

Title: Black hole decay and Fast Radio Bursts

Date: Jun 04, 2015 02:30 PM

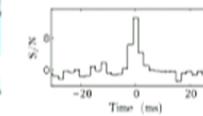
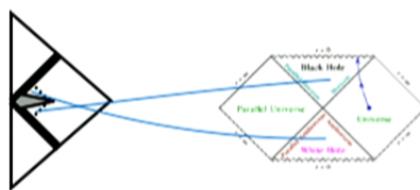
URL: <http://pirsa.org/15060006>

Abstract: <p>Quantum effects render black holes unstable. Besides Hawking radiation there is another, genuinely quantum gravitational, source of instability: the Hajicek-Kiefer explosion via tunnelling to a white hole. A recent result in classical general relativity makes this decay channel plausible: there is an exact external solution of the Einstein equations locally (but not globally) isometric to extended Schwarzschild, which describes an object collapsing into a black hole and then exploding out of a white hole. The tunnelling time can in principle be computed using Loop Quantum Gravity. If it is sufficiently short, present explosions of primordial black hole could be observables, opening a new observational window on quantum gravity. I discuss the possibility that these explosions could be related to the recently observed and mysterious Fast Radio Bursts. </p>

Black hole decay and Fast Radio Bursts

carlo rovelli

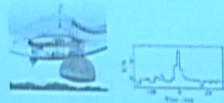
hal haggard, aurelien barrau, francesca vidotto



Black hole decay and Fast Radio Bursts

carlo rovelli

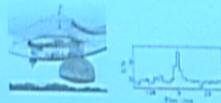
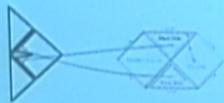
hal haggard, aurelien barrau, francesca vidotto



Black hole decay and Fast Radio Bursts

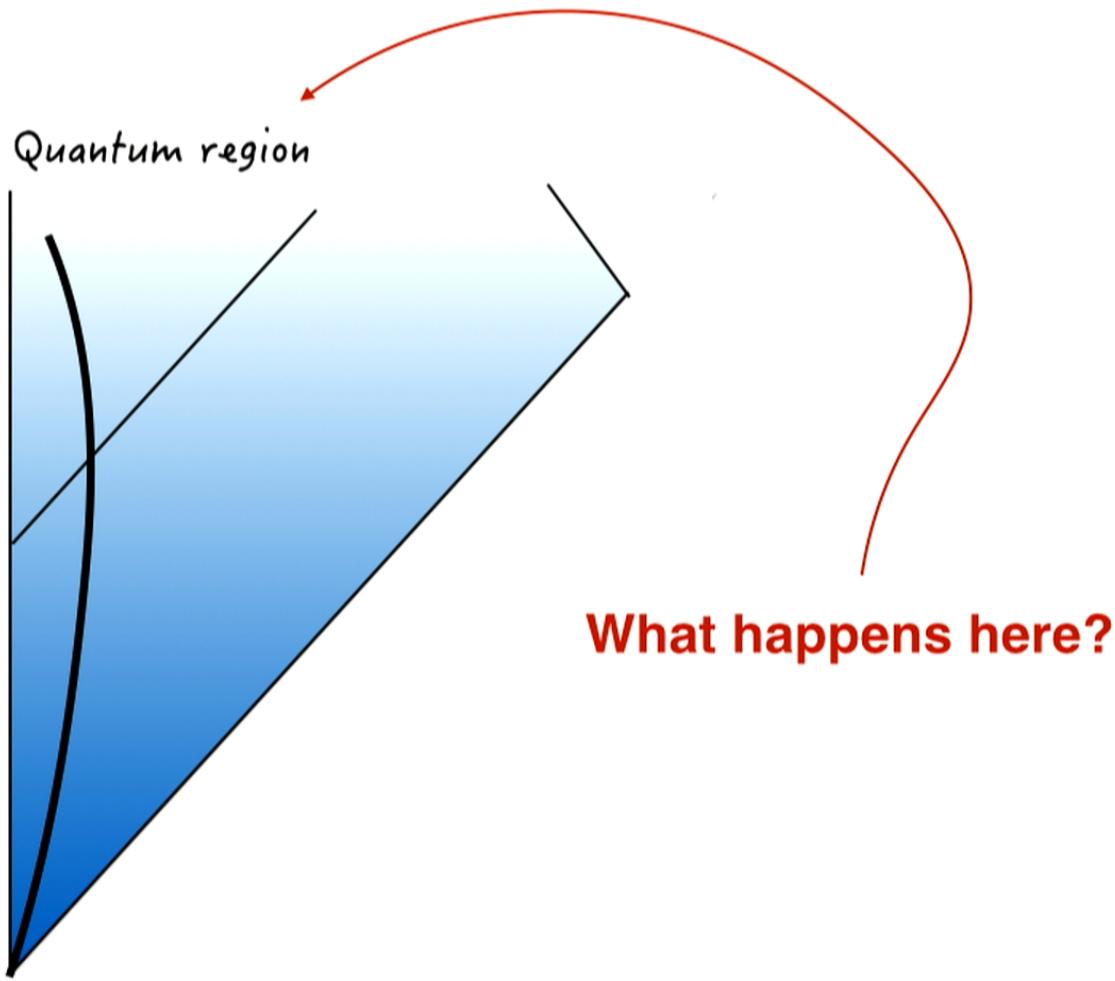
carlo rovelli

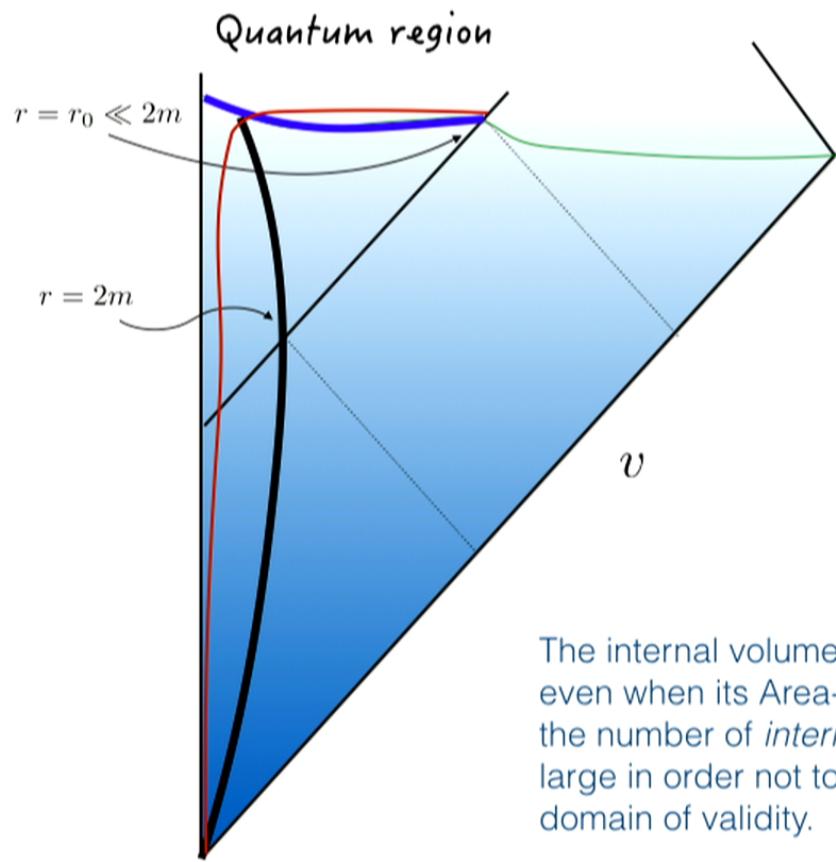
hal haggard, aurelien barrau, francesca vidotto



Quantum gravity: we need observations

- Early universe : CMB
- Short scale : Lorentz violations
- Black holes : Observable quantum effects





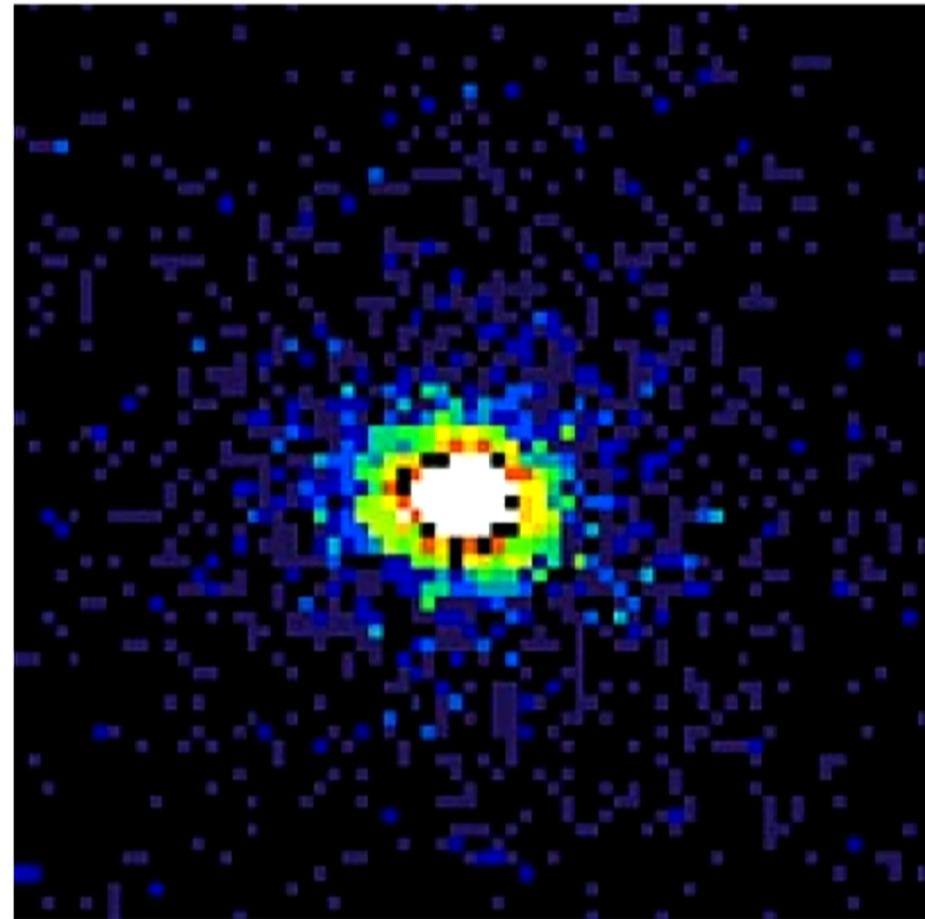
$$V \sim 3\sqrt{3}\pi m^2 v$$

How big is a black hole?
 Marios Christodoulou, Carlo Rovelli,
Phys.Rev. D91 (2015) 6, 064046

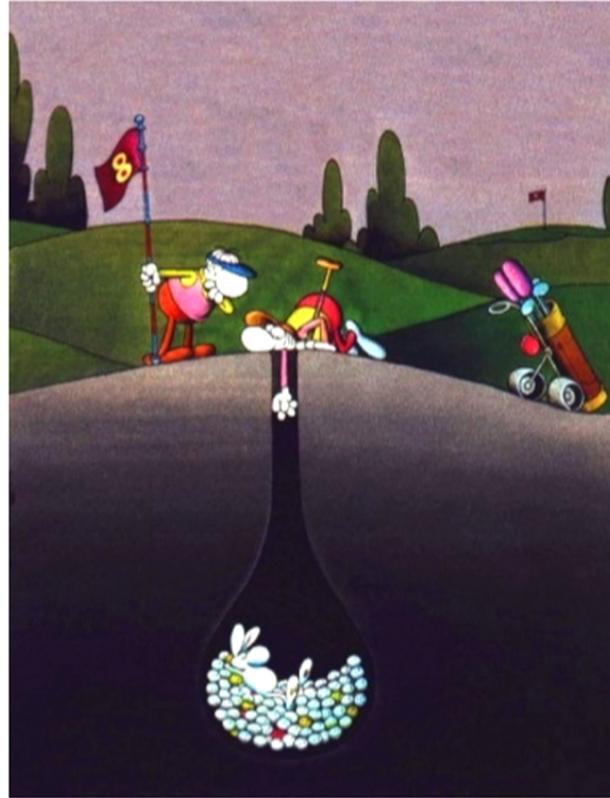
After a time m^3 : $V \sim m^5$

At the end of the evaporation: $V \sim m^{\frac{7}{2}}$

The internal volume of a black hole is large even when its Area-Entropy is small:
 the number of *internal* degrees of freedom must be large in order not to violate local qft in its own domain of validity. ... holography?

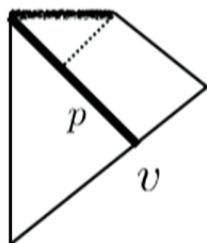
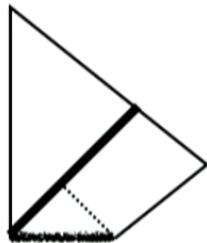






What happens to the matter falling into black holes?

- It disappears (?)
- It creates "another universe" (Lee)
- It stays there forever (nothing is forever)
- It comes out.



The Hajicek-Kiefer bounce

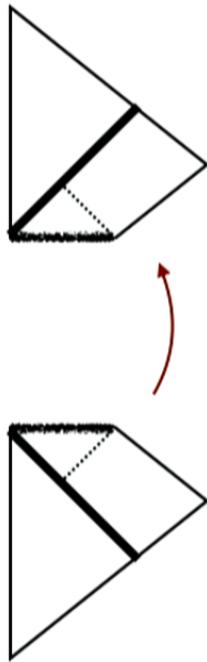
- Spherical symmetry
- Null shell of matter
- Classically: Finite dimensional phase space (v,p) separated in two disconnected components:
 - $p>0$: shell collapsing into white hole (future singularity)
 - $p<0$: shell emerging from a white hole (past singularity)

- Can a black hole truly tunnel into a white hole?

Singularity avoidance by collapsing shells in quantum gravity

Petr Hájíček, Clauss Kiefer.

IJMP D, (2001), 775.

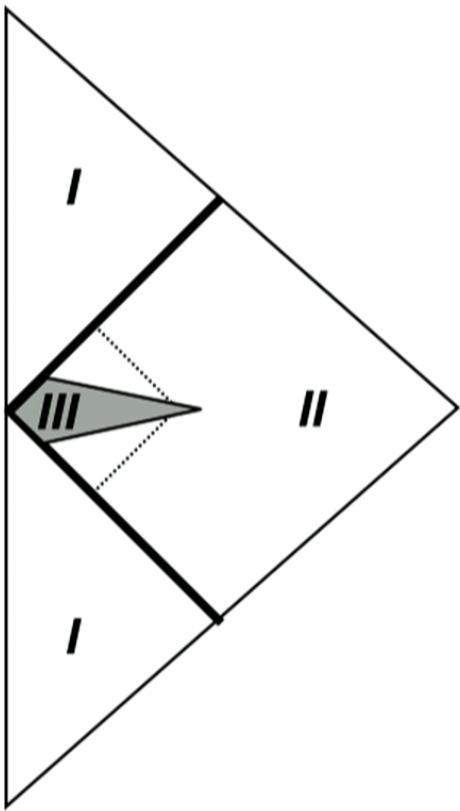


Intriguing aspects:

- Classical stable objects often unstable via quantum effects (cfr. nuclear decay).
- Hawking radiation takes a huge amount of time ($t \sim m^3$) and is not a full quantum gravitational phenomenon.

Difficulties:

- Two distinct asymptotic regions.
- A quantum jump involving the entire universe.
- What could determine the tunneling time?



A technical result in classical GR:

The following metric is an exact vacuum solution, plus an ingoing and outgoing shell, of the Einstein equations outside a finite spacetime region (grey).

$$ds^2 = -F(u, v)dudv + r^2(u, v)(d\theta^2 + \sin^2\theta d\phi^2)$$

Region I $F(u_I, v_I) = 1, \quad r_I(u_I, v_I) = \frac{v_I - u_I}{2}.$
 $v_I < 0.$

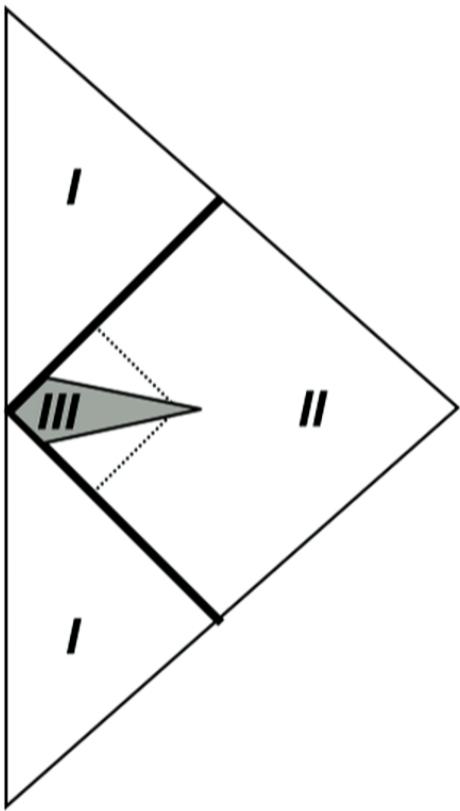
Region II $F(u, v) = \frac{32m^3}{r}e^{\frac{r}{2m}} \quad \left(1 - \frac{r}{2m}\right)e^{\frac{r}{2m}} = uv.$

Matching: $r_I(u_I, v_I) = r(u, v) \rightarrow u(u_I) = \frac{1}{v_o} \left(1 + \frac{u_I}{4m}\right) e^{\frac{u_I}{4m}}.$

Region III $F(u_q, v_q) = \frac{32m^3}{r_q}e^{\frac{r_q}{2m}}, \quad r_q = v_q - u_q.$

Black hole fireworks: quantum-gravity effects outside
the horizon spark black to white hole tunneling
Hal M. Haggard, Carlo Rovelli
arXiv:1407.0989

The metric is determined by three constants: m, ϵ, δ



A technical result in classical GR:

The following metric is an exact vacuum solution, plus an ingoing and outgoing shell, of the Einstein equations outside a finite spacetime region (grey).

$$ds^2 = -F(u, v)dudv + r^2(u, v)(d\theta^2 + \sin^2\theta d\phi^2)$$

Region I $F(u_I, v_I) = 1, \quad r_I(u_I, v_I) = \frac{v_I - u_I}{2}.$
 $v_I < 0.$

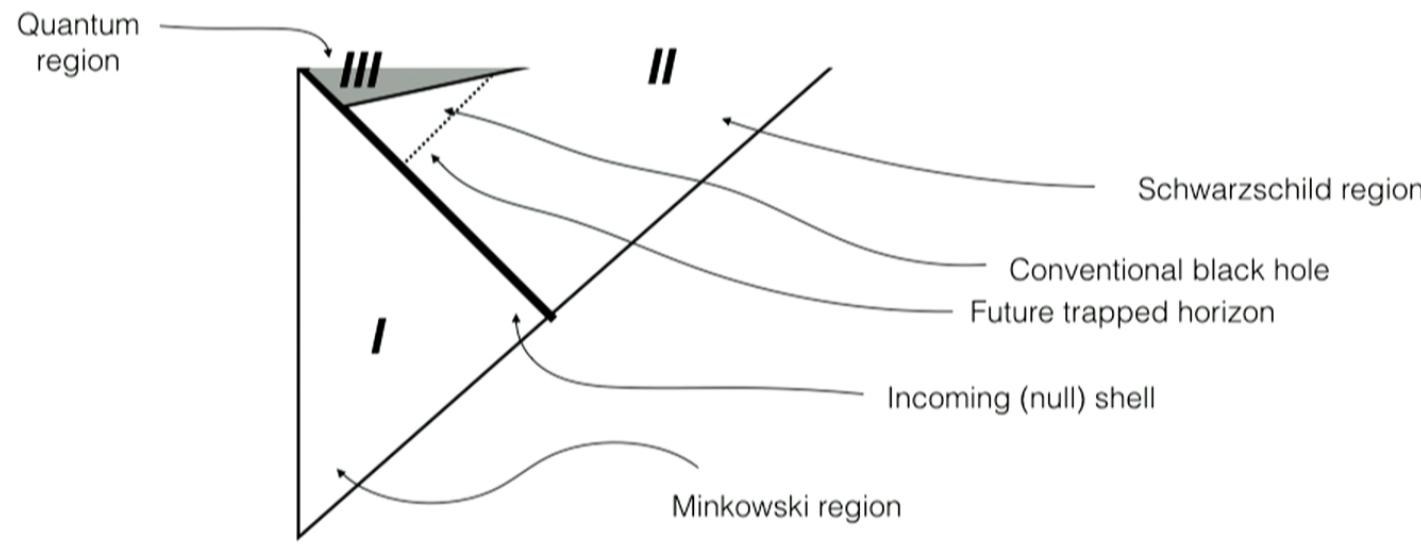
Region II $F(u, v) = \frac{32m^3}{r}e^{\frac{r}{2m}} \left(1 - \frac{r}{2m}\right)e^{\frac{r}{2m}} = uv.$

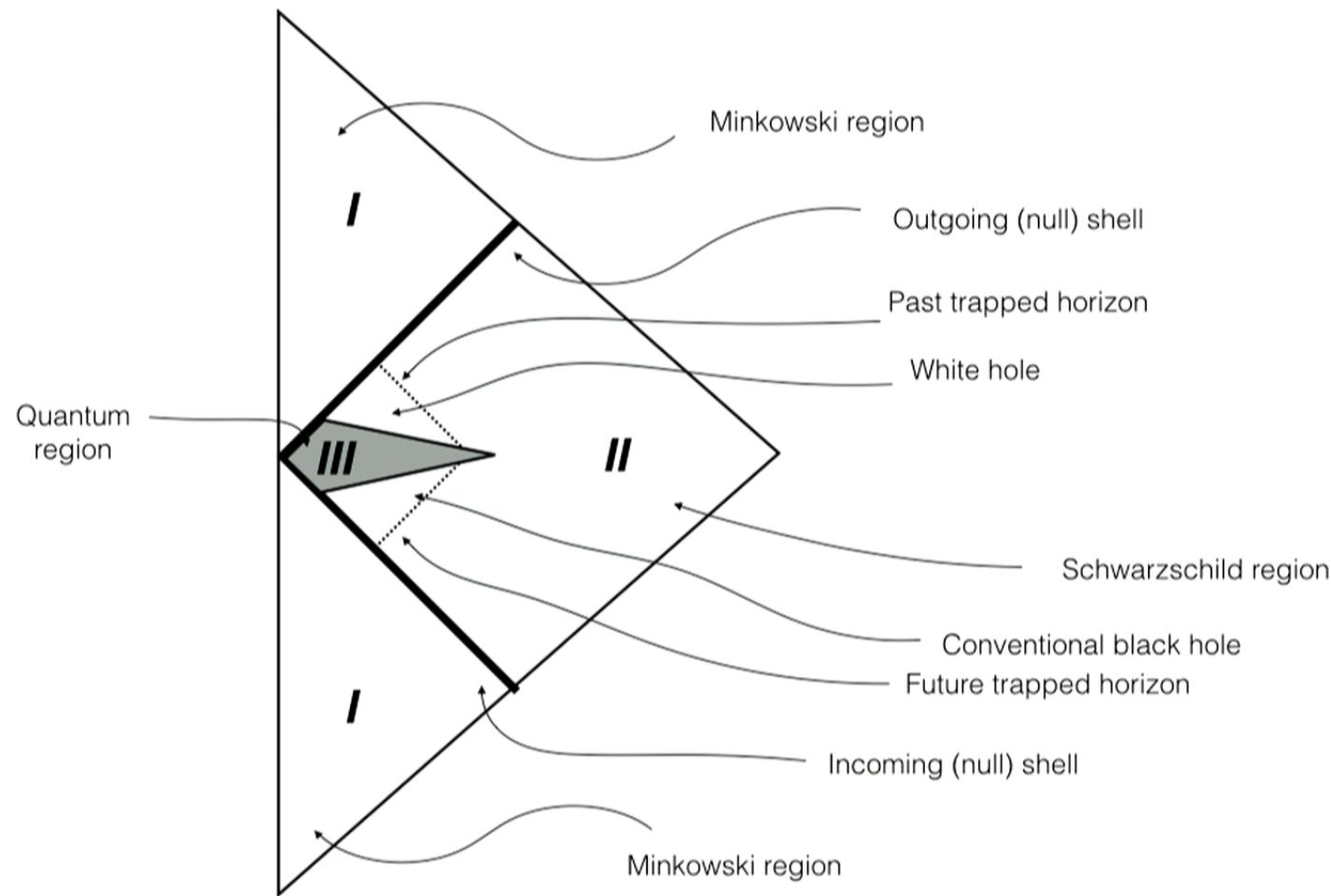
Matching: $r_I(u_I, v_I) = r(u, v) \rightarrow u(u_I) = \frac{1}{v_o} \left(1 + \frac{u_I}{4m}\right) e^{\frac{u_I}{4m}}.$

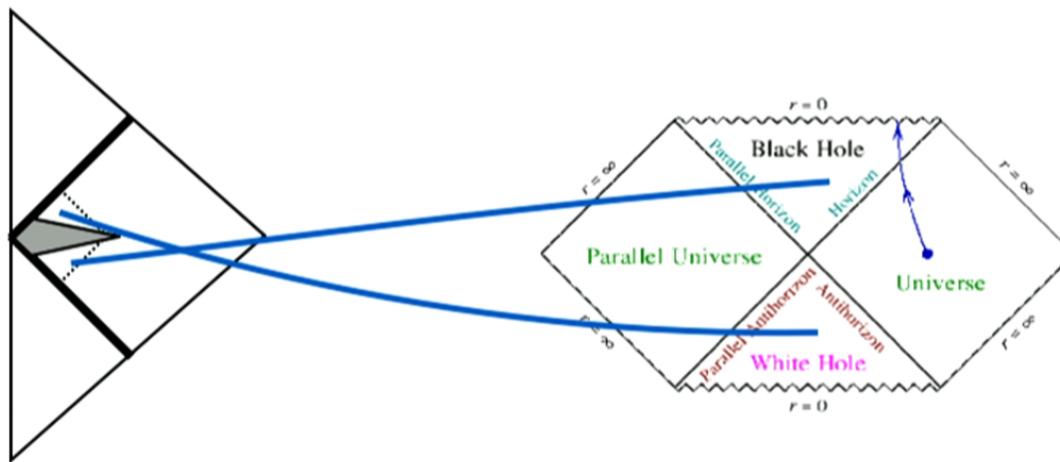
Region III $F(u_q, v_q) = \frac{32m^3}{r_q}e^{\frac{r_q}{2m}}, \quad r_q = v_q - u_q.$

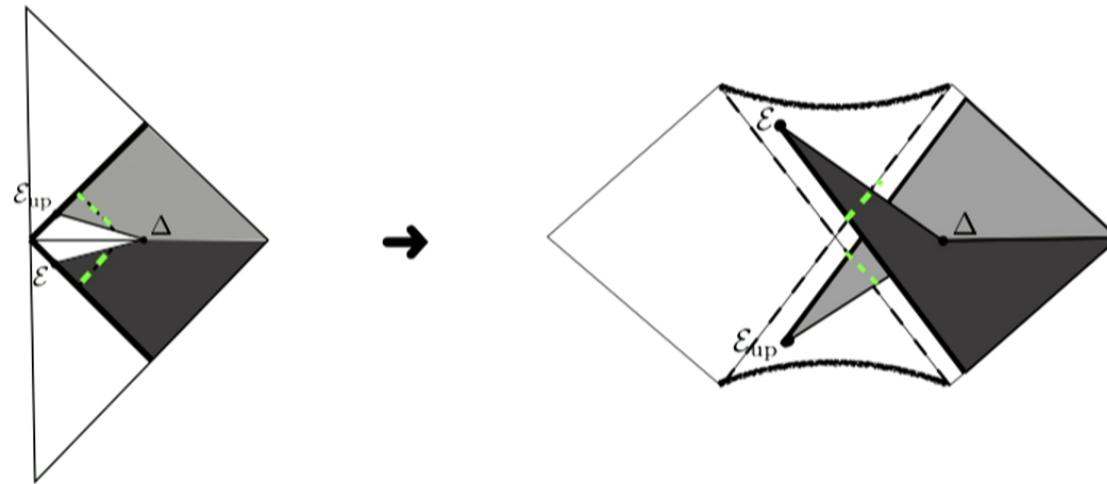
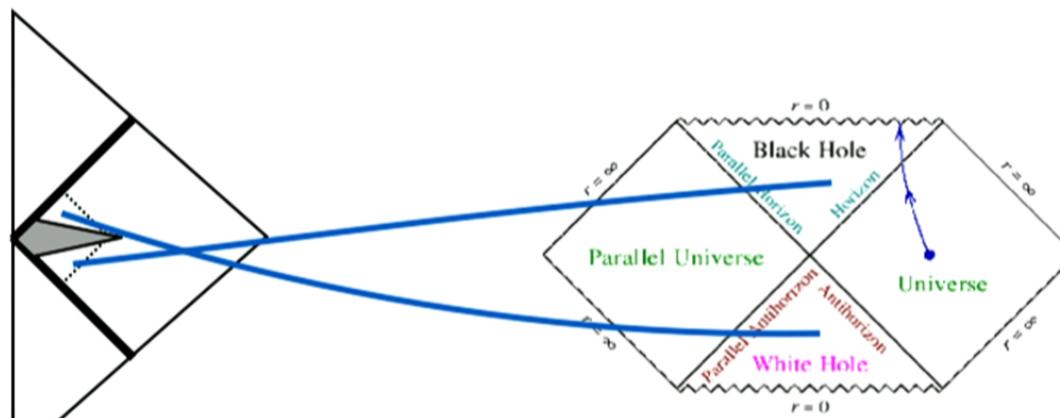
Black hole fireworks: quantum-gravity effects outside
the horizon spark black to white hole tunneling
Hal M. Haggard, Carlo Rovelli
arXiv:1407.0989

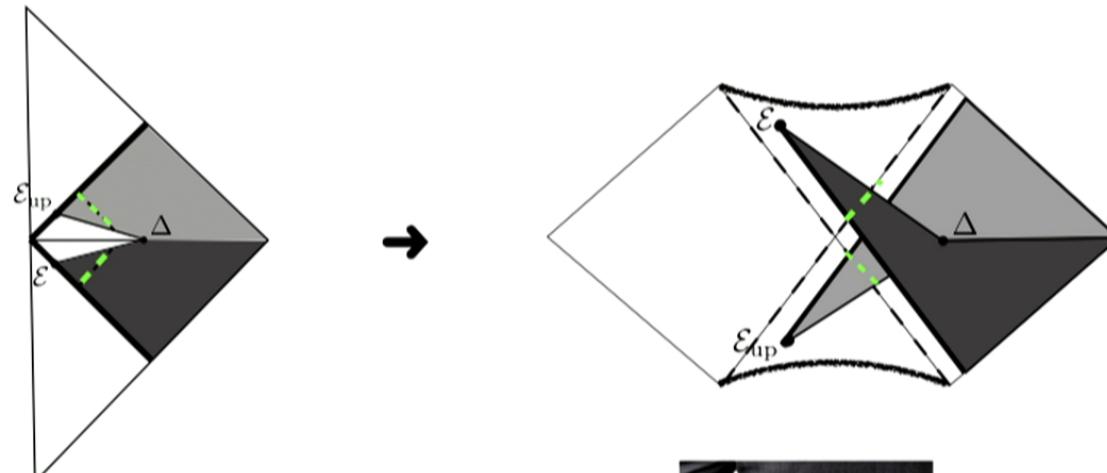
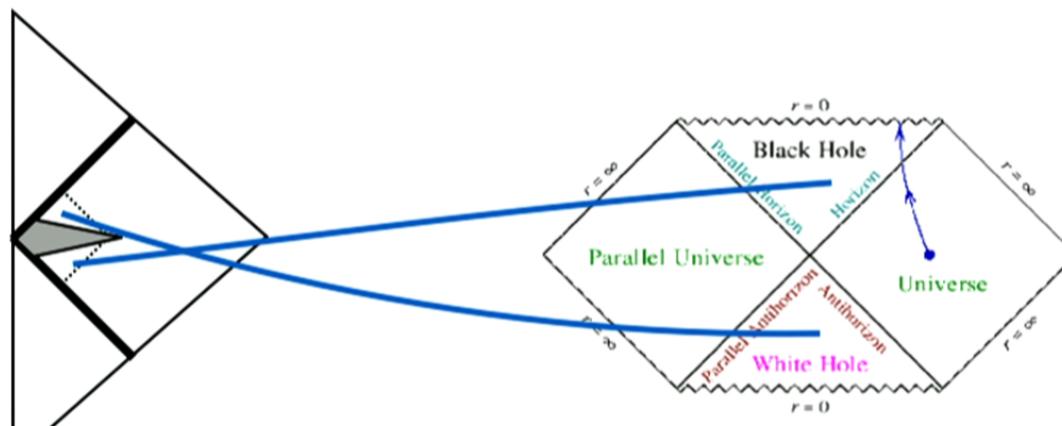
The metric is determined by three constants: m, ϵ, δ



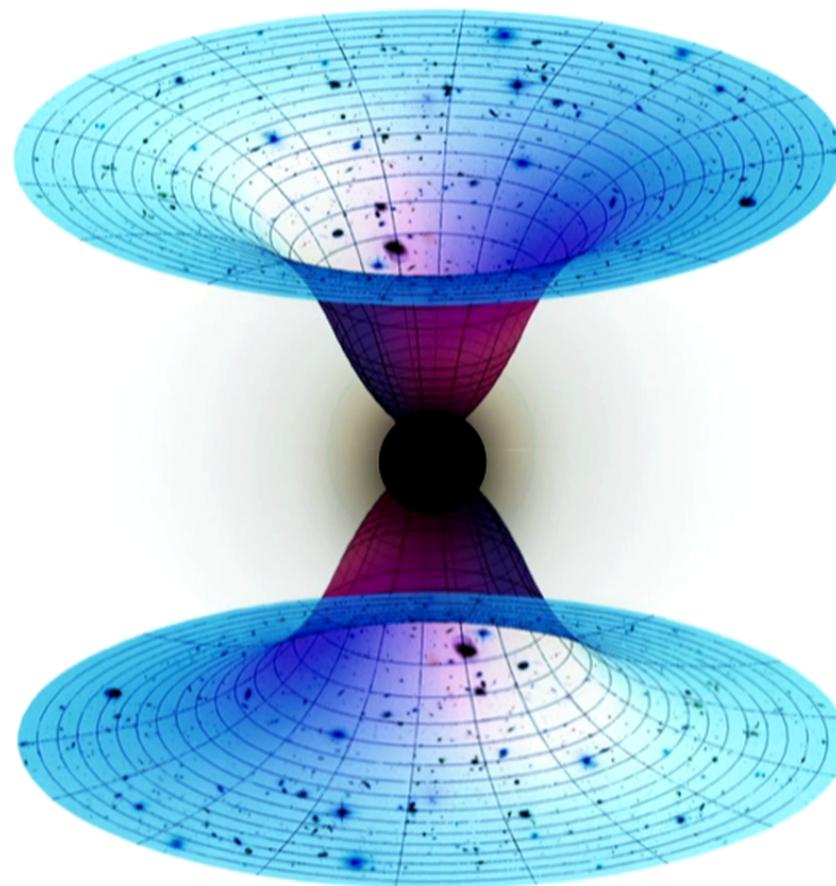








The Fingers Crossed



The metric is determined by three constants:

- m is the mass of the collapsing shell.

- ϵ is the radius where quantum effect start on the shell: $\epsilon \sim \left(\frac{m}{m_P^3}\right)^{\frac{1}{3}} l_P$.

Quantum pressure causes the bounce

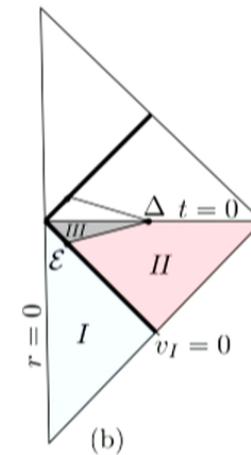
Planck stars

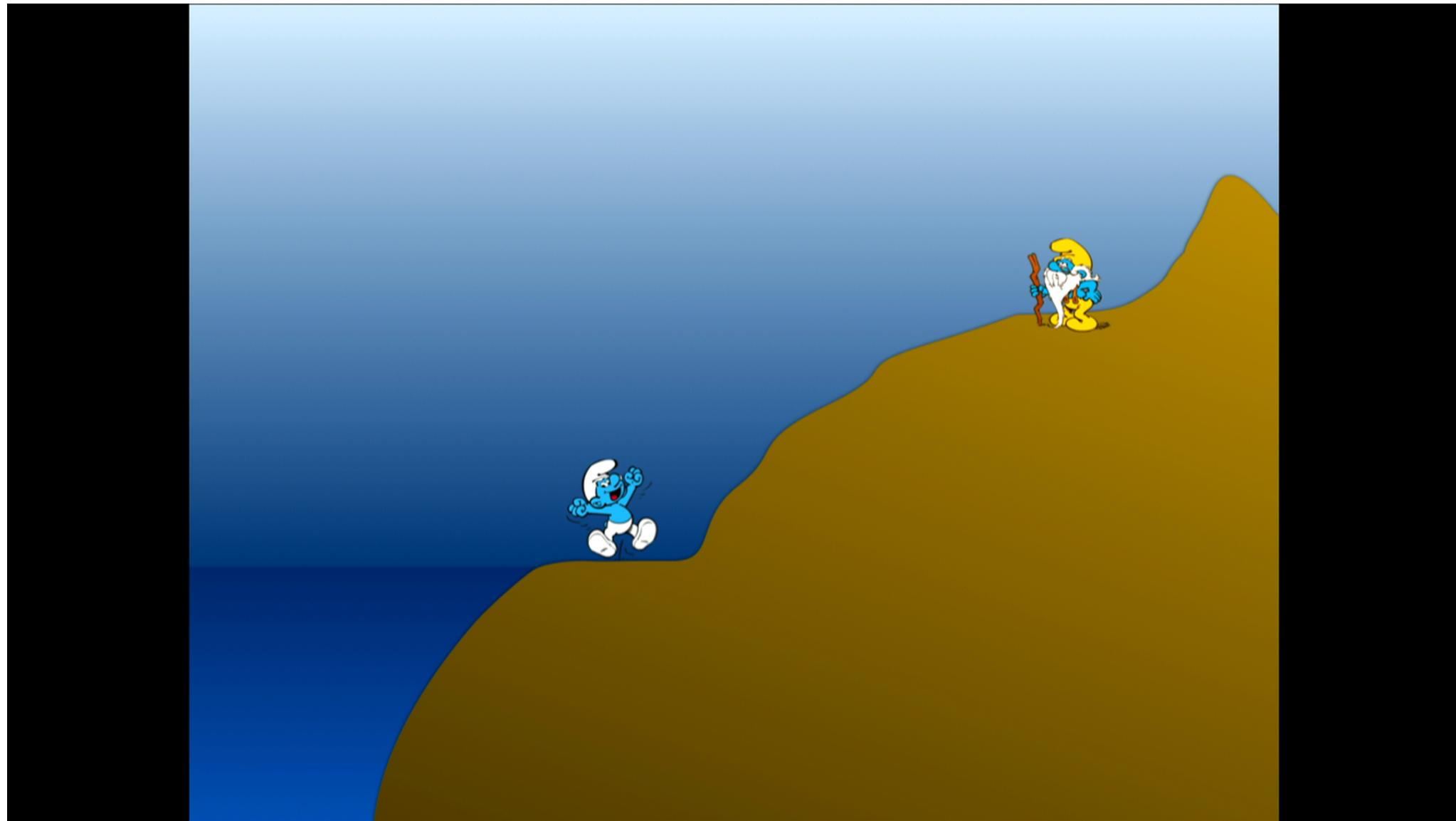
Carlo Rovelli, Francesca Vidotto

Int.J.Mod.Phys. D23 (2014) 12, 1442026

- δ is related to the distance from the horizon where the theory is entirely classical

What does δ represent and what determines it?





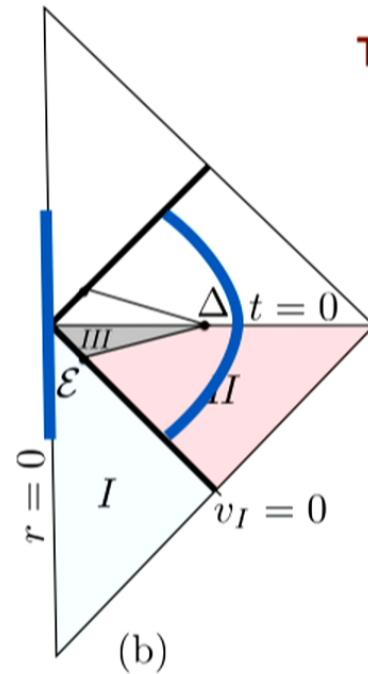
Time dilation

$$\tau_R = 2R - m \ln(\delta/m)$$

T: bounce time (very large)

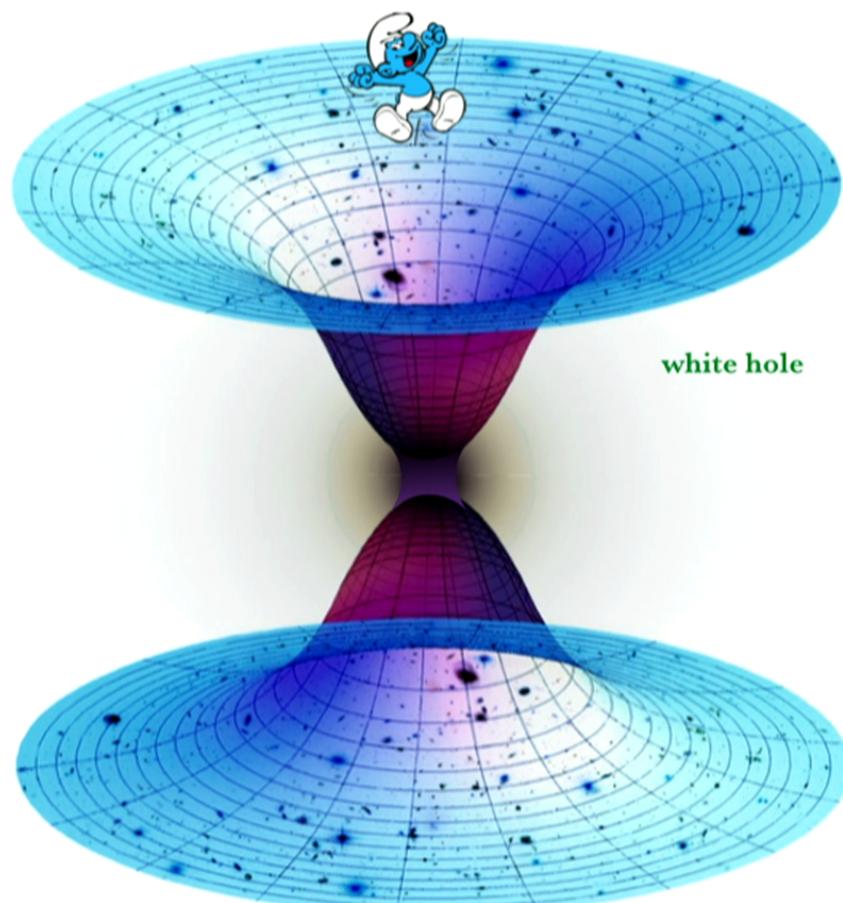
$$\tau_{internal} \sim m \sim 1ms$$

$$\tau_{external} \sim m^2 \sim 10^9 years$$



"A black hole is a short cut to the future"

Time inside: 1 ms

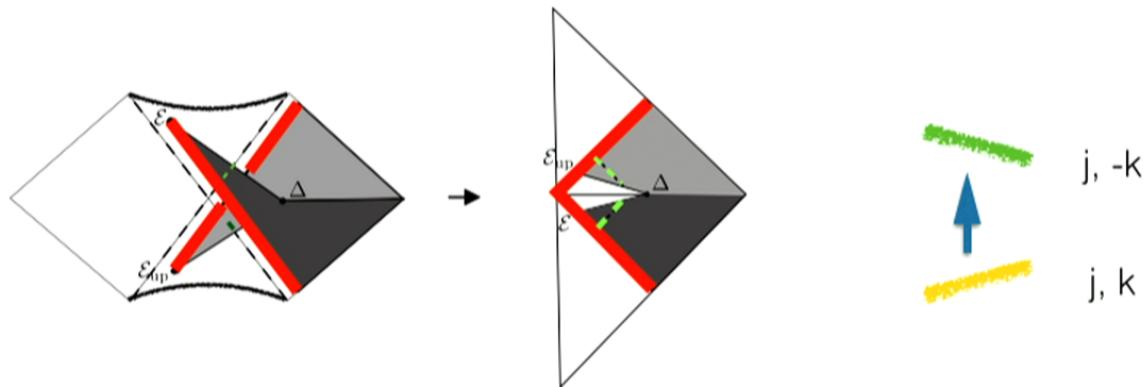


What do we expect of the bouncing time:

$$T = \begin{cases} \sim e^m & \text{Naive expectation from analogy with tunnelling in space} \\ \sim m^3 & \text{Page time. Requiring that AMPS firewall are avoided} \\ \sim m^2 & \text{Minimal failure of local qft: } RT > L_{Planck}^{-1} \\ \sim m \ln m & \text{Calculation from LQG, first contribution (too short!)} \end{cases}$$

Fast Radio Bursts and White Hole Signals
Aurélien Barrau, Carlo Rovelli, Francesca Vidotto.
Phys.Rev. D90 (2014) 12, 127503

- Covariant loop quantum gravity. Calculation of $T(m)$.

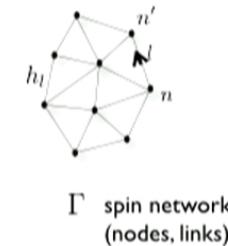


- Covariant loop quantum gravity. Full definition.

Kinematics Boundary

State space $\mathcal{H}_\Gamma = L^2[SU(2)^L / SU(2)^N]_\Gamma \ni \psi(h_l) \quad \mathcal{H} = \lim_{\Gamma \rightarrow \infty} \mathcal{H}_\Gamma$

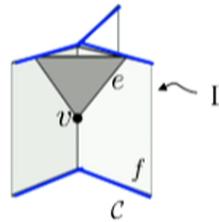
Operators [cfr: Steve]: $\vec{L}_l = \{L_l^i\}, i = 1, 2, 3 \quad L^i \psi(h) \equiv \frac{d}{dt} \psi(h e^{t \tau_i}) \Big|_{t=0}$



Dynamics Bulk

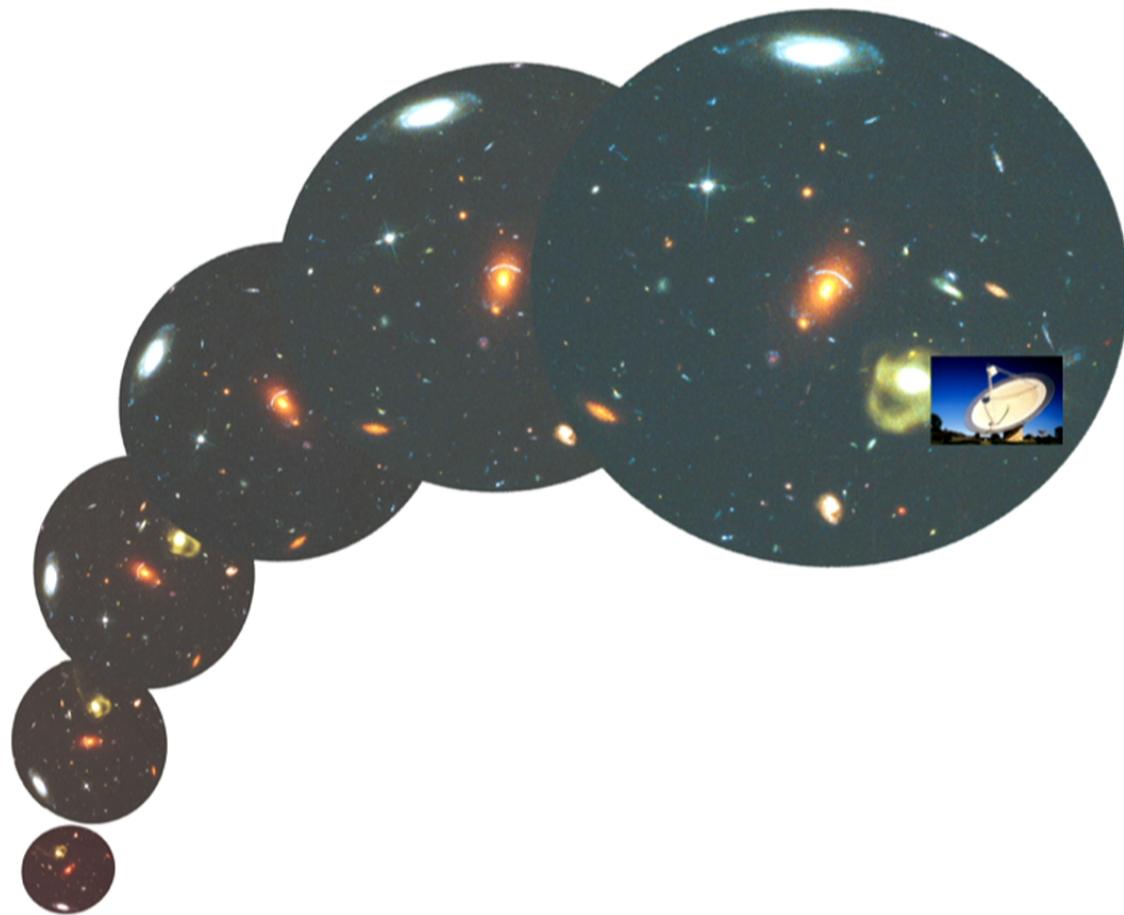
Transition amplitudes $W_C(h_l) = N_C \int_{SU(2)} dh_{vf} \prod_f \delta(h_f) \prod_v A(h_{vf}) \quad W = \lim_{C \rightarrow \infty} W_C \quad h_f = \prod_v h_{vf}$

Vertex amplitude $A(h_{vf}) = \int_{SL(2, \mathbb{C})} dg'_e \prod_f \sum_j (2j+1) D_{mn}^j(h_{vf}) D_{jmjn}^{\gamma(j+1), j}(g_e g_e^{-1}) \quad 8\pi\gamma\hbar G = 1$



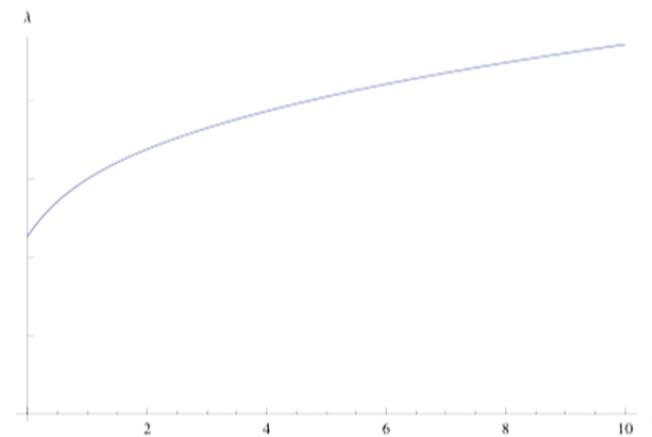
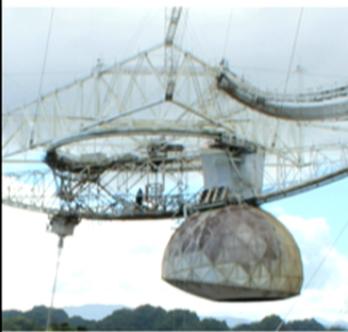
spin foam
(vertices, edges, faces)

SL(2,C) Chern-Simons Theory, a non-Planar Graph Operator, and 4D Loop Quantum Gravity with a Cosmological Constant
Hal M. Haggard, Muxin Han, Wojciech Kamiński, Aldo Riello.
[arXiv:1412.7546](https://arxiv.org/abs/1412.7546)



Signature: peculiar wavelength versus distance relation

$$\lambda_{obs} \sim \frac{2Gm}{c^2} (1+z) \sqrt{\frac{H_0^{-1}}{6k\Omega_\Lambda^{1/2}}} \sinh^{-1} \left[\left(\frac{\Omega_\Lambda}{\Omega_M} \right)^{1/2} (z+1)^{-3/2} \right].$$



Summary

- Technical results: black holes may tunnel to white holes locally.
- The tunnelling time can in principle be computed with LQG.
- Gamma-rays or Fast Radio Bursts possible phenomenology.
- Peculiar wavelength to distance signature.

