

Title: Limitations of statistical mechanics for quantum thermodynamics.

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Abstract: How should we describe the thermodynamics of extreme quantum regimes, where features such as coherence and entanglement dominate?

I will discuss possible limitations of a traditional statistical mechanics approach, and then describe work that applies modern techniques from the theory of quantum information to the foundations of thermodynamics. In particular I discuss recent progress in quantum resource theories and argue that to properly encapsulate the thermodynamic structure of quantum coherence and entanglement we must make use of concepts beyond free energies.

# Limitations of statistical mechanics for quantum thermodynamics.

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# Trajectory

- Motivations + background.
- Possible limitations of traditional approaches.
- Resource theories & symmetry principles.
- Quantum coherence in thermodynamics.

(Work with M. Lostaglio, K Korzekwa, T. Rudolph, J. Oppenheim, M. Frenzel.)

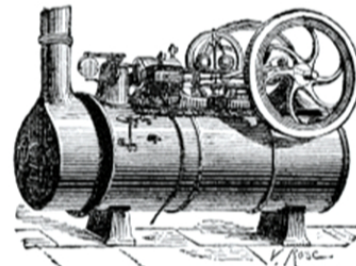
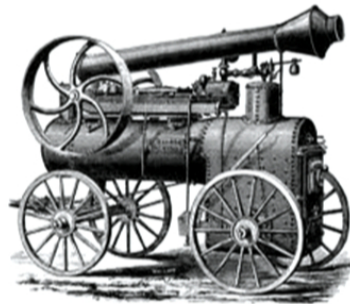
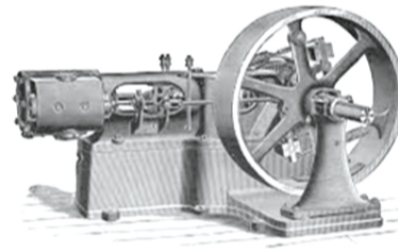
# 2nd Law of Thermodynamics

- “It is impossible to construct a device who’s sole effect is the extraction of work from heat.”
- “It is impossible to construct a device who’s sole effect is the erasure of a bit.”
- “It is impossible to see inside a furnace, solely by the light of the furnace.”\*



\* Bennett (1987)

# Limitations of existing thermodynamics



# The Thermodynamic Limit



- “*Thermodynamics means the thermodynamic limit.*”
- (Except it doesn't)

(a) Structure of kinesin

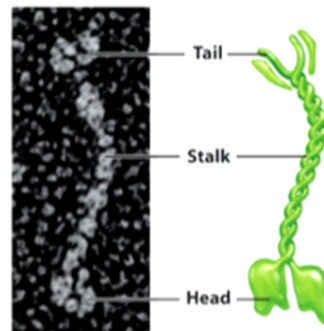
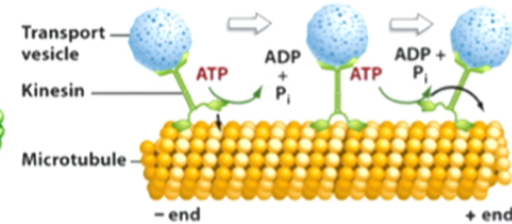


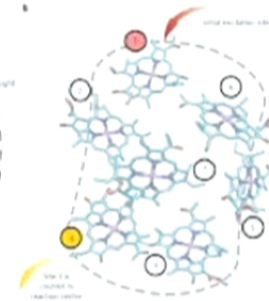
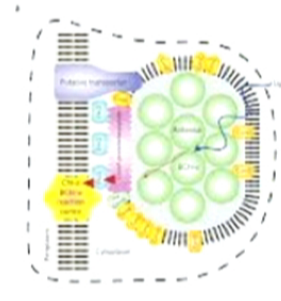
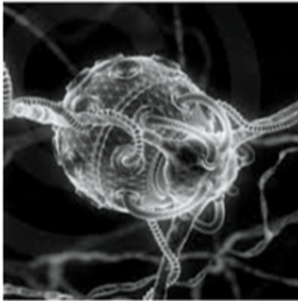
Figure 7-37 Biological Science, 2/e

(b) Kinesin “walks” along a microtubule track.



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# Non-asymptotic thermodynamics

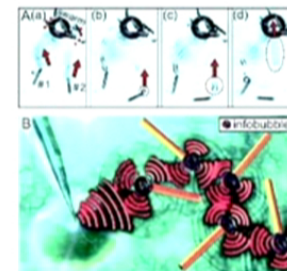


- **Q: What thermodynamic laws operate at micro/nano/pico/... scales?**
- **Q: How do coherent superpositions extend thermodynamic processes?**

# Applications

- Active work to develop nanoscale thermodynamic machines.
- Nanotechnology ~\$6 billion (currently)

***What are the fundamental thermodynamic laws beyond the thermodynamic limit?***

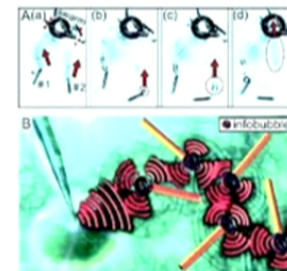




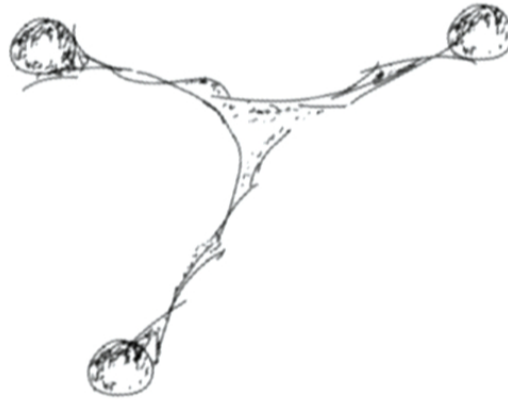
# Applications

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***What are the fundamental thermodynamic laws beyond the thermodynamic limit?***

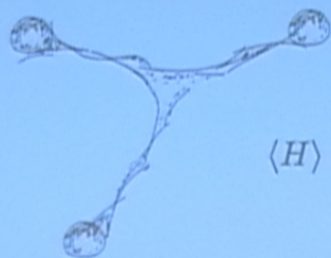


# Extreme Regimes



- **Determine the thermodynamics of highly entangled quantum systems in extreme regimes.**

# Extreme Regimes



$\langle H \rangle$  and "ensemble of microstates" ??

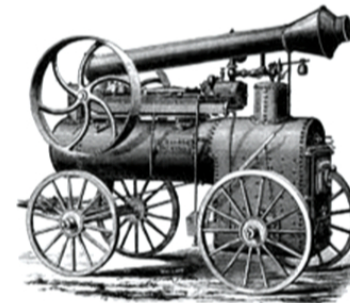
- Determine the thermodynamics of highly entangled quantum systems in extreme regimes.

*Potential limitations of  
more traditional analysis*

# Traditional 2nd Law

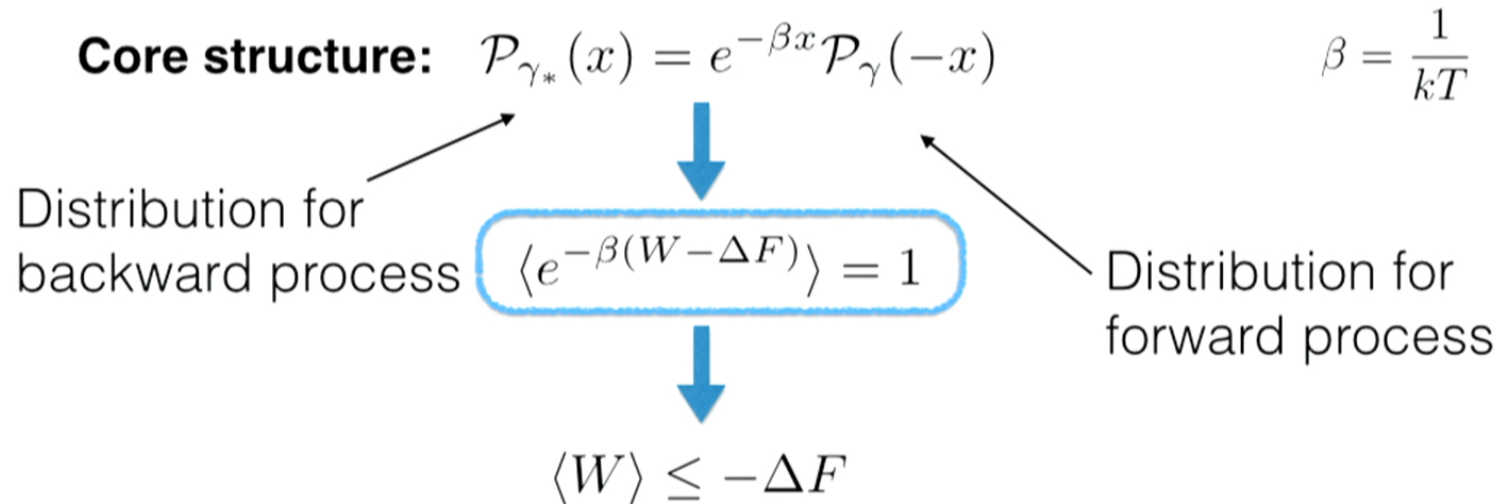
- “It is impossible to construct a device who’s sole effect is the extraction of work from heat.”

$$\langle W \rangle \leq -\Delta F$$



# Fluctuation Theorems?

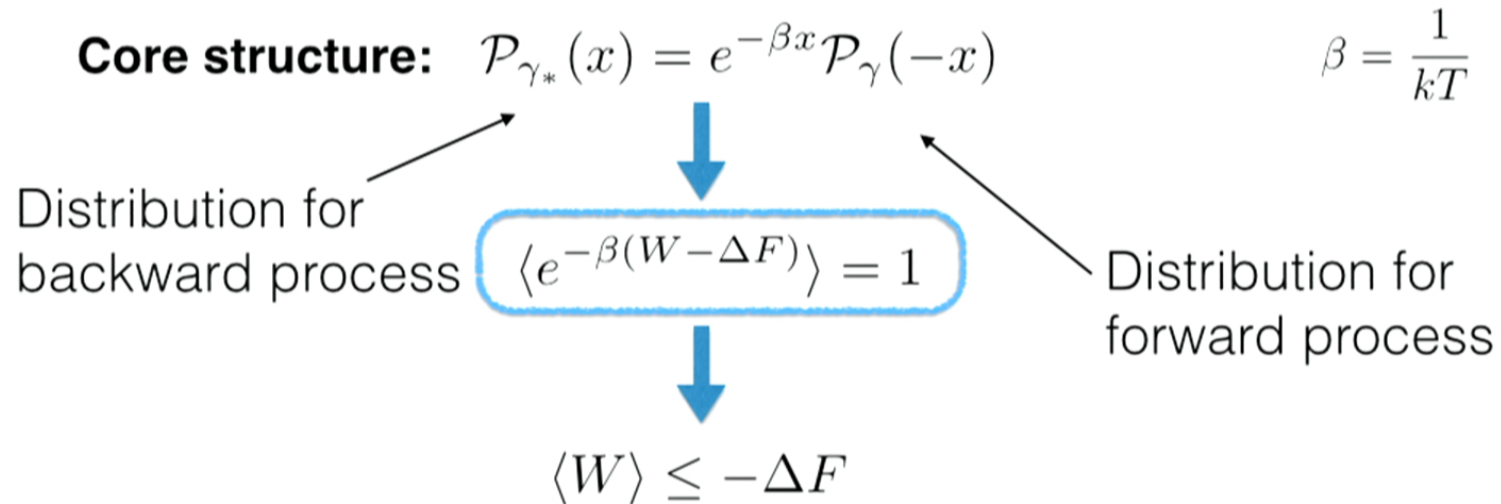
- **Arbitrarily violent dynamics** on thermal state.
- Sharpening of 2nd Law to an **equality**.



\* C. Jarzynski, *Phys. Rev. Lett.* (1997)

# Fluctuation Theorems?

- **Arbitrarily violent dynamics** on thermal state.
- Sharpening of 2nd Law to an **equality**.



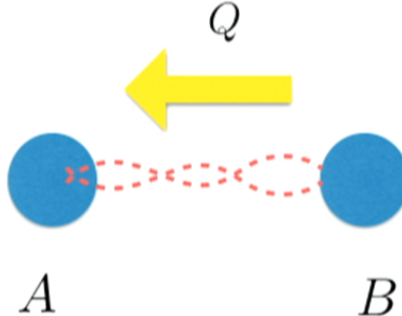
\* C. Jarzynski, *Phys. Rev. Lett.* (1997)

# Fluctuation Theorems

$$\langle e^{-\Delta\beta Q} \rangle = \gamma := \langle e^{-\Delta\mathcal{I}} \rangle$$

Energy exchange  $\nearrow$   $\langle e^{-\Delta\beta Q} \rangle$

$\nwarrow$  Initial correlations  $\langle e^{-\Delta\mathcal{I}} \rangle$



$$\langle \Delta\mathcal{I} \rangle = I_c(A : B)$$

$$\rho_{AB} \rightarrow \sum_{j,k} \Pi_{j,k} \rho_{AB} \Pi_{j,k}$$

Only **classical** correlations contribute.

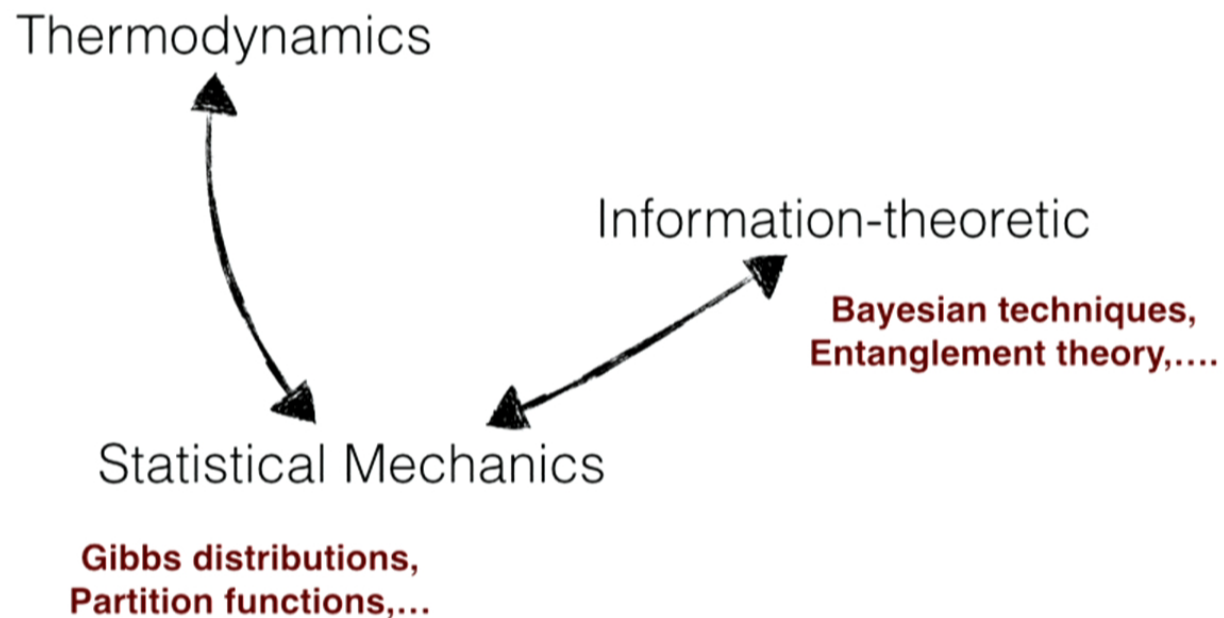
Poorly suited to handling **coherence** and **entanglement**.

\* DJ, et al, Phys. Rev. E (2015)

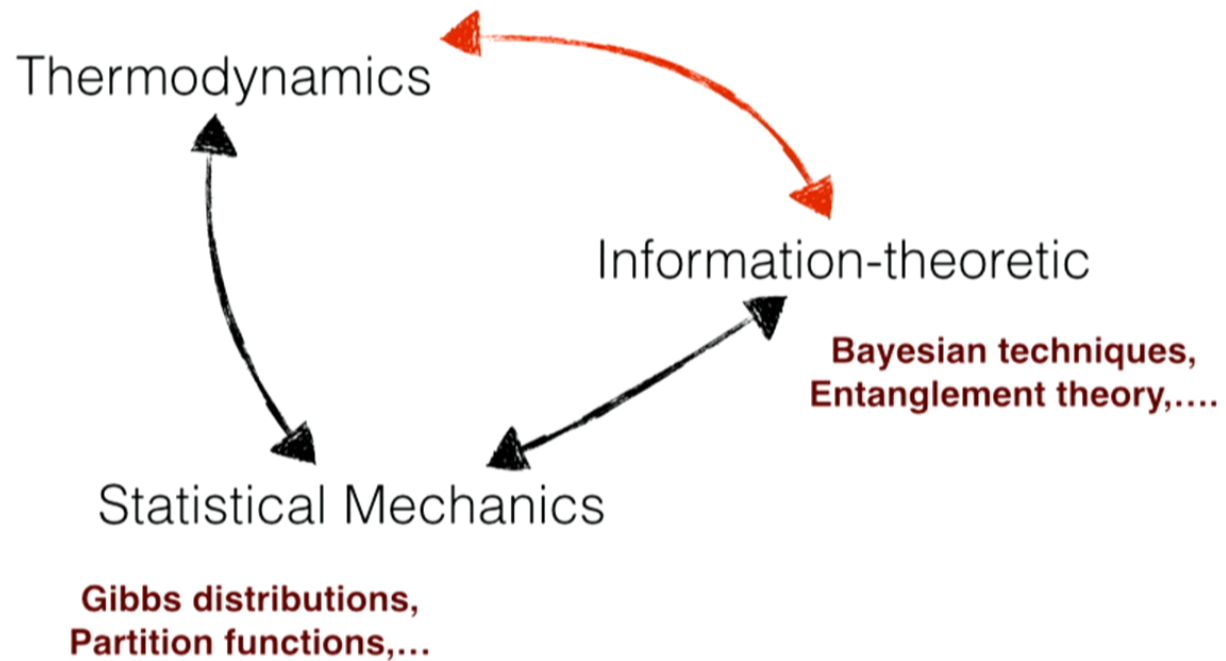


how *should* we handle  
intrinsically quantum-  
mechanical properties?

# Maximum entropy principles?



# Maximum entropy principles?



# A different kind of analysis

- **“Heat”, “temperature”** — ambiguous/complex/indirect...or maybe not even definable.
- **Giles (1964)**: *thermodynamics ultimately concerns the accessibility/inaccessibility of one physical state from another.*

*\*“The mathematical foundations of thermodynamics”, R. Giles (1964)*

# A different kind of analysis

- **Thermodynamics**: *the accessibility/inaccessibility of one physical state from another under particular operations.*
- **Entanglement theory**: *the accessibility/inaccessibility of one quantum state from another under Local operations and Classical Communications.*

# Primitive measures

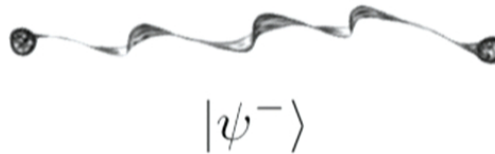
A single unit of “entropy”



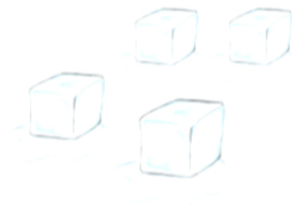
1g of ice @ 0 Celsius

1g of water @ 0 Celsius

A single unit of entanglement

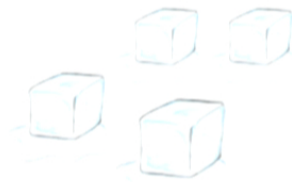


# An adiabatic process



Star

# An adiabatic process



Maximum # ice cubes  
that can be melted?



Star



Supernova



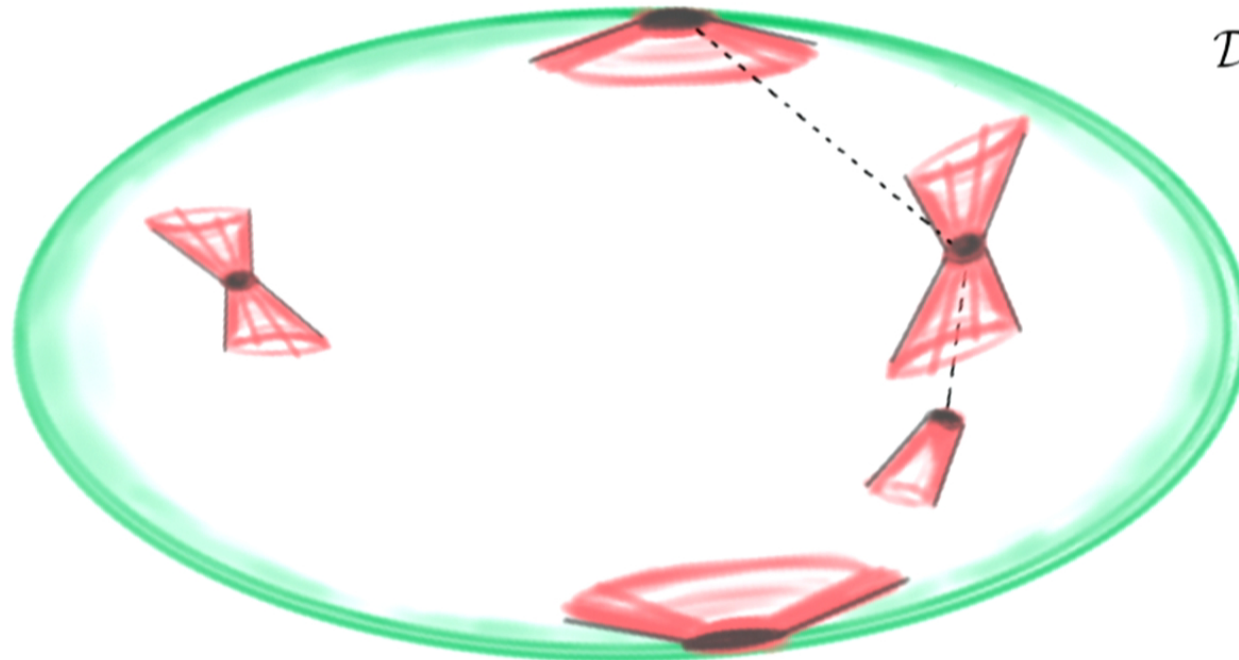
Remnant



# Thermodynamics

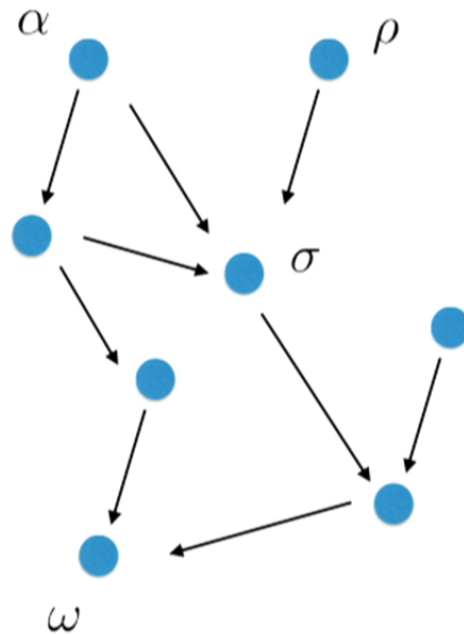
maximally ordered states

$\mathcal{D}(\mathcal{H})$



maximally disordered states

# Thermodynamic Entropy



When does thermodynamic ordering have a unique entropic formulation?

$$\rho \rightarrow \sigma$$
$$\Leftrightarrow S(\rho) \leq S(\sigma)$$

**A unique additive entropy exists  
if and only if the following 7 conditions hold:**

<b>Reflexivity</b>	$\rho \rightarrow \rho$
<b>Transitivity</b>	$\rho \rightarrow \sigma$ and $\sigma \rightarrow \tau$ implies $\rho \rightarrow \tau$
<b>Consistency</b>	$\rho_1 \rightarrow \sigma_1$ and $\rho_2 \rightarrow \sigma_2$ then $(\rho_1, \rho_2) \rightarrow (\sigma_1, \sigma_2)$
<b>Scale invariance</b>	$\rho \rightarrow \sigma$ then $\rho^{\otimes t} \rightarrow \sigma^{\otimes t}$ for $t \geq 0$
<b>Splitting</b>	$\rho \leftrightarrow (\rho^{\otimes t}, \rho^{\otimes(1-t)})$
<b>Stability</b>	$(\rho, \epsilon_1) \rightarrow (\sigma, \epsilon_2)$ then $\rho \rightarrow \sigma$
<b>Comparability</b>	if $\alpha \rightarrow \rho$ and $\beta \rightarrow \rho$ then $\alpha \rightarrow \beta$ or $\beta \rightarrow \alpha$

\* E. Lieb & J. Ingvason *The physics and mathematics of the second law* (1999)

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# MINING FOR DEEPER INSIGHTS

I THINK WE BROUGHT  
THE WRONG TOOLS.



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THE WRONG TOOLS.



*Learn from Entanglement Theory*

# The Quantum Resource Theory Recipe:

**Step 1: Define “free states”**

**Step 2: Define “free operations”**



1. **Ordering** of accessible/inaccessible states
2. **Resource states.**
3. Resource **measures.**

# Entanglement



1. **Free states: separable states**  $\sum_k p_k \rho_k \otimes \sigma_k$

2. **Free quantum operations:**

**Local Operations & Classical  
Communications**

**—Entanglement is non-increasing—**



# *Some very recent resource directions*

- Weilenmann, Kramer, Faist, Renner (2015)
- Fritz (2015)
- Brandao, Gour (2015)
- Coecke, Fritz, Spekkens (2014)
- and others.....

# Resource Theory of Thermodynamics

1. **Free states: thermal states**  $\gamma = e^{-\beta H} / Z$

2. **Free quantum operations:**

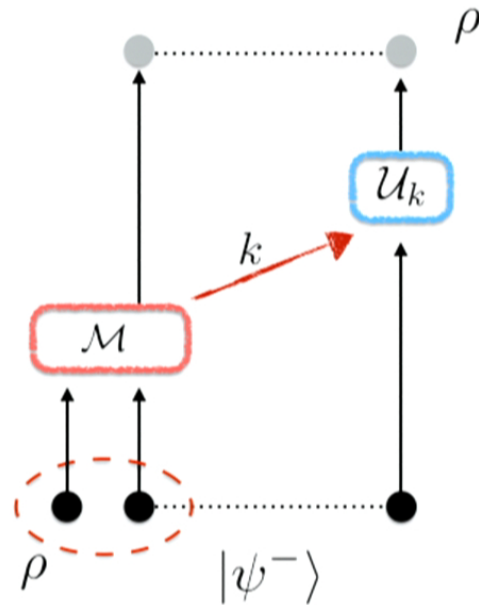
$$\mathcal{E}(\rho) = \text{tr}_b[U(\rho \otimes \gamma_b)U^\dagger]$$

$$[U, H_{\text{tot}}] = 0$$

—athermality is non-increasing—

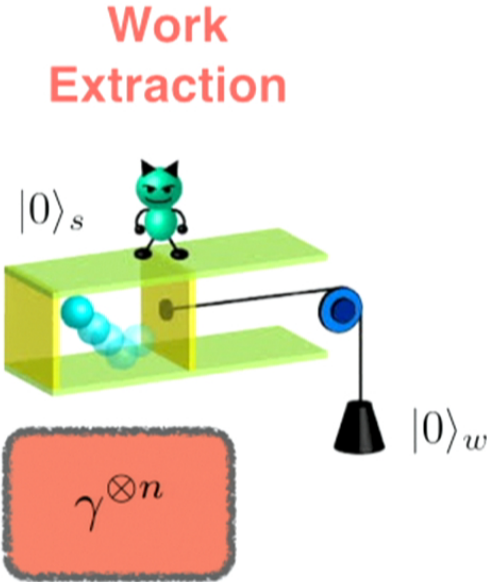
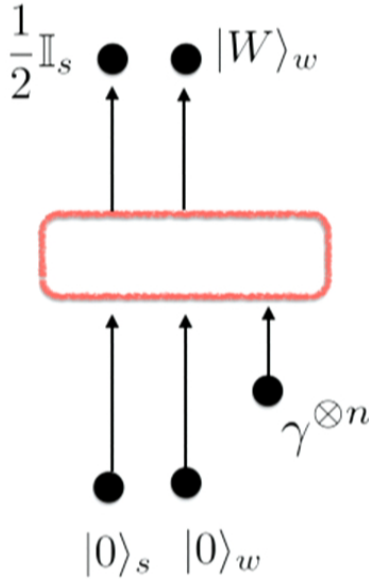
\* Brandao et al, Phys. Rev. Lett. (2013)

# Use of resources: Entanglement

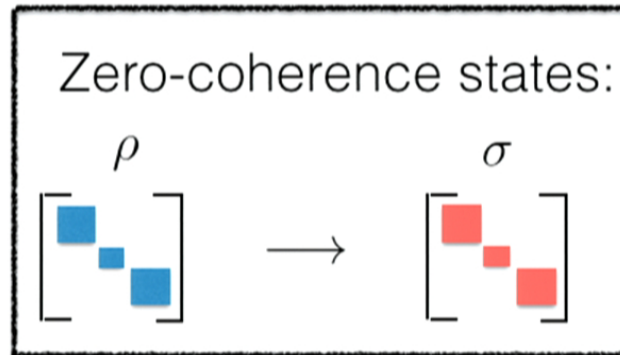
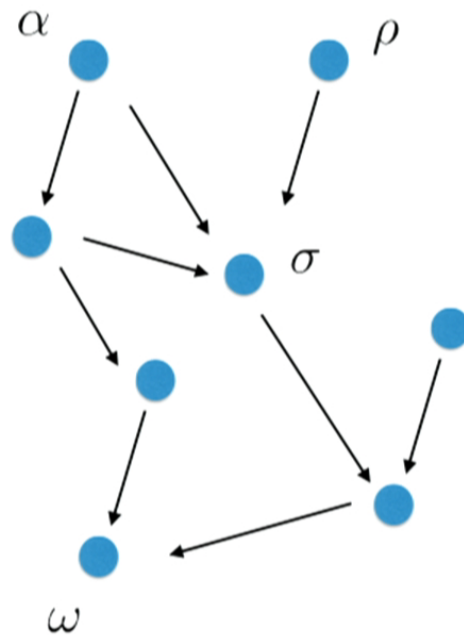


**Quantum  
Teleportation**

# Use of resources: Thermodynamics



# Mapping the general thermal structure



**Q: Does the ordering of states admit an entropic formulation?**

# “A Set of Second Laws” (Jonathan's talk!)

**Theorem:** For zero coherence states, the transformation  $\rho \rightarrow \sigma$  is possible

if and only if  $F_\alpha(\rho) \geq F_\alpha(\sigma) \quad \forall \alpha$

A **set** of entropy conditions

**Renyi-divergences:**  $D_\alpha(\rho||\sigma) = \frac{1}{\alpha - 1} \log [\text{tr}(\sigma^\kappa \rho \sigma^\kappa)^\alpha]$

$$F_\alpha(\rho) := D_\alpha(\rho||\gamma)$$

$$\kappa = \frac{1 - \alpha}{2\alpha}$$

\* Brandao et al, PNAS (2014)

**But what about  
Quantum Coherence  
& Thermodynamics?**

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Quantum Coherence  
& Thermodynamics?**



# Symmetry & the 1st Law of Thermodynamics

- Traditional form:  $dE = dQ + dW$
- Microscopic energy conservation (system+bath).

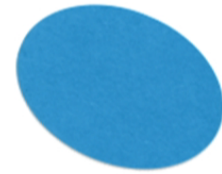
## Quantum Mechanical Symmetry:

$$[U, H_{\text{tot}}] = 0$$

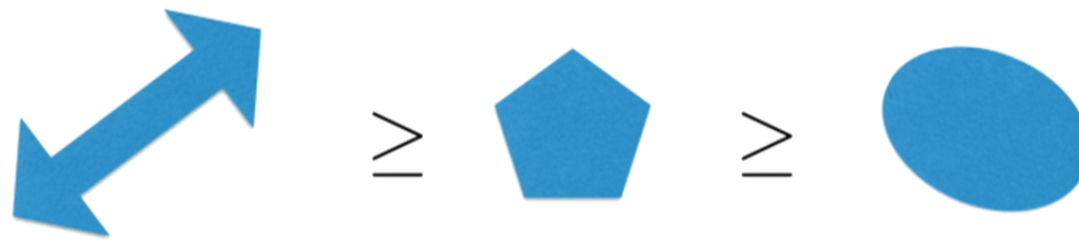
$$t \mapsto e^{-itH_{\text{tot}}}$$

Constrains **non-conservation** of **two** quantities:

- (a) System energy
- (b) System “coherence”



*When is A is more asymmetric than B?*



# A theory of asymmetry

*Related papers:*

- *I. Marvian, R. Spekkens Phys. Rev. A 90, (2014)*
- *I. Marvian, R. Spekkens, New J. Phys. 15, (2013)*
- *G. Gour, I Marvian, R. Spekkens Phys. Rev. A 80, 012307 (2009)*
- *M. Ahmadi, DJ, T. Rudolph, New. J. Phys. 15 (2013)*

# A theory of asymmetry

- “Group-theoretic Anna Karenina Principle”:

*“all symmetric objects are alike; each asymmetric object can be asymmetric in its own way.”*

# A theory of asymmetry

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# Resource Theory of Asymmetry

$$U : G \rightarrow U(g)$$

1. **Free states: symmetric states**  $U\rho U^\dagger = \rho$

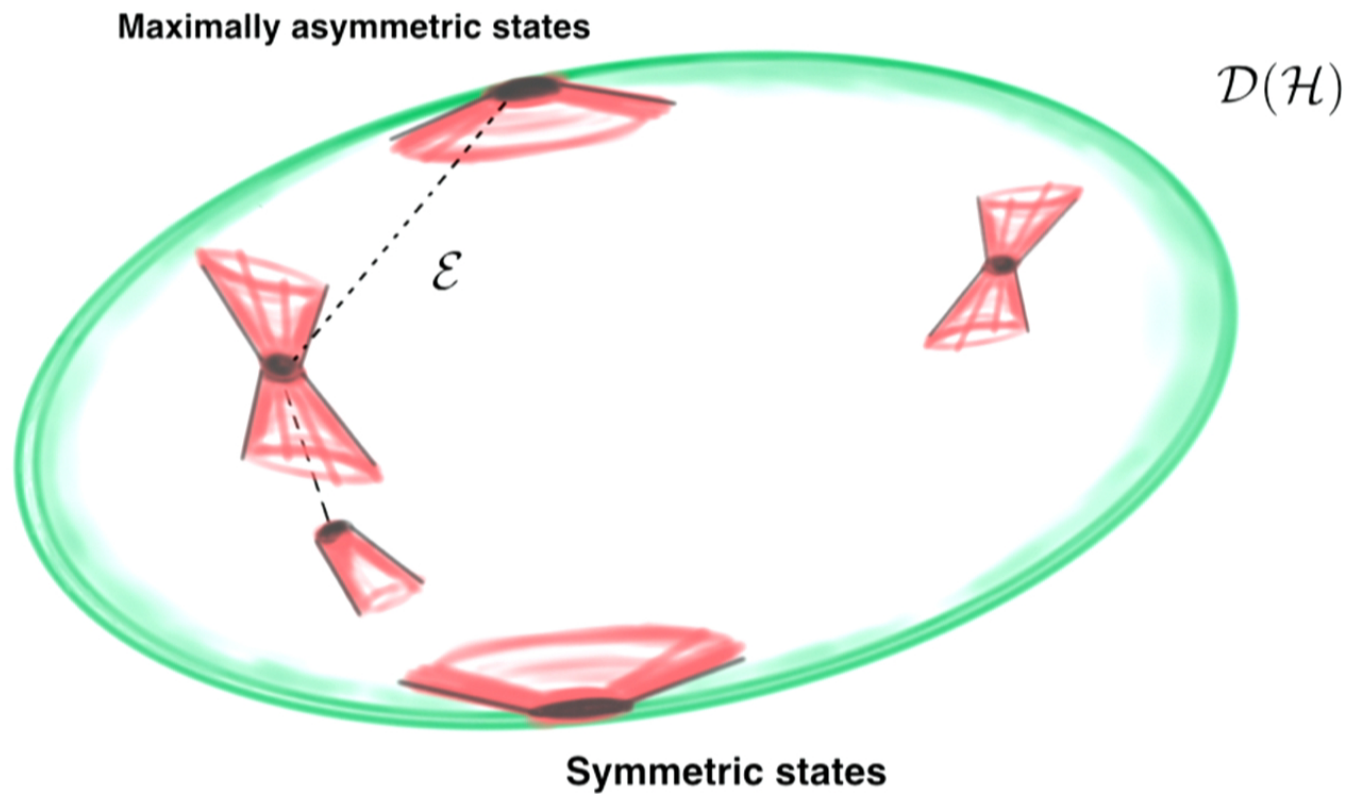
2. **Free quantum operations:**

**symmetric/covariant operations**

$$U\mathcal{E}(\rho)U^\dagger = \mathcal{E}(U\rho U^\dagger)$$

—**asymmetry is non-increasing**—

# A theory of asymmetry

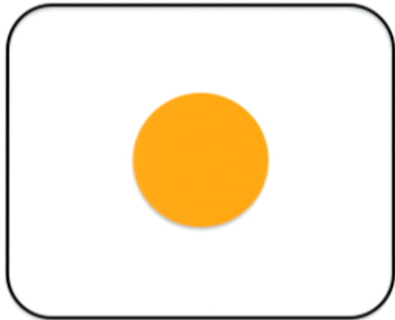




# Example, $G = SU(2)$

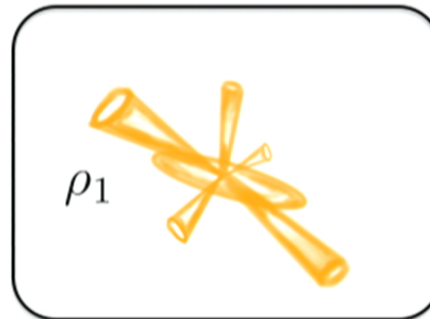
**Rotationally  
invariant states**

$$|\psi^-\rangle, \rho_s = \frac{1}{2}\mathbb{I}$$

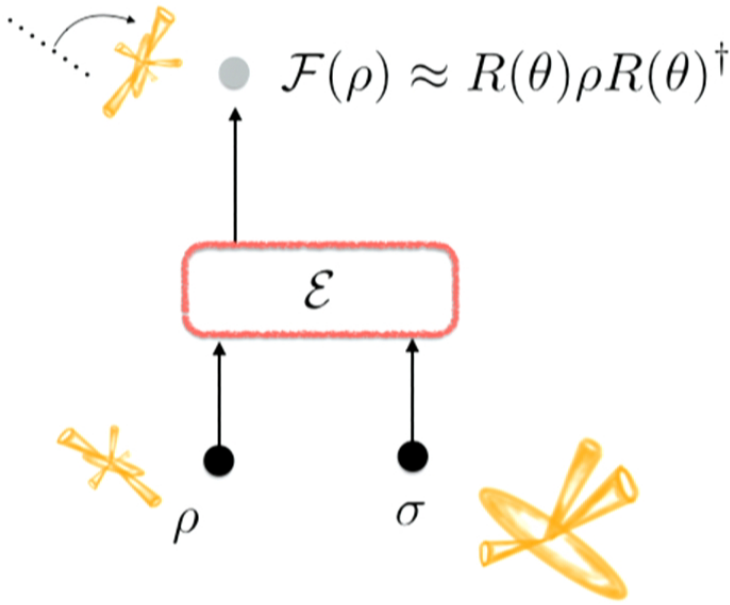


**Pointy/Asymmetric  
states**

$$\rho_s = |\uparrow\uparrow\rangle, \quad \frac{1}{3}|l,l\rangle\langle l,l| + \frac{2}{3}|l,0\rangle\langle l,0|$$



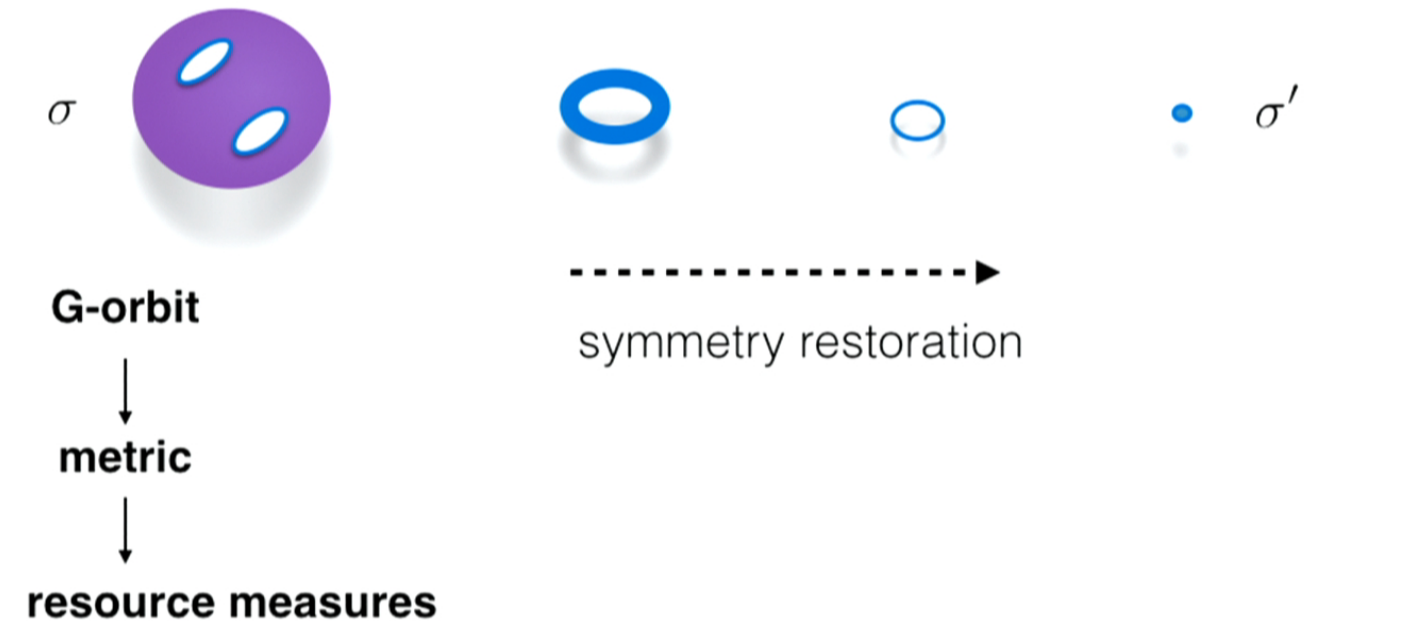
# Use of resources: asymmetry



**Spatial  
Rotation**

# Asymmetry structure

## Resource state



\*C. Cirstoiu, DJ, (2015)

# The 1st Law of Thermodynamics

- Traditional form:  $dE = dQ + dW$
- Microscopic energy conservation.  $[U, H_{\text{tot}}] = 0$

$$t \mapsto e^{-itH_{\text{tot}}}$$

$$\sqrt{p}|E_0\rangle + \sqrt{1-p}|E_1\rangle \rightarrow e^{-itE_0}\sqrt{p}|E_0\rangle + e^{-itE_1}\sqrt{1-p}|E_1\rangle$$

# Core connection

$$[U, H_{\text{tot}}] = 0$$

Global energy  
conservation  
+  
Thermal states



Thermodynamics is a  
**particular** asymmetry  
theory

$$\gamma = \frac{e^{-\beta H}}{Z}$$

Coherence



U(1)-asymmetry

## THERMAL $\subset$ $U(1)$ COVARIANT

Thermal operations:  $\mathcal{E}(\rho) = \text{tr}_b[\mathbf{U}(\rho \otimes \gamma)\mathbf{U}^\dagger]$

Covariant operations:  $\mathcal{F}(U_1(g)\rho U_1(g)^\dagger) = U_2(g)\mathcal{F}(\rho)U_2(g)^\dagger$   
 $\forall g \in G$

\* M. Lostaglio, DJ, T. Rudolph, *Nature Comm.* (2015)

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 $\forall g \in G$

Covariant  
Stinespring dilation:  $\mathcal{F}(\rho) = \text{tr}_2[\mathbf{V}(\rho \otimes \sigma)\mathbf{V}^\dagger]$

**symmetric auxiliary state**

\* M. Keyl, R. Werner *J.Math.Phys.* (1998)

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**symmetric auxiliary state**

\* M. Keyl, R. Werner *J.Math.Phys.* (1998)



## Applications of Framework:

- 1 The insufficiency of free energy relations.
- 2 Coherence “work-locking”.
- 3 General thermodynamic bounds on coherence.
- 4 Intrinsically-quantum 2nd law constraints.

### Based on:

- Lostaglio, Korzekwa, Oppenheim, DJ, *arxiv:05.xx* (2015)
- Lostaglio, DJ, Rudolph, *Nature Comm.* (2015)
- Lostaglio, Korzekwa, DJ, Rudolph, *Phys. Rev. X* (2015)

# (1). Insufficiency of free energies in thermodynamics.

Consider any set of functions  $\{D_\alpha(\cdot)\}_\alpha$  that  
**“behave like free energies”**:

If  $\rho \rightarrow \sigma$  then we have  $\{D_\alpha(\rho) \geq D_\alpha(\sigma)\}_\alpha$   
and  $D_\alpha(\rho) \geq c\|\rho - \gamma\|$

then

$\{D_\alpha(\cdot)\}_\alpha$  cannot provide a complete set of thermodynamic constraints.

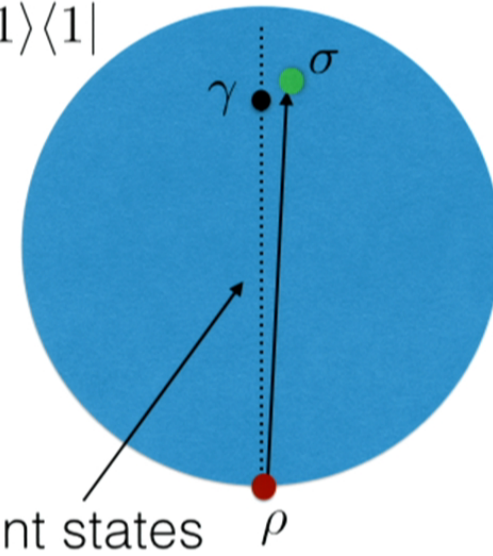
# (1). Insufficiency of free energies in thermodynamics.

## Proof:

$D_\alpha$  say “get closer to  $\gamma$ .”

Symmetry says:  
“asymmetry non-increasing.” ■

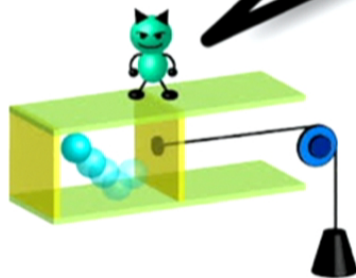
$$H = |1\rangle\langle 1|$$



Symmetric/incoherent states  $\rho$

$$\sigma = \begin{bmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{bmatrix}$$

*Is a qubit worth  
 $kT \ln 2$  of energy?*

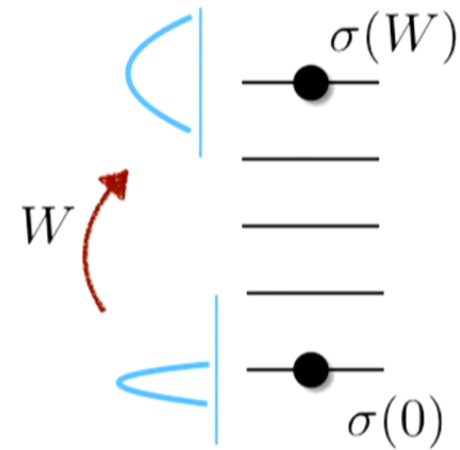


# Work / Ordered Energy

**Allow a broad work definition:**  
*“raising a weight up a ladder by height  $W$ ”*

$$\sigma(0) \rightarrow \sigma(W)$$

$$W = f(\langle H \rangle, \langle H^2 \rangle, \langle H^3 \rangle, \dots)$$



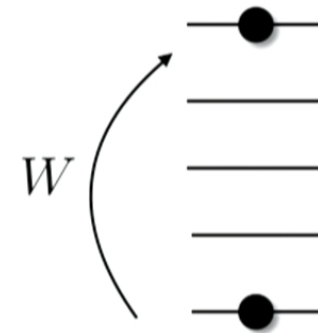
## (2). Work-locked in coherence

**Theorem:**

$$[\rho \rightarrow W] \Leftrightarrow [\mathcal{D}(\rho) \rightarrow W]$$

where

$$\mathcal{D}(\rho) = \mathcal{G}_H(\rho) = \int dt e^{-itH} \rho e^{itH}$$

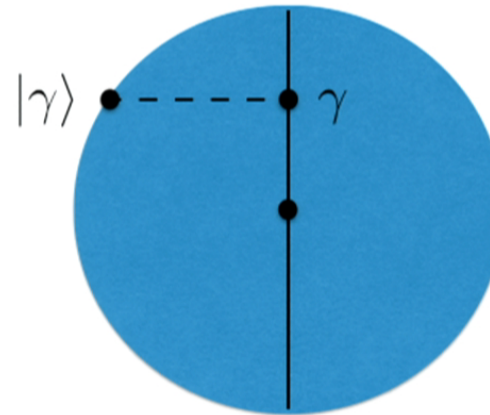


# Szilard and coherence

Pure state  $|\gamma\rangle$

$$\mathcal{D}(|\gamma\rangle\langle\gamma|) = \gamma$$

⇒ **No work** can be extracted from  $|\gamma\rangle$  **on its own.**



Value of a qubit ? Non-trivial.  
Requires careful “resource counting”.

# Unlocking coherence for work.

- Use a second coherent source:

$$\mathcal{D}(|\gamma\rangle\langle\gamma|) = \gamma$$

**(relational  
coherence  
protected)**

$$\mathcal{D}(|\gamma\rangle\langle\gamma| \otimes \sigma_R) \neq \mathcal{D}(|\gamma\rangle\langle\gamma|) \otimes \mathcal{D}(\sigma_R)$$

$\sigma_R$  acts as quantum reference frame for  $|\gamma\rangle$

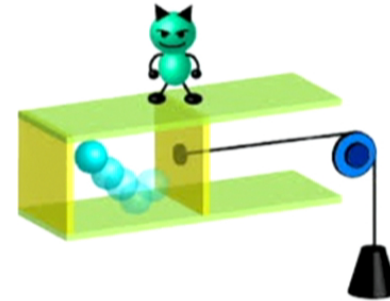
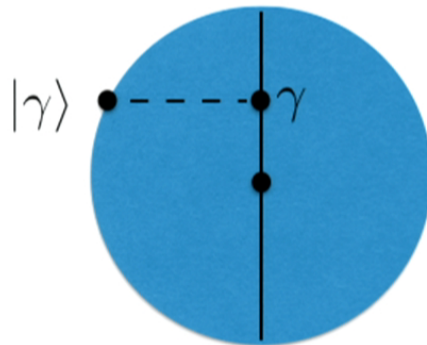
E.g. 
$$|\gamma\rangle \otimes |\gamma\rangle \rightarrow W \leq Z^{-1} e^{-\frac{E}{kT}} (E - 2kT \ln Z)$$
$$= kT \ln 2 \text{ (for } E = 0)$$



# Quantum Szilard engine

- It is only for a particular “classical” regime that we can associate the free energy to every qubit state.

$$|\Psi\rangle \longrightarrow W = -\Delta F$$



*M. Lostaglio, K. Korzekwa, J. Oppenheim, DJ,  
“Extracting work from quantum coherence” arXiv 05.xx (2015)*

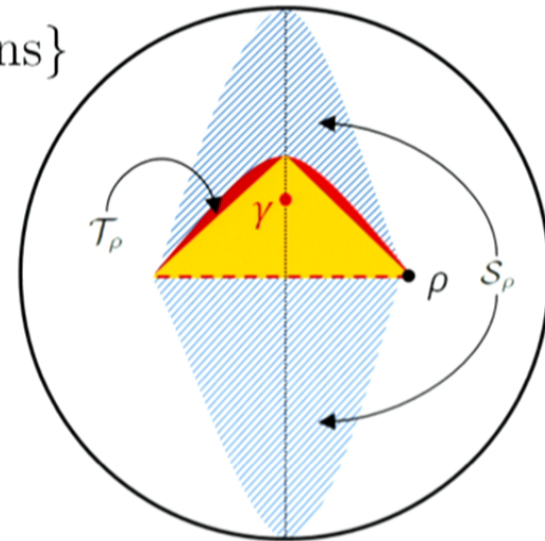
## Is free energy all there is to thermo?

$\mathcal{T}_\rho := \{\sigma : \rho \rightarrow \sigma \text{ under thermal operations}\}$

Now allow **arbitrarily** many  $|w\rangle^{\otimes n}$

Extends  $\mathcal{T}_\rho$  only to:

$\mathcal{S}_\rho := \{\sigma : \rho \rightarrow \sigma \text{ under covariant maps}\}$



---

From Covariant-Stinespring:

$$\mathcal{E}(\rho) = \text{tr}_b[\mathbf{U}(\rho \otimes \gamma_b)\mathbf{U}^\dagger]$$

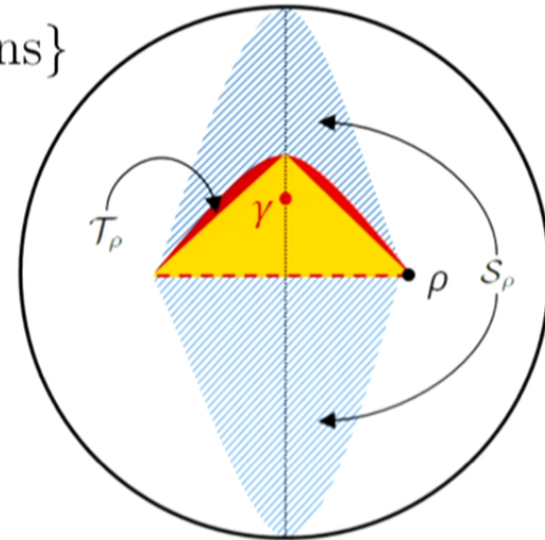
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# Non-trivial Coherence

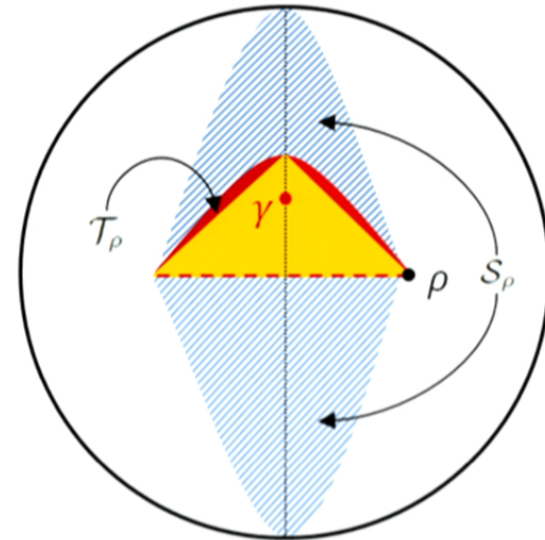
Yellow region:

$$\{\sigma = p\rho + (1 - p)\mathcal{E}(\mathcal{D}(\rho))\}$$



(thermal map on dephased state)

**Suggests at least two  
resources at play:  
(purity, asymmetry)**



# Mode operators

- Apply harmonic analysis to operators: irreps of group action.

$$\mathcal{B}(\mathcal{H}) = \bigoplus_{\nu} V_{\nu}$$

$$U(t)\rho^{(\nu)}U(t)^{\dagger} = e^{-i\nu t}\rho^{(\nu)}$$

$$\rho = \sum_{\nu=-d}^d \rho^{(\nu)}$$

Thermal	$[\mathcal{E}(\rho)]^{(\nu)} = \mathcal{E}(\rho^{(\nu)})$
operations	$\ \mathcal{E}(\rho)^{(\nu)}\ _1 \leq \ \rho^{(\nu)}\ _1$

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# (3). General Bounds on Coherence

Lower bound:

$$\sigma^{(\nu)} = \lambda_{\star} \rho^{(\nu)}$$

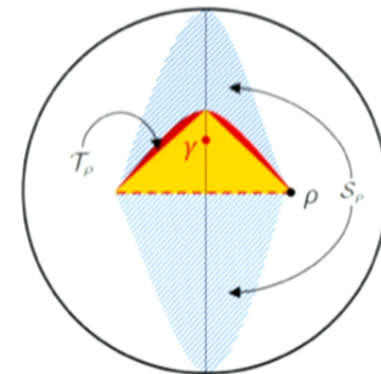
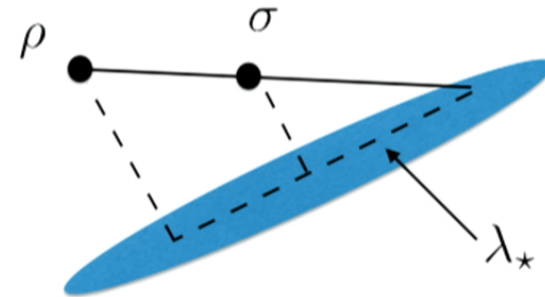
Upper bound:

$$|\sigma_k^{(\nu)}| \leq \sum_{c: \omega_c \leq \omega_k} |\rho_c^{(\nu)}| e^{-\beta \hbar (\omega_k - \omega_c)} + \sum_{c: \omega_c > \omega_k} |\rho_c^{(\nu)}|$$

Previous bound:

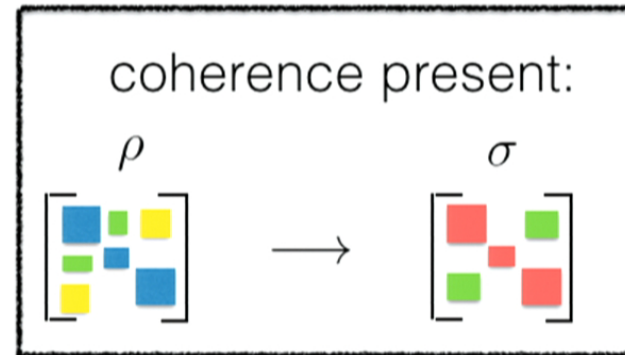
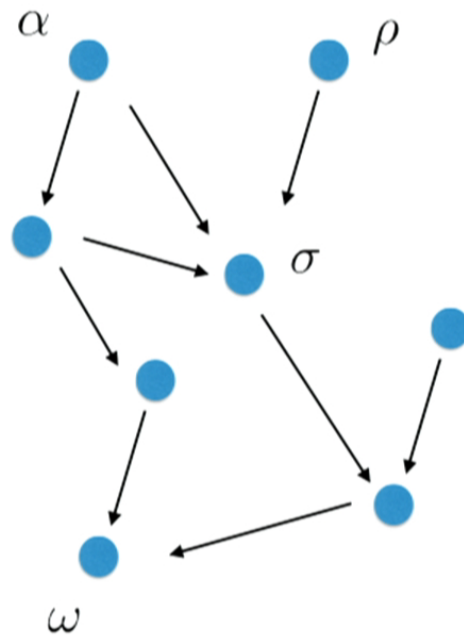
$$|\sigma_{nm}| \leq |\rho_{nm}| \sqrt{p_n |n| p_m |m|}$$

\* Cwiklinski, Studzinski, Horodecki, Oppenheim, arxiv (2014)



\*M. Lostaglio, K. Korzekwa, D.J. T. Rudolph, Phys. Rev. X (2015)

# (4). The full thermodynamic ordering of states?



**Q: Does the ordering of states admit an entropic formulation?**



# Thermodynamic structure

- Entanglement theory ~ non-locality monotones.
- Asymmetry theory ~ asymmetry monotones.
- Thermodynamics ~ **ordered energy + asymmetry**

**Work**

**Quantum  
coherence**

# Distillation/Formation

**Work:**

$$|w\rangle = |1\rangle^{\otimes w}$$

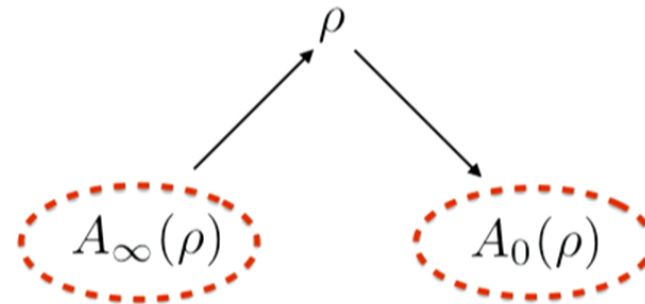
$$W_{\text{form}} \rightarrow \rho \rightarrow W_{\text{distill}}$$

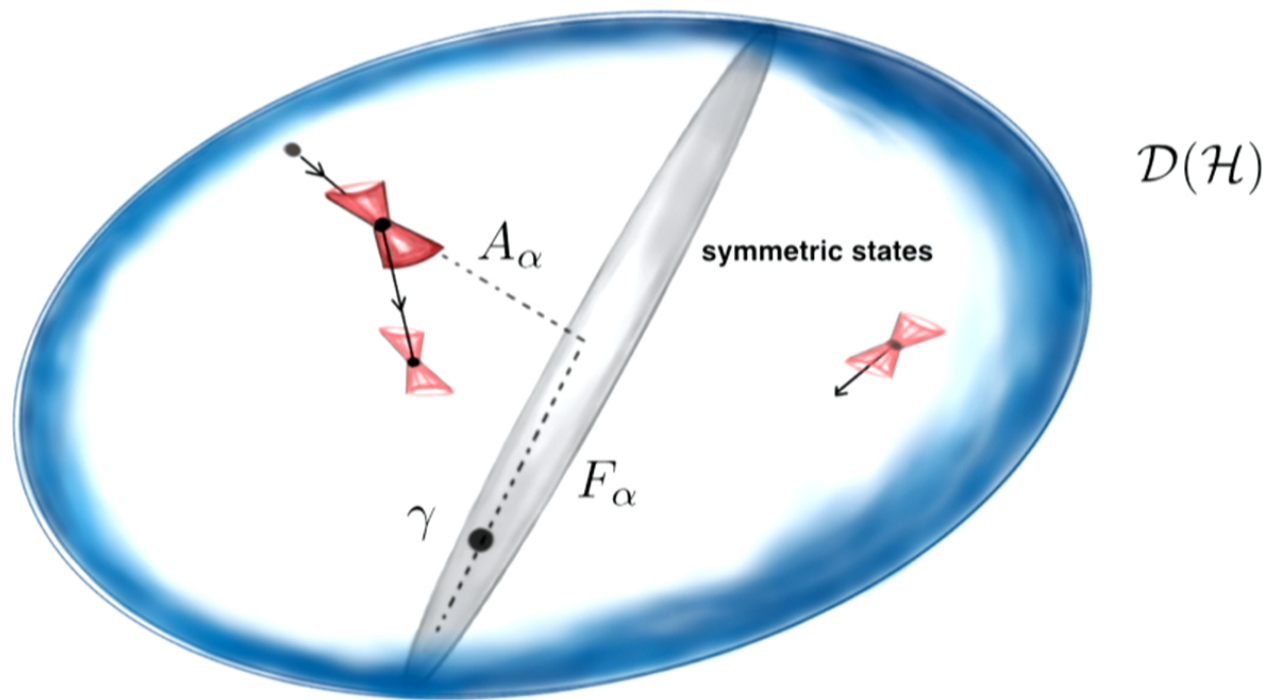
$$D_{\infty}(\rho||\gamma)$$

$$D_0(\rho||\gamma)$$

**Coherence:**

$$|+\rangle = (|0\rangle + |1\rangle)/\sqrt{2}$$





# (4). Necessary entropic constraints

**Theorem:** For arbitrary quantum states, the thermodynamic transformation  $\rho \rightarrow \sigma$  is possible provided

$$F_\alpha(\rho) \geq F_\alpha(\sigma)$$

$$A_\alpha(\rho) \geq A_\alpha(\sigma) \quad \forall \alpha \geq 0$$

**Asymmetry  
monotones:**

$$A_\alpha(\rho) := D_\alpha(\rho || \mathcal{G}(\rho))$$

$$\mathcal{G}(\rho) = \int_G dg U(g)\rho U(g)^\dagger$$

$$F = U - TS$$

$$F(\rho) = F(\mathcal{G}(\rho)) + A_1(\rho)$$

**incoherent component**

**coherent component**

# Macroscopic regime

- **Theorem:** for any  $\rho \in \mathcal{B}(\mathcal{H})$  we have

$$\lim_{n \rightarrow \infty} \frac{1}{n} \begin{bmatrix} F_\alpha(\rho^{\otimes n}) \\ A_\alpha(\rho^{\otimes n}) \end{bmatrix} = \begin{bmatrix} F(\rho) - F(\gamma) \\ 0 \end{bmatrix}$$

$$F = \langle H \rangle - TS$$



# Review



Essentially unique entropy.

$$\rho \rightarrow \sigma \Leftrightarrow S(\rho) \leq S(\sigma)$$

---



$$\langle e^{-\beta(W-\Delta F)} \rangle = 1 \quad (\text{provably incomplete})$$

$$\rho \rightarrow \sigma \Leftrightarrow D_\alpha(\rho||\gamma) \leq D_\alpha(\sigma||\gamma)$$

---



$$\begin{bmatrix} F_\alpha(\rho) \\ A_\alpha(\rho) \end{bmatrix} \leq \begin{bmatrix} F_\alpha(\sigma) \\ A_\alpha(\sigma) \end{bmatrix} + ?$$

# Outlook

- Applications: coherent transport systems (conductivity, energy transport scenarios) ?
- Non-equilibrium state states? Feedback control?
- New entropic conditions — are they both necessary and sufficient?

