Title: Black Holes and Reversibility

Date: Apr 13, 2015 02:00 PM

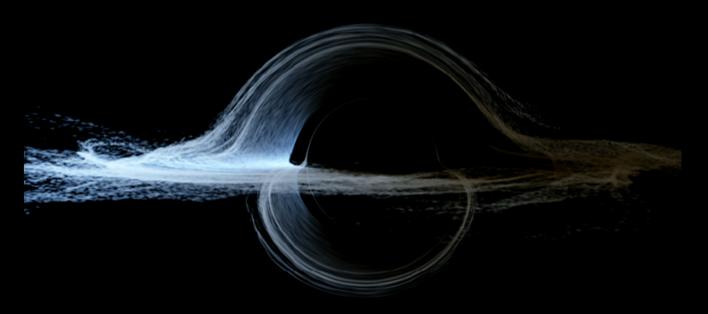
URL: http://pirsa.org/15040140

Abstract: It has become a platitute to say that black holes are fascinating objectsâ€"but they really are, in part because they challenge our understanding of the fundamental reversibility of physical processes.

In the first part of the talk, I will review some of the classical ways black holes behave as dissipative systems, such as the "hair loss― phenomenon and the monotonic growth of horizon area. In the second part, I will explain how quantum mechanics (more precisely, the coupling of black holes to the quantum vacuum) affects the classical picture at late times, notably through particle creation and evaporation. I will argue that techniques from two-dimensional field theory can help bring clarity to the associated "information loss―

problem, and perhaps also point to new, unexpected predictions. My approach will be as model-independent as possible; that is, rather than investigating a particular scenario for black hole evaporation, I will aim to derive generic consequences from basic assumptions regarding the reversibility of black hole evaporation.

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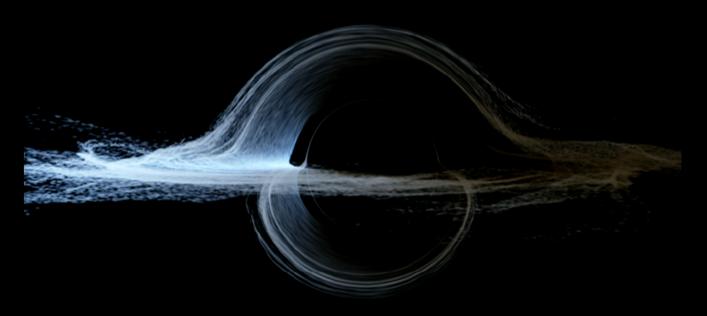


THORNE ET AL. 15

"[Black holes] are the most perfect macroscopic objects in the universe: the only elements in their construction are our concepts of space and time."

Chandrasekhar

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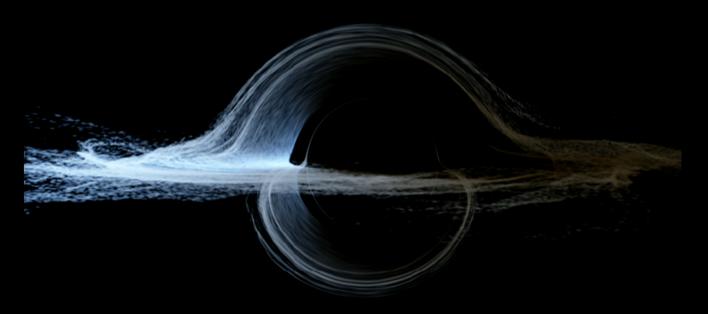


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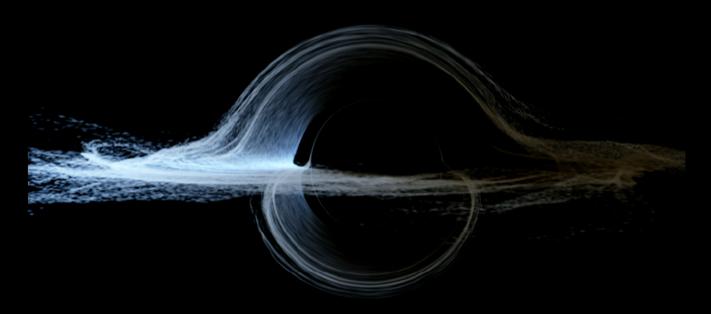


THORNE ET AL. 15

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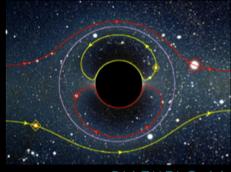
THORNE ET AL. 15

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A SPACE ODDITY



can go in, cannot come out

RIAZUELO 14

- hole is black, but see images of the entire sky near it
- a trip near a black hole is a trip to the future

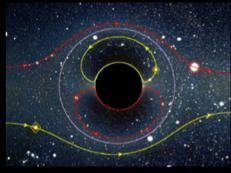
$$\Delta t_r = (1 - r_s/r)^{1/2} \, \Delta t$$

black hole are small outside, large inside

 $V(t) \sim 3\sqrt{3}M^2t$

CHRISTODOULOU, ROVELLI 15

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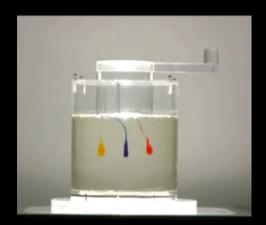
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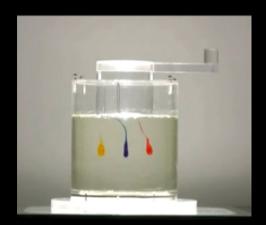
"Everything that can be done, can be undone*."



* in principle, if not in practice

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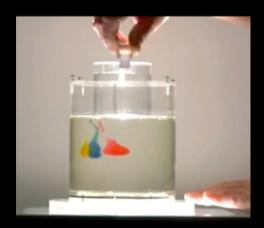
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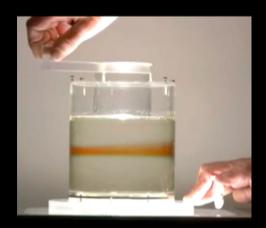
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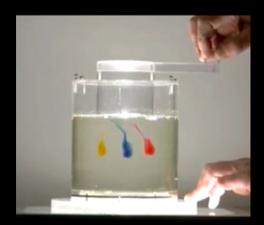
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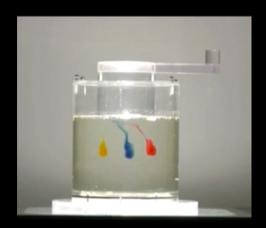
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Pirsa: 15040140 Page 13/64

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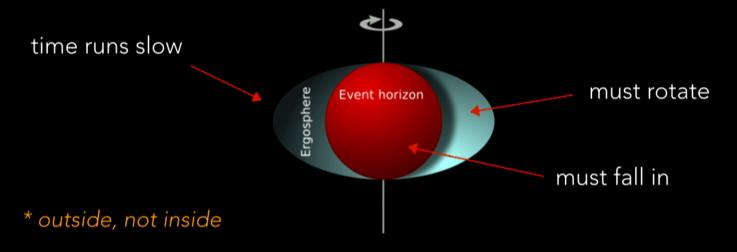
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1. CLASSICAL IRREVERSIBILITY

$$R_{ab} - \frac{1}{2}Rg_{ab} = \kappa T_{ab}$$

NO HAIR THEOREM: STATIONARY BLACK HOLES HAVE ONLY 2 PARAMETERS: M AND J*

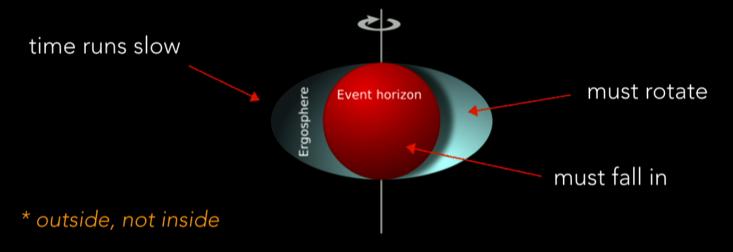
This means that all stationary black holes can be described exactly, with no approximation: they are indeed "perfect".



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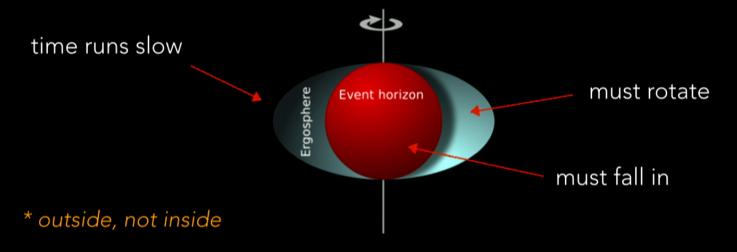
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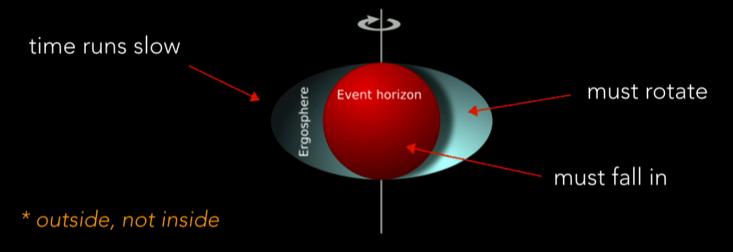
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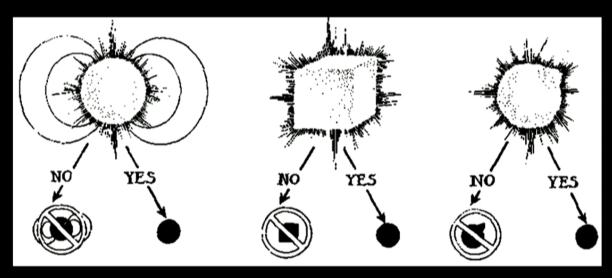
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BLACK HOLE BALDING - OLD INSIGHTS

PRICE'S LAW: WHATEVER CAN BE RADIATED IS RADIATED*



©THORNE

PRICE 72

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^{*} albeit slowly, i.e. with power law tails

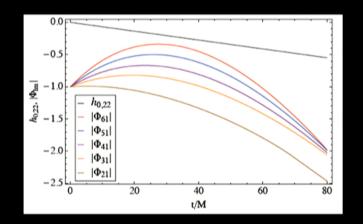
BLACK HOLE BALDING - NEW INSIGHTS

Turbulent balding

• parametric instability

$$\ddot{x} + \gamma \dot{x} + \omega(t)^2 x = 0$$

 inverse energy cascade small scales → large scales



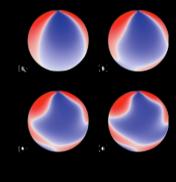
Gravitational Reynolds number:

$$Re_g = \frac{h_0}{m\gamma}$$
 amplitude

azimuthal frequency

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decay rate



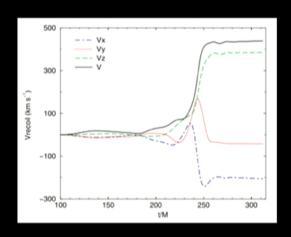
HUANG ET AL. 14

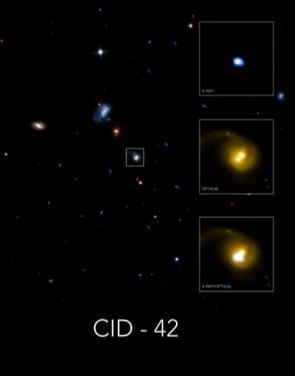
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BLACK HOLE KICKS

Large recoil velocity ($\approx 1\% c$) after merger:

- large spin orbit coupling
- focusing of gravitational waves





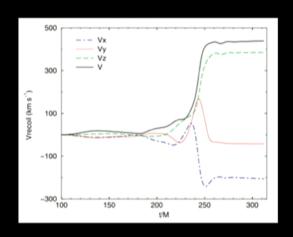
CAMPANELLI ET AL. 07

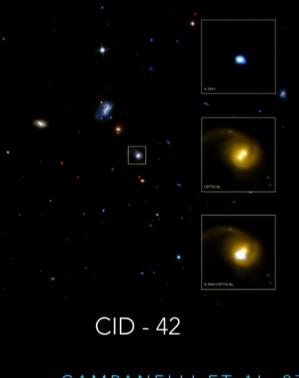
Pirsa: 15040140 Page 22/64

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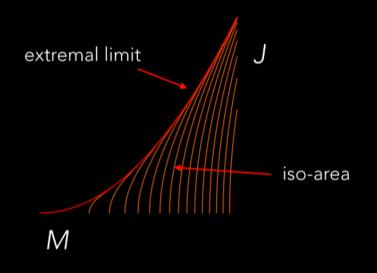


CAMPANELLI ET AL. 07

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BLACK HOLE TRANSFORMATIONS

AREA LAW: THE AREA OF A BLACK HOLE NEVER DECREASES*



* when interacting with normal matter

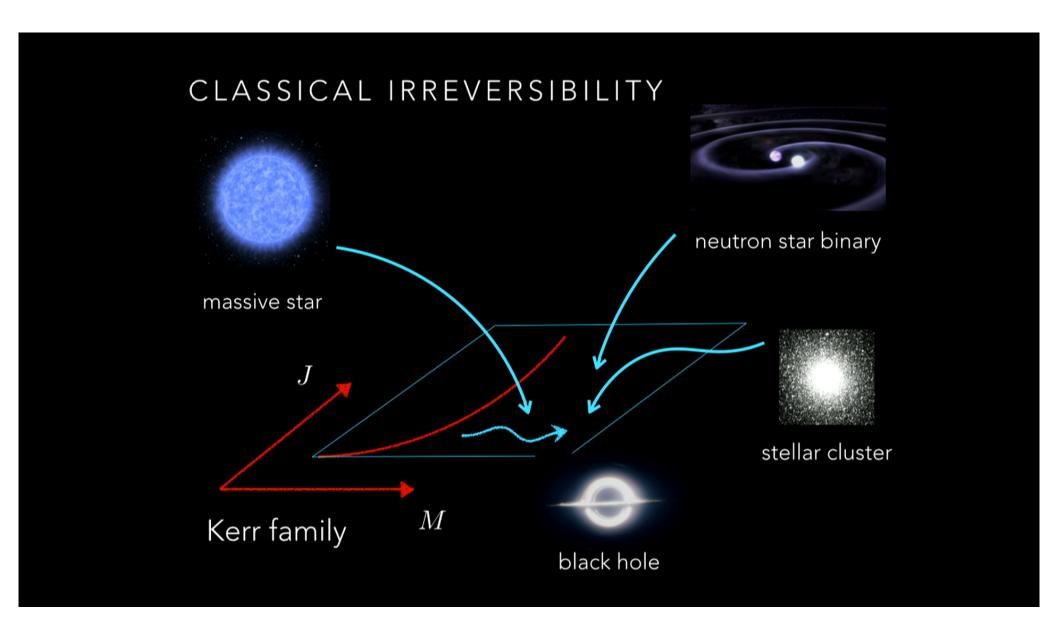
HAWKING 71

Two types of black hole transformations:

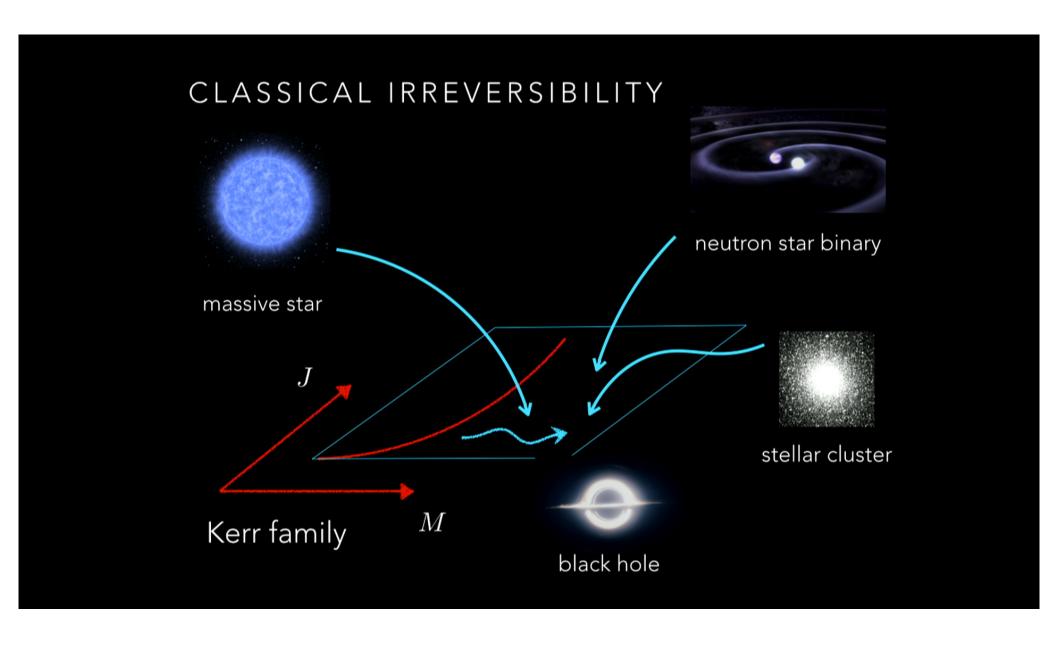
- reversible if $\delta A = 0$ (\approx add/remove spin)
- irreversible if δA > 0
 (≈ add mass)

CHRISTODOULOU 70

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SINGULARITIES - OLD INSIGHTS

SINGULARITY THEOREM: ALL BLACK HOLES CONTAIN SINGULARITIES*

PENROSE 65; HAWKING, PENROSE 70

An infalling observer will run into "trouble" in finite time.

The nature of this "trouble" is left unspecified.

* if the black hole contains normal matter

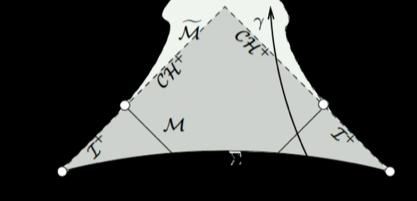
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SINGULARITIES - NEW INSIGHTS

Old idea: time and space end at the (spacelike) singularity.

More careful computations suggest weak null singularities:

- spacetime can be extended continuously
- curvature blows up but not tidal disruptions
- unpredictable other universe?



ISRAEL, POISSON 89; DAFERMOS 13

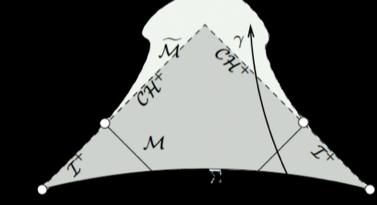
Pirsa: 15040140 Page 28/64

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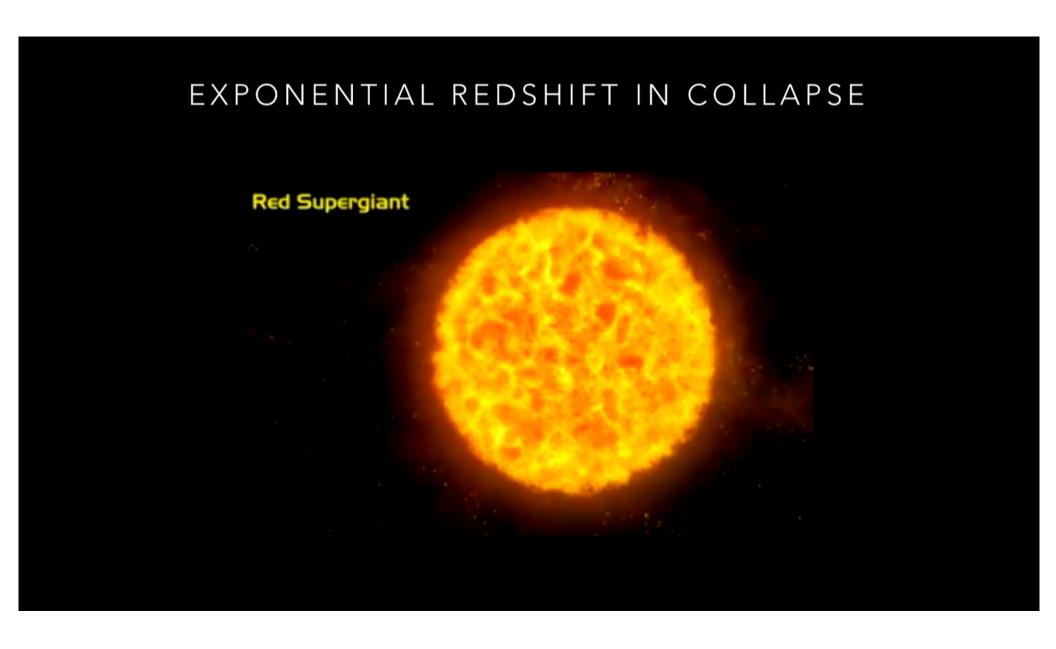


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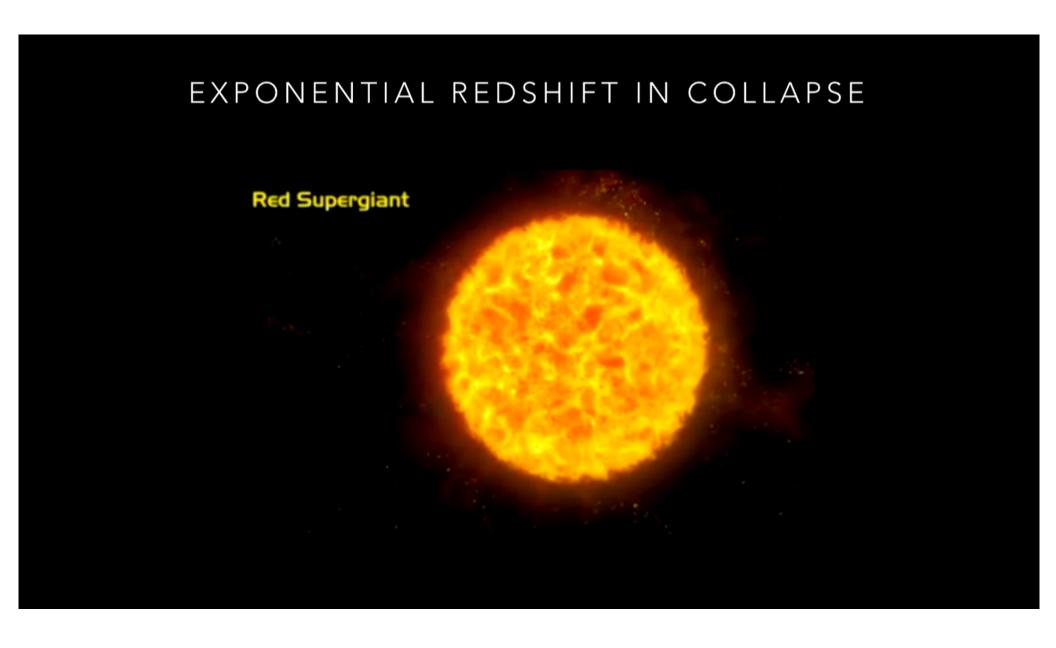
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2. QUANTUM REVERSIBILITY?

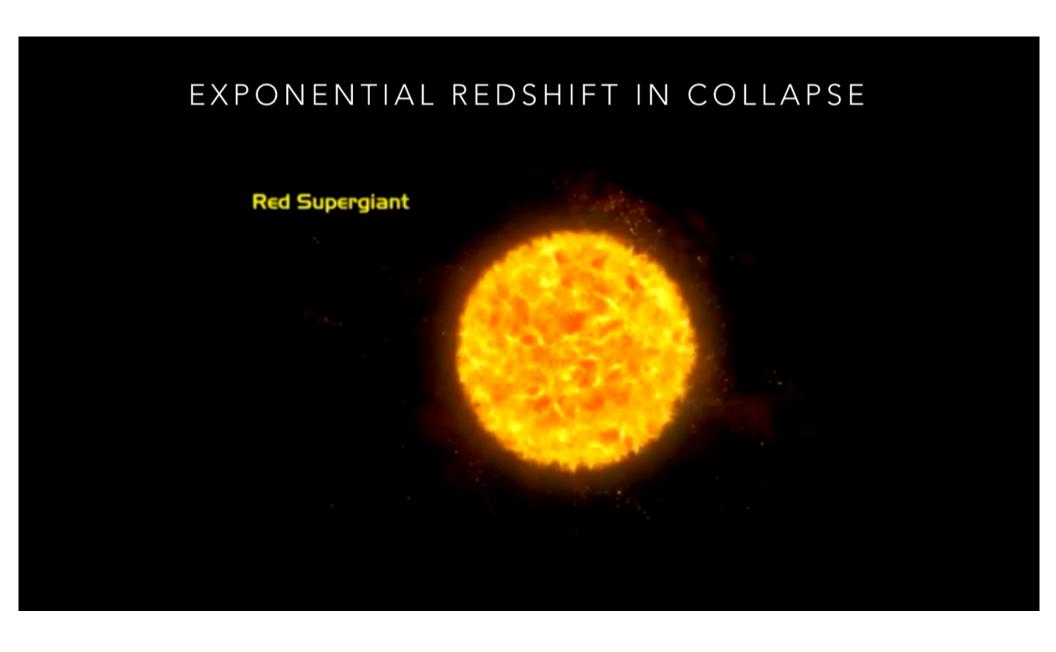
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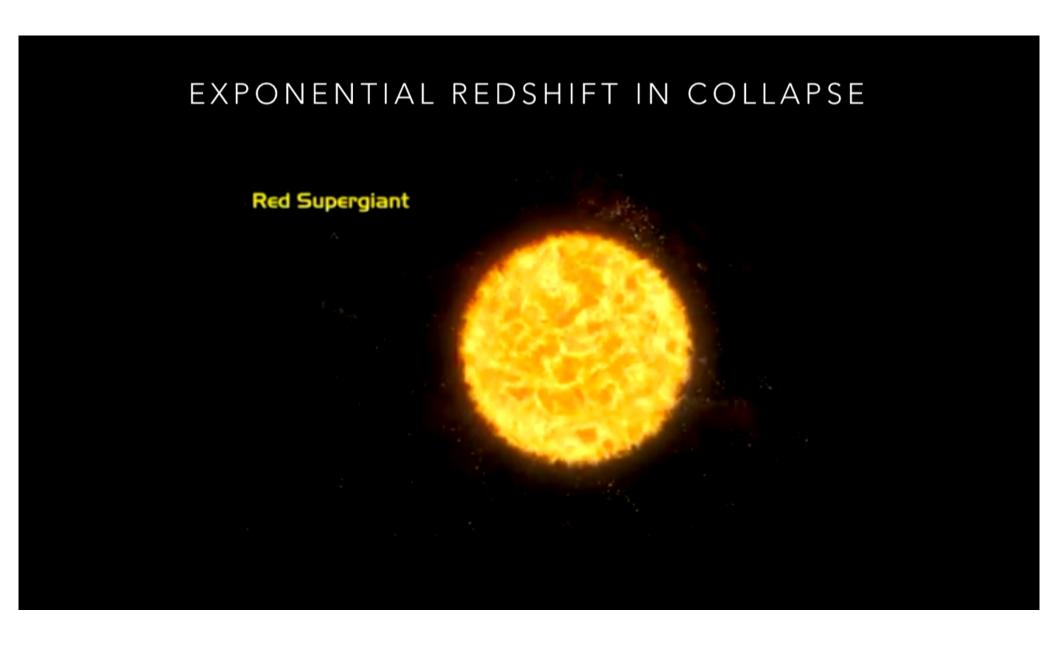
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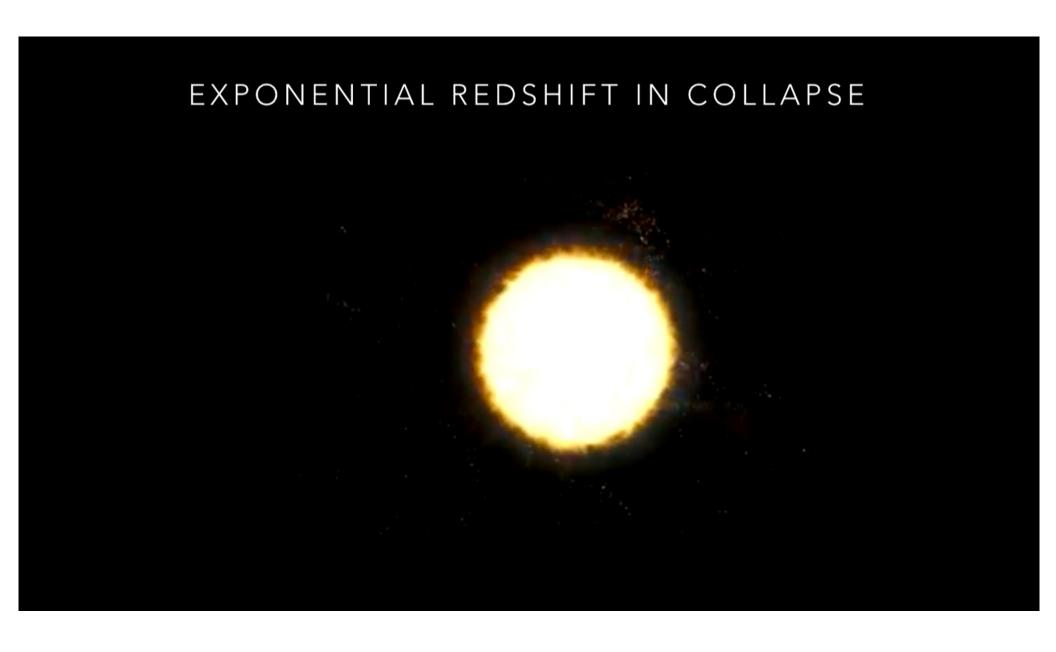
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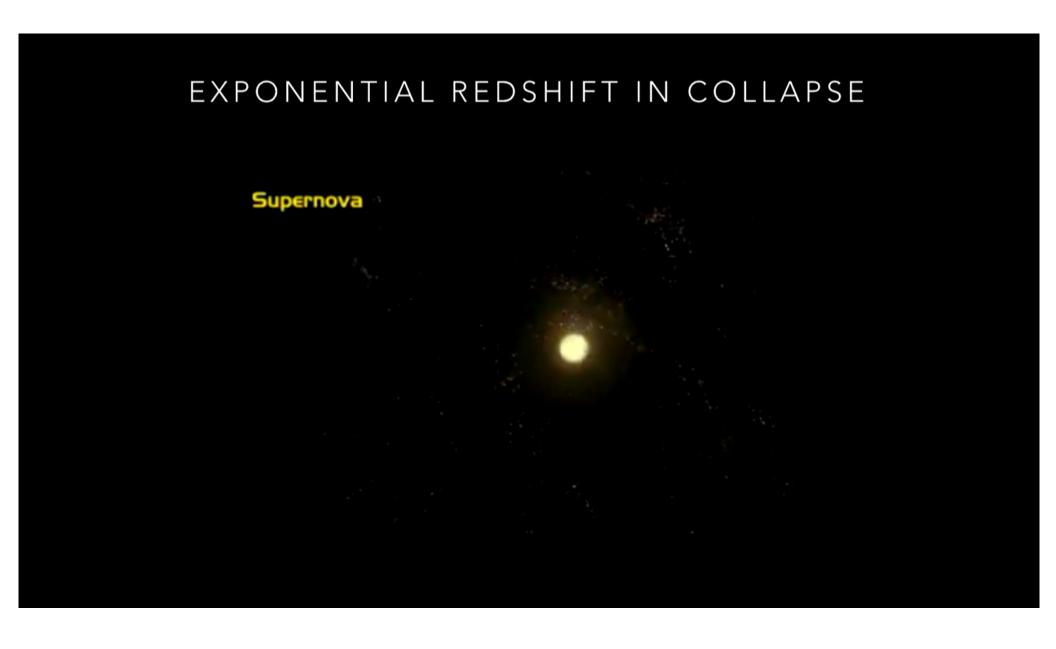
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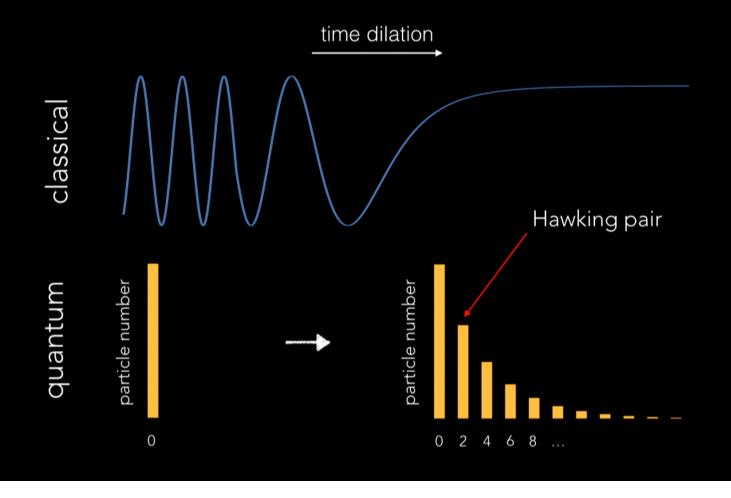


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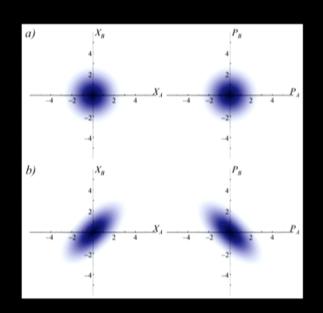
Pirsa: 15040140 Page 36/64

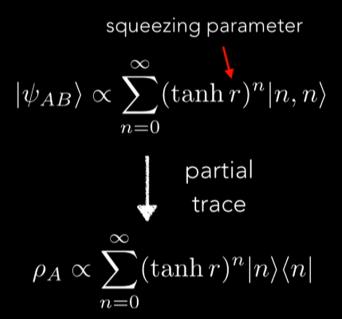
THE HAWKING EFFECT - PAIR CREATION



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TWO-MODE SQUEEZED VACUUM STATE

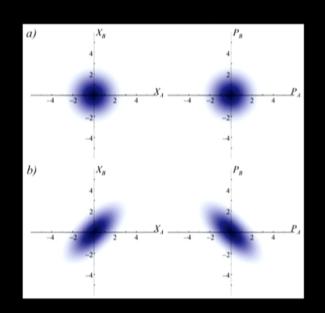




Higher squeezing, higher temperature.

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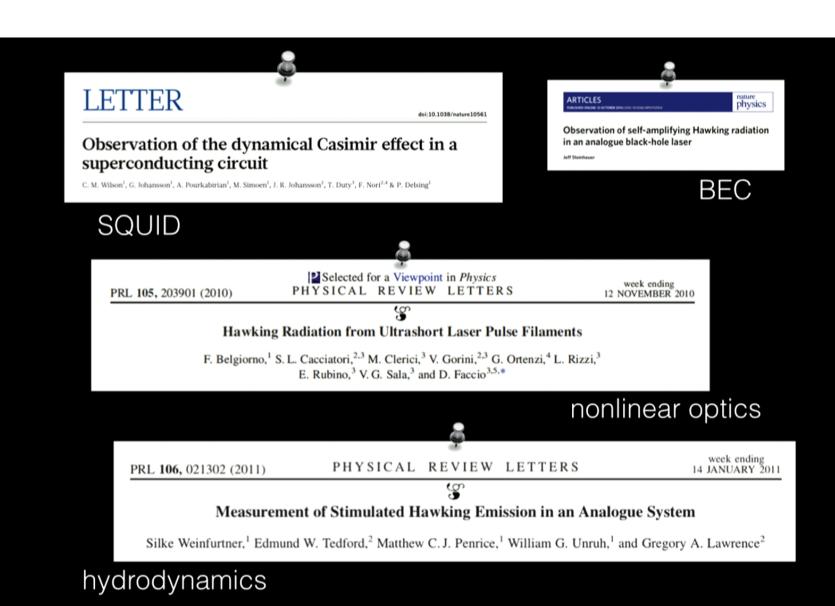
TWO-MODE SQUEEZED VACUUM STATE



squeezing parameter
$$|\psi_{AB}
angle \propto \sum_{n=0}^{\infty} (\tanh r)^n |n,n
angle$$
 partial trace $ho_A \propto \sum_{n=0}^{\infty} (\tanh r)^n |n
angle \langle n|$

Higher squeezing, higher temperature.

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- Black hole excites radiation from the vacuum.
- The outgoing part of the flux is thermal:

$$T_H = \frac{\hbar c^3}{8\pi G M k_B}$$

Mass loss rate:

$$\frac{dM}{dt} = -\frac{\hbar c^4}{15360\pi G^2 M^2}$$

Final fate?

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EVAPORATION AND REVERSIBILITY

Quantum mechanics contradicts all classical lessons:

- No hair: Hawking radiation is "hair"
- 2. Area law: area shrinks in time ("evaporation")
- 3. Singularity: quantum repulsion may resolve it



Does quantum mechanics restore reversibility?

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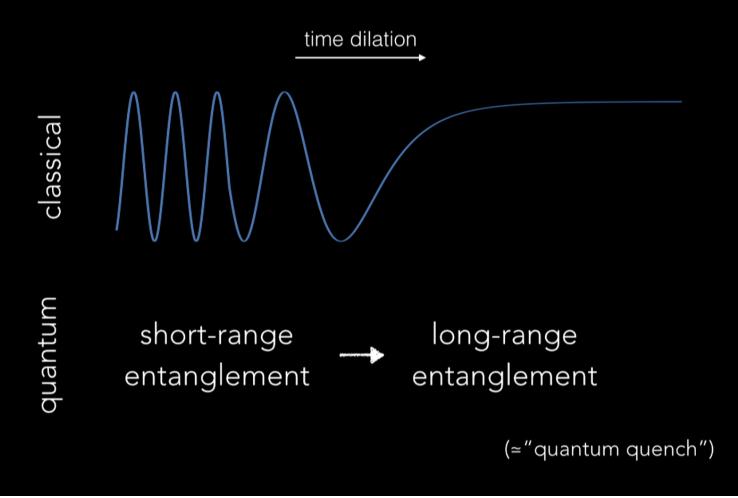
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Does quantum mechanics restore reversibility?

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THE HAWKING EFFECT - ENTANGLEMENT



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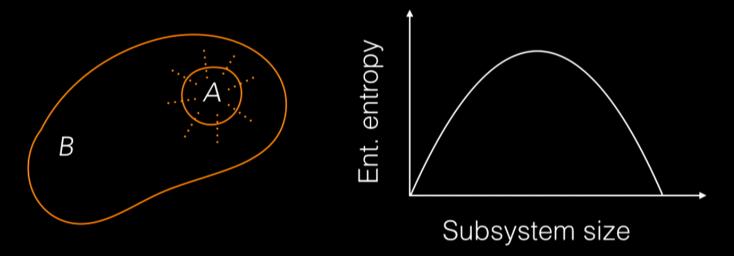
ENTANGLEMENT ENTROPY

In quantum mechanics, entanglement entropy is

$$S[\rho_A] \equiv -\operatorname{tr}_A[\rho_A \ln \rho_A]$$
 with $\rho_A \equiv \operatorname{tr}_B[\rho_{AB}]$

with triangle inequality

$$|S[\rho_A] - S[\rho_B]| \le S[\rho_{AB}] \le S[\rho_A] + S[\rho_B]$$



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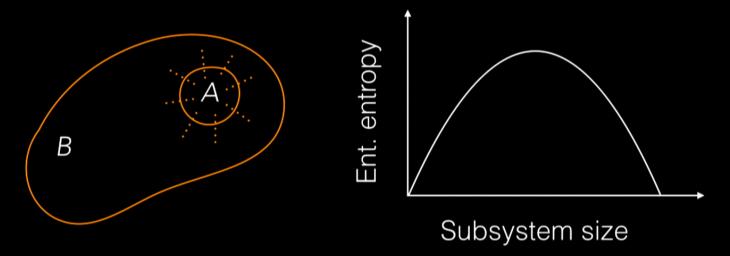
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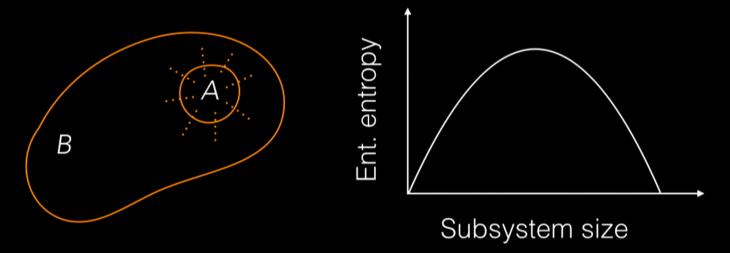
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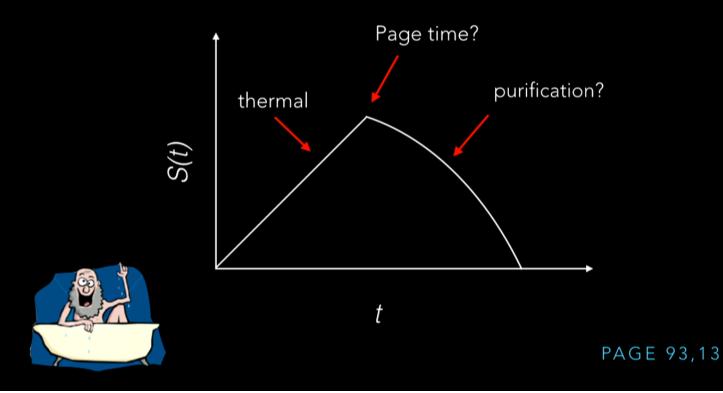
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THE PAGE CURVE

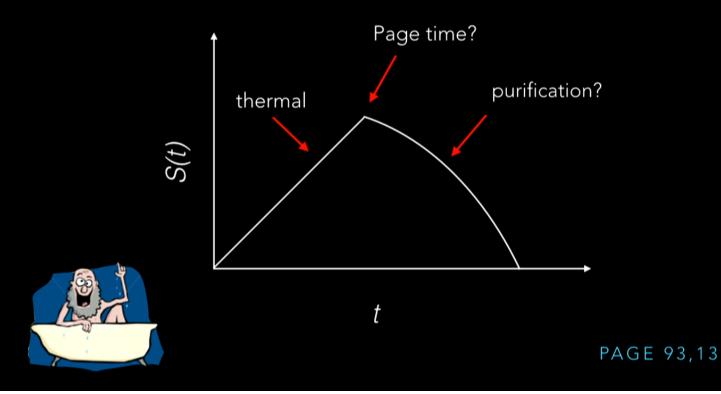
Outstanding open conjecture in semiclassical gravity:



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THE PAGE CURVE

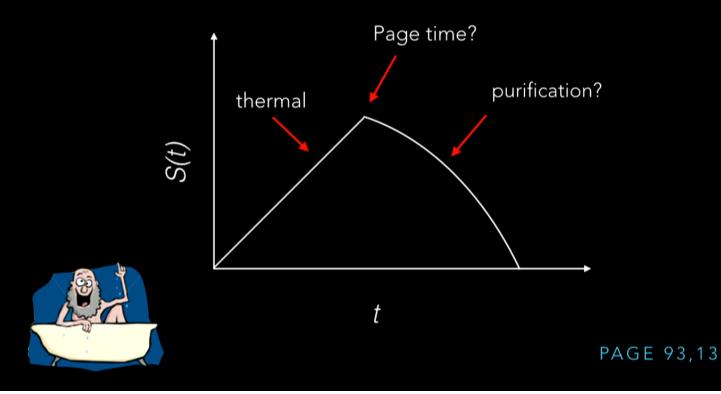
Outstanding open conjecture in semiclassical gravity:



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THE PAGE CURVE

Outstanding open conjecture in semiclassical gravity:



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RADIATION ENTROPY - THE 2D APPROXIMATION

Approximations of the field dynamics

- 1. neglect multipole moments l > 0
- 2. neglect "backscattering"

yields the geometric formula:

renormalized entropy

gravitational redshift

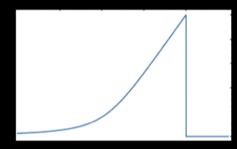
$$S(t) = \frac{1}{12} \ln \frac{\omega(t)}{\omega_0}$$

HOLZHEY ET AL. 94; BIANCHI, MS 14

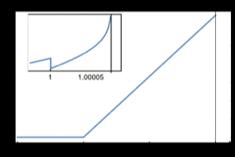
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FROM GEOMETRY TO ENTROPY

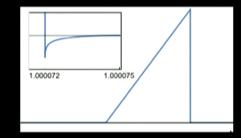
Direct probe of geometric vacuum deformations:



 ϵ -star



monotonic evaporation



nonsingular hole



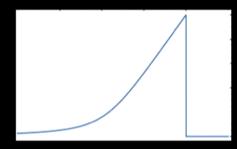
First computations of the Page curve.

BIANCHI, DE LORENZO, MS 14

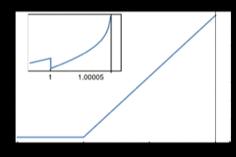
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FROM GEOMETRY TO ENTROPY

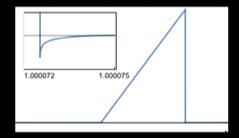
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BIANCHI, DE LORENZO, MS 14

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REVERSIBILITY: THREE QUESTIONS

Is the evaporation process:

1. unitary, i.e. is purity of the vacuum preserved?

"information loss problem"

2. cyclic, i.e. is entanglement restored to its original value?

$$S(+\infty) = S(-\infty)$$
?

(and conservative, i.e. do input and output energy match?)



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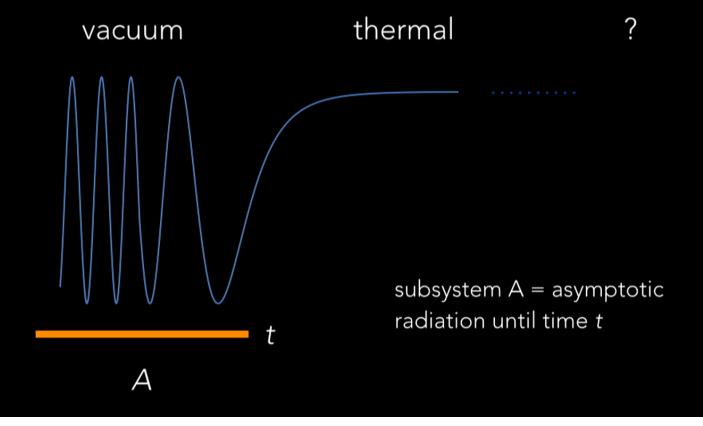
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More importantly: what difference would it make?

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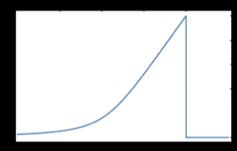
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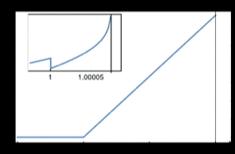
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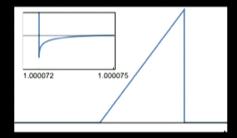
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First computations of the Page curve.

BIANCHI, DE LORENZO, MS 14

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IT FROM BIT - AND THE GENERALIZED SECOND LAW



spacetime quantum field

Entropy-flux identity:

$$2\pi F = 6\dot{S}^2 + \ddot{S}$$

BIANCHI, MS 14

Generalizes the "generalized second law":

- includes a non-adiabatic term (hence identity)
- does not require special causal structure (event horizon)
- implies the GSL in spherical symmetry



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IT FROM BIT - CONSEQUENCES OF CYCLICITY

$$2\pi F = 6\dot{S}^2 + \ddot{S}$$

LAST GASP

If evaporation is cyclic, then

- 1. the black hole must "gasp": $F(t_{
 m Page}) < 0$
- 3. purification must be slow:



$$au_P \ge \xi \; rac{(M_0^2 - M_1^2)^2}{M_1 m_P^2}$$

initial mass

mass at the end of thermal phase

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CONCLUSION: A TALE OF TWO THEORIES

Classical relativity suggests black holes break reversibility.

Quantum effects may resolve this, by enforcing

unitarity?

cyclicity?

2d field theory helps us imagine what this might mean physically:

- black hole gasping
- slow purification



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FOR MORE DETAILS

- International Loop Quantum Gravity Seminar (Mar. 15)
- References:
 - Bianchi, MS, PRD 90 (4), 041904(R) (2014)
 - Bianchi, de Lorenzo, MS, arXiv: 1409.0144
 - Bianchi, MS, GRG 46, 1809 (2014)

Thank you!

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