

Title: Some combinatorial comments on amplitudes

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Abstract: I will begin with the perspective that the perturbative expansion is an augmented generating function and then discuss some of the results which follow from this perspective.

segmented
generating functions
and combinatorics!

Dyson-Schwinger
equations.

I like excuses to sum
over graphs

Precisely

\mathcal{C} comb class

and form

$$\sum_{c \in \mathcal{C}} c x^{|c|} \in \mathbb{Q}(\mathcal{C})[[x]]$$

$$\sum_{c \in \mathcal{C}} c x^{|c|} \in \mathbb{Q}[\mathcal{C}][[x]]$$

and then I want
an evaluation map

$$\phi: \mathcal{C} \rightarrow \mathbb{Q}$$

$$c \mapsto 1$$

get usu. gen. fn.

a Feynman graph
to its integral

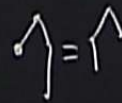
more to the point

$$\phi: \mathcal{C} \rightarrow \mathbb{Q}[[L]](\dots) \text{ or formal integral expressions or } \dots$$

Feynman rule taking

How to specify an
interesting carb class
with a focus on rooted
trees.

I want rooted trees
with no plane structure
no valence restrictions.
call this \mathcal{T}

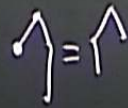


define $B_+ : \mathcal{Q}[\mathcal{T}] \rightarrow \mathcal{Q}[\mathcal{T}]$
as a forest = monomial in
rooted trees

$$B_+(t_1 \dots t_n) = \begin{array}{c} \wedge \\ \vdots \\ t_1 \dots t_n \end{array}$$

interesting comb class
with a focus on rooted
trees.

I want rooted trees
with no plane structure
no vertex restrictions.
call them \tilde{T}



can a forest = monomial in
rooted trees

$$B_+(t_1, \dots, t_n) = \begin{array}{c} \triangle \\ \vdots \\ t_1 \dots t_n \end{array}$$

use B_+ to build
interesting classes

$$Q(\tilde{T}) \rightarrow Q(\tilde{T})$$

eg $B_+(1) =$ 

$$(1 - x + 2x^2)^2$$

$$= 1 - 2x + 4x^2$$

Spec. eg.  empty tree

$$T(x) = 1 + xB_+(T(x)^2)$$

$$= 1 + x^0 + 2x^2 + x^3(1+4) + x^4(8) + 2x^4 + 4x^4 + \dots$$

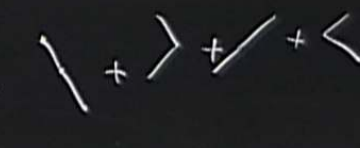
eg $B_+(1) =$ 

Spec. eg.  empty tree

$$\begin{aligned} & (1 + x + 2x^2)^2 \\ &= 1 + 2x + \underbrace{4x^2 + 4x^3}_{\text{circled}} + 4x^4 \end{aligned}$$

$$T(x) = 1 + xB_+(T(x)^2)$$

$$= 1 + x^0 + 2x^2 + x^3(1 + 4) + x^4(8 + 2 \cdot 1 + 4 \cdot 1) + \dots$$

what do these coeffs count: binary trees eg $4 =$ 

$$\begin{aligned}
 \text{eg } T(x) &= \underline{1} - x B_4 \left(\frac{1}{T(x)} \right) & \frac{1}{1-x-x^2} &= 1 + \overset{\circ}{x} + x^2 + \dots \\
 & & & + x^3 + \dots \\
 &= \underline{1} - x \cdot -x^2 - x^3 (\lambda + 1) - x^4 (\lambda + 2\lambda + \lambda + 1) \\
 & & & + \dots
 \end{aligned}$$

no value restriction,
call these $\}$

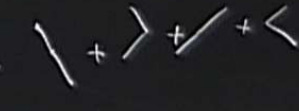
eg $B_+(\cdot) =$ 

Spec. eg. \swarrow empty tree

$$(1 + x + 2x^2)^2 = 1 + 2x + 4x^2 + 4x^3 + x^4$$

$$T(x) = 1 + x B_+(T(x)^2)$$

$$= 1 + x \cdot 0 + x^2 \cdot 1 + x^3 (1 + 4) + x^4 (8 + 2 \cdot 1 + 4 \cdot 1) + \dots$$

what do these coeffs count: binary trees eg $4 \}$ = 

$$\text{eg } T(x) = \mathbb{1} - xB + \left(\frac{1}{T(x)}\right) \quad \frac{1}{1-x-x^2} = 1 + \frac{x}{1-x^2} + x^2 + \dots$$

$$= \mathbb{1} - x \cdot -x^2 - x^3(\Lambda+1) - x^4(\Lambda+2\Lambda+1) + \dots$$

Call these kinds of things
Combinatorial Dyson-Schwinger
equations.

Combinatorial DSEs
have alg. gen. fns
as solns.

Apply Feynman rules ϕ - to be Comb DSE
get analytic DSE
- to be solns get
Green's functions.

② Tree hook weights
(joint with Brad Jones)

The enumer. comb. community
studied this indep

A simple case of the preceding ϕ just
gives numbers (leading
log of more realistic cases)

specifically

$$\phi(t) = \prod_{\text{vet}} B_{|t_v|}$$

where t_v is the subtree rooted
at v .

a Feynman graph

A simple case of the preceding ϕ just gives numbers (leading log of more realistic cases)

specifically

$$\phi(t) = \prod_{v \in t} B_{|t_v|}$$

where t_v is the subtree rooted at v .

eg tree factorial

$$t! = \prod_{v \in t} |t_v|$$

eg  $= 4 \cdot 3 \cdot 1 \cdot 1 = 12$

Then the hook length series ("Greens function")

is $F_{T,B}(x) = \sum_{t \in T} \phi(t) x^{|t|}$

$$F_{T,B}(x) = \sum_{t \in T} \phi(t) x^{|t|}$$

Hook weight formula - nice if $F_{T,B}$ has a closed form or interpretation.

Call these)

Go on
nice comb. DSE's
and nice Eyring rules
with nice Greens functions

$$T(x) = 1 - x \beta_+ \left(\frac{1}{T(x)^2} \right)$$

eg



$$\phi(t) = \frac{L^{|t|}}{t!}$$

$$\text{then } G(x, L) = (1 - 3xL)^{\frac{1}{3}}$$

③ More physical DSEs

there is something closer to what you might think of as a Dyson-Schwinger eqn.

(Broadhurst - Kreizler)

$$G(\alpha, L) = 1 - \left(\frac{\alpha}{g^2} \right) \left(\int^{d^4 k} \frac{k \cdot q}{|k|^2 G(\alpha, \log \frac{|k|^2}{\mu^2}) |k+q|^2} \dots \int^{d^2 p} \frac{1}{g^2 \mu^2} \right)$$



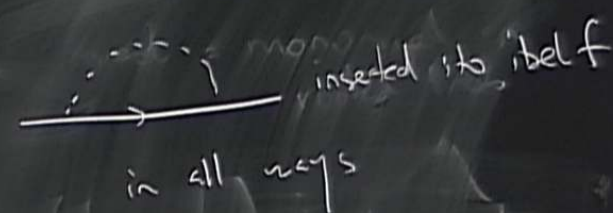
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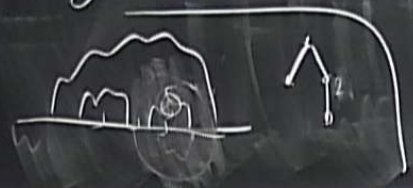
$$G(\alpha, L) = 1 - \left(\frac{\alpha}{g^2} \right) \int d^4k \frac{k \cdot q}{|k|^2 G(\alpha, \log \frac{|k|^2}{\mu^2}) |k+q|^2} \dots \left(\frac{g^2}{\mu^2} \right)$$

like in Yukawa theory



ee rooted

eg



$$T(x) = 1 - x \beta_+ \left(\frac{1}{T(x)} \right)^2$$

$$\phi(t) = \frac{L^{1/t}}{t!}$$

$$\text{then } G(x, L) = (1 - 3xL)^{1/3}$$

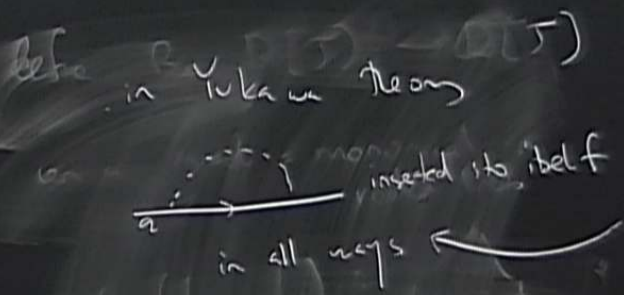
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$$G(\alpha, L) = 1 - \left(\frac{\alpha}{g^2} \right) \int d^4k \frac{k \cdot q L}{|k|^2 G(\alpha, \log \frac{1+t}{\mu}) |k+q|^2} \dots \left(\frac{\alpha}{g^2 \mu^2} \right)$$

$L = \log \frac{\Lambda^2}{\mu^2}$



function")

$$\phi(t) \times |t|$$

if $\bar{F}_{T,B}$

fisher.

eg

Call



eg $T(x) = 1 - xB_+ \left(\frac{1}{T(x)} \right)$

Now do the following to

① expand $G(x, L) = 1 - \sum_{k \geq 1} \gamma_k(x) L^k$
 plug in and expand geom. series

② swap \sum, \int

③ use $\frac{\partial^i}{\partial p^i} y^p \Big|_{p=0} = (-1)^i \log^i y$

④ swap $\int, \frac{\partial}{\partial p}$

⑤ undo ①

glt

- ② swap $(1, 2)$
 ③ use $\frac{\partial^i}{\partial(p)^i} y^p \Big|_{p=0} = (-1)^i \log^i y$

get $G(\alpha, L) = 1 - \alpha G(\alpha, \frac{\partial}{\partial(p)})^{-1} (e^{-Lp} - 1) F(p) \Big|_{p=0}$



- ② swap $(z, 1/z)$
 ③ use $\frac{\partial^i}{\partial p^i} y^p \Big|_{p=0} = (-1)^i \log^i y$

get $G(\alpha, L) = 1 - \alpha \underbrace{G(\alpha, \frac{\partial}{\partial p})^{-1}}_{\text{series } \frac{\partial}{\partial p} \text{ in } \alpha}$ $\underbrace{(e^{-Lp} - 1) F(p)}_{\text{series } L, p} \Big|_{p=0}$

integral for
 with \int regularized by p .

CAUTION
 DO NOT TOUCH THE BOARD SURFACE
 IT IS COVERED BY A PROTECTIVE FILM
 WHICH CAN BE DAMAGED BY THE BOARD SURFACE

- ② swap $(z, 1)$
 ③ use $\frac{\partial^i}{\partial p^i} y^p \Big|_{p=0} = (-1)^i \log^i y$

get $G(\alpha, L) = 1 - \alpha \underbrace{G(\alpha, \frac{\partial}{\partial p})^{-1}}_{\text{series in } \frac{\partial}{\partial p}} \underbrace{(e^{-Lp-1} F(p)) \Big|_{p=0}}_{\text{series } L, p}$

Slightly more generally

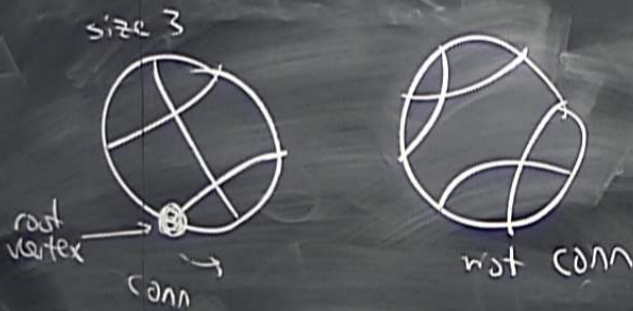
$$G(\alpha, L) = 1 - \sum_{k \geq 1} \alpha^k G(\alpha, \frac{\partial}{\partial p})^{1-sk} (e^{-Lp-1} F_k(p)) \Big|_{p=0}$$

where $F_k(p) = \sum_{i \geq 0} f_{ki} p^{i-1}$

integral for
 with regularization by p .



where the sum runs
over rooted connected
chord diagrams



$|C| = \# \text{ chords}$

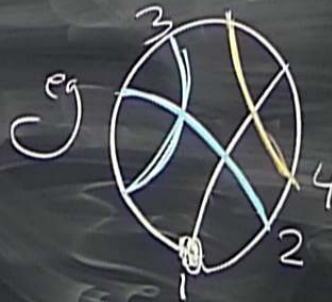
need recursive chord order

root chord is 1

remove root chord. get

connected components C_1, C_2, \dots, C_k

order C_1 next to C_2 etc.



and in general Markus Hihn

then $G(x, L) = (1 - x)L$

$$F(p) = f_0 p^{-1} + f_1 + f_2 p + \dots$$

then $b(c) = t_1$

$$f_c = f_{t_1 - t_0} \cdot f_{t_3 - t_2} \cdot f_{t_2 - t_1} \cdot f_0$$

say a chord is terminal
if everything it crosses occurs
before it in the recursive
chord order

let $t_1 < t_2 < \dots < t_k$ be the terminal
chords