

Title: The Cosmological Constant Problem (and its sequester)

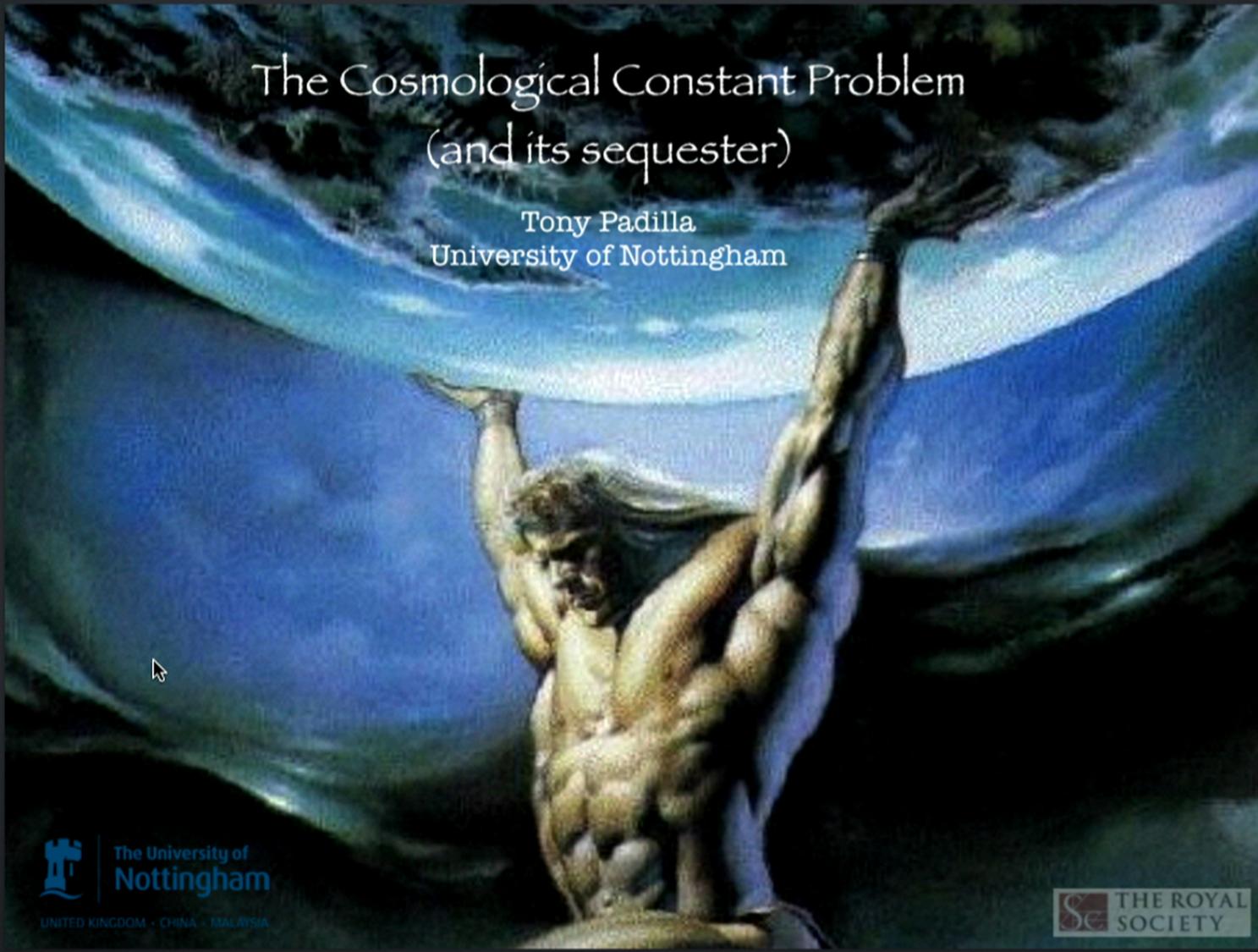
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Abstract: I will review the notorious cosmological constant problem, sometimes described as the worst fine tuning problem in Physics. I will explain the true nature of the problem, which is one of radiative instability against any change in the effective description. I will recall Weinberg's venerable no-go theorem that prohibits certain attempts to "solve" this problem before going on to explain a new mechanism that circumvents Weinberg. This is the vacuum energy sequester, a global modification of GR that results in the cancellation of large vacuum energy contributions from a protected matter sector (taken to include the Standard Model) at each and every order in the perturbative loop expansion. Cosmological consequences are a Universe which has finite space-time volume, will ultimately crunch, and for which dark energy can only be a transient. Furthermore, using a linear scalar potential within the sequestering set-up, I will show that dark energy today can be intimately related to the trigger that brings about cosmological collapse in the not too distant future, at the same time providing a possible solution to the "Why Now?" problem.

The Cosmological Constant Problem (and its sequester)

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with Nemanja Kaloper

1309.6562

1406.0711

1409.7073

see also

1502.05296

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The Cosmological Constant Problem

General Covariance & Equivalence Principle \Rightarrow Vacuum Energy Gravitates

$$-V_{vac} \int \sqrt{-g} d^4x \implies T_{\mu\nu} = -V_{vac} g_{\mu\nu}$$

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The Cosmological Constant Problem

General Covariance & Equivalence Principle \Rightarrow Vacuum Energy Gravitates

$$-V_{vac} \int \sqrt{-g} d^4x \implies T_{\mu\nu} = -V_{vac} g_{\mu\nu}$$

...add a bare cosmological constant...

$$\frac{1}{4} (V_{vac} + \Lambda_{bare}) \int \sqrt{-g} d^4x \implies T_{\mu\nu} = -\Lambda_{tot} g_{\mu\nu}$$

$$\text{where } \Lambda_{tot} = V_{vac} + \Lambda_{bare} \lesssim (meV)^4$$

Estimating the vacuum energy

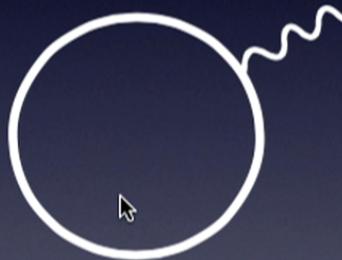
$$V_{vac} \supset \sum_m \int d^3k \frac{1}{2} \hbar \sqrt{k^2 + m^2}$$
$$\sim c_\nu m_\nu^4 + c_e m_e^4 + c_\mu m_\mu^4 + \dots + M_{\text{cut-off}}^4$$

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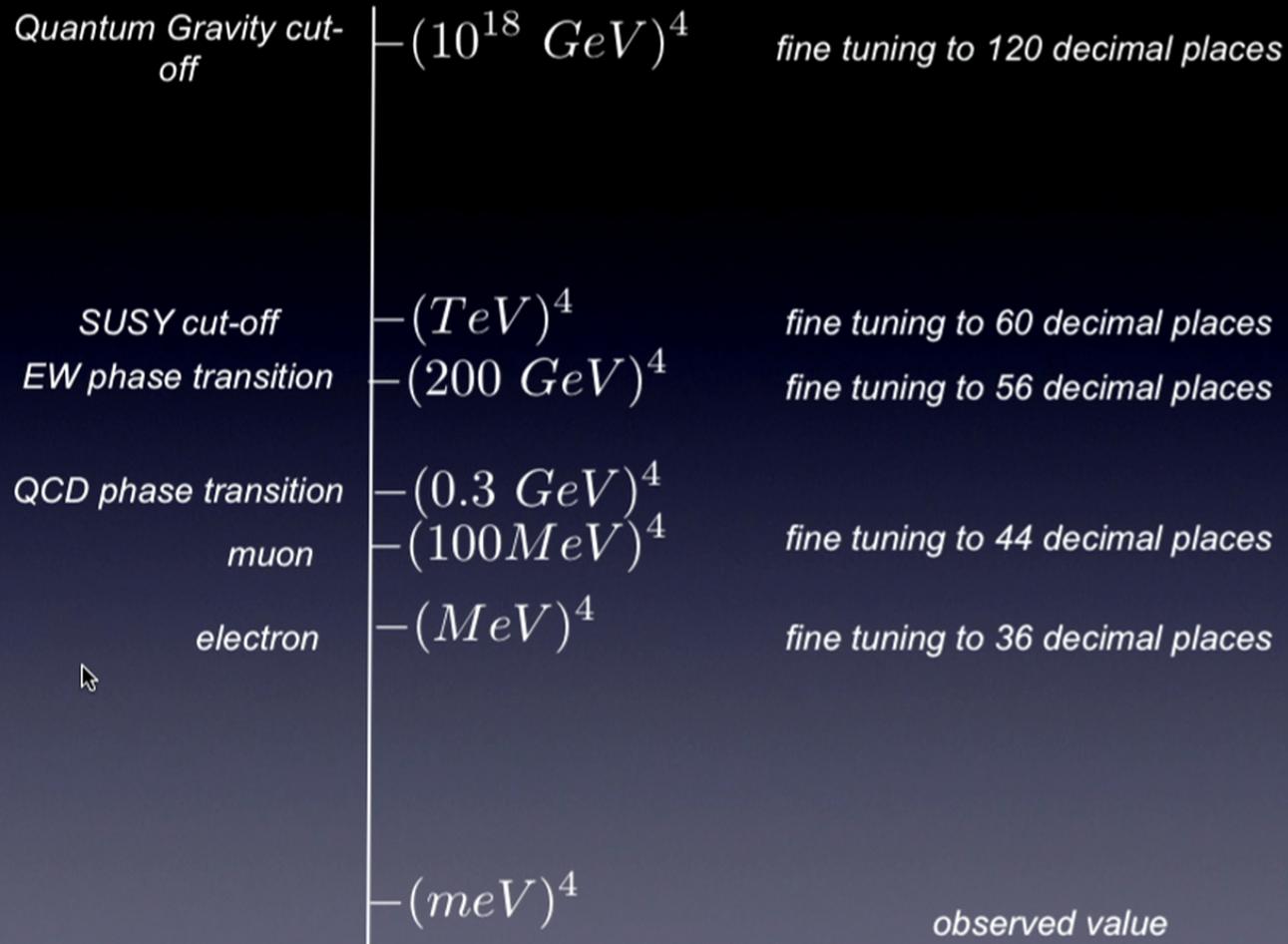
$$\sim \frac{1}{2} \int \frac{d^4k}{(2\pi)^4} \left[\frac{k^\mu k^\nu - \frac{1}{2} \eta^{\mu\nu} (k^2 + m^2)}{k^2 + m^2} \right] = -\frac{i}{2} \eta^{\mu\nu} V_{vac}^{\phi, 1\text{-loop}}$$

$$V_{vac}^{\phi, 1\text{-loop}} = -\frac{m^4}{(8\pi)^2} \left[\frac{2}{\epsilon} + \text{finite} + \ln \left(\frac{M_{UV}^2}{m^2} \right) \right]$$

What the cosmological constant problem is not....

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Why is the cosmological constant so small?



“Fine tuning” and renormalization

$$V_{vac}^{\phi, 1\text{-loop}} = -\frac{m^4}{(8\pi)^2} \left[\frac{2}{\epsilon} + \text{finite} + \ln \left(\frac{M_{UV}^2}{m^2} \right) \right] \quad \text{Divergent vacuum energy}$$

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“Fine tuning” and renormalization

$$V_{vac}^{\phi, 1\text{-loop}} = -\frac{m^4}{(8\pi)^2} \left[\frac{2}{\epsilon} + \text{finite} + \ln \left(\frac{M_{UV}^2}{m^2} \right) \right] \quad \text{Divergent vacuum energy}$$

$$\Lambda_{bare} = \frac{m^4}{(8\pi)^2} \left[\frac{2}{\epsilon} + \ln \left(\frac{M_{UV}^2}{\mathcal{M}^2} \right) \right] \quad \text{Bare counterterm}$$

Renormalized vacuum energy:

$$\Lambda_{ren} = V_{vac}^{\phi, 1\text{-loop}} + \Lambda_{bare} = \frac{m^4}{(8\pi)^2} \left[\ln \left(\frac{m^2}{\mathcal{M}^2} \right) - \text{finite} \right]$$

Depends on arbitrary subtraction scale so cannot be predicted...must be measured!



goglobal

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Radiative instability of Λ

At one loop, $V_{vac} = V_{vac}^{tree} + V_{vac}^{1loop} \sim M_{UV}^4 \gtrsim (TeV)^4$

tune finite part of Λ_{bare} to great precision

At two loops, $V_{vac} = V_{vac}^{tree} + V_{vac}^{1loop} + V_{vac}^{2loop}$, $V_{vac}^{2loop} \sim V_{vac}^{1loop}$

REtune finite part of Λ_{bare} to same precision

At three loops,

MISSION: IMPOSSIBLE™ MISSION: IMPOSSIBLE

“to make the cosmological constant radiatively stable”

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..but why they hardly gravitate.

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The Sequester

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$$S = \int d^4x \sqrt{-g} \left[\frac{M_{pl}^2}{2} R - \Lambda - \mathcal{L}(g^{\mu\nu}, \Psi) \right]$$

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Introduce global dynamical variables Λ

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Introduce global dynamical variables Λ, λ

$$S = \int d^4x \sqrt{-g} \left[\frac{M_{pl}^2}{2} R - \Lambda - \lambda^4 \mathcal{L}(\lambda^{-2} g^{\mu\nu}, \Psi) \right]$$

λ sets the hierarchy between matter scales and M_{pl}

λ

$$\frac{m_{phys}}{M_{pl}} = \frac{\lambda m}{M_{pl}}$$

Introduce global dynamical variables Λ, λ

$$S = \int d^4x \sqrt{-g} \left[\frac{M_{pl}^2}{2} R - \Lambda - \lambda^4 \mathcal{L}(\lambda^{-2} g^{\mu\nu}, \Psi) \right] + \sigma \left(\frac{\Lambda}{\lambda^4 \mu^4} \right)$$

λ sets the hierarchy between matter scales and M_{pl}

I

$$\frac{m_{phys}}{M_{pl}} = \frac{\lambda m}{M_{pl}}$$

Equations of motion

$$\Lambda \text{ equation} : \frac{\sigma'}{\lambda^4 \mu^4} = \int d^4x \sqrt{g}$$

$$\lambda \text{ equation} : 4\Lambda \frac{\sigma'}{\lambda^4 \mu^4} = \int d^4x \sqrt{g} \lambda^4 \tilde{T}^\mu{}_\mu$$

$$g_{\mu\nu} \text{ equation} : M_{pl}^2 G_\nu^\mu = -\Lambda \delta_\nu^\mu + \lambda^4 \tilde{T}_\nu^\mu$$

$$T_{\mu\nu} = \lambda^4 \tilde{T}_{\mu\nu}$$

$$\tilde{T}_{\mu\nu} = \frac{2}{\sqrt{-g}} \frac{\delta}{\delta \tilde{g}^{\mu\nu}} \int d^4x \sqrt{-\tilde{g}} \mathcal{L}(\tilde{g}^{\mu\nu}, \Psi)$$

Equations of motion

$$\begin{aligned} \Lambda \text{ equation} & : \quad \Lambda = \frac{1}{4} \langle T^\alpha{}_\alpha \rangle, & \langle Q \rangle &= \frac{\int d^4x Q \sqrt{g}}{\int d^4x \sqrt{g}} \\ \lambda \text{ equation} & : \\ g_{\mu\nu} \text{ equation} & : \quad M_{pl}^2 G^\mu{}_\nu = -\Lambda \delta^\mu{}_\nu + T^\mu{}_\nu \end{aligned}$$

I

Equations of motion

$$M_{pl}^2 G^\mu{}_\nu = T^\mu{}_\nu - \frac{1}{4} \delta^\mu{}_\nu \langle T^\alpha{}_\alpha \rangle$$

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Equations of motion

$$M_{pl}^2 G^\mu{}_\nu = T^\mu{}_\nu - \frac{1}{4} \delta^\mu{}_\nu \langle T^\alpha{}_\alpha \rangle$$

$$T^\mu{}_\nu = -V_{vac} \delta^\mu{}_\nu + \tau^\mu{}_\nu$$

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$$M_{pl}^4 G_{\nu}^{\mu} = -\frac{1}{4} \langle \tau_{\alpha}^{\alpha} \rangle \delta_{\nu}^{\mu} + \tau_{\nu}^{\mu}$$

Vacuum energy drops out at each and every loop order

No hidden equations — this is everything!

Residual cosmological constant $\Lambda_{eff} = \frac{1}{4} \langle \tau_{\alpha}^{\alpha} \rangle$

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Residual cosmological constant $\Lambda_{eff} = \frac{1}{4} \langle \tau_{\alpha}^{\alpha} \rangle$

Λ_{eff} has nothing to do with vacuum energy

It is radiatively stable

OK to fix it empirically

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How big is Λ_{eff} ?

For standard matter, space-time integrals dominated by time when universe is largest

$$\int d^4x \sqrt{-g} \sim \frac{1}{H_{\text{age}}^4} \quad \text{where lifetime } t_{\text{age}} \sim \frac{1}{H_{\text{age}}} \gtrsim 13.7 \text{ Gyrs}$$

$$\langle \tau_\alpha^\alpha \rangle \sim \rho_{\text{age}} \sim \text{energy density at largest size} < \rho_c$$

$\Rightarrow \Lambda_{\text{eff}}$ is not dark energy ... too small!

Universe has finite spacetime volume

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$$\frac{\sigma'}{\lambda^4 \mu^4} = \int d^4 \sqrt{g}$$

space-time volume must be finite or else $\lambda \rightarrow 0$

$$\frac{m_{phys}}{M_{pl}} = \frac{\lambda m}{M_{pl}}$$

if $\lambda \rightarrow 0$ particle masses go to zero

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$w=-1$ is transient

$\Omega_k > 0$

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 $w = -1$ is transient
 $\Omega_k > 0$

circles in the sky?
possible correlation
between $1+w$ and Ω_k

COLLAPSE TRIGGER = DARK ENERGY

Linear potential $V=m^3\phi$

*form protected by shift symmetry,
size of m^3 technically natural*

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Radiatively stable choice of collapse time?

Radiatively stable choice of φ_{in} ?

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Yes, thanks to m^3

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But its not even a "choice".... $\langle R \rangle = 0$ picks out precisely those solutions with $\varphi_{in} > M_{pl}$!!!!!!!

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(Non) Issues

Fine tuning Λ_{eff}

Silly — loops are gone!

Acausality

*Some acausality essential for solving CCP
...need to scan all of spactime to measure true vacuum energy*

We fix one number, there are no local pathologies

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Summary

Found a way to cancel SM vacuum energy at any order in loops

Locally theory looks just like GR with a small Λ

Small Λ is radiatively stable & determined by a global boundary condition

Universe is doomed

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