

Title: How does extended supersymmetry affect the elliptic genus?

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Abstract: The elliptic genus of K3 and its decomposition into characters of the $N=4$ superconformal algebra of associated conformal field theories can be viewed as the outset of Mathieu Moonshine. Thus, extended supersymmetry induces additional properties of the elliptic genus, which so far lack a satisfactory geometric interpretation. We investigate the implications of this decomposition on geometric structures that underlie the elliptic genus.

How does extended supersymmetry affect the elliptic genus?

(MOCK) MODULARITY, MOONSHINE AND STRING THEORY
Perimeter Institute, Waterloo, Canada, April 13-17, 2015

Katrin Wendland
Albert-Ludwigs-Universität Freiburg

The elliptic genus under extended supersymmetry

Introduction

- 1 The elliptic genus of Calabi-Yau manifolds
- 2 Effects of extending the $N = (2, 2)$ superconformal algebra
- 3 Interpretation and symmetry-surfing

[Creutzig/W15], [Carnahan/Grimm/W15] *work in progress*

[W15] *K3 en route from geometry to conformal field theory*;
to appear in: Proceedings of the 2013 Summer School "Geometric, Algebraic and Topological Methods for Quantum Field Theory", Villa de Leyva, Colombia;
arXiv:1503.08426 [math.DG]

[W14] *Snapshots of conformal field theory*;
in "Mathematical Aspects of Quantum Field Theories", Mathematical Physics Studies, Springer 2015, pp. 89-129; arXiv:1404.3108 [hep-th]

[Taormina/W13] *A twist in the M_{24} moonshine story*, to appear in Confluentes Mathematici;
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[Taormina/W12] *Symmetry-surfing the moduli space of Kummer $K3$ s*, to appear in: Proceedings of String-Math 2012; arXiv:1303.2931 [hep-th]

[Taormina/W11] *The overarching finite symmetry group of Kummer surfaces in the Mathieu group M_{24}* , JHEP **1308**:152 (2013); arXiv:1107.3834 [hep-th]

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1. The elliptic genus of Calabi-Yau manifolds

Let M denote a compact Calabi-Yau D -fold, $T := T^{1,0}M$,

$$\mathbb{E}_{q,-y} := y^{-\frac{D}{2}} \Lambda_{-y} T^* \otimes \bigotimes_{n=1}^{\infty} [\Lambda_{-yq^n} T^* \otimes \Lambda_{-y^{-1}q^n} T \otimes S_{q^n} T^* \otimes S_{q^n} T],$$

where for any bundle $E \rightarrow M$, $\Lambda_x E := \bigoplus_{m=0}^{\infty} x^m \Lambda^m E$, $S_x E := \bigoplus_{m=0}^{\infty} x^m S^m E$

Definition

With $q := e^{2\pi i\tau}$, $y := e^{2\pi iz}$ for $\tau, z \in \mathbb{C}$, $\text{Im}(\tau) > 0$,

$$\mathcal{E}_M(\tau, z) := \chi(\mathbb{E}_{q,-y}) = \int_M \text{Td}(M) \text{ch}(\mathbb{E}_{q,-y})$$

is the **ELLIPTIC GENUS** of M .

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$$E_H = \lambda(\mathbb{E}_{q, y}) \quad c(\tau) = \prod_j (1 + x_j)$$

$$\downarrow$$

$$\int \prod_j x_j \frac{\mathcal{D}_1(\tau, z - x_j)}{\mathcal{D}_1(\tau, z)}$$

$$\mathcal{D}_1(\tau, z) = i q^{1/8} y^{-1/2} \prod_{k=1}^{\infty} (1 - q^k)(1 - y^{-1} q^{k-1})(1 - y q^k)$$

$$D_1(\tau, z) = i q^{1/8} y^{-1/2} \prod_{k=1}^{\infty} (1 - q^k) (1 - y^{-1} q^{k-1}) (1 - y q^k)$$

$q \rightarrow 0$
 "no-loop
 limit"

$$\chi(y^{-D/2} \Lambda_y T^+) = \text{Td}(M) \leftarrow_{k=1} \frac{\Lambda_{\cdot} T^+}{ST_{(k \neq 1)}} \frac{\Lambda T}{ST}$$

"graded Euler"

expect: $E_M(\tau, z) = \text{str}_{\text{HNS}} \left(y^{\tau_0 - \frac{\epsilon}{6}} \dots q^{L_0^{\text{top}}} \right)$

for SCFT = NLSM(M), $L_0^{\text{top}} = L_0 - \frac{1}{2} \tau_0$
 $c = 3D$

CAUTION

DO NOT TOUCH THE BOARD WHEN IT IS HOT
 IT IS HOTTER THAN YOU THINK
 PLEASE BE CAREFUL

expect: $E_M(\tau, z) \underset{\tau \rightarrow \infty}{\sim} \text{str HNS} \left(y^{\tau_0 - \frac{c}{6}} L_0^{\text{top}} \right)$

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ok: $M = \mathbb{C}^D / \mathcal{N} + \text{orbifolds}$

CAUTION

THE BOARD IS LOANED FROM RESEARCH GROUPS
 PLEASE REPORT TO THE OFFICE OF THE DEAN
 IF IT IS DAMAGED OR STOLEN
 3000 UNIVERSITY AVENUE
 ANN ARBOR MI 48106-1500

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Chiral de Rham complex

Definition [Malikov/Schechtman/Vaintrob99]

For open $U \subset M$ with local holomorphic coordinates z^1, \dots, z^D :

$\Omega_M^{ch}(U) :=$ Fock space for the fields $\phi^j, p_j, \psi^j, \rho_j, j \in \{1, \dots, D\}$,
(D copies of a $(bc - \beta\gamma)$ -system)

where $\phi^j \leftrightarrow z^j, p_j \leftrightarrow \frac{\partial}{\partial z^j}, \psi^j \leftrightarrow dz^j, \rho_j \leftrightarrow \frac{\partial}{\partial(dz^j)}$.

This yields a sheaf of vertex algebras over M ,

the CHIRAL DE RHAM COMPLEX Ω_M^{ch} .

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Theorem [Malikov/Schechtman/Vaintrob99; Borisov/Libgober00]

$H^*(M, \Omega_M^{ch})$ carries the structure of a topological $N = 2$ superconformal vertex algebra, induced from globally well-defined fields on M ,

$$L^{top} = - : p_j \partial \phi^j : - : \rho_j \partial \psi^j :, \quad J = : \rho_j \psi^j :, \quad Q = - : \psi^j p_j :, \quad G = : \rho_j \partial \phi^j :.$$

The elliptic genus $\mathcal{E}_M(\tau, z)$ is the bigraded Euler characteristic of Ω_M^{ch} .

$$D_1(\tau, z) = i q^{1/8} y^{-1/2} \prod_{k=1}^{\infty} (1 - q^k)(1 - y^{-1} q^{k-1})(1 - y q^k)$$

$q \rightarrow 0$
 "no-loop
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$$\frac{\Lambda_{T^+}}{\Lambda_T} \rightsquigarrow \rho_{j|k}$$

$$T d \ln \left| \prod_{k=1}^{\infty} S T^+_{(k \neq 1)} \right|$$

\downarrow
 $\varphi_{j|k-1}$

$$S T \rightsquigarrow \rho_{j|k}$$

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for SCFT = NLSM(M), $L_0^{\text{top}} = L_0 - \frac{1}{2}c$
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ok: $M = \mathbb{C}^D / \mathcal{N}$ + orbifolds

$M \rightsquigarrow H^+(\mathcal{M}, \Omega^{\text{ch}}) \rightsquigarrow \mathcal{E}_\mathcal{M}(\tau, z)$
VOA

model: $\mathbb{E}_{g_1 - y}$

CAUTION

DO NOT TOUCH THE BOARD
IF IT IS NECESSARY TO Wipe
PLEASE CONTACT THE STAFF
ATTENDING BOARD



calculate VOA
 H^+ (M , Ω^{cm})

ok.
 C^0/A ,

kurve
 K_3

CAUTION

DO NOT TOUCH THE CONTROL PANEL
OR THE POWER SUPPLY UNIT
OR THE POWER SUPPLY UNIT
OR THE POWER SUPPLY UNIT
OR THE POWER SUPPLY UNIT



calculate VOA
 $H^+ (17, \Omega^{(4)})$

- dependence on
cpx str.?

relation to NLSM
Γ Nitten '92 top half twist
+ Kapustin 06 (large vol. limit)

ok.
CP/A, ✓

Kummer
K3, ✓

[Carrollian]

Grimm/W

work

in progress

CAUTION

DO NOT TOUCH THE BOARD WHEN
POWER IS ON. ALWAYS USE THE
FRONT PANEL TO OPEN
AND CLOSE THE BOARD.
OTHER PERSONS MUST

Chiral de Rham complex

Definition [Malikov/Schechtman/Vaintrob99]

For open $U \subset M$ with local holomorphic coordinates z^1, \dots, z^D :

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The elliptic genus $\mathcal{E}_M(\tau, z)$ is the bigraded Euler characteristic of Ω_M^{ch} .

if $\mathcal{E}_H(\tau, z) = \text{str}_{\text{HNS}} (\dots)$
with extended SUSY : decompose!

CAUTION

DO NOT TOUCH THE SURFACE OF THE BOARD.
IT IS NECESSARY TO OPEN
THE BOARD PROPERLY.

2. Extending to $N = (4, 4)$ supersymmetry

Assume: $D = 2$, i.e. $c = 6$ and $N = (4, 4)$ supersymmetry.

3 types of $N = 4$ irreps \mathcal{H}_\bullet with $\chi_\bullet(\tau, z) = \text{sTr}_{\mathcal{H}_\bullet} (y^{J_0} q^{L_0 - 1/4})$:
VACUUM \mathcal{H}_0 , MASSLESS MATTER $\mathcal{H}_{m.m.}$, MASSIVE MATTER $\mathcal{H}_{h>0}$.

$$\mathcal{E}_{K3}(\tau, z) = -2\chi_0(\tau, z) + 20\chi_{m.m.}(\tau, z) + \sum_{n=1}^{\infty} A_n \chi_n(\tau, z)$$

Proof (with T. Creutzig)

┌ Creutzig (Höhn '13)

Proof (with T. Creutzig) Γ Creutzig (Höhn '13)

T : $SU(2)$ principal bundle

$$\begin{array}{c} \mathbb{E}_{q_1-y} \\ \xrightarrow[\text{dist.}]{\text{holonomy}} \end{array} \mathbb{E}_{q_1-y} = \bigoplus_{n=1}^{\infty} W_n \otimes M_n \leftarrow \text{multipl.}$$

\uparrow
 n -dim
indep. of $SU(2)$
 $S^{n-1}(T)$

Proof (with T. Creutzig)

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T: $SU(2)$ principal bundle

\mathbb{F}_{q_1-y}
holonomy
dist.

— " —

$$\mathbb{F}_{q_1-y} = \bigoplus_{n=1}^{\infty} W_n \otimes M_n \leftarrow \text{multipl.}$$

\uparrow
n-dim
irrep. of $SU(2)$

$$= \bigoplus_{n=1}^{\infty} S^{n-1}(T) \otimes_n(\mathbb{C}^2)$$

\uparrow
 $\mathbb{Z}(N=4 \text{ char.})$

\uparrow
 $N=4$
action
(since $SU(2)$ -
inv.)

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Conjecture [W13]

Let $M=K3$. There are polynomials p_n for every $n \in \mathbb{N}$, such that

$$\mathbb{E}_{q,-y} = -\mathcal{O}_{K3} \chi_0(\tau, z) - T \chi_{m.m.}(\tau, z) + \sum_{n=1}^{\infty} p_n(T) \chi_n(\tau, z),$$

where $A_n = \int_{K3} \text{Td}(K3) p_n(T) = \chi(p_n(T))$ for all $n \in \mathbb{N}$.

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if $E_n(\tau, z) = \text{str}_{\text{HNS}}$ (...)
 with extended SUSY : decompose!

$$\chi(\mathbb{P}^n(\mathbb{T})) \rightarrow \dim \mathbb{P}^n = A_n \longleftrightarrow \text{Rep}(\mathbb{T}^2)$$

$$\downarrow$$

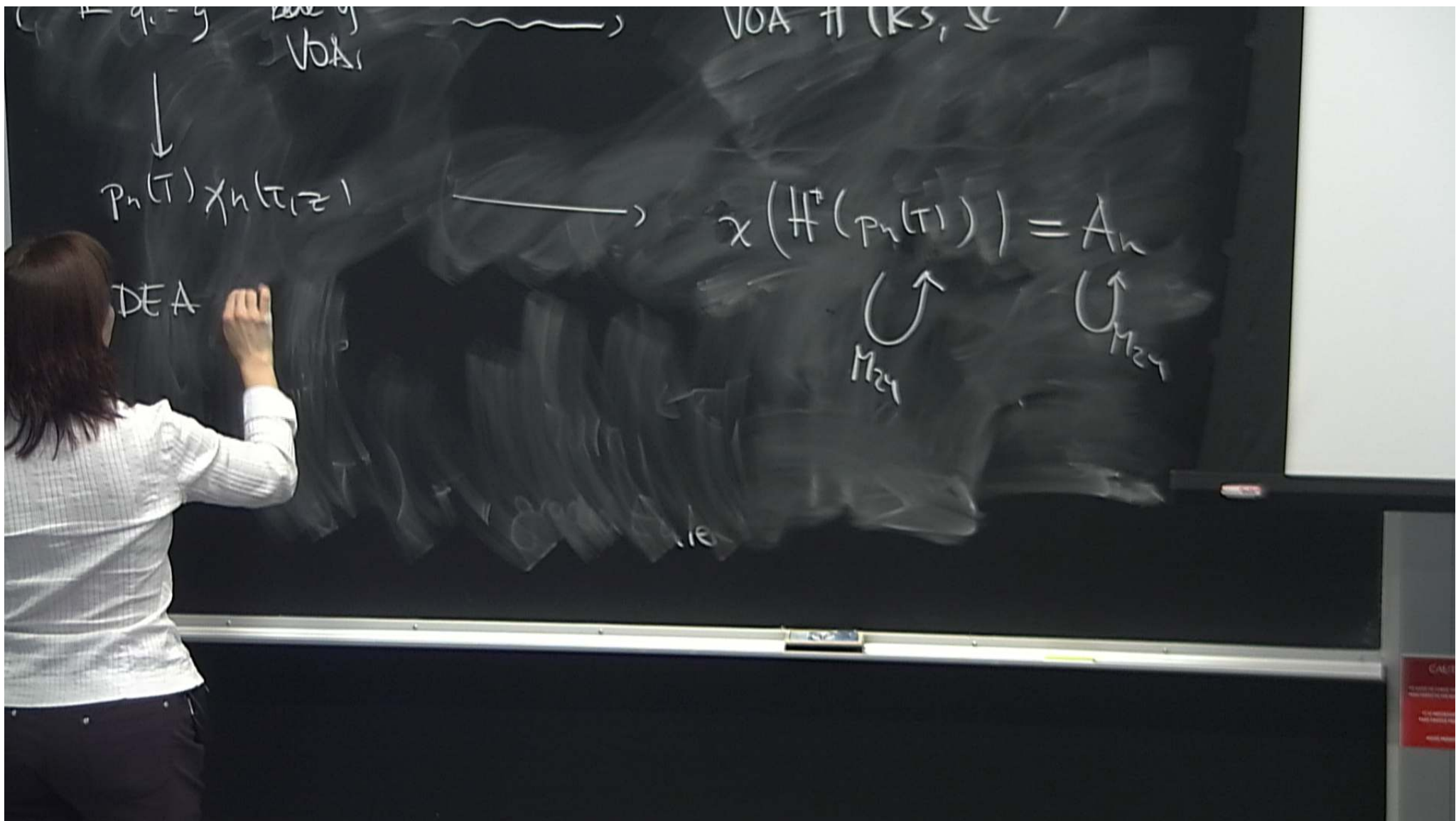
$$H^+(\mathbb{P}^n(\mathbb{T}))$$

Γ Eguchi / Ooguri / Tachikawa '10
 the audience
 (annou'12)

! decompose $\#_{g, y}$

$$\mathbb{P}^n \mathbb{T} = S^2 \mathbb{T}$$





work a case \mathcal{M} of moduli space of $MSp(K3)$

e.g. for \mathbb{Z}_2 -orbifold CFTs:

choose geom. interpr. on $T_{11}/\mathbb{Z}_2, T_{11} = \mathbb{C}^2/\Lambda$

$$\Lambda/2\Lambda \cong \mathbb{F}_2^4$$

↑
choose!

work a coset $\tilde{\mathcal{M}}$ of moduli space of $MSFT(K3)$

e.g. for \mathbb{Z}_2 -orbifold CFTs:

choose geom. interpr. on $\tilde{T}_1/\mathbb{Z}_2, T_1 = \mathbb{C}^2/\Lambda$

$$\wedge^{1/2}\Lambda \cong \mathbb{F}_2^4$$

↑
choose!

(\rightsquigarrow common marking

$$H^{\tilde{z}}(\tilde{T}_1/\mathbb{Z}_2, \mathbb{C}) \cong \text{standard lattice}$$

work a coset $\tilde{\mathcal{M}}$ of moduli space of $MSM(K3)$

e.g. for \mathbb{Z}_2 -orbifold CFTs:

choose geom. interpr. on $T_{1/2\Lambda} \sim T_{\Lambda} = \mathbb{C}^2/\Lambda$

$$\uparrow_{1/2\Lambda} \cong \mathbb{F}_2^4$$

↑
choose!

(\rightsquigarrow common marking

$$H^{\pm}(\tilde{T}_{1/2\Lambda}, \mathbb{C}) \cong \text{standard lattice}$$

1. [TW11] : $M_1, M_2 \in \tilde{\mathcal{M}}$
 $H^*(\pi_{k_1}, \mathbb{Z}) \xrightarrow[\cong]{\theta} \boxed{N A_1^{24}}$
 $\uparrow \cong_{G_k}$ $\uparrow \cong_{G_k}$
 $\mathbb{Z}_2^4 \Delta A_7$

2. [TW12] \mathbb{F}_2^4
 $\uparrow \cong_{G_k}$
 $\mathbb{Z}_2^4 \Delta A_8 \in \pi_{13}$
 [TW13] \mathcal{R}_1 : space of states common to all \mathbb{Z}_2 -orbifold CFT on k_1

Symmetry-surfing: Symmetries of \mathbb{Z}_2 -orbifold CFTs on K3

G : **geometric** symmetry group of a \mathbb{Z}_2 -orbifold CFT
with **geometric interpretation** on $X = T_\Lambda/\mathbb{Z}_2$, $T_\Lambda = \mathbb{C}^2/\Lambda$

$$\implies G = (\mathbb{Z}_2)^4 \rtimes G_\Lambda,$$

using [Fujiki88] \implies $(\mathbb{Z}_2)^4$ acts by shifts on $\Lambda/2\Lambda$, $G_\Lambda \subset SO(3)$

- $G_\Lambda \subset G_{\Lambda_k}$, one of three **maximal finite groups**

$$G_{\Lambda_1} = A_4, \quad G_{\Lambda_2} = S_3, \quad G_{\Lambda_0} = \mathbb{Z}_2^2$$

- our choice of **lattice generators** $\Lambda = \text{span}_{\mathbb{C}}\{\lambda_1, \lambda_2, \lambda_3, \lambda_4\}$
for the corresponding **maximally symmetric** lattices Λ_k is

$$\Lambda_1 = \text{span}_{\mathbb{Z}}\{(1,0,0,0), (0,1,0,0), (0,0,1,0), \frac{1}{2}(1,1,1,1)\},$$

$$\Lambda_2 = \text{span}_{\mathbb{Z}}\{(1,0,0,0), \frac{1}{2}(-1, \sqrt{3}, 0, 0), (0,0,1,0), \frac{1}{2}(0,0, -1, \sqrt{3})\},$$

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 $\implies G = (\mathbb{Z}_2)^4 \rtimes G_\Lambda \subset (\mathbb{Z}_2)^4 \rtimes \text{GL}_4(\mathbb{F}_2) \stackrel{[\text{Jordan1870}]}{\cong} (\mathbb{Z}_2)^4 \rtimes A_8 \subset M_{24}$,
 $\mathbb{F}_2^4 \cong \Lambda/2\Lambda$, $G_\Lambda \subset \text{SO}(3)$
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- $G_\Lambda \subset G_{\Lambda_k}$, one of three **maximal finite groups**
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- $G_\Lambda \subset G_{\Lambda_k}$, one of three **maximal finite groups**
 $G_{\Lambda_1} = A_4$, $G_{\Lambda_2} = S_3$, $G_{\Lambda_0} = \mathbb{Z}_2^2$
- our choice of **lattice generators** $\Lambda = \text{span}_{\mathbb{C}}\{\lambda_1, \lambda_2, \lambda_3, \lambda_4\}$
 for the corresponding **maximally symmetric** lattices Λ_k is
 - $\Lambda_1 = \text{span}_{\mathbb{Z}}\{(1,0,0,0), (0,1,0,0), (0,0,1,0), \frac{1}{2}(1,1,1,1)\}$,
 - $\Lambda_2 = \text{span}_{\mathbb{Z}}\{(1,0,0,0), \frac{1}{2}(-1, \sqrt{3}, 0, 0), (0,0,1,0), \frac{1}{2}(0,0, -1, \sqrt{3})\}$,
 - $\Lambda_0 = \text{span}_{\mathbb{Z}}\{(1,0,0,0), (0,1,0,0), (0,0,1,0), (0,0,0,1)\}$.

Symmetry surfing the moduli space of Kummer K3s

Results [Taormina/W11&12&13]

- For the \mathbb{Z}_2 -orbifold CFTs on K3 with geometric interpretation on some $X = \widetilde{T_\Lambda/\mathbb{Z}_2}$ ($T_\Lambda = \mathbb{C}^2/\Lambda$), the **joint action** of **all symmetry groups** yields the **maximal** subgroup $(\mathbb{Z}_2)^4 \rtimes A_8 \subset M_{24}$.

Note: $(\mathbb{Z}_2)^4 \rtimes A_8$ is **not** a subgroup of M_{23} .

- There is a **90-dim.** space \mathcal{R}_1 of states **common** to all K3-theories that are \mathbb{Z}_2 -orbifolds of toroidal theories. As **common representation space** of **all geometric symmetry groups** of Kummer K3s, \mathcal{R}_1 carries an action of $(\mathbb{Z}_2)^4 \rtimes A_8$ induced from $\mathcal{R}_1 \cong \mathbf{45} \oplus \overline{\mathbf{45}}$ with **irreps** $\mathbf{45}$, $\overline{\mathbf{45}}$ of M_{24} .

Note: There is a **twist** in this action of the groups $G_{\Lambda_1}, G_{\Lambda_2}, G_{\Lambda_0}$ on $\mathcal{R}_1 \cong \mathbf{15} \otimes (\mathbf{3} \oplus \overline{\mathbf{3}})$.

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Note: There is a twist in this action of the groups $G_{\Lambda_1}, G_{\Lambda_2}, G_{\Lambda_0}$ on $\mathcal{R}_1 \cong 15 \otimes (3 \oplus \overline{3})$.