

Title: The Outside Story: String spreading in the S-matrix and black holes

Date: May 02, 2015 09:00 AM

URL: <http://pirsa.org/15050140>

Abstract:

Gary has had an enormous impact on the field in general. For me personally his work has been a huge inspiration. First some great collaborations:

HEP 4 records found Search took 0.15 seconds.

1. Inflation Physics from the Cosmic Microwave Background and Large Scale Structure

K.N. Abazajian, K. Arnold, J. Austermann, B.A. Benson, C. Bischoff, J. Bock, J.R. Bond, J. Borrill, I. Buder, D.L. Burke (Chicago U., KICP) *et al.*, Sep 20, 2013. 11 pp.

Published in *Astropart.Phys.* 63 (2015) 55-65

FERMILAB-PUB-13-442-A

DOI: [10.1016/j.astropartphys.2014.05.013](https://doi.org/10.1016/j.astropartphys.2014.05.013)

e-Print: [arXiv:1309.5381](https://arxiv.org/abs/1309.5381) [astro-ph.CO] | [PDF](#)

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)

[ADS Abstract Service](#); [Fermilab Library Server](#) (fulltext available); [Link to Fulltext](#)

[Detailed record](#) - Cited by 31 records

$\langle \gamma \gamma \rangle$

GH: t -dep background cosmology

2. Insightful D-branes

Gary Horowitz (UC, Santa Barbara), Albion Lawrence (Brandeis U. & Santa Barbara, KITP), Eva Silverstein (SLAC & Stanford U., Phys. Dept. & Santa Barbara, KITP). Apr 2009. 30 pp.

Published in *JHEP* 0907 (2009) 057

NSF-KITP-09-53, SLAC-PUB-13584, SU-ITP-09-16, BRX-TH-607

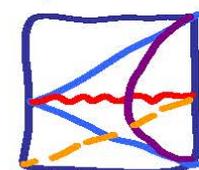
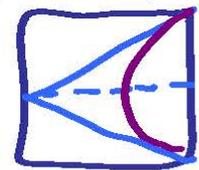
DOI: [10.1088/1126-6708/2009/07/057](https://doi.org/10.1088/1126-6708/2009/07/057)

e-Print: [arXiv:0904.3922](https://arxiv.org/abs/0904.3922) [hep-th] | [PDF](#)

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[CERN Document Server](#); [ADS Abstract Service](#); [SLAC Document Server](#)

[Detailed record](#) - Cited by 47 records



3. The Inside story: Quasilocal tachyons and black holes

Gary T. Horowitz (UC, Santa Barbara), Eva Silverstein (SLAC & Stanford U., Phys. Dept.). Jan 2006. 33 pp.

Published in *Phys.Rev.* D73 (2006) 064016

SLAC-PUB-11616, SU-ITP-06-01

DOI: [10.1103/PhysRevD.73.064016](https://doi.org/10.1103/PhysRevD.73.064016)

e-Print: [hep-th/0601032](https://arxiv.org/abs/hep-th/0601032) | [PDF](#)

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[ADS Abstract Service](#); [Phys. Rev. D Server](#); [SLAC Document Server](#)

[Detailed record](#) - Cited by 60 records 50+

4. Clean time dependent string backgrounds from bubble baths

Ofer Aharony (Weizmann Inst.), Michal Fabinger (Stanford U., Phys. Dept. & SLAC), Gary T. Horowitz (UC, Santa Barbara), Eva S

Published in *JHEP* 0207 (2002) 007

SLAC-PUB-9203, WIS-16-02-DPP

DOI: [10.1088/1126-6708/2002/07/007](https://doi.org/10.1088/1126-6708/2002/07/007)

e-Print: [hep-th/0204158](https://arxiv.org/abs/hep-th/0204158) | [PDF](#)

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[CERN Document Server](#); [ADS Abstract Service](#); [JHEP Electronic Journal Server](#); [SLAC Document Server](#)

[Detailed record](#) - Cited by 120 records 100+

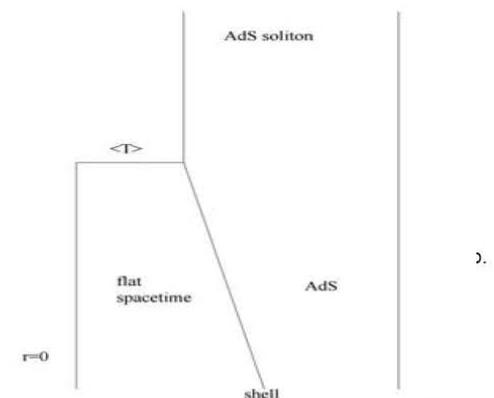


Fig. 1: A shell of D3-branes slowly contracts. The spacetime outside is approximately $AdS_5 \times S^5$, while the spacetime inside is approximately flat. The branes are wrapped around a Spherically-Symmetric circle, and when this circle reaches the string scale, the winding tachyons condense. The exterior geometry becomes a bubble which settles down to the AdS soliton (cross S^5). We will be interested in the fate of excitations in the $\langle T \rangle$ region.

Gary's paradigms like Calabi-Yau compactification, singularity resolution (or not) in various contexts, holographic SC, etc. have been extraordinarily stimulating

Vacuum Configurations for Superstrings

P. Candelas (Texas U. & Santa Barbara, KITP), Gary T. Horowitz (UC, Santa Barbara), Andrew Strominger (Princeton, Inst. Advanced Study), Edward Witten (Princeton U.). Mar 1985. 28 pp.
Published in *Nucl.Phys.* **B258** (1985) 46-74
NSF-ITP-84-170
DOI: [10.1016/0550-3213\(85\)90602-9](https://doi.org/10.1016/0550-3213(85)90602-9)

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)
[CERN Document Server](#); [Science Direct](#)

Special geometry

Lifshitz Singularities

Gary T. Horowitz, Benson Way (UC, Santa Barbara). Nov 2011. 14 pp.
Published in *Phys.Rev.* **D85** (2012) 046008
DOI: [10.1103/PhysRevD.85.046008](https://doi.org/10.1103/PhysRevD.85.046008)
e-Print: [arXiv:1111.1243](https://arxiv.org/abs/1111.1243) [hep-th] | [PDF](#)

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)
[CERN Document Server](#); [ADS Abstract Service](#)

Two matters for would-be tidal force singularity

• • •

THE NATURE OF SPACE AND TIME - PROFESSOR GARY T. HOROWITZ.

To se mi líbí 0



Of course Gary's main interest, now that he's 60, is:

3) Time machines are not obviously impossible

There are space-times in general relativity which describe time travel. But space-times can satisfy the equations of general relativity and still be "unphysical". One must examine them carefully.

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The Nature of Space and Time - Professor Gary T. Horowitz.



I am not suggesting that Gary is a homicidal maniac. However:

Time travel need not cause paradoxes

Famous question: What if you travel into the past and kill your grandmother?

The Outside Story: string spreading in the S-matrix and black holes

with M. Dodelson

2015 thought-expt data release

I. Longitudinal nonlocality in the string S-matrix

II. String spreading and

Outline

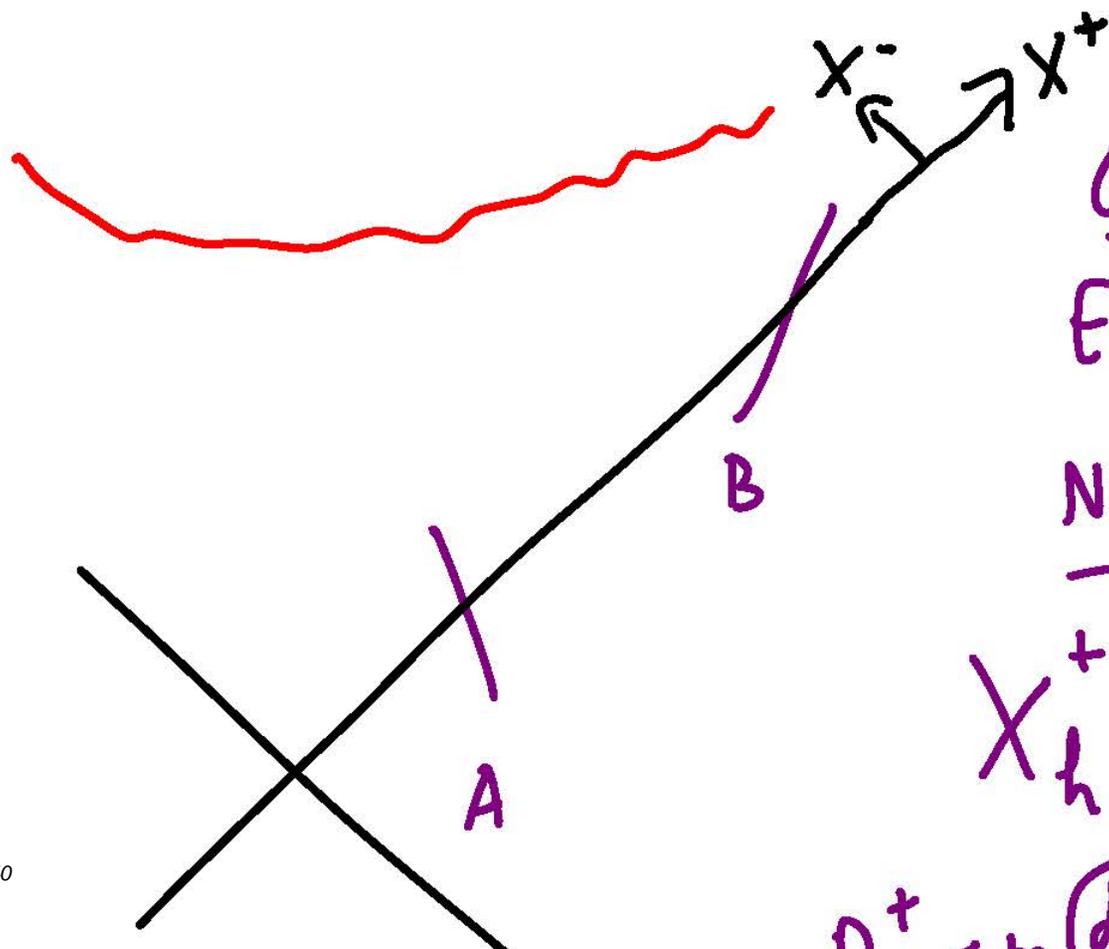
- I. Intro { Black Hole as accelerator
string spreading { Susskind '94
Brower-Pokhinski
Strassler-Tan '06
- II. Longitudinal Nonlocality in tree-level
String amplitudes from phases & trajectories
in { Veneziano
Bardakci-Ruegg
- III. Quantitative estimates for

Black Holes : conflict between EFT and Quantum Mechanics

... Hawking ... Mathur ... **AMPS** ...
↓
single observer can
see the conflict

→ Check level of violations of
EFT in string theory, with

$$ds^2 = -\frac{2r_s}{r} e^{\frac{1-r}{r_s}} dx^+ dx^- + r^2 d\Omega^2$$



Outside :
 E, m fixed

Near Horizon :

$$X_h^+ = 2r_s \sqrt{e} \frac{E}{m} e^\eta$$

$$d^+ \left(\sqrt{dx^+} \right) = m e^\eta$$

Near horizon: huge Energy, but
separated along X^+ .

String Spreading

- Susskind '94
- Brown Polchinski
- Strassler Tan '06

Light Cone gauge $X^- \sim p^- \tau$,

Constraint determines X^+ in terms of X^\perp

- Apparent asymmetry between X_{\perp} and X^{\pm} directions?

No: the RMS longitudinal spreading is detectable for

$X^{\pm} \approx$ direction of relative motion

More precisely: Brick wall frame

Light cone time resolution:

$$\Delta X^- \sim \frac{1}{P_{\text{detector}}^+} \quad N_{\text{max}} \sim \frac{g'}{\Delta T} \quad Y \sim \frac{X^-}{P_s^-}$$

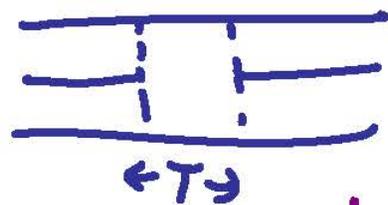
$$\Rightarrow N_{\text{max}} \sim P_s^- P_d^+$$

$$\left(\text{for } t \sim -k_{\perp}^2 \sim \mathcal{O}\left(\frac{1}{g'}\right) \right)$$

More generally,



☆ This physical idea is confirmed
explicitly in BPST '06 calculation
 of 4-point Regge amplitude in
 light-cone gauge.



in brick wall frame: $P_{\perp r} \sim \frac{k_{\perp}}{2}, m=0$

Note that the fact that X^+ is constrained in terms of X_{\perp} does not make it unphysical, there are many familiar examples (perhaps most basic being the expansion of the universe, in spacetime or on a string worldsheet)

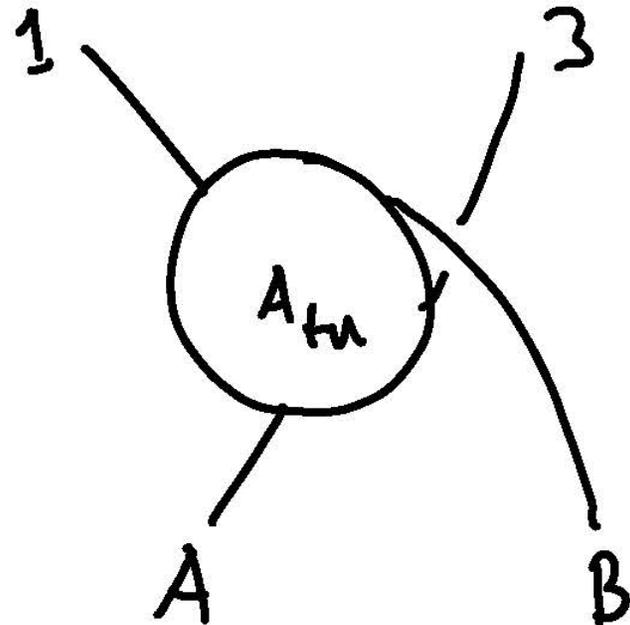
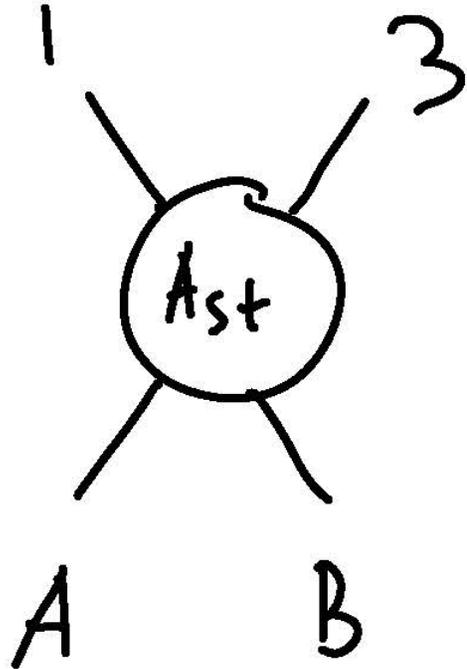
Let us take as given the transverse
spreading

$$\Delta X_{\perp} \sim \sqrt{r'} \log N_{\max}$$

$$\sim \sqrt{r'} \log \frac{S}{-t}$$

- Can be seen from impact parameter transform in forward scattering
- well-established in BPST

4 pt)
Veneriano

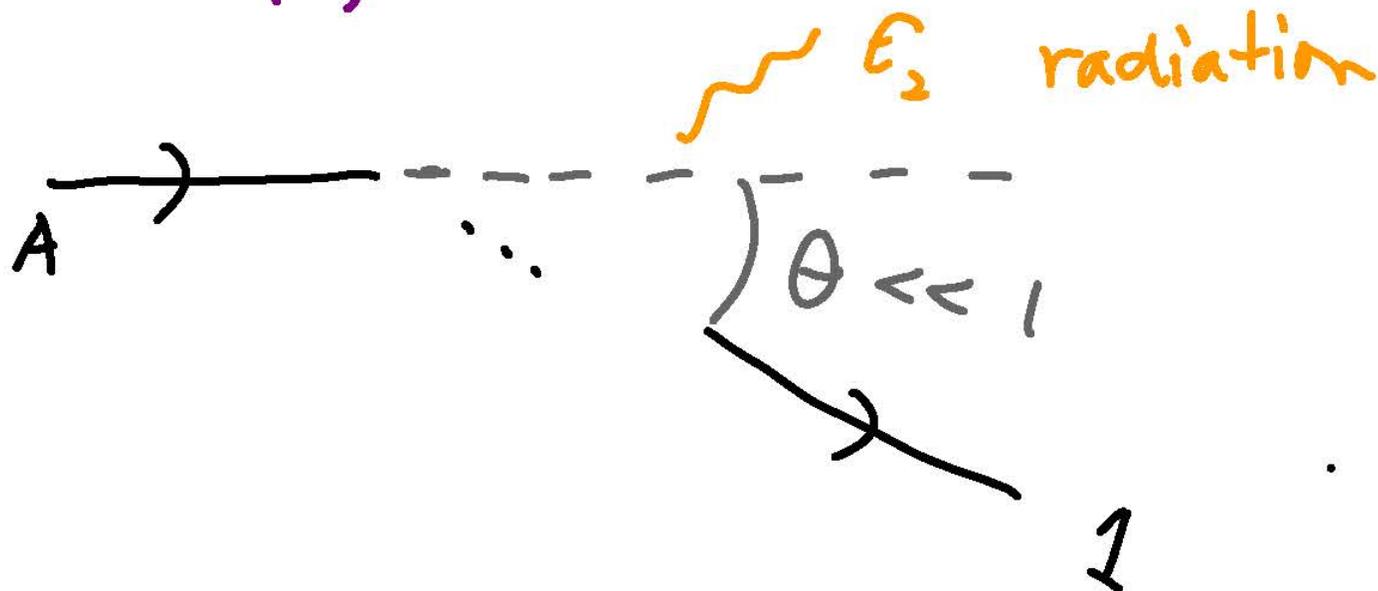


5 pt)
Bardaki - Rwegg
Bialas - Pokorski

S-matrix 'data' analysis

- Work in Regge or double Regge

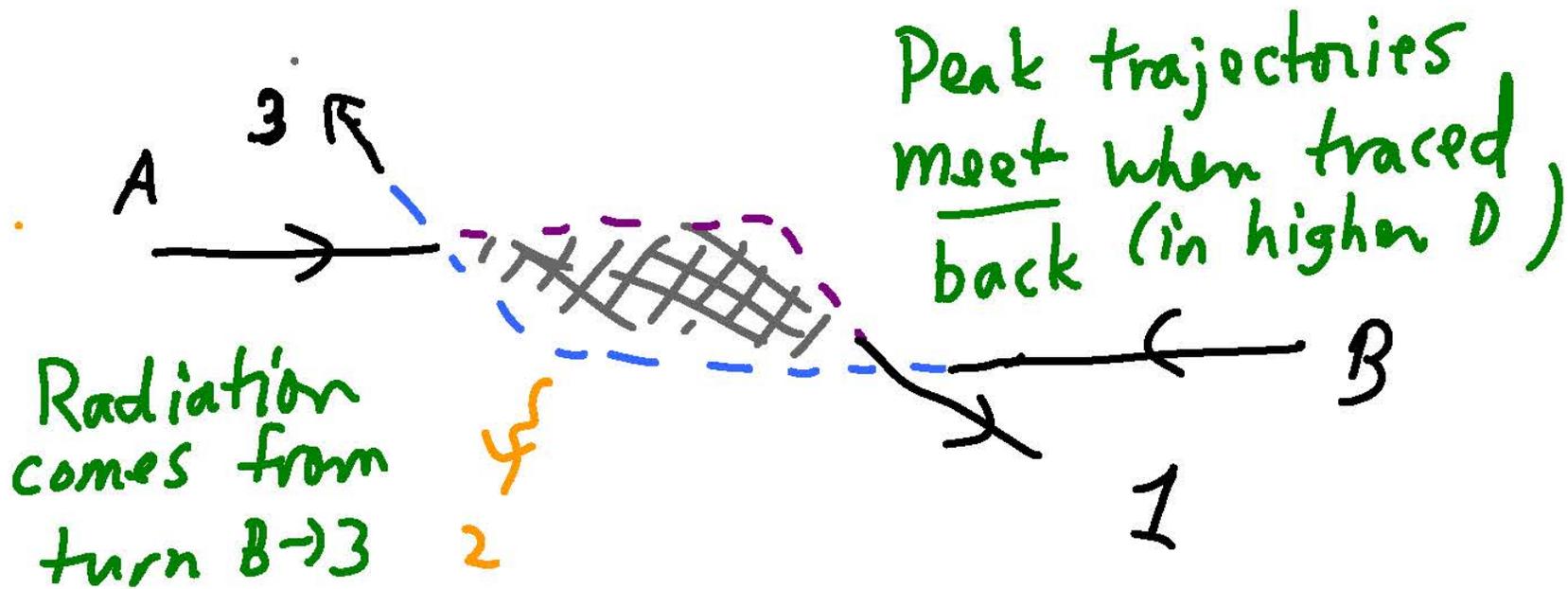
$$E\theta_1, E_2 \ll E$$



(1) Keep track of phases in
amplitudes $A \sim e^{i\sigma(k_I \cdot k_J)} A_{\text{slow}}$

in tractable regime, and

(2) convolve with wavepackets
to determine peak trajectories
peak impact parameter b_{AB}
time delays/advances T_{delay}



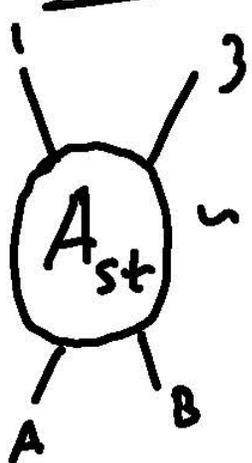
(3) Explicit & simple string solutions for intermediate S -channel states \leftrightarrow ^{imag.} parts & quantitative agreement with peak $b_2 T$

(1) ... and limited \perp

2 Main examples: 4 point:

$$A_{st} \sim g_0^2 \underbrace{\Gamma(-1 - g't)}_{e^{g't} \left(\frac{-s}{t}\right)^{g't}} (g's)^{1+g't} e^{-i\pi t} + \dots$$

$\frac{\sin(\pi g'(s+t))}{\sin(\pi g's)}$
 ← phase

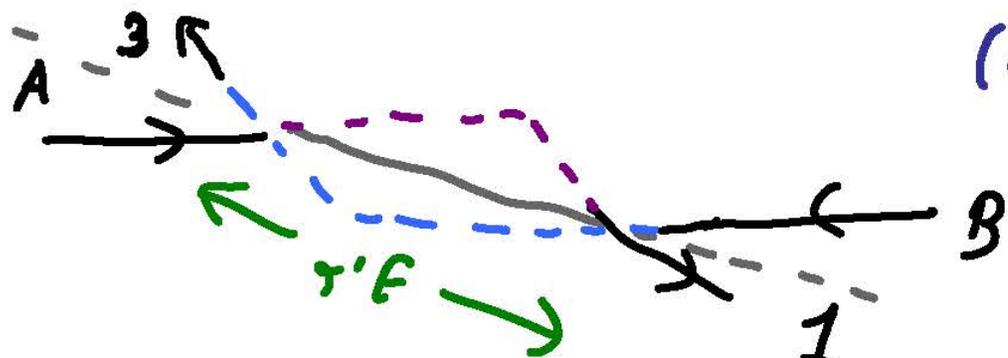


→ Peak trajectories

$$b_{AB} = -2\pi g' E \sin \theta_1 \neq 0!$$

$$T_1 = 2\pi g' E (1 - \cos \theta_1) \approx \pi E g'^2$$

time delay
(cf attractive potential)



The interaction timescale^{*} from
 BPST/GM is $\Delta X_{int}^- \sim \frac{g't}{P_A^+}$

During this time, A travels
 along X^+ a distance

$$\Delta X^+ \sim \Delta X^- P_A^{+2} \sim \Delta X^- P_A^{+2} g't^+$$

Alternate transverse effect?

(see appendix D)

Fourier transform

$$\int dq_{\perp} s^{-\alpha' q_{\perp}^2} e^{i\pi\alpha' q_{\perp}^2} e^{iq_{\perp} b} = \frac{\exp\left(-\frac{b^2}{4\alpha'(\log s - i\pi)}\right)}{\sqrt{\frac{\log s}{\pi} - i}},$$

• Not a real density

cf $A(\vec{q}) = A_{pt}(\vec{q}) F(\vec{q})$

source $\rho(\vec{r}) = |T_s(\vec{r})|^2$

$$F(\vec{q}) = \int d\vec{r} e^{i\vec{q}\cdot\vec{r}} \rho(\vec{r})$$

real



all directions

Interpret Fourier transform as wave function?

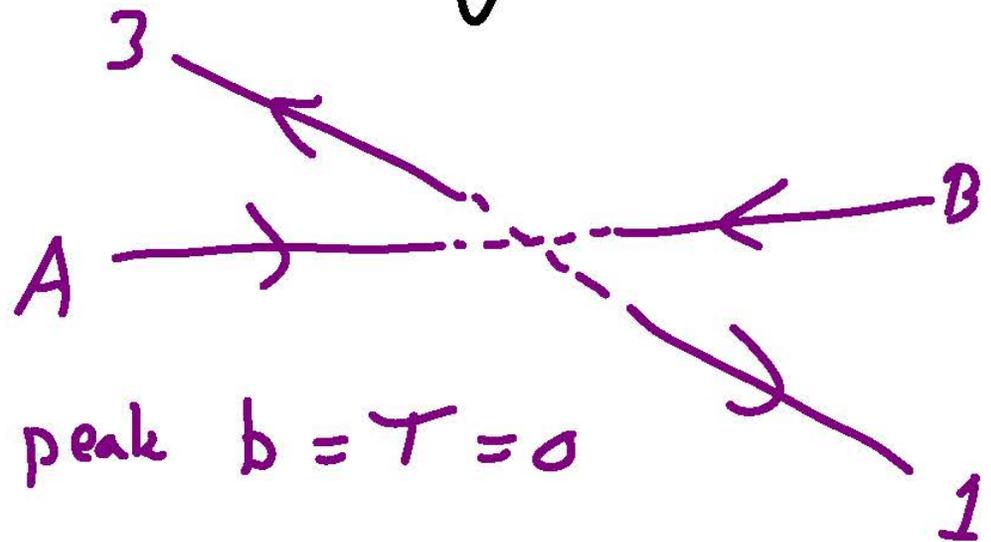
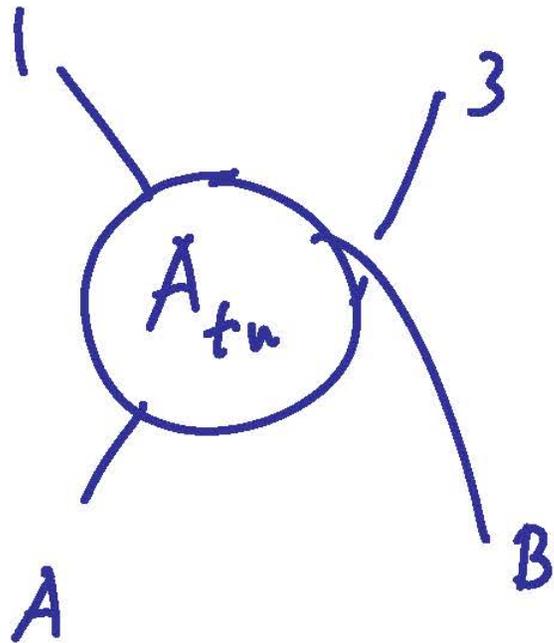
- Known ground state wave function

$$\psi_0(\Delta y) = N_0 e^{-\frac{(\Delta y)^2}{\sigma' \log \frac{n_{\max}}{n_0}}}$$

- Alternate hypothesis:

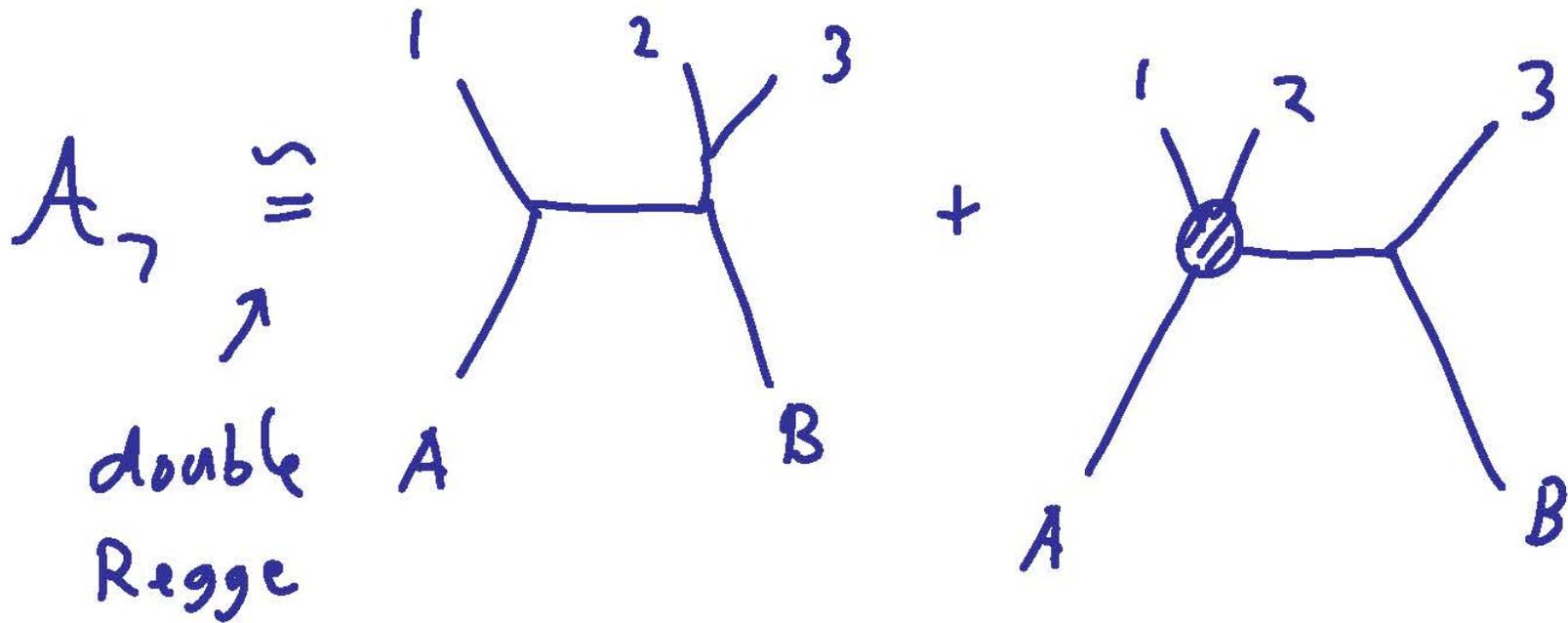
$$\psi_{\pi}(\Delta y) = N_{\pi} e^{-\frac{(\Delta y)^2}{\sigma' (\log n_{\max} - i\pi)}}$$

One of the 4 pt diagrams has zero time delay / advance



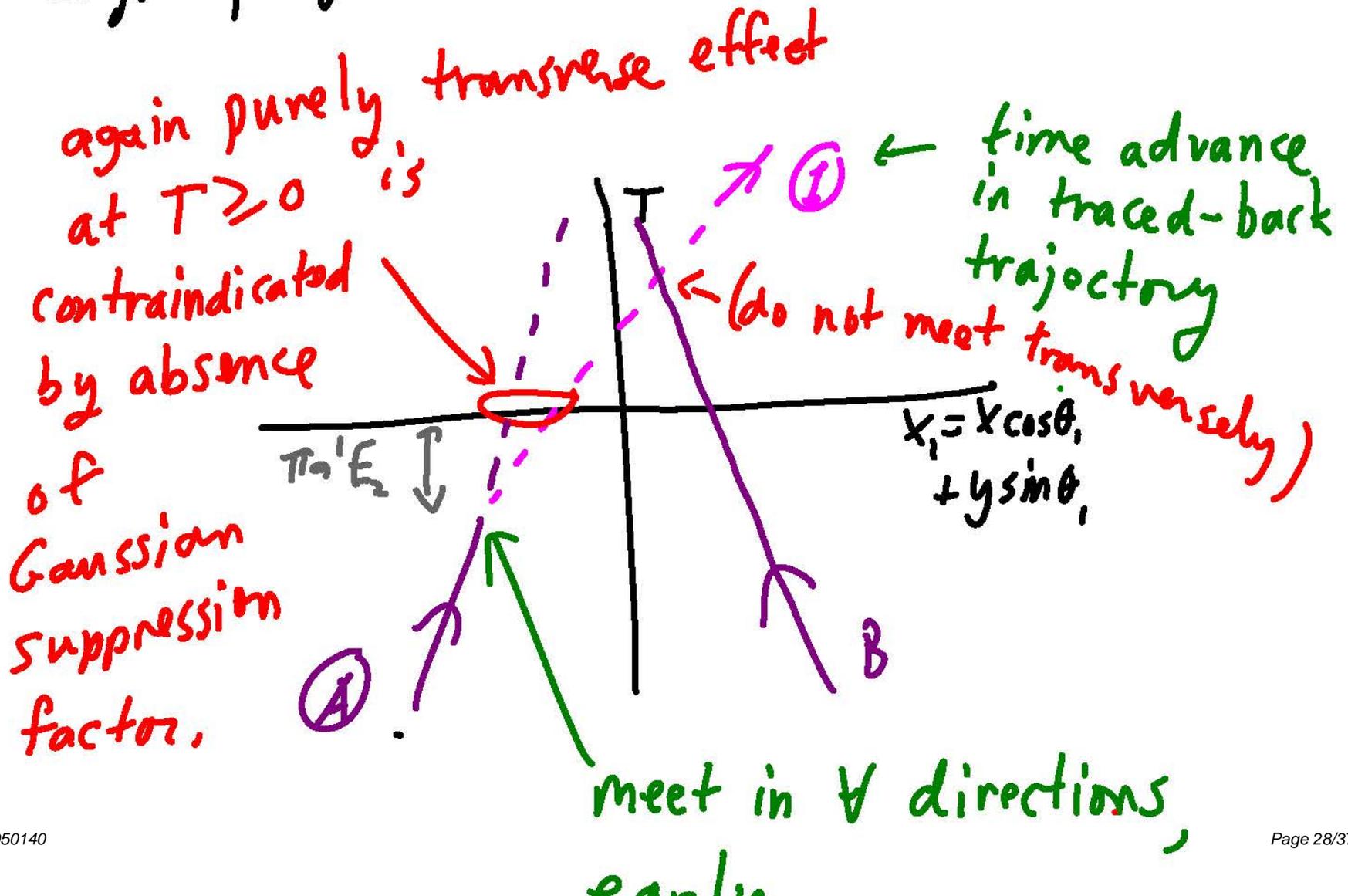
peak $b = T = 0$

In this one, when we upgrade



In a concrete regime with $A_7 \approx e$ $i\pi(k_{33} - k_{11}) \hat{A}_{slow}$,
 peak trajectories, traced back,
 such that

e.g. projected to 1+1 dimensions



Remarks

- These S-matrix amplitudes admit a simple interpretation in terms of long spreading; not captured by $\rho(X_{\perp})$.
- However, BH problem has additional complications: (weak) curvature; strings appear in near-horizon

Horizon physics



- In Schwarzschild BH, find (including decays/secondary probes)

$$\Delta X^+ \sim p_{\text{detector}}^+ q' \Rightarrow \text{detectable}$$

spreading for $m_{\text{det}} > \frac{r_s}{q'} \frac{E_{\text{det}}}{m}$

→ breakdown of EFT for late infalling detector given

Remarks

★ despite $Rg' \ll 1$, BH accelerates trajectories to generate large near-horizon relative boost

(Not simply Minkowski/Rindler dynamics)

- Causal (source string always spread, detector develops the.

- relative boost sets in outside

horizon; neither strict 'firewall'
AMPS

nor low-energy 'non violent nonlocality'
Giddings

but has elements of each.

Also timescale $\Delta t \gg 2r_s$
short compared to Page time

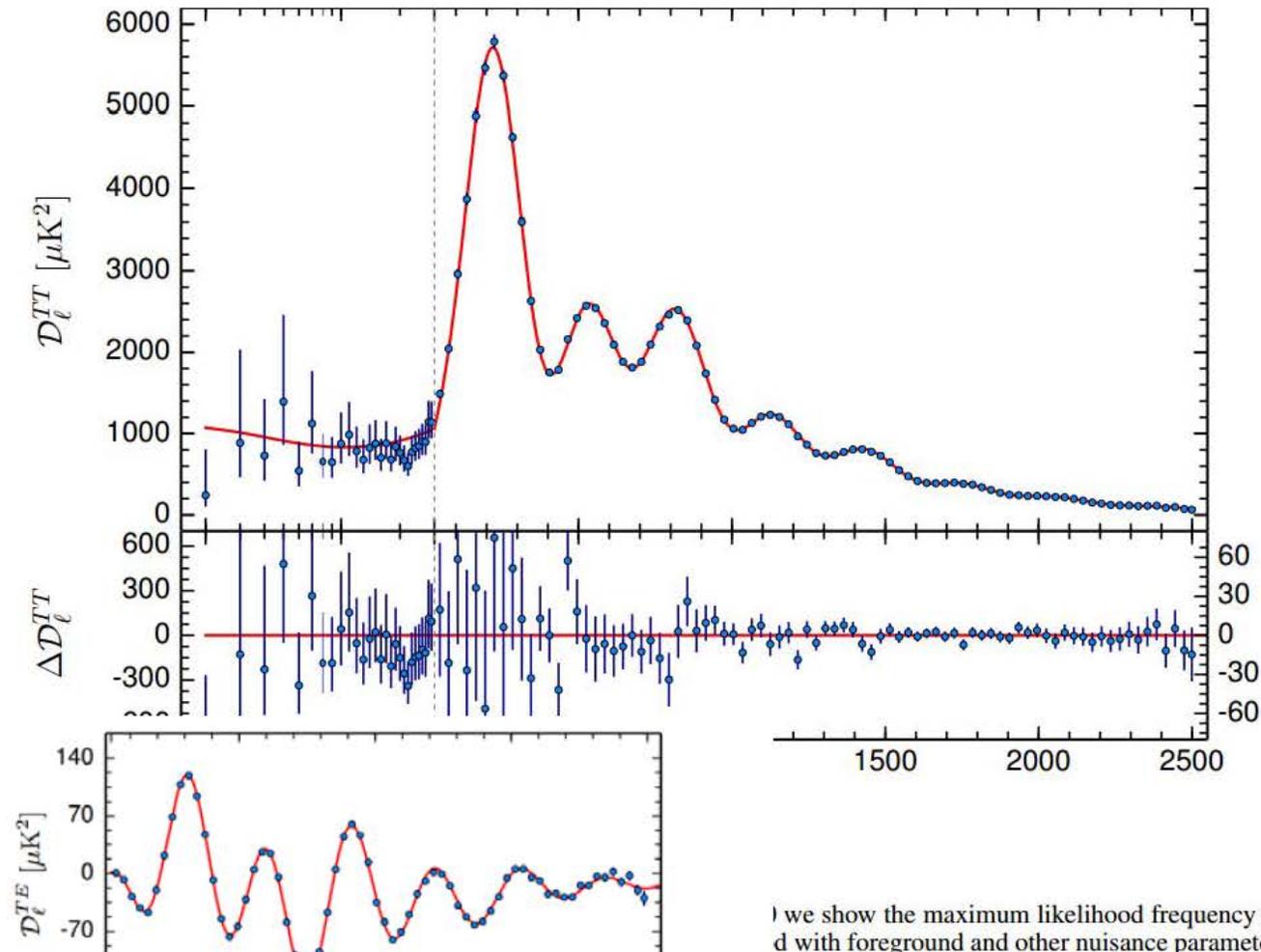
- Observational tests ??

Cosmological Horizons?

observer -
dependent

Planck Collaboration: Cosmological parameters

...



we show the maximum likelihood frequency averaged
d with foreground and other nuisance parameters deter-

Cosmological horizons

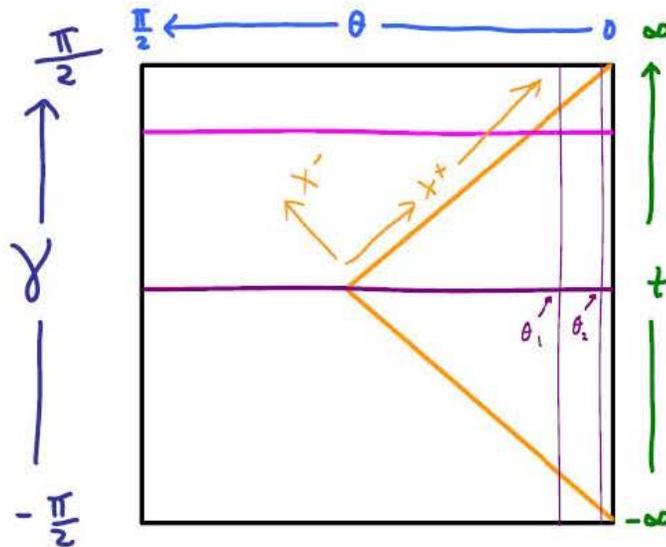


Figure 6: Trajectories 1 and 2 in the late de Sitter universe, as described in the text. For small values of the global spatial coordinate θ , the trajectories fall across the indicated observer horizon at a late global time, so that the spatial slices are nearly flat as in our observed universe. Within that regime, the hierarchy $\frac{\theta_2}{\theta_1} \ll 1$ leads to a large relative boost at the horizon, generated by the cosmological background.

l. i. ... : At of order L_{ds}

Future directions

- Other regimes of string and QG effects, e.g. single-Regge, 6-pt, ...
- background fields (linear dilaton, tachyon wall, AdS^* , electromagnetic flds, etc.) may further help tease out/test the longitudinal nonlocality
- 5-pt function of Shenker-Stanford

Happy Birthday Gary!

3) Time machines are not obviously impossible

There are space-times in general relativity which describe time travel. But space-times can satisfy the equations of general relativity and still be “unphysical”. One must examine them carefully.

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