

Title: Displaced vertices from dark matter freeze-in

Date: May 08, 2015 12:30 PM

URL: <http://pirsa.org/15050042>

Abstract: <p>

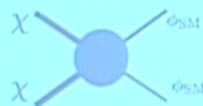
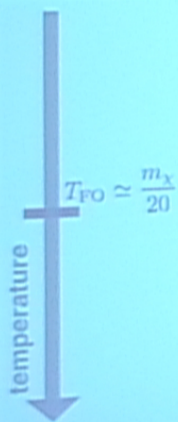
Freeze-in is a general and calculable mechanism for dark matter production in the early universe. Assuming a standard cosmological history, such a framework predicts metastable particles with a lifetime generically too long to observe their decays in a collider environment. In this talk I will report work in progress where we consider alternative cosmologies, where entropy dumped in the primordial Standard Model plasma leads to shorter lifetime for the metastable particles in order to reproduce the observed dark matter density. Famous examples are moduli decays in SUSY theories and inflationary reheating. Remarkably, for a large region of the parameter space the decay lengths are in the displaced vertex range and they can be observable at present and future colliders.</p>

Thermal Freeze-out

$$\frac{dn_X}{dt} = -3Hn_X - \langle \sigma v_{\text{rel}} \rangle (n_X^2 - n_X^{\text{eq}2})$$

Number density changes due to universe expansion and reactions

Lee and Weinberg, *Phys. Rev. D* 17, 339 (1977)



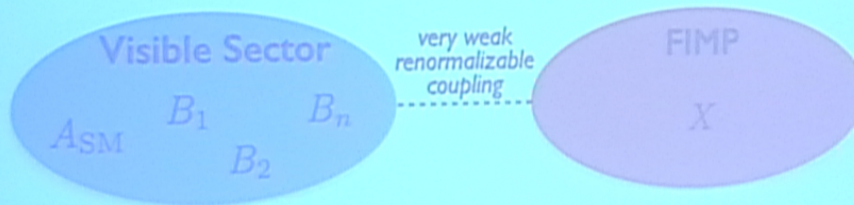
$$n_X^{\text{eq}} \sigma v_{\text{rel}} \simeq H$$

$$Y_{\text{FO}} = \frac{n_{\text{FO}}}{s} \simeq \frac{1}{m_X M_{\text{Pl}} \sigma v_{\text{rel}}}$$

- Efficient annihilations, WIMP full T^3 abundance
- Expansion takes over (freeze-out)
- Number density scales only with the expansion

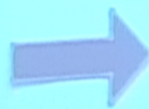
FIMP Dark Matter

Freeze-in of Feeble Interacting Massive Particles (FIMPs)



- Visible Sector: SM fields and bath particles B_i in thermal equilibrium
 - FIMP interactions too weak to ever bring X to equilibrium
- Hall, Jedamzik, March-Russell, West, JHEP 1003 (2010)

Assume negligible initial X abundance



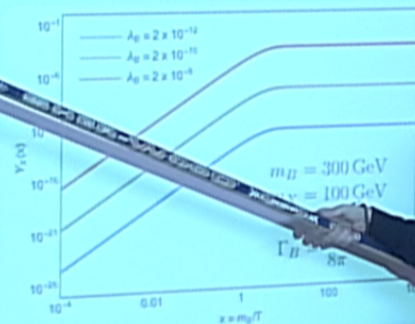
FIMP can be produced from bath particles collisions or decays

Freeze-in

$$B \rightarrow A_{SM} X \quad \frac{dn_X}{dt} = -3Hn_X + \Gamma_B n_B^{\text{eq}} \frac{K_1[m_B/T]}{K_2[m_B/T]}$$

Effective
temperature $T \sim m_B$

$$Y_{\text{FI}} = \frac{n_{\text{FI}}}{s} \simeq \Gamma_B \frac{M_{\text{Pl}}}{m_B^2}$$



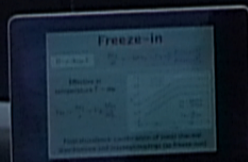
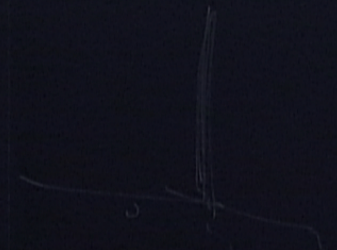
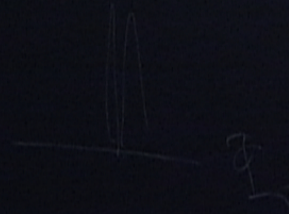
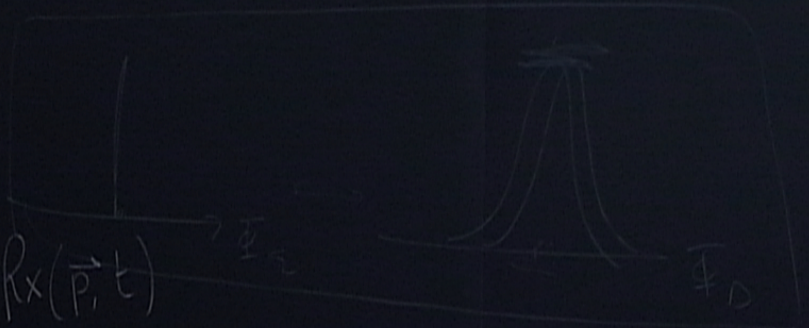
Final abundance: combination of initial thermal distributions and masses/couplings (as freeze-out)

CLASSICAL

R_x

$$n_x = \int \frac{d^3 p}{(2\pi)^3} R_x(\vec{p}, t)$$

$$R_x \ll R_B^{29}$$

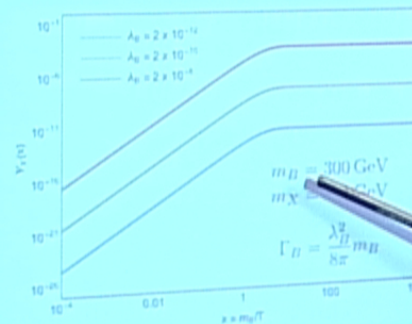


Freeze-in

$$B \rightarrow A_{SM} X \quad \frac{dn_X}{dt} = -3Hn_X + \Gamma_B n_B^{eq} \frac{K_1[m_B/T]}{K_2[m_B/T]}$$

Effective at temperature $T \sim m_B$

$$Y_{FI} = \frac{n_{FI}}{s} \simeq \Gamma_B \frac{M_{Pl}}{m_B^2}$$



Final abundance: combination of initial thermal distributions and masses/couplings (as freeze-out)

Freeze-in Scenarios

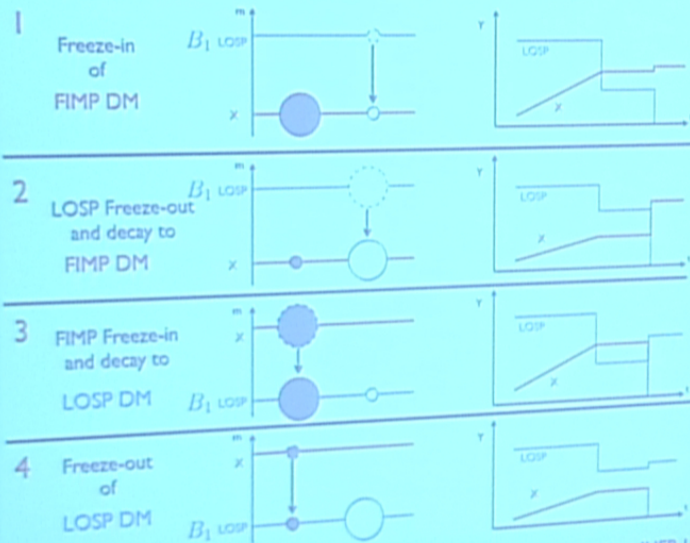
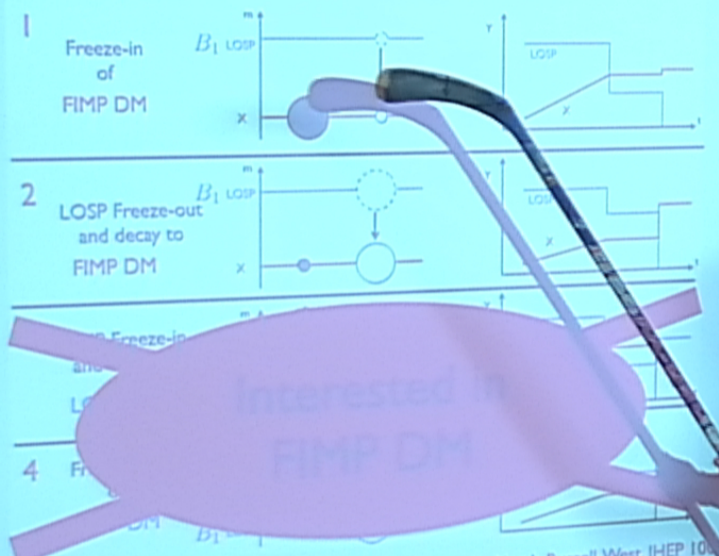


Figure from Hall, Jedamzik, March-Russell, West, JHEP 1003 (2010)

$$f_X = \left(\frac{dX}{dt} \right)_{LOSP}$$

$$R_X \ll f_X$$

Freeze-in Scenarios



Handwritten notes on a blackboard:

$$f_X = \left(\frac{d^3}{d^3 p} \right) f_X(p, t)$$

$$\ll f_B$$

$$E_1$$

$$E_2$$

$$T_1$$

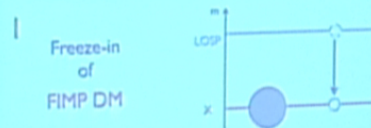
$$T_2$$

Freeze-in Signals

Signal depends on the LOSP nature and whether DM is LOSP or FIMP

Collider Signals

LOSP can be produced at colliders with subsequent decays



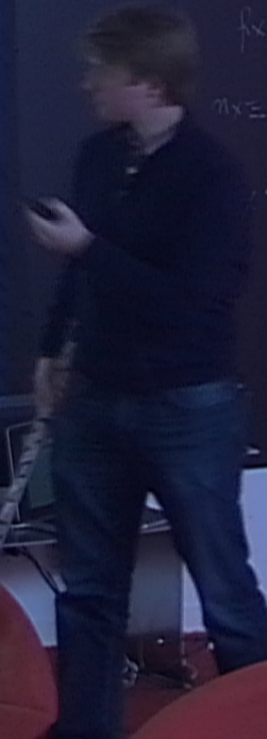
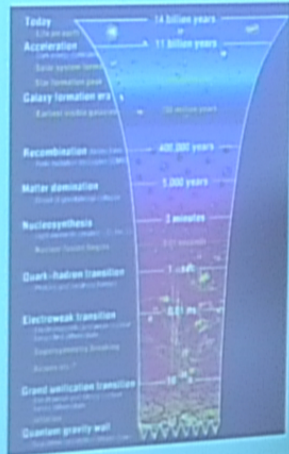
$$\tau_B = \Gamma_B^{-1} \simeq 3.7 \times 10^8 \text{ cm} \left(\frac{m_X}{100 \text{ GeV}} \right) \left(\frac{300 \text{ GeV}}{m_B} \right)^2$$

Decay length much bigger than detector size
(decays of stopped particles still possible)



Radiation Era?

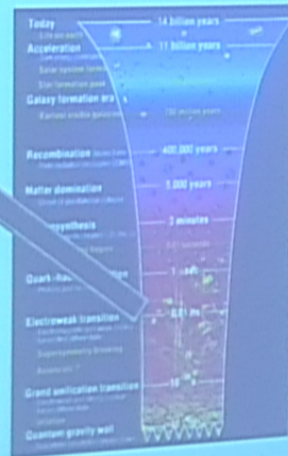
Freeze-out and Freeze-in active when $T \sim m_{\text{weak}}$



Radiation Era?

Freeze-out and Freeze-in active when $T \sim m_{\text{weak}}$

Assuming standard cosmology
this happens during a
radiation dominated era

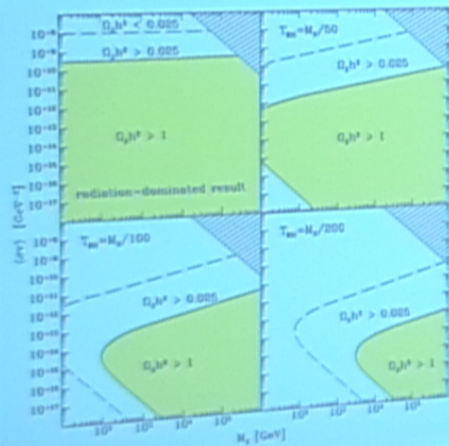


Reheating

Particle production during inflationary reheating

WIMPs Axions Neutrinos Baryogenesis

Chung, Kolk, Riotto, PRD 60 (1999)
Giudice, Kolk, Riotto, PRD 64 (2001)



WIMPs

General message:
smaller cross section allowed
(dilution and $Y_{FO} \propto (\sigma_{\text{rel}})^{-1}$)

$$R_X = \int \frac{d^3 p}{(2\pi)^3} R_X(p, t)$$

$$R_X \ll R_{\text{FO}}$$

Reheating

Particle production during inflationary reheating

WIMPs

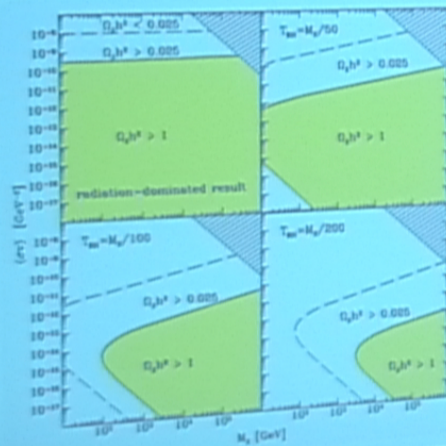
Axions

Neutrinos

Baryogenesis

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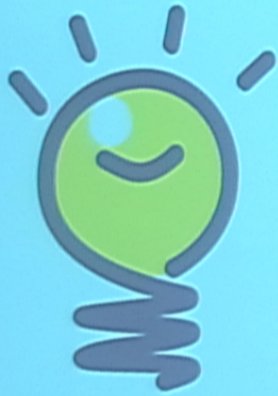
Giudice, Kolk, Riotto, PRD 64 (2001)



WIMPs

General message:
smaller cross section allowed
(dilution and $Y_{FO} \propto (\sigma_{Vrel})^{-1}$)

Reheating



$$Y_{FI} \propto \Gamma_B$$

FIMP production would
require smaller lifetime

Decay length may be in the
displaced vertex region!

[GeV]

M_{ν} [GeV]

Plan for Today's Talk

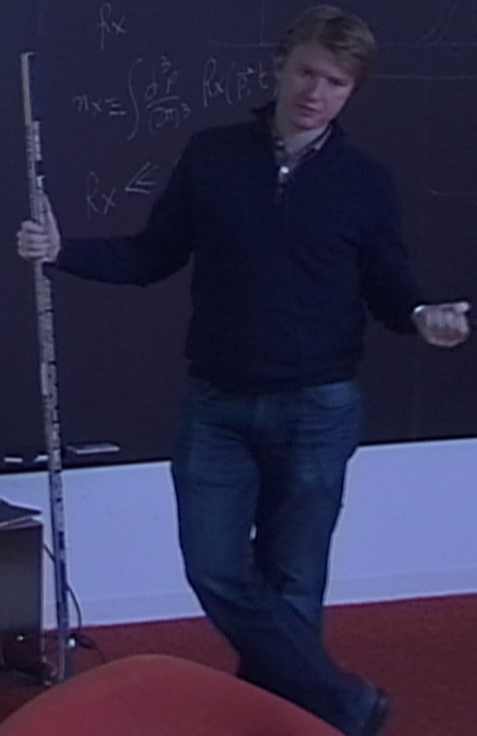
Freeze-in: general analysis

Freeze-in and inflationary reheating

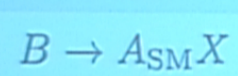
Freeze-in and "moduli"

A motivated candidate

Outlook

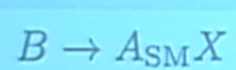


Relic Calculation



$$\frac{dn_X}{dt} = -3Hn_X + \Gamma_B n_B^{\text{eq}} \frac{K_1[m_B/T]}{K_2[m_B/T]}$$

Relic Calculation



$$\frac{dn_X}{dt} = -3Hn_X + \Gamma_B n_B^{\text{eq}} \frac{K_1[m_B/T]}{K_2[m_B/T]}$$

Hubble Parameter

$$H = \left(\frac{8\pi}{3}\right)^{1/2} \frac{\sqrt{\rho(t)}}{M_{\text{Pl}}}$$

$\rho(t)$ depends on the cosmological history

Thermal Distributions

$$T(t) = \left(\frac{30}{g_* \pi^2}\right)^{1/4} \rho_{\text{rad}}(t)^{1/4}$$

Radiation temperature $T(t)$ depends on the cosmological history

General Cosmology

Scale factor a
as "time" Variable

$$\frac{d}{dt} = H \frac{d}{d \ln a}$$

Comoving Density

$$\mathcal{X} = n_X a^3$$

Inflationary Cosmology

INFLATION

$$a \propto \exp(H_I t) \quad T \simeq 0$$

$$\rho_r \simeq \text{const} \quad \rho_{\text{rad}} \simeq 0$$

OSCILLATIONS

$$a \propto t^{2/3} \quad T \propto a^{-4/3}$$

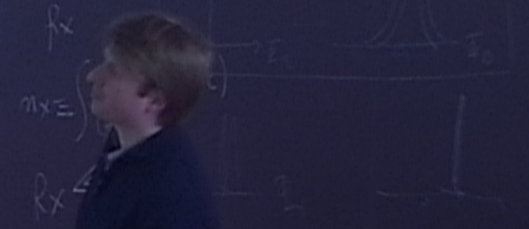
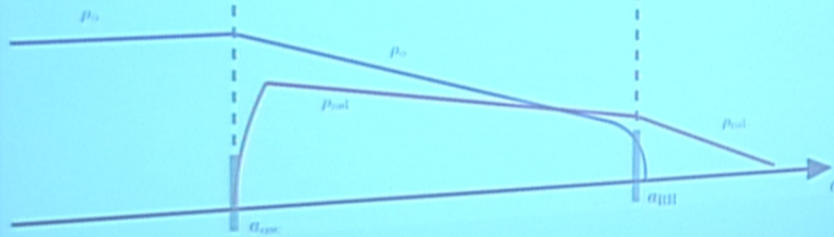
$$\rho_r \propto a^{-4} \quad \rho_{\text{rad}} \propto a^{-4/2}$$

RD

$$a \propto t^{1/2}$$

$$T \propto a^{-1}$$

$$\rho_{\text{rad}} \propto a^{-4}$$

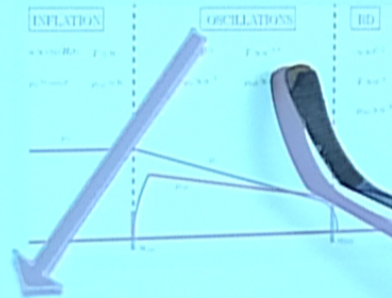


Inflationary Cosmology

INFLATION

$$\rho \simeq \rho_\phi \simeq \text{const} \simeq E_I^4$$

$$\rho_{\text{rad}} \simeq 0$$



INFLATON OSCILLATIONS

$$\rho \simeq \rho_\phi \propto a^{-3}$$

$$T_{\text{max}} \simeq (E_I^2 M_{\text{Pl}} \Gamma_\phi)^{1/4}$$

$$T \propto a^{-3/8}$$

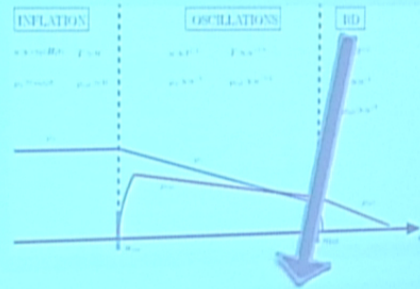
f_R
 $\gamma = \int \frac{\partial \rho}{\partial T} dT$

Inflationary Cosmology

INFLATION

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INFLATON OSCILLATIONS

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REHEATING

$$\rho \simeq \rho_{\text{rad}} \propto a^{-4}$$

$$T_{\text{RH}} \simeq (M_{\text{Pl}} \Gamma_\phi)^{1/2}$$

$$T \propto a^{-1}$$

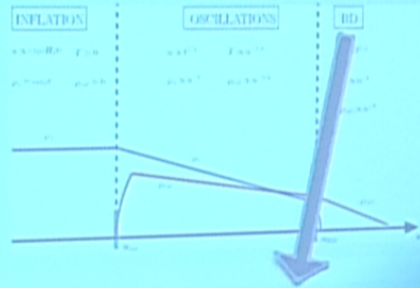


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REHEATING

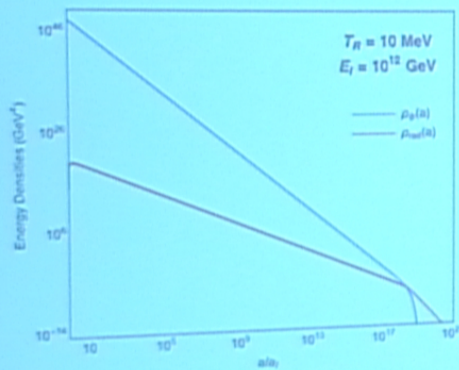
$$\rho \simeq \rho_{\text{rad}} \propto a^{-4}$$

$$T_{\text{RH}} \simeq (M_{\text{Pl}} \Gamma_\phi)^{1/2}$$

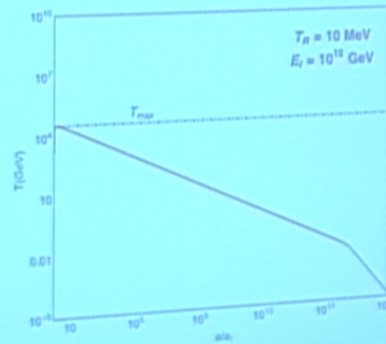
$$T \propto a^{-1}$$

f_X
 $\pi_X = \int \frac{d^3 p}{(2\pi)^3} R_X(p, t)$
 $R_X \ll R_B^{\text{eff}}$

Inflationary Cosmology



$$\frac{d\rho_\phi}{dt} + 3H\rho_\phi = -\Gamma_\phi\rho_\phi$$



$$\frac{d\rho_{\text{rad}}}{dt} + 4H\rho_{\text{rad}} = \Gamma_\sigma\rho_\sigma$$

f_X
 $\pi_X = \int \frac{d^3p}{(2\pi)^3} R_X(p, t)$
 $R_X \ll R_B$

Inflationary Cosmology

INFLATION

$$a \propto \exp(H_I t) \quad T \approx 0$$

$$\rho_{\text{tot}} \approx \text{const} \quad \rho_{\text{rad}} \approx 0$$

OSCILLATIONS

$$a \propto t^{2/3} \quad T \propto a^{-3/2}$$

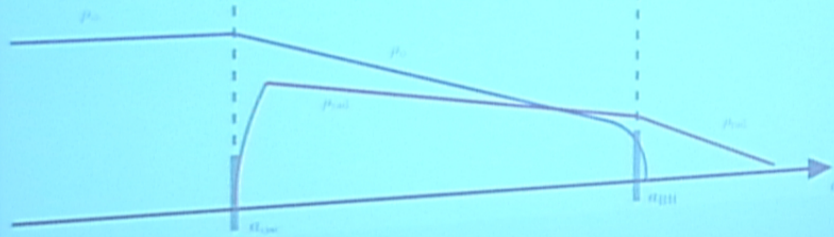
$$\rho_{\text{tot}} \propto a^{-3} \quad \rho_{\text{rad}} \propto a^{-4}$$

RD

$$a \propto t^{1/2} \quad T \propto a^{-2}$$

$$\rho_{\text{tot}} \propto a^{-3}$$

Freeze-in and inflationary reheating



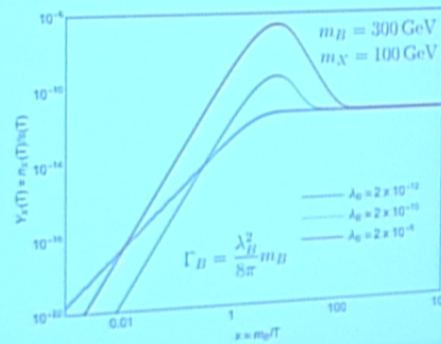
$$r_X = \left(\frac{d^3 \rho}{dt^3} \right) \frac{R_X(\rho, t)}{R_X(\rho, t)}$$

$$R_X \ll R_B$$

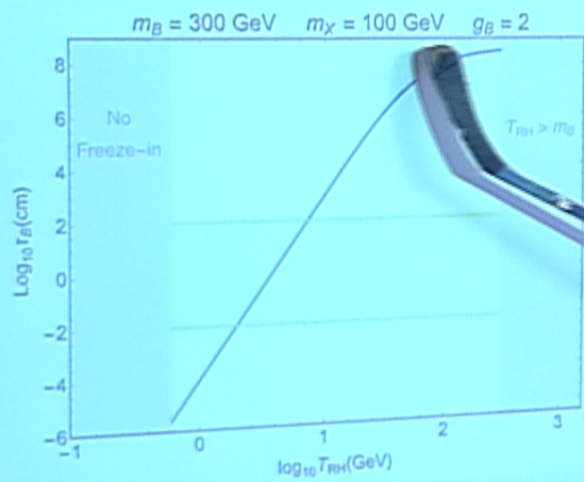
Freeze-in

Freeze-in density depends on the reheat temperature
 FIMP more coupled to the visible sector for low T_{RH}

$$\lambda_B = \begin{cases} 2 \times 10^{-12} & T_{RH} = 10^6 \text{ GeV} \\ 2 \times 10^{-10} & T_{RH} = 17.4 \text{ GeV} \\ 2 \times 10^{-8} & T_{RH} = 4.6 \text{ GeV} \end{cases}$$



Displaced Vertices

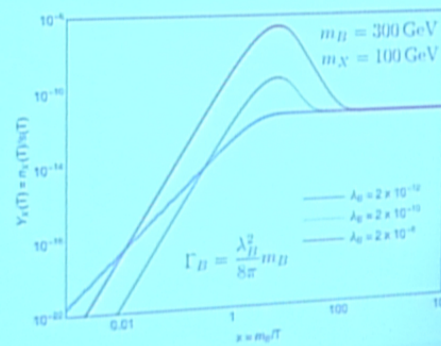


Full overlap freeze-in and displaced vertices region!

Freeze-in

Freeze-in density depends on the reheat temperature
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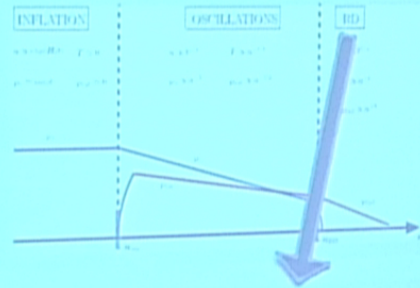
f_X
 $n_X = \int \frac{d^3p}{(2\pi)^3} f_X(p)$
 $R_X \ll R_B$

Inflationary Cosmology

INFLATION

$$\rho \simeq \rho_\phi \simeq \text{const} \simeq E_I^4$$

$$\rho_{\text{rad}} \simeq 0$$



INFLATON OSCILLATIONS

$$\rho \simeq \rho_\phi \propto a^{-3}$$

$$T_{\text{max}} \simeq (E_I^2 M_{\text{Pl}} \Gamma_\phi)^{1/4}$$

$$T \propto a^{-3/8}$$

REHEATING

$$\rho \simeq \rho_{\text{rad}} \propto a^{-4}$$

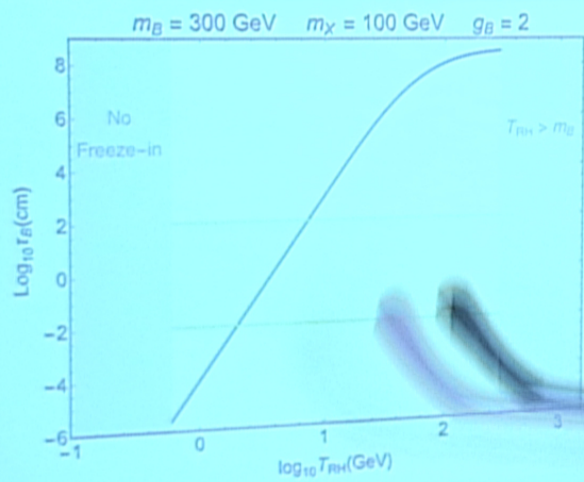
$$T_{\text{RH}} \simeq (M_{\text{Pl}} \Gamma_\phi)^{1/2}$$

$$T \propto a^{-1}$$

$$n_s = \frac{\partial \ln P(k)}{\partial \ln k}$$

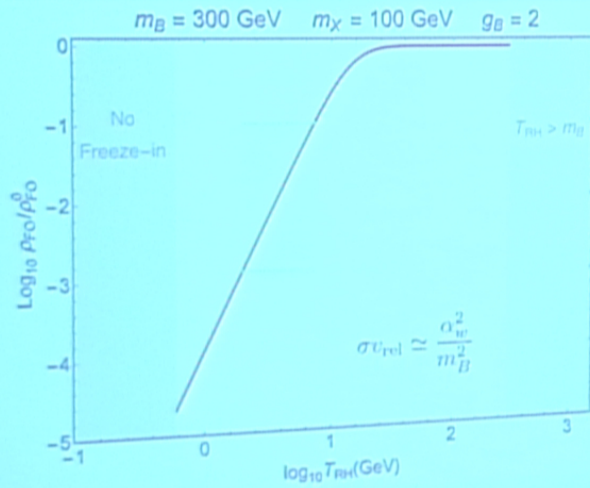
$$R_x \ll R_B$$

Displaced Vertices



Full overlap freeze-in and displaced vertices region!

Freeze-out of B?



Freeze-out contribution subdominant

$$f_X = \int \frac{d^3p}{(2\pi)^3} f_X(p, T)$$

$$R_X \ll R_B$$

Plan for Today's Talk

Freeze

Freeze-in

Freeze-in and "moduli"

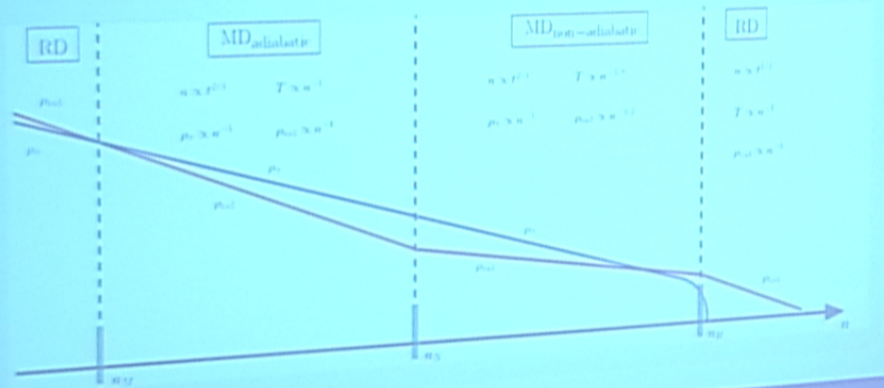
f_x

$\pi_x = \int_{\mathcal{P}^3} R_x(\vec{p}, \vec{c})$

$R_x \leftarrow \rightarrow$

“Moduli” Cosmology

“Modulus”: metastable massive particles which dominated the energy density of the universe and decayed out of equilibrium

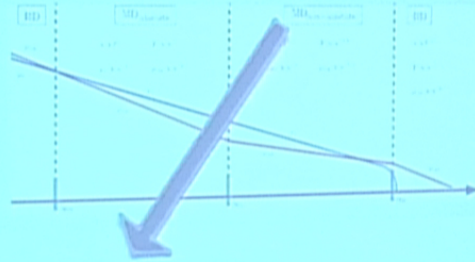


“Moduli” Cosmology

EARLY RD

$$\rho \simeq \rho_{\text{rad}} \propto a^{-4}$$

$$T \propto a^{-1}$$



MA adiabatic

starts at $T = T_M$

$$\rho \simeq \rho_\phi \propto a^{-3}$$

$$T \propto a^{-1}$$

MA non-adiabatic

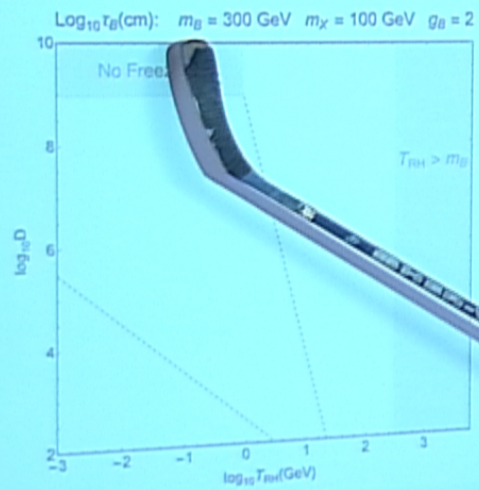
starts at $T = T_N$

$$\rho \simeq \rho_\phi \propto a^{-3}$$

$$T \propto a^{-3/8}$$



The (T_{RH}, D) plane

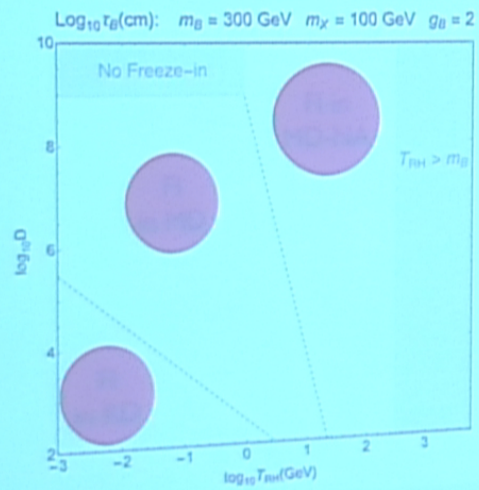


$$f_X$$

$$\eta_X = \int \frac{d^3p}{(2\pi)^3} R_X(p, t)$$

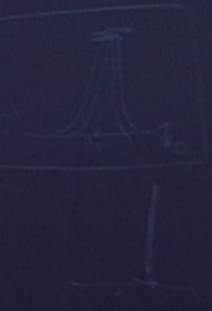
$$R_X \ll R_B$$

The (T_{RH}, D) plane

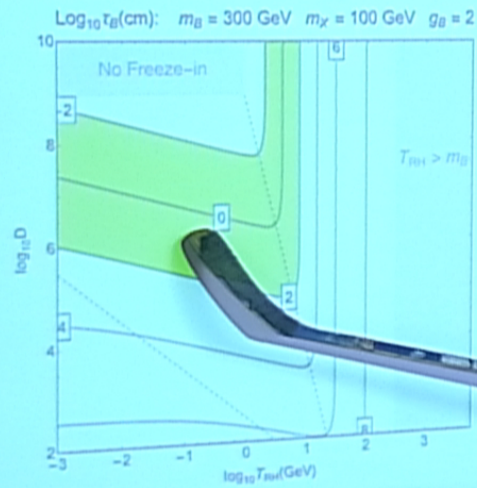


$$f_X = \int \frac{d^3p}{(2\pi)^3} \frac{d^4p}{(2\pi)^4} \dots$$

$$R_X \ll \dots$$



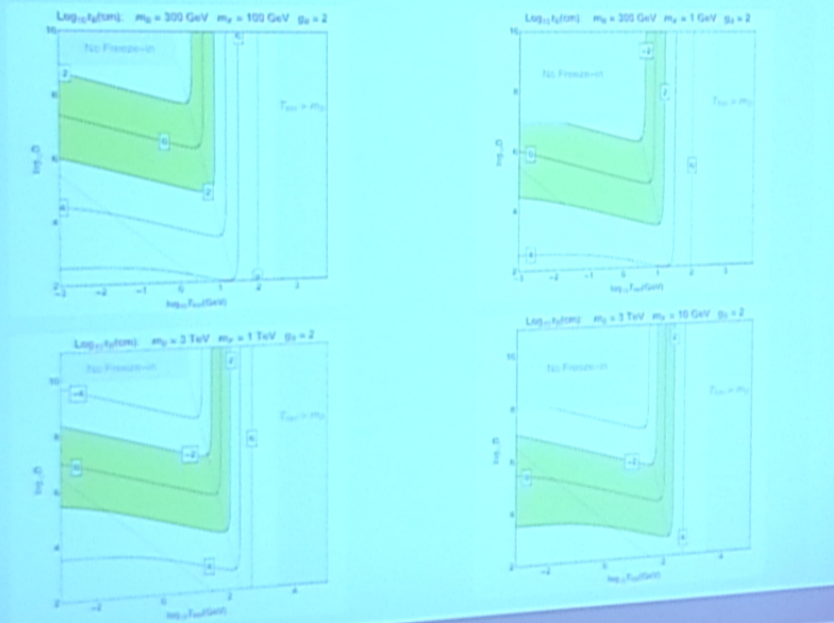
Displaced Vertices



Significant region of the plane has displaced vertices!

f_X
 $n_X = \int \frac{d^3 p}{(2\pi)^3} f_X(p, t)$
 $R_X \ll R_B$

Displaced Vertices

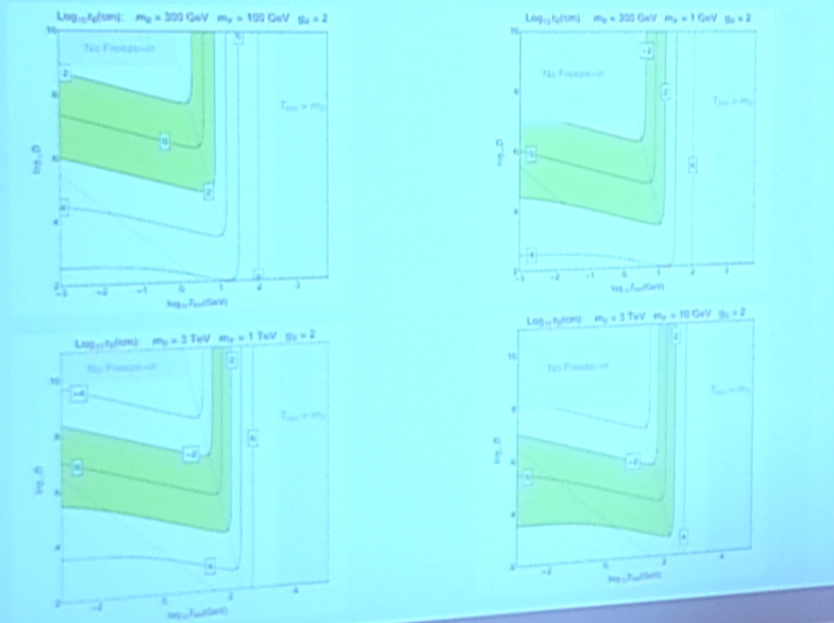


$$f_x$$

$$\eta_x = \int \frac{d^3p}{(2\pi)^3} R_x(p^0)$$

$$R_x \ll R_B$$

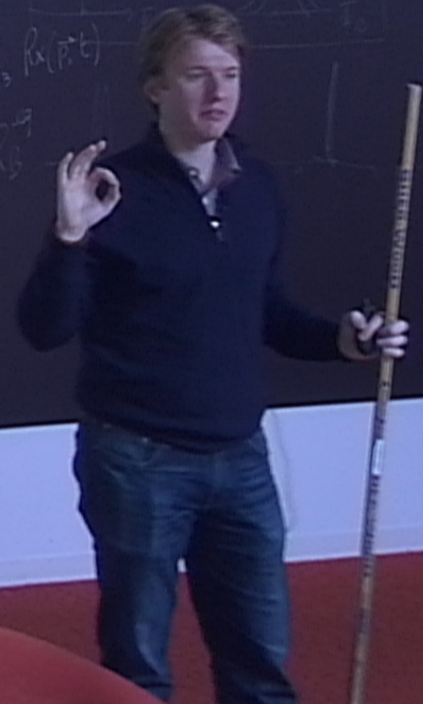
Displaced Vertices



$$f_x$$

$$n_x = \int \frac{d^3p}{(2\pi)^3} f_x(p, t)$$

$$R_x \ll R_B^{\text{dec}}$$

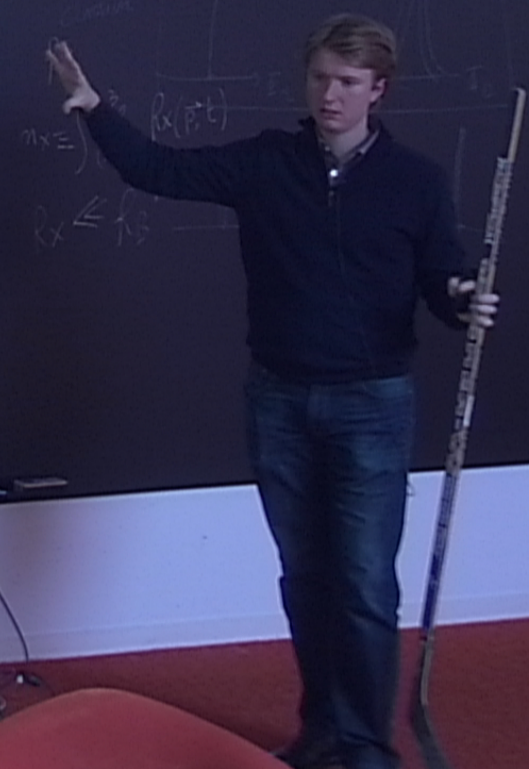


Plan for Today's Talk

Freeze

Freeze-in

A motivated candidate



Axino

$$W = \frac{S^2 H_u H_d}{M_*} \longrightarrow \mu \simeq \frac{f^2}{M_*} \simeq m_{\text{weak}}$$

Axino FIMP

$$\mathcal{L} \supset \frac{\mu}{f} \bar{h} \tilde{a} h$$

Cosmological LSP Axino Problem

$$\frac{\mu}{f} \simeq 10^{-8} \left(\frac{\mu}{1 \text{ TeV}} \right) \left(\frac{10^{11} \text{ GeV}}{f} \right) \longrightarrow \text{Bad, our universe}$$

Handwritten notes on a chalkboard:

$$n_x = \int \frac{d^3 p}{(2\pi)^3} f_x(p, t)$$
$$R_x \ll \frac{f^2}{M_*^2}$$

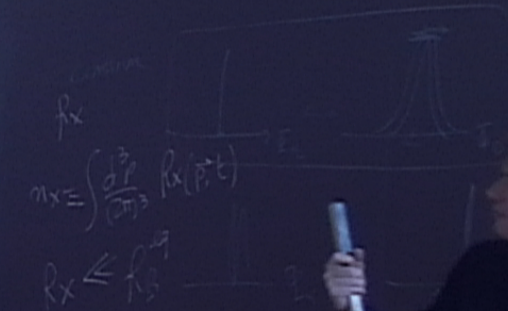
Displaced Axino

$$W = \frac{S^2 H_u H_d}{M_*} \quad \mathcal{L} \supset \frac{H \tilde{h} \tilde{a} h}{f} \quad \mu \simeq \frac{f^2}{M_*} \simeq \frac{m_{\text{weak}}}{\sqrt{2}}$$

Axino FIMP \tilde{N}_1 $\tilde{h} \tilde{a}$ $\mathcal{L} \supset \frac{\tilde{\mu}}{f} \tilde{h} \tilde{a} h \quad Z \tilde{a}$

Cosmological LSP Axino Problem $\left(\frac{f}{10^{11} \text{ GeV}}\right)^2$

$$\frac{\mu}{f} \simeq 10^{-8} \left(\frac{\mu}{1 \text{ TeV}}\right) \left(\frac{10^{11} \text{ GeV}}{f}\right) \rightarrow \text{Badly overcloses our universe}$$



Displaced Axino

$$\mathcal{L} \supset \frac{\mu}{f} \tilde{h} \tilde{a} h$$



$$\tilde{N}_1 \rightarrow h \tilde{a}$$

$$\tilde{N}_1 \rightarrow Z \tilde{a}$$

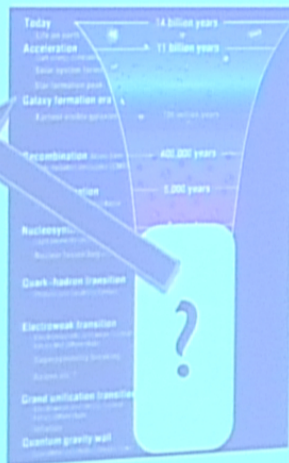
$$\frac{1}{\Gamma_{\tilde{N}_1}} \simeq \delta_{\text{mix}}^{-2} \left(\frac{\mu^3}{8\pi f^2} \right)^{-1} \simeq 5 \text{ cm} \left(\frac{1}{\delta_{\text{mix}}} \right)^2 \left(\frac{1 \text{ TeV}}{\mu} \right)^3 \left(\frac{f}{11 \text{ GeV}} \right)^2$$

Fine with alternative cosmology:

Outlook

Exploring freeze-in production for DM FIMP during a early matter dominated era (inflation, SUSY moduli, etc.)

- LOSP freeze-out suppressed
- Displaced vertices at colliders



Work in progress in collaboration with Raymond Co, Lawrence Hall and Duccio Pappadopulo

$$f_x$$
$$\eta_x = \int \frac{d^3p}{(2\pi)^3} f_x(p, t)$$
$$R_x \ll R_B$$