

Title: Subluminal Vainshtein Screening in Massive Gravity

Date: Apr 09, 2015 03:15 PM

URL: <http://pirsa.org/15040102>

Abstract: I will discuss the Vainshtein mechanism in massive gravity. I will show that the spherically symmetric backgrounds that were believed to have superluminal sound speed are in fact unstable. Instead, there is a new class of phenomenologically relevant solutions with stable and subluminal perturbations.

$$m \rightarrow 0$$

$$M_{Pl} \rightarrow \infty$$

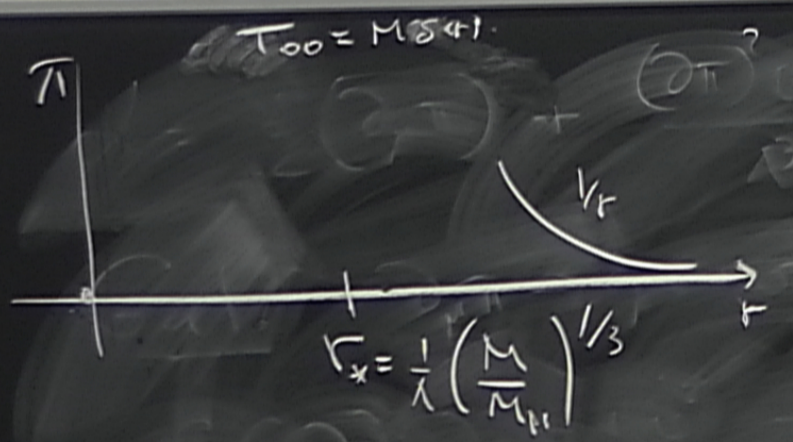
$$\lambda = (m^2 M_{Pl})^{1/3}$$

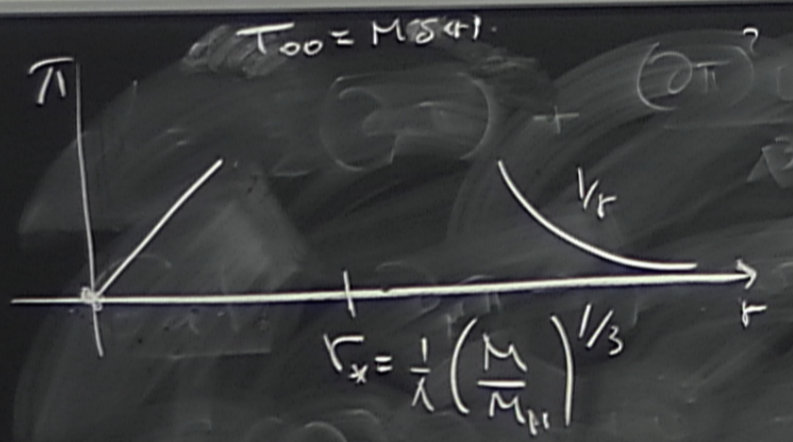
$$\mathcal{L} = h \xi h + h \left( \partial \partial \pi + \alpha \frac{(\partial \partial \pi)^2}{\lambda^3} + \beta \frac{(\partial \partial \pi)^3}{\lambda^6} \right) + \frac{h T}{M_{Pl}}$$

$$\beta = 0.$$

$$h \rightarrow h + \eta \pi + \partial_\mu \pi \partial_\mu \pi$$

$$\mathcal{L}_\pi = + \pi \square \pi + \alpha \frac{(\partial \pi)^2 \square \pi}{\lambda^3} + \alpha^2 \frac{(\partial \pi)^2 (\partial \partial \pi)^2}{\lambda^6} + \frac{\pi T}{M_{Pl}} + \alpha \frac{\partial_\mu \pi \partial_\nu \pi T_{\mu\nu}}{\lambda^3 M_{Pl}}$$





$$m \rightarrow 0$$

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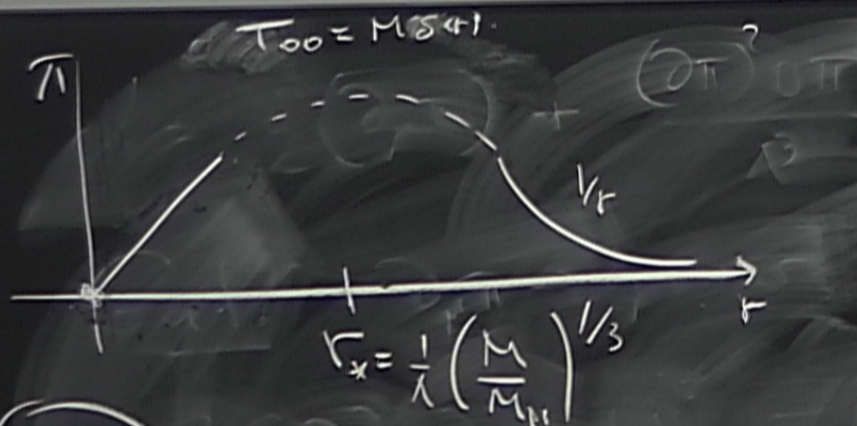
$$\lambda = (m^2 M_{Pl})^{1/3}$$

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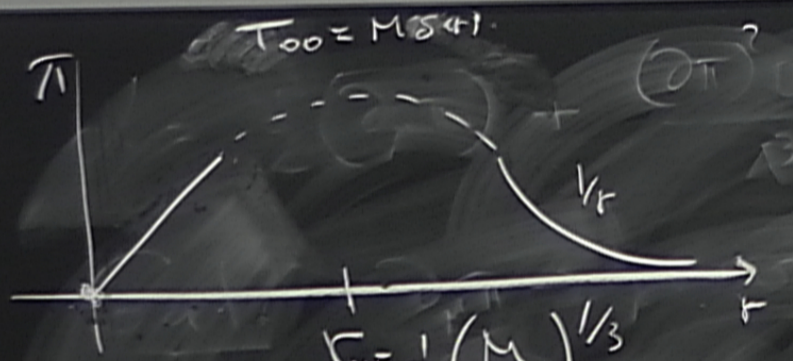
$$\beta = 0.$$

$$h \rightarrow h + \eta \pi + \partial_\mu \pi \partial_\mu \pi$$

$$\mathcal{L}_\pi = + \pi \square \pi + \alpha \frac{(\partial \pi)^2 \square \pi}{\lambda^3} + \alpha^2 \frac{(\partial \pi)^2 (\partial \partial \pi)^2}{\lambda^6} + \frac{\pi T}{M_{Pl}} + \alpha \frac{\partial_\mu \pi \partial_\nu \pi T_{\mu\nu}}{\lambda^3 M_{Pl}}$$

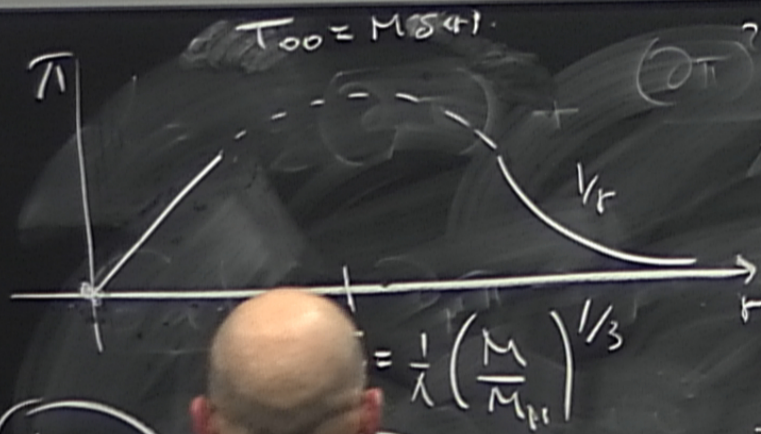


$\alpha < 0$



$$r_x = \frac{1}{\lambda} \left( \frac{M}{M_{pl}} \right)^{1/3}$$

$\alpha < 0$



$\alpha < 0$



$\pi = \bar{\pi}_1 + \delta\pi$

$\mathcal{L}_2(\delta\pi) = (\delta\pi)^2 \left( \frac{\partial^2 \bar{\pi}_1}{\lambda^3} \right) +$

$+ \frac{\alpha P}{\lambda^3 M_{Pl}} (\delta\pi)^2$

$\frac{\partial^2 \bar{\pi}_1}{\lambda^3} \sim$

$\frac{1}{r^2}$

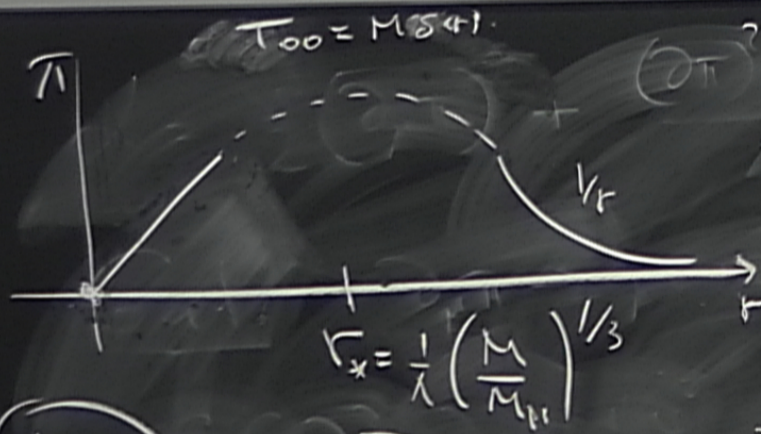
$P \sim$

$\frac{1}{r^3}$

$\frac{P}{r^3}$

$\left( \frac{r_+}{R} \right)^2$





$\alpha < 0$



$\pi = \bar{\pi}_1 + \delta\pi$

$\mathcal{L}_2(\delta\pi) = (\delta\pi)^2 \left( \frac{\partial^2 \bar{\pi}_1}{\lambda^3} \right) + \left( \frac{r_y}{R} \right)^2$

$\frac{\alpha \rho}{\lambda^3 M_{Pl}} (\delta\pi)^2$

$\frac{\partial^2 \pi}{\partial r^2}$

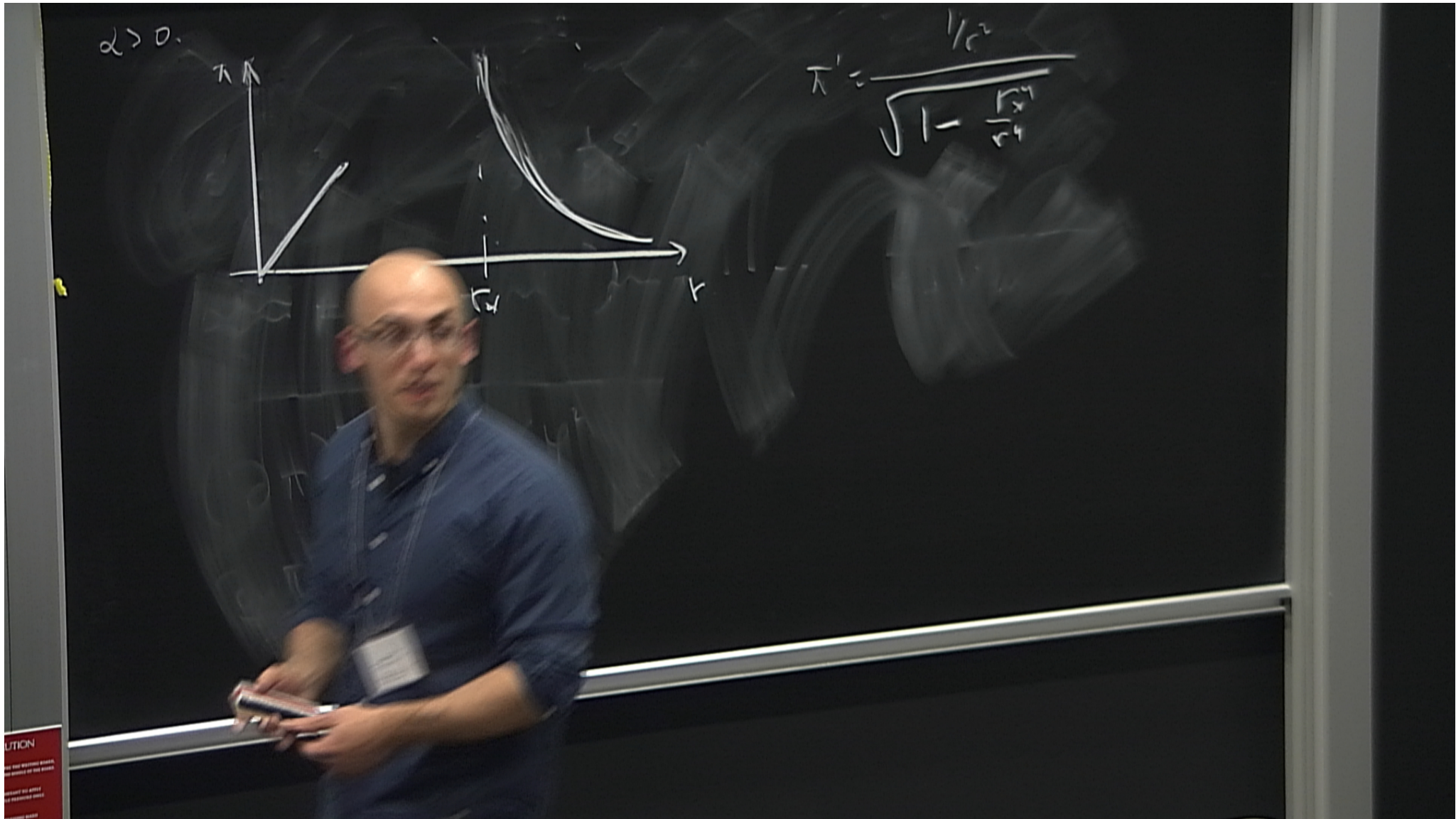
$\frac{1}{r^2}$

$\rho$

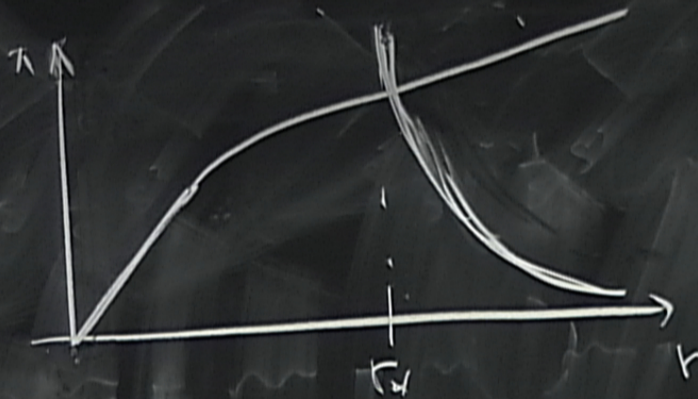
$\frac{1}{r^2}$

$\frac{r}{R}$





$\alpha > 0.$



$$\lim_{r \rightarrow \infty} \pi \sim \Lambda^3 r^2$$

$$W \approx -0.1.$$

$$\pi = \Lambda^3 \chi^2$$

$$\mathcal{E}h = \partial \pi + \partial$$

$$m \rightarrow 0$$

$$M_{pl} \rightarrow \infty$$

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$$h \rightarrow h + \eta \pi + \partial_\mu \pi \partial_\mu \pi$$

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$$-\frac{(p-2p)(\partial_r \delta\pi)^2}{q}$$

$$\lim_{r \rightarrow \infty} \pi \sim \Lambda^3 r^2$$

$$W \approx -0.1$$

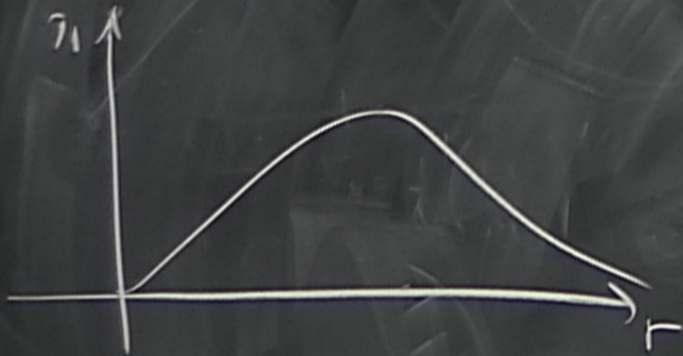
$$\pi = \Lambda^3 \chi^2$$

$$C_5 < 1$$

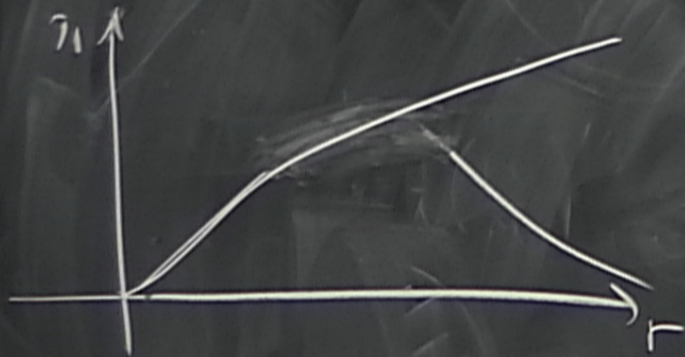
$$\delta\pi^2 \left( d^{2/3} \left( \frac{r_+}{r} \right)^2 + d^{1/3} \frac{r_+}{r} \right)$$

$$(\partial_r \delta\pi)^2 \left( d^{2/3} \left( \frac{r_+}{r} \right)^2 \right)$$

$$\Delta L = \pi T$$



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$$\lim_{r \rightarrow \infty} \pi = \lambda^2 r$$

$$\partial_r \pi = \lambda^2$$



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