

Title: Double perovskite materials: From Chern bands to unusual Mott insulators

Date: Apr 30, 2015 11:30 AM

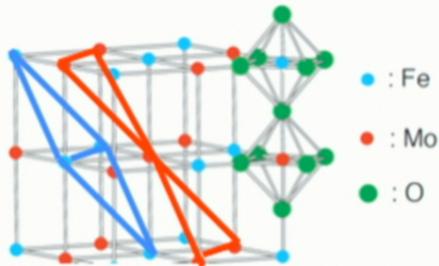
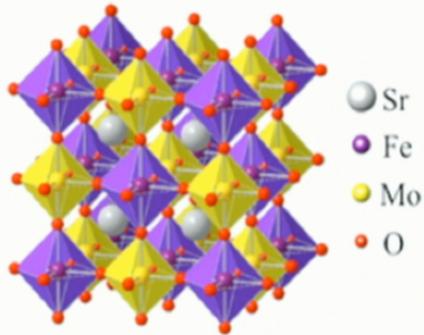
URL: <http://pirsa.org/15040159>

Abstract: Double perovskites, a class of oxide materials with 3d and 5d transition metal ions, can realize a wide variety of interesting phases. Here we focus on our recent theoretical work suggesting the appearance of Chern insulators, and possible emergent nematic phases at Chern transitions in such systems. We will also discuss Mott insulating double perovskites with iridium moments which appear to host unusual magnetic interactions on the 3-dimensional fcc lattice, and present comparisons with experimental data.

# Ordered Double Perovskites

General formula:  $A_2BB'O_6$  ( $B, B' = 3d, 4d, 5d$ )

Double perovskite lattice



## Metallic systems

B: Magnetism and B': Conduction electrons

- . Half metallic ferrimagnets (eg:  $Sr_2FeMoO_6$ ,  $T_C = 420K$ )
- . Large polarization: good for spin injection
- . Interplay of magnetism, **SOC**, metallic transport

## Mott insulators

B=magnetism, B'=inert or magnetism

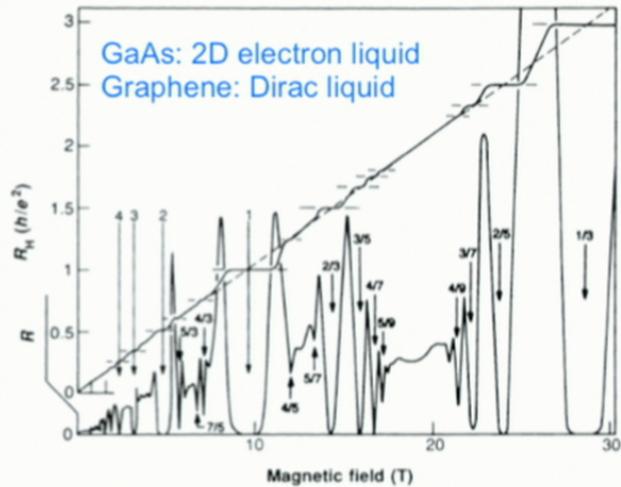
- . Get fcc lattice of spins (eg:  $Ba_2YReO_6$ )
- . Unusual spin-orbit coupled liquids?
- . Insulating ferrimagnets (eg:  $Sr_2CrOsO_6$ ,  $T_C = 725K$ )

## Outline of the talk

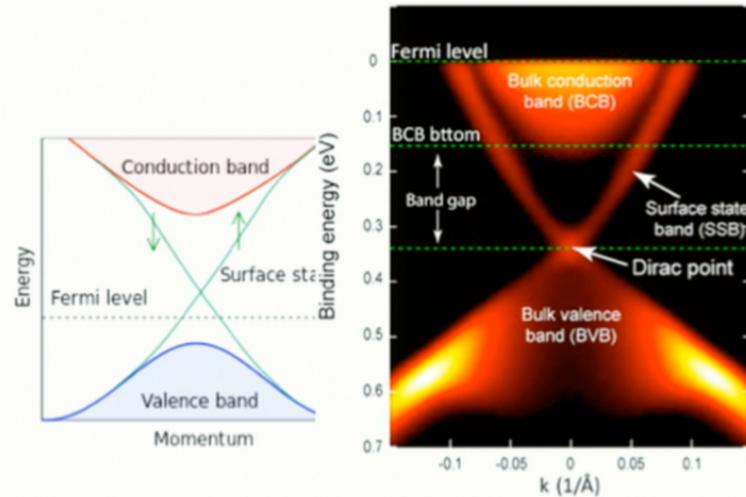
- I. Bulk metallic double perovskites with SOC  
- application to  $\text{Ba}_2\text{FeReO}_6$
- II. Spin dynamics in bulk double perovskites  
- application to  $\text{Ba}_2\text{FeReO}_6$
- III. Topological states in “engineered” double perovskite films  
- application to  $\text{Sr}_2\text{FeMoO}_6$
- IV. Exotic spin Hamiltonians in Mott insulating double perovskites  
- application to  $\text{La}_2\text{MgIrO}_6$  /  $\text{La}_2\text{ZnIrO}_6$

# Topological phases of quantum matter

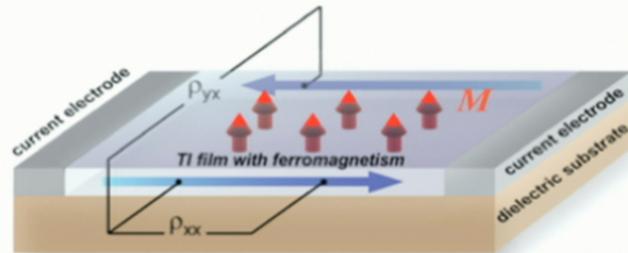
## Quantum Hall effects



## Topological insulators, top. crystalline insulators



## Quantum anomalous Hall insulators



$(\text{Bi,Sb})_2\text{Te}_3$  film doped with Cr or V atoms

- C.Z. Chang et al, Science 2013 (Xue group, Tsinghua)
- C.Z. Chang et al, arXiv (M. Chan + J. Moodera groups, PSU/MIT)
- A. J. Bestwick et al, arXiv (Goldhaber-Gordon group, Stanford)
- A. Kandala, et al, arXiv (N. Samarth and C.X. Liu groups, PSU)

# How to “engineer” correlated topological states in solids?

## (1) Strong correlations

- Brink of Mott localization or deep Mott insulator regime
- Common in 3d oxides: Kinetic energy  $\sim$  Interactions
- “High” energy/temperature scales for correlations/magnetism

Periodic Table of the Elements

Legend:

- Alkali Metal
- Alkaline Earth
- Transition Metal
- Post-Transition Metal
- Nonmetal
- Halogen
- Noble Gas
- Lanthanide
- Actinide

The periodic table shows elements from Hydrogen (H) to Oganesson (Og). A blue box highlights the row of transition metals from Titanium (Ti) to Zinc (Zn). The legend at the bottom identifies various groups of elements by color.

## How to “engineer” correlated topological states in solids?

- Q: Can we tie together these two disparate ingredients?  
Nontrivial topological bands + strong electronic correlations
  - Correlated 3d oxides, interaction effects, magnetism, etc strong!  
Spin orbit coupling  $\sim 10\text{meV}$  (typically negligible effects)
  - Simple 5d oxides - 5d orbitals strongly overlap, SOC large  $\sim 500\text{ meV}$   
 $\text{ReO}_3$ ,  $\text{IrO}_2$  (electrocatalyst): metallic (no magnetism, no strong correlations)
  - Ternary 5d oxides - 5d atoms further apart  
 $\text{Sr}_2\text{IrO}_4$ : Magnetism, but Slater-Mott insulator, so no conduction fluid
- Beyond simple and ternary oxides?
- Can we tailor SOC and correlations independently?

# “Oxide engineering”

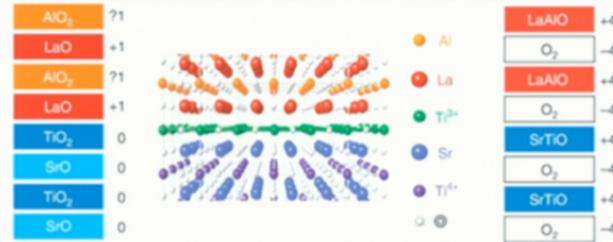
## Controlled layer-by-layer growth using MBE/PLD

### LaAlO<sub>3</sub>/SrTiO<sub>3</sub> interface [100]

H. Y. Hwang, J.M. Triscone, J. Mannhart,  
R. Ashoori, K. A. Moler, ...

**Basic physics:** Magnetism + Superconductivity

**Applications:** Write/Erase circuits using electric field



### Prediction of topological insulators in simple TMO bilayers

ARTICLE

Received 20 Jun 2011 | Accepted 18 Nov 2011 | Published 20 Dec 2011

DOI: 10.1038/nmat1324

Interface engineering of quantum Hall effects in digital transition metal oxide heterostructures

Di Xiao<sup>1</sup>, Wenguang Zhu<sup>1,2</sup>, Ying Ran<sup>3</sup>, Naoto Nagaosa<sup>4,5</sup> & Satoshi Okamoto<sup>1</sup>

Rapid Communication

Topological insulators from complex orbital order in transition-metal oxides heterostructures

Andreas Rüegg and Gregory A. Fiete  
Phys. Rev. B **84**, 201103(R) – Published 14 November 2011

### 3d-3d Superlattices along [111]

nature materials

LETTERS

PUBLISHED ONLINE: 22 JANUARY 2012 | DOI: 10.1038/nmat1324

#### Exchange bias in LaNiO<sub>3</sub>-LaMnO<sub>3</sub> superlattices

Marta Gibert<sup>1\*</sup>, Pavlo Zubko<sup>1</sup>, Raoul Scherwitzl<sup>1</sup>, Jorge Íñiguez<sup>2</sup> and Jean-Marc Triscone<sup>1</sup>

### 3d/5d Superlattices

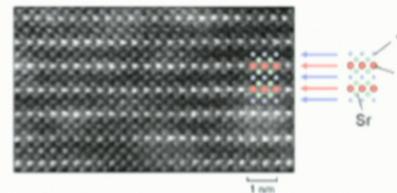
H. Takagi group (2014, 2015)



Local electronic and magnetic studies of an artificial La<sub>2</sub>FeCrO<sub>6</sub> double perovskite

Benjamin Gray, Ho Nyung Lee, Jian Liu, J. Chakhalian, and J. W. Freeland

Citation: Applied Physics Letters **97**, 013105 (2010); doi: 10.1063/1.3455323



# “Engineering” spin-orbit coupled topological states in solids

nature  
materials

LETTERS

PUBLISHED ONLINE: 22 JANUARY 2012 | DOI: 10.1038/NMAT3224

## Exchange bias in $\text{LaNiO}_3$ - $\text{LaMnO}_3$ superlattices

Marta Gibert<sup>1\*</sup>, Pavlo Zubko<sup>1</sup>, Raoul Scherwitzl<sup>1</sup>, Jorge Íñiguez<sup>2</sup> and Jean-Marc Triscone<sup>3</sup>

$(\text{LaNiO}_3)_m$ - $(\text{LaMnO}_3)_n$  superlattices along  $\{111\}$   
 $(\text{LaFeO}_3)_m$ - $(\text{LaCrO}_3)_n$  superlattices along  $\{111\}$   
Infinite  $m=1, n=1$  superlattice: Double Perovskite

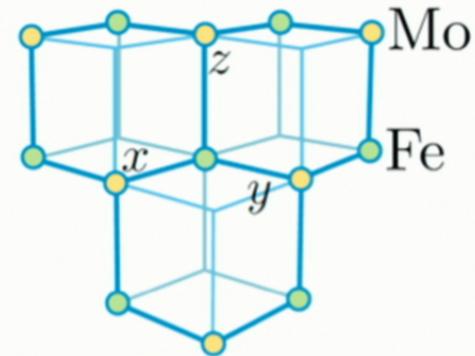
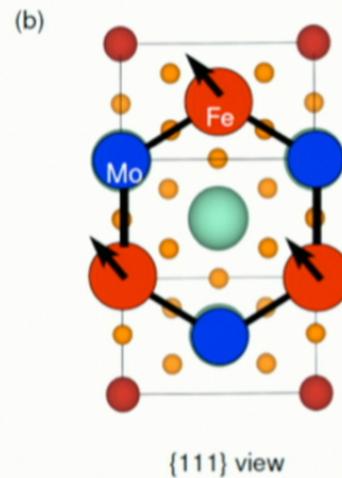
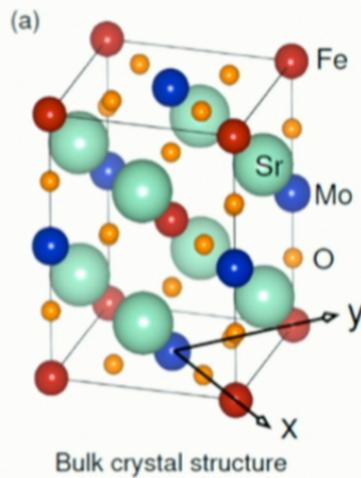


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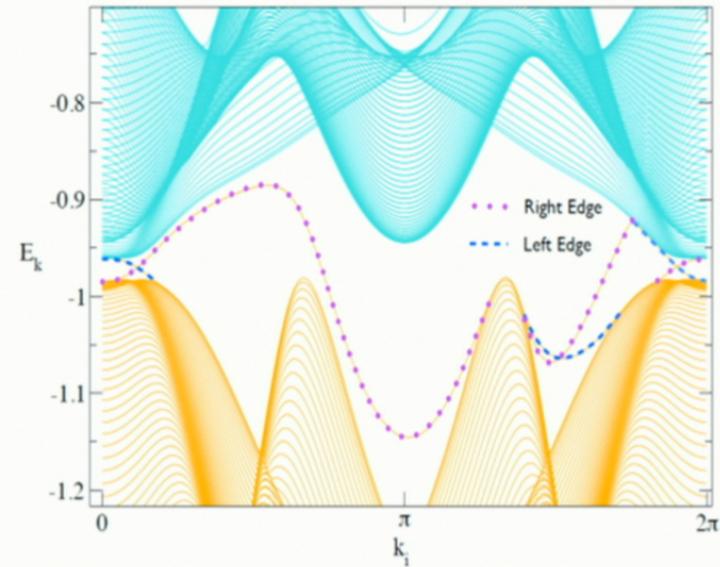
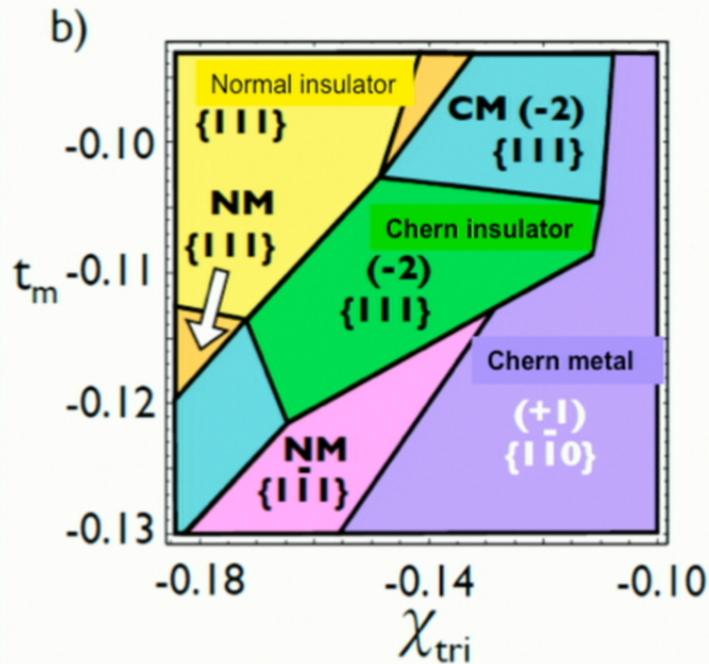
Citation: *Applied Physics Letters* 97, 013105 (2010); doi: 10.1063/1.3455323

## SFMO: $\{111\}$ bilayer



Fe and Mo/Re on honeycomb lattice  
Fe: Local moments  
Mo/Re: Itinerant  $t_{2g}$  electrons

## Phase diagram of SFMO in {111} bilayer



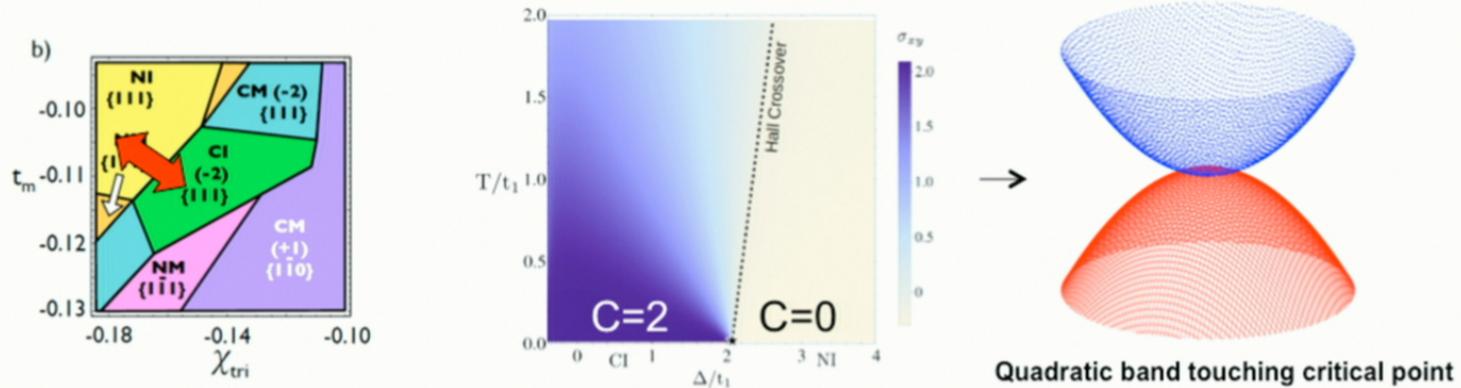
Edge state spectrum in  
C=2 Chern insulator

Quantum Hall gap  $\sim 75K$

**SFMO bilayer supports  
topologically nontrivial metals  
and C=2 Chern insulator**

. A. Cook, AP (PRL 2014)

## Can interactions drive new phases at topological critical points?



$$H = -\sum_{\langle ij \rangle} t_{ij}^{\alpha\beta} c_{i\alpha}^\dagger c_{j\beta} + \Delta \sum_i (n_{i\uparrow} - n_{i\downarrow}) + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

$$H_\Gamma^{\text{low}} = \left( \begin{array}{cc} \frac{3}{2}(t_2+t_1)k^2 + r & \frac{3}{4}t_3(k_x - ik_y)^2 \\ \frac{3}{4}t_3(k_x + ik_y)^2 & \frac{3}{2}(t_2-t_1)k^2 - r \end{array} \right)$$

$$S_{\text{int}} = u \int_0^\beta d\tau \sum_i \bar{\psi}_{i\uparrow}(\tau) \bar{\psi}_{i\downarrow}(\tau) \psi_{i\downarrow}(\tau) \psi_{i\uparrow}(\tau)$$

Flow

$$\frac{dr}{d\ell} = 2r + \frac{u\Lambda^2}{4\pi} \frac{t_1}{\sqrt{t_1^2 + t_3^2/4}}$$

$$\frac{du}{d\ell} = \frac{u^2}{6\pi} \frac{1}{\sqrt{t_1^2 + t_3^2/4}}$$

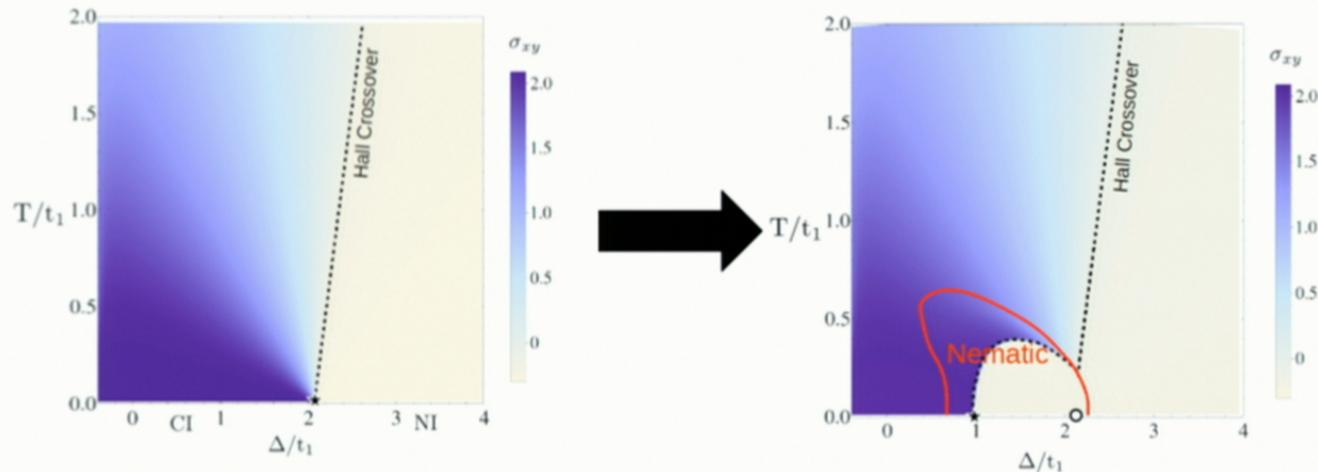
- Single "hot spot"
- Single marginally relevant coupling

## Chern transition: Interactions can drive “emergent” phases

$$\vec{m} = \frac{1}{2N} \sum_{\mathbf{k}} \langle c_{\mathbf{k}\alpha}^\dagger \vec{\sigma}_{\alpha\beta} c_{\mathbf{k}\beta} \rangle$$

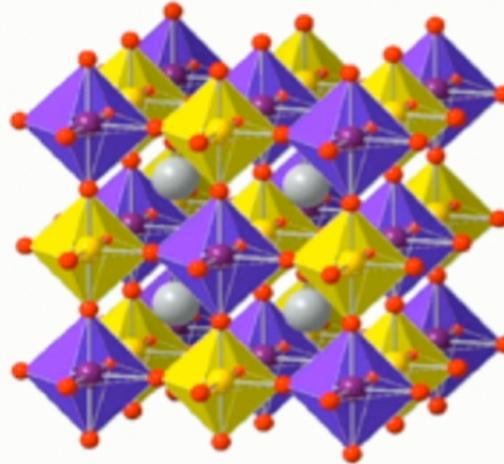
$$H_{\text{mf}}(\mathbf{k}) = \begin{pmatrix} A_{\mathbf{k}} - \mu - U m_z & D_{\mathbf{k}} - U(m_x - i m_y) \\ D_{\mathbf{k}}^* - U(m_x + i m_y) & B_{\mathbf{k}} - \mu + U m_z \end{pmatrix}$$

Breaks rotational symmetry  
- Emergent “nematic order”



## Can SOC lead to unusual magnetism in Mott insulators?

### Double perovskite lattice



### Mott insulating double perovskites

A sublattice: Inert closed shell atom

B sublattice: Magnetic ion

- Atoms farther apart, suppresses delocalization
- Rich variety of fcc lattice Mott insulators

$\text{Ba}_2\text{YMoO}_6$ ,  $\text{Sr}_2\text{YReO}_6$ ,  $\text{Sr}_2\text{YIrO}_6$ ,  $\text{Ba}_2\text{YRuO}_6$ ,...

T. Aharen et al, Phys. Rev. B 81, 224409 (2010)

J. P. Carlo, et al, Phys. Rev. B 88, 024418 (2013)

G. Cao, et al, Phys. Rev. B 87, 155136 (2013)

G. Chen, L. Balents, PRB 84, 094420 (2011)

T. Dodds, T.P.Choy, Y.B. Kim, PRB 84, 104439 (2011)

## Iridate double perovskites

$\text{La}_2\text{ZnIrO}_6$ ,  $\text{La}_2\text{MgIrO}_6$  [G. Cao, et al, PRB 2013]

- Small CW temperature,  $\Theta_{\text{CW}} \sim -24\text{K}$  in  $\text{La}_2\text{MgIrO}_6$
- Deep in the Mott insulator phase: Consistent with larger Ir-Ir spacing compared to perovskites or honeycomb iridates
- A-type magnetic ordering at  $\sim 10\text{K}$  (FM layers, AFM stacking)
- Opportunity to explore magnetism of SOC moments on the frustrated fcc lattice

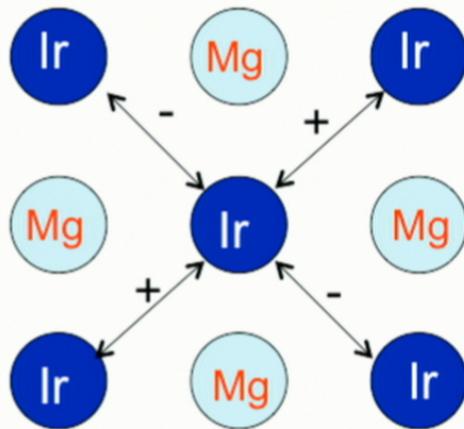
# Magnetic Hamiltonian: Symmetry analysis

## Honeycomb $\alpha$ - $\text{Na}_2\text{IrO}_3/\text{Li}_2\text{IrO}_3$

Jackeli, Khaliullin, Chaloupka (2011,2012) – **Kitaev Hamiltonian**  
J. Rau, H. Y. Kee (PRL+PRB 2014) – **Beyond Kitaev**

## Hyper-honeycomb $\beta/\gamma$ - $\text{Li}_2\text{IrO}_3$

Takagi group (2014), J. Analytis group (2014) – **3D Kitaev generalization**  
I.Kimchi, R. Coldea, A. Vishwanath (2014); E.K.H.Lee, Y.B.Kim (2014)



$$H_{xy} = J\vec{S}_1 \cdot \vec{S}_2 + J_I S_1^z S_2^z \pm \Gamma(S_1^x S_2^y + S_1^y S_2^x)$$

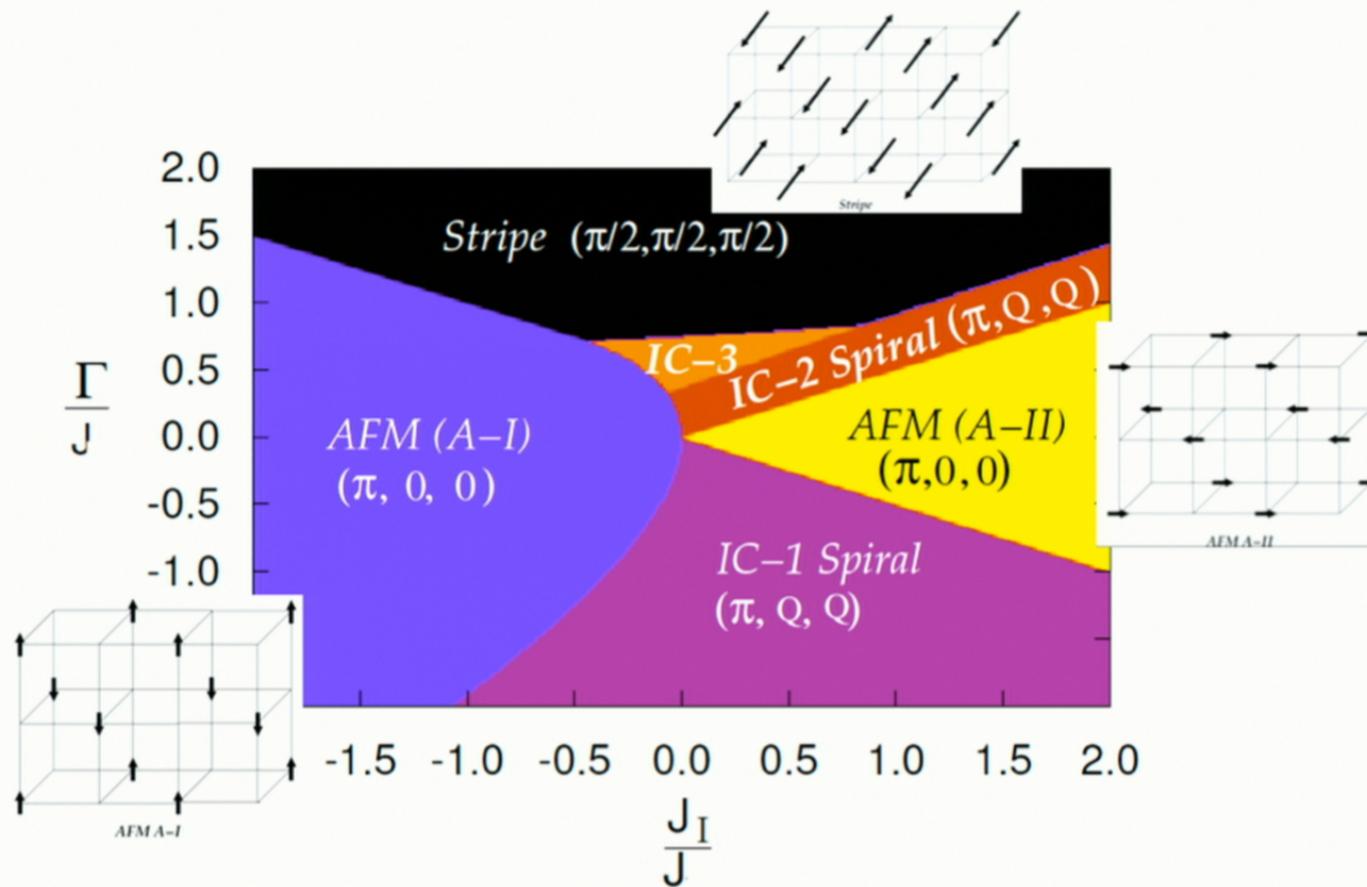
$$H_{yz} = J\vec{S}_1 \cdot \vec{S}_2 + J_I S_1^x S_2^x \pm \Gamma(S_1^y S_2^z + S_1^z S_2^y)$$

$$H_{xz} = J\vec{S}_1 \cdot \vec{S}_2 + J_I S_1^y S_2^y \pm \Gamma(S_1^x S_2^z + S_1^z S_2^x)$$

H. Ishizuka, L. Balents (PRB 2014)

A. Cook, S. Matern, C. Hickey, A. A. Aczel, AP (arXiv:1502.01031)

# Luttinger-Tisza phase diagram



A. Cook, S. Matern, C. Hickey, A. A. Aczel, AP (arXiv:1502.01031)

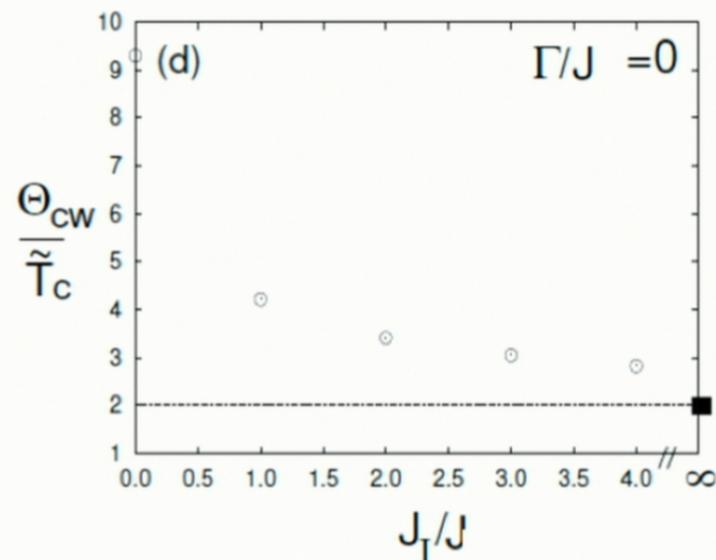
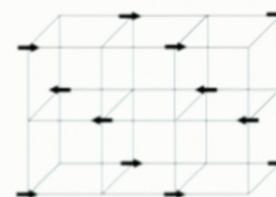
## Comparison with materials

### $\text{La}_2\text{MgIrO}_6$

- .  $\Theta_{\text{CW}} = -24\text{K}$ , orders with  $T_c = 12\text{K}$
- . Ordering type: A-II  
(neutron diffraction + ab initio)
- .  $\Theta_{\text{CW}} = - (3 J + J_I)$
- . Getting  $\Theta_{\text{CW}}/T_c$  to be small needs big SOC – expect  $J_I$  dominant

### $\text{La}_2\text{ZnIrO}_6$

- .  $\Theta_{\text{CW}} = -5\text{K}$ , orders with  $T_c = 7.5\text{K}$
- . Ordering type: A-II  
(neutron diffraction + ab initio)
- . Experimental uncertainties in  $\Theta_{\text{CW}}$ ?



Material has octahedral rotations which are staggered between layers  
Do we expect a ferromagnetic moment in the ordered AFM state?

**Experiment:** Observed in  $\text{La}_2\text{ZnIrO}_6$  (ok) but not in  $\text{La}_2\text{MgIrO}_6$  (puzzle)

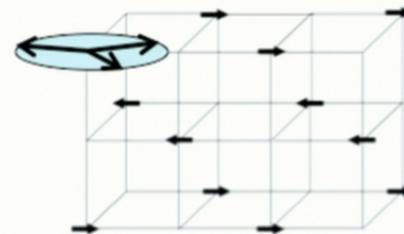
**Prediction:** Staggered octahedral rotations do not coincide with stacking direction in  $\text{La}_2\text{MgIrO}_6$  – test using Xrays

A. Cook, S. Matern, C. Hickey, A. A. Aczel, AP (arXiv:1502.01031)

## Spin dynamics

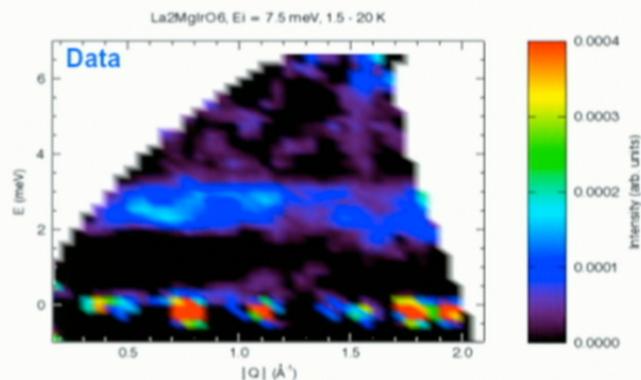
- Minimal Hamiltonian: "Quantum Ising model"

$$H = J_I \sum_{\langle ij \rangle_{XY}} S_i^z S_j^z + J_I \sum_{\langle ij \rangle_{YZ}} S_i^x S_j^x + J_I \sum_{\langle ij \rangle_{XZ}} S_i^y S_j^y$$



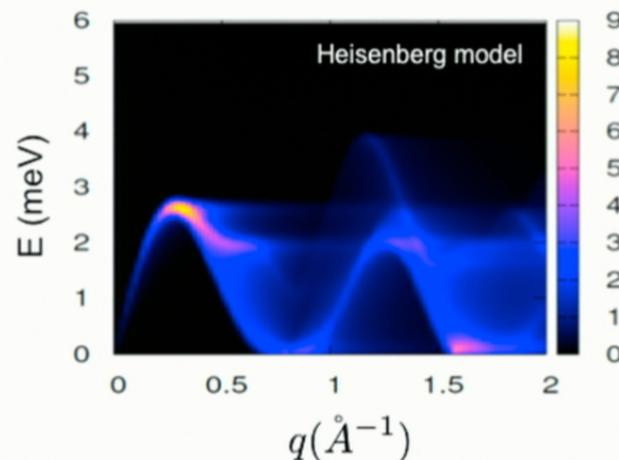
Classically: Spins can point anywhere in XY plane

- Thermodynamic studies:  $J_1 \sim 24\text{K}$
- Quantum order by disorder**
- Pins moments** to point along Ir-O bond direction
- Leads to a gap** in the spin wave spectrum



Note: Ir is a strong neutron absorber

- Data much less dispersive, lower in energy
- Impurity physics? Spin-phonon coupling?



A. A. Aczel, A. Cook, AP (manuscript in preparation)

# Summary

- **3D metallic double perovskites:  $\text{Ba}_2\text{FeReO}_6$** 
  - band structure with Weyl nodes
  - spin and orbital magnetization from strong correlations
  - spin dynamics in agreement with neutron data
- **2D double perovskite films along {111}:  $\text{Sr}_2\text{FeMoO}_6$** 
  - Topological phases including emergent Chern bands
  - C=2 quantum anomalous Hall insulators
- **Mott insulating double perovskites:  $\text{La}_2\text{ZnIrO}_6$ ,  $\text{La}_2\text{MgIrO}_6$** 
  - Spin Hamiltonian with exotic non-Heisenberg interactions
  - magnetic order in agreement with experimental data
  - resolve differences between  $\text{La}_2\text{ZnIrO}_6$  and  $\text{La}_2\text{MgIrO}_6$
  - Qualitative understanding of spin dynamics

## References

1. K. Plumb, A. Cook, J.P. Clancy, A. Kolesnikov, B.C. Jeon, T.W. Noh, AP, Y.J. Kim, PRB 87, 184412 (2013)
2. A. Cook, AP, PRB 88, 235102 (2013)
3. A. Cook, AP, PRL 113, 077203 (2014)
4. A. Cook, C. Hickey, AP, PRB 90, 085145 (2014)
5. A. Cook, S. Matern, C. Hickey, A. A. Aczel (arXiv:1502.01031)

THANK YOU!