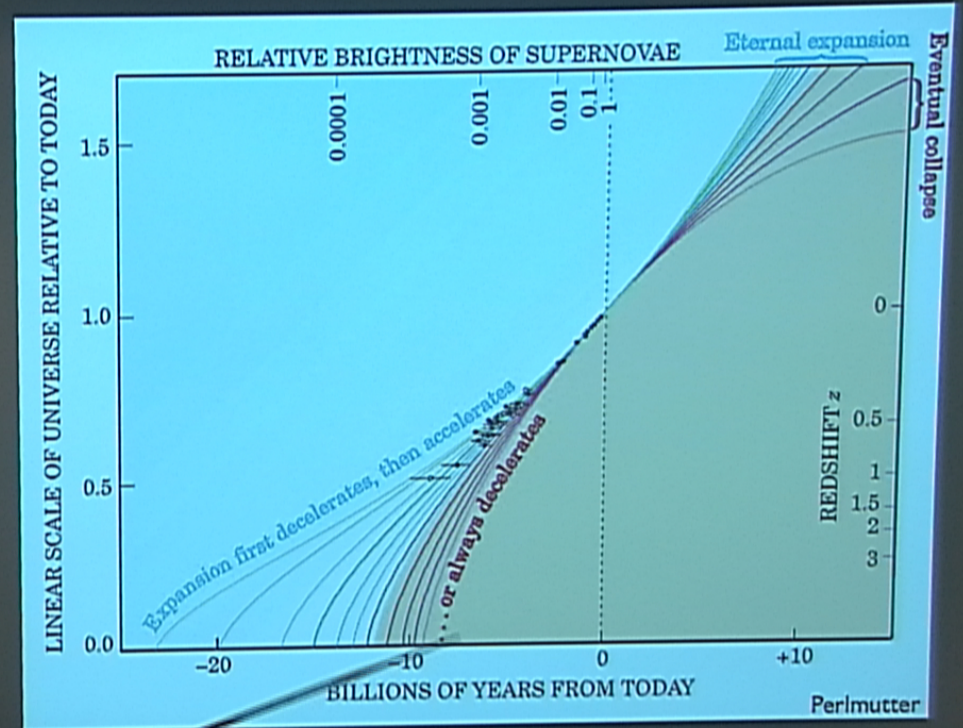


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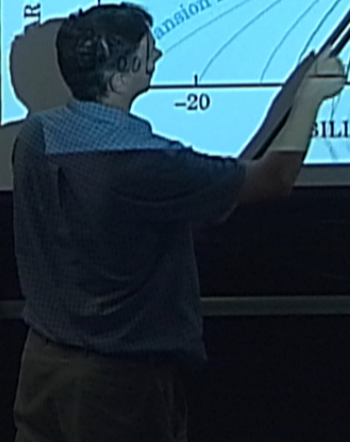
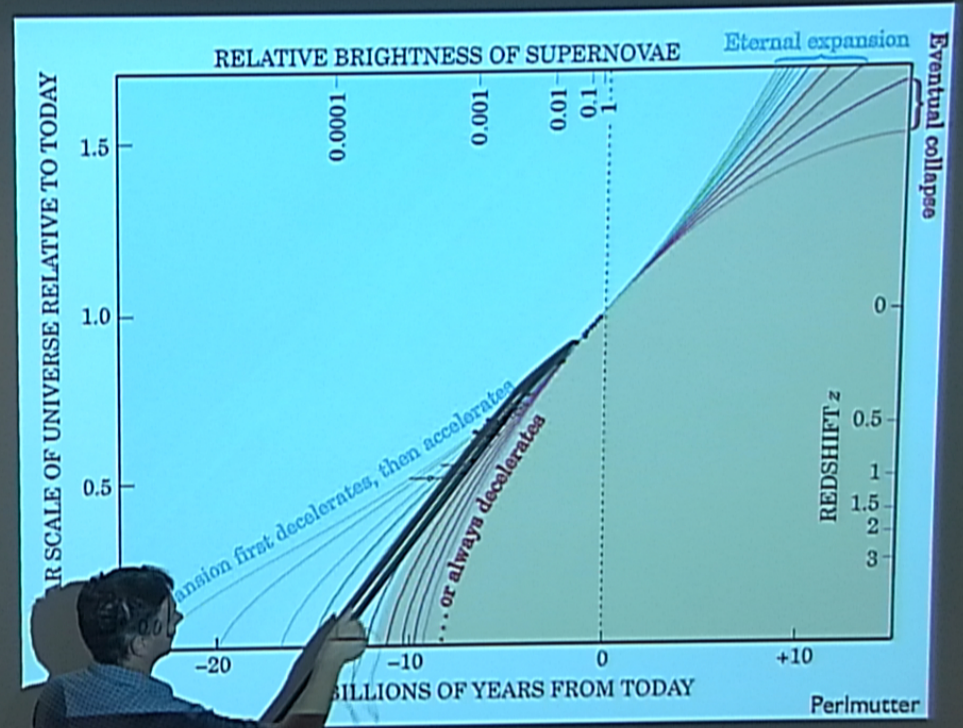


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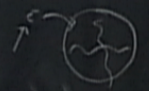
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Quantum

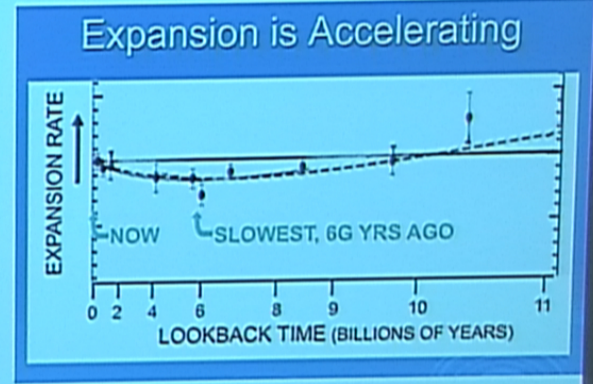


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The Expansion is Speeding Up!

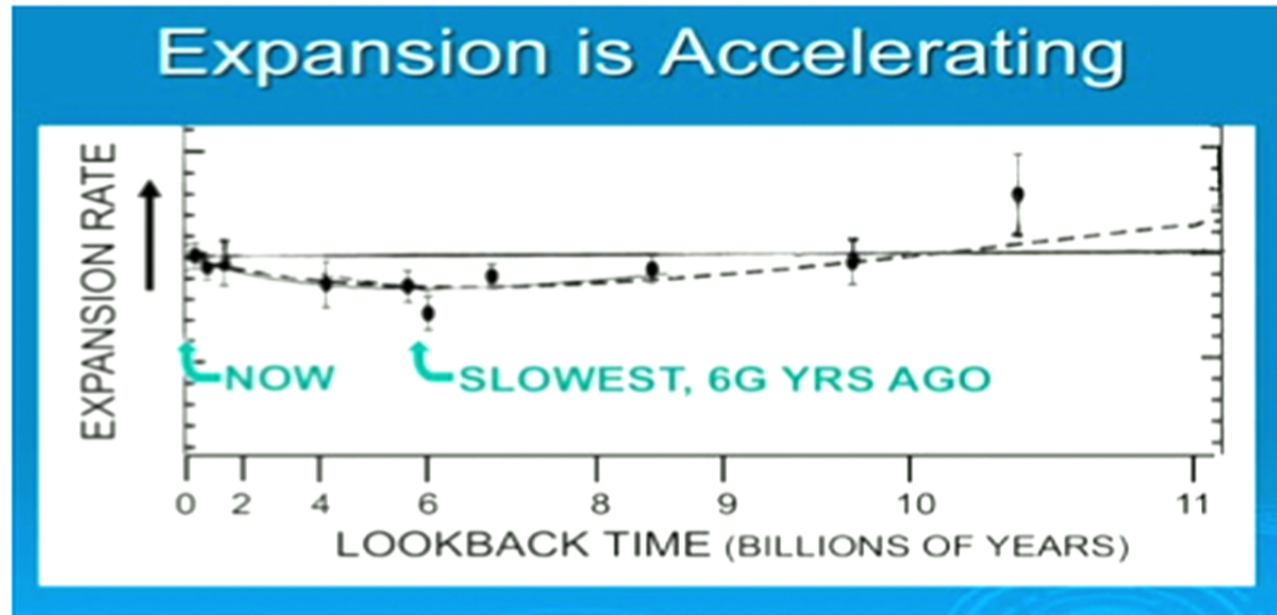
- It slowed down, as expected, but has recently begun to accelerate
- The universe is apparently full of some kind of "dark energy" that stays roughly constant in density as the universe expands



<http://www.guidetothecosmos.com/newsletter-Accel-Universe.html>

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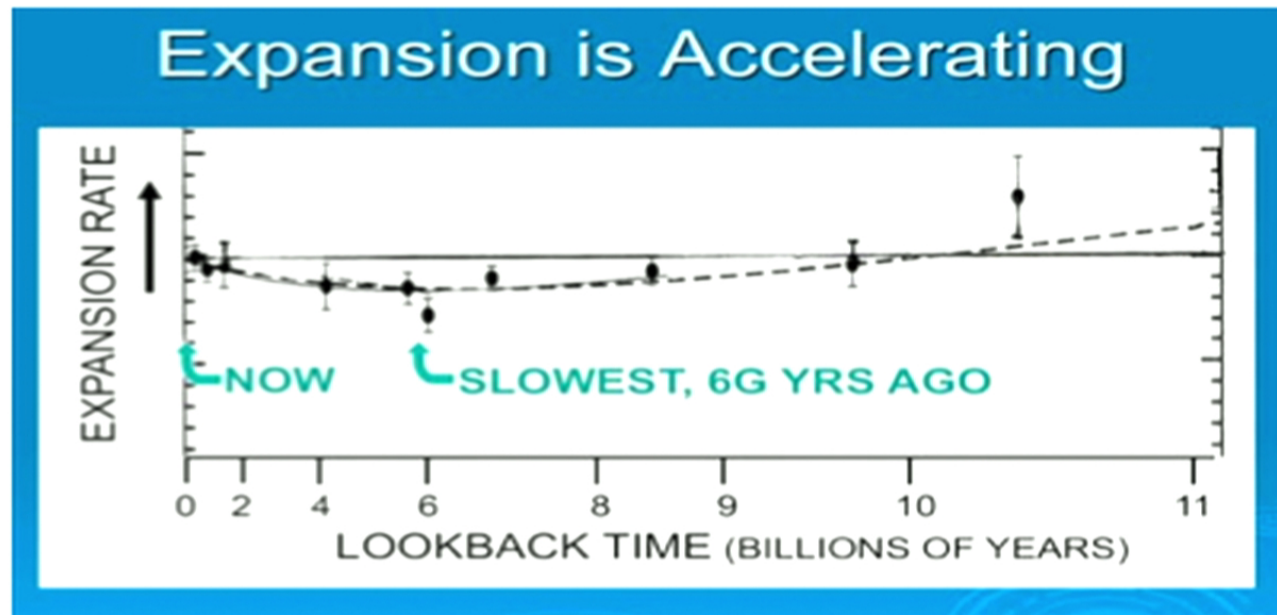
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Example of why this is even harder than it looks:
"K-corrections"

- Objects at different distances appear at different wavelengths
- Requires accurate understanding of instrument response as a function of wavelength to sources in the field of optics

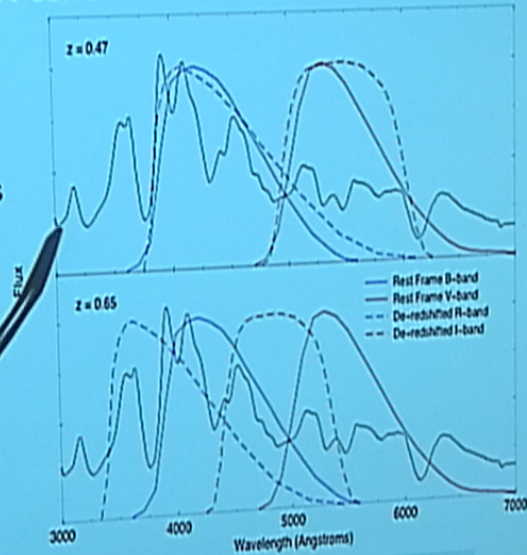


Fig from Nugent et al 2002

$\rho_m = \rho_{m0} \left(\frac{a}{a_0}\right)^{-3}$ $w=0$
 $\rho_r = \rho_{r0} \left(\frac{a}{a_0}\right)^{-4}$ $w=\frac{1}{3}$
 $\rho_x = \rho_{x0} \left(\frac{a}{a_0}\right)^{-2(1+w_x)}$ $\left\{ \begin{array}{l} \rho_x w_x = P_x \end{array} \right.$

$H^2 = H_0^2 \left[\Omega_m (1+z)^3 + \Omega_{rad} (1+z)^4 + \Omega_\Lambda (1+z)^{-2} \right]$
 $r = \int \frac{c dz'}{H(z')} \approx \frac{c H_0^{-1}}{(3000 \text{ Mpc})} \left(1 - \frac{1}{1+z} \right)$
 $t = \int \frac{dz}{H(z)} \left(\frac{1}{1+z} \right) \approx \frac{H_0^{-1}}{2} \left(1 - \frac{1}{(1+z)^2} \right)$
 $\lambda_0 z \rightarrow \lambda \rightarrow r \rightarrow c H_0^{-1} \left(1 - \frac{1}{1+z} \right)$
 $t \rightarrow \frac{c}{2} \left(1 - \frac{1}{(1+z)^2} \right)$

Summary

- The universe is expanding
 - Hubble (1929)
- Expansion or contraction is the norm for any uniform self-gravitating medium
- Expansion mucks around with time and distance measures
- Universe used to be hot and dense
- Currently dominated by something which
 - ~constant in density ("dark energy")
 - evidence from Type IA supernovae

$$\rho_m = \rho_{m0} \left(\frac{a}{a_0} \right)^{-3}$$

$$w=0$$

$$\rho_r = \rho_{r0} \left(\frac{a}{a_0} \right)^{-4}$$

$$w=\frac{1}{3}$$

$$\rho_x = \rho_{x0} \left(\frac{a}{a_0} \right)^{-3(1+w_x)}$$

$$\rho_x w_x = p_x$$

$$H^2 = H_0^2 \left[\Omega_m (1+z)^3 + \Omega_{rad} (1+z)^4 + \frac{\Omega_\Lambda (1+z)^{-3}}{(-\Omega_\Lambda - w_\Lambda)} \right]$$

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$$t = \int \frac{dz}{H(z) (1+z)} \approx \frac{H_0^{-1}}{2} \left(1 - \frac{1}{(1+z)^2} \right)$$

$$\Lambda \rightarrow \rho_\Lambda = r \rightarrow c H_0^{-1} \left. \begin{array}{l} t \rightarrow \frac{1}{2} H_0^{-1} \\ ct = \frac{d}{2} c \end{array} \right\}$$