

Title: Symmetries in large scale structure

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Abstract:

Symmetries in large scale structure

Lam Hui
Columbia University

Outline:

1. Equivalence principle: a generic test of modified gravity
- with Alberto Nicolis.
2. Parity: symmetry in the measurement of LSS
- with Camille Bonvin & Enrique Gaztanaga.
3. Dilation & beyond: symmetry in the theory of LSS
- with Kurt Hinterbichler & Justin Khoury,
Walter Goldberger & Alberto Nicolis,
Cremineilli, Gleyzes, Simonovic & Vernizzi,
Bart Horn, & Xiao Xiao.

Summary 1:

Test for presence of extra (scalar) forces by looking for off-centered black holes.

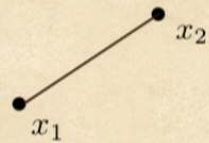
Footnote 1: No hair theorem for galileons (LH, Nicolis).

Footnote 2: The case of massive gravity (Gruzinov & Mirbabayi, Berezhiani, Chkareuli, de Rahm, Gabadadze, Tolley).

Footnote 3: Analogs for chameleon mechanism (Khoury, Weltman; Hu; Jain, Vanderplas; Pourhasan, Afshordi, Mann, Davis; Cabre, Vikram, Zhao, Jain, Koyama; LH, Nicolis, Stubbs).

Idea 2: parity in the measurement of LSS

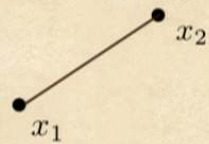
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$$\langle \delta(x_1) \delta(x_2) \rangle$$

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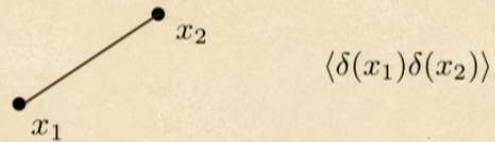
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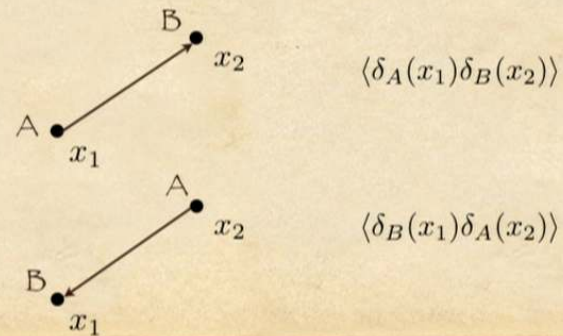
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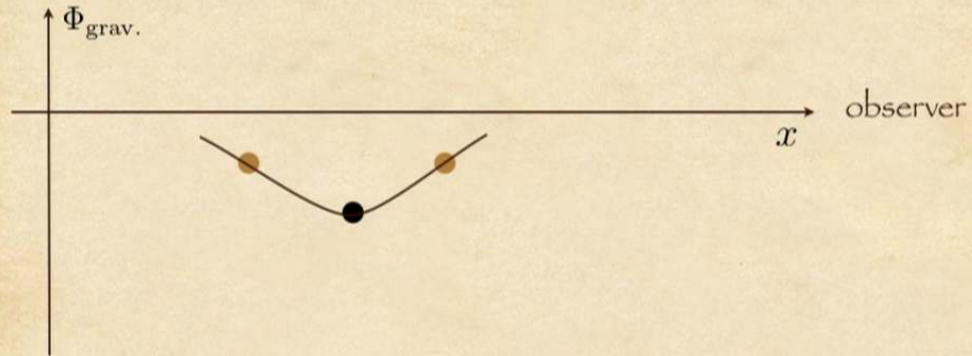
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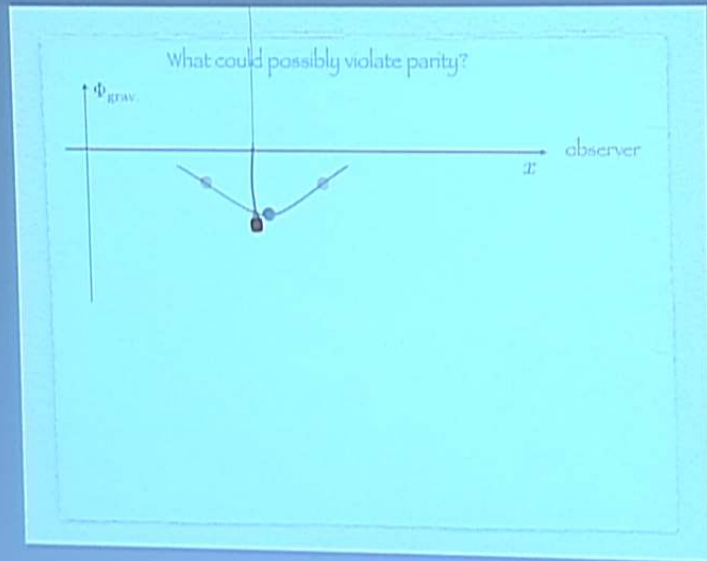


- But how about cross-correlation between 2 different kinds of galaxies, A & B?

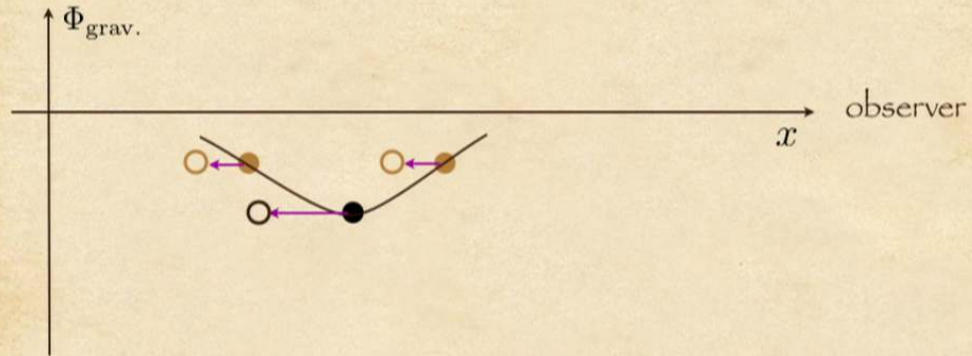


What could possibly violate parity?

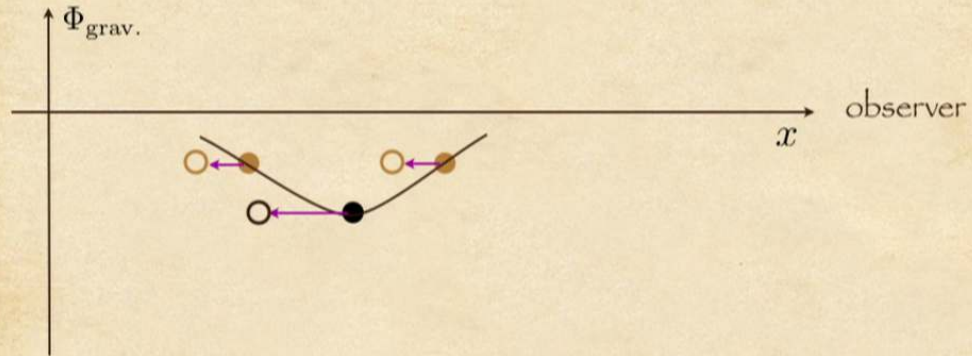




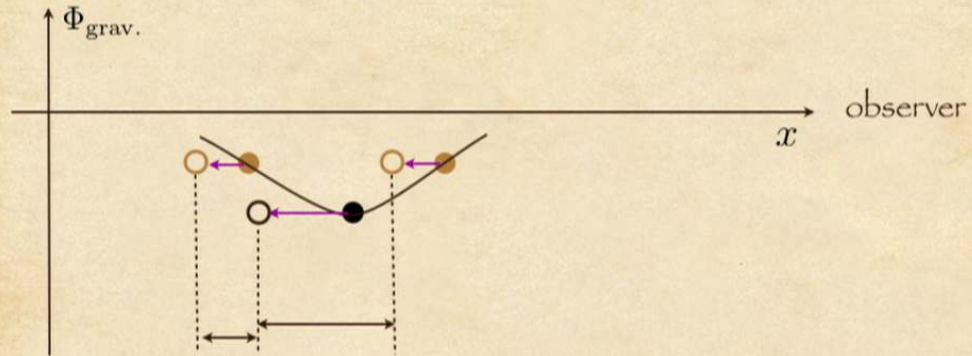
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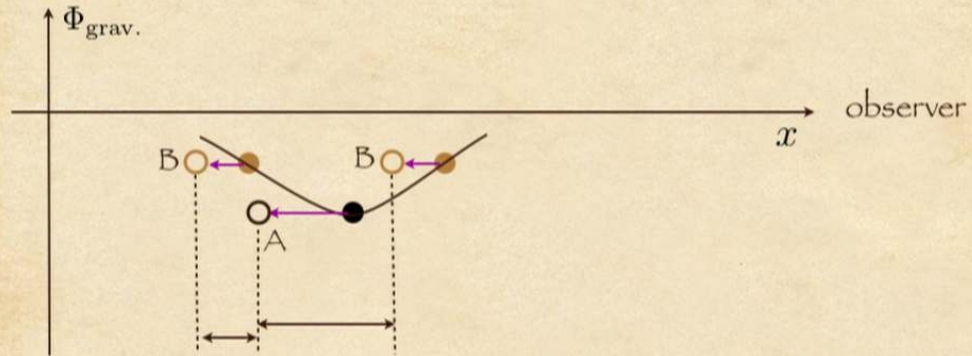
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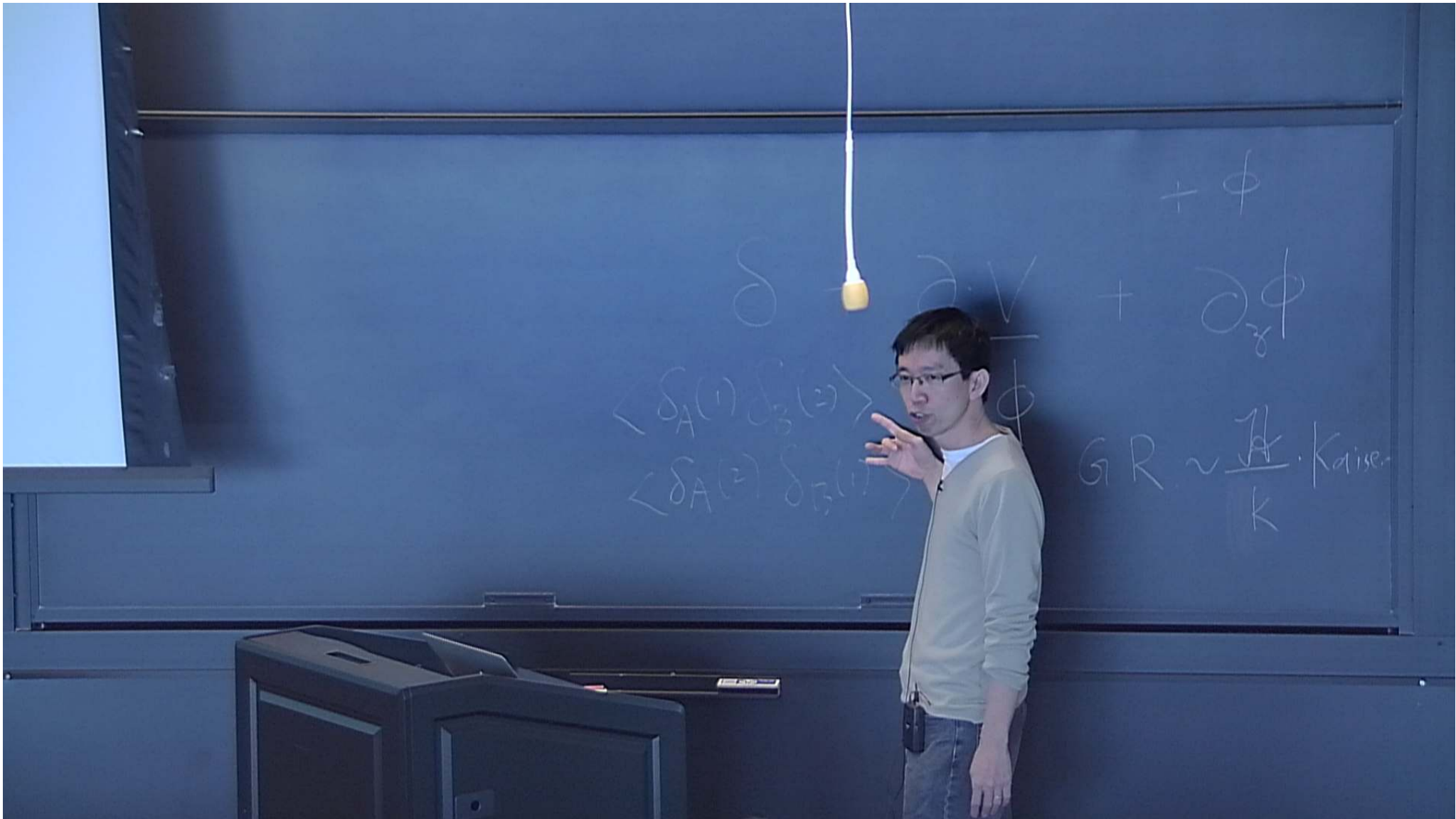


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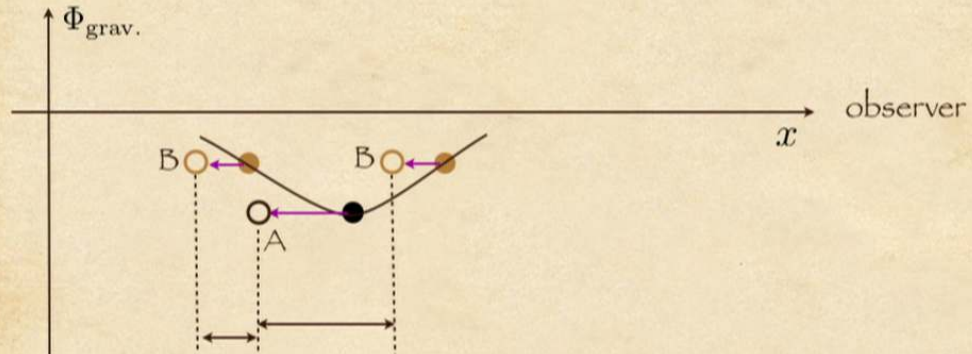




$$\delta \left(\frac{\partial V}{\partial \phi} + \nabla^2 \phi \right) \sim \frac{1}{H^2}$$



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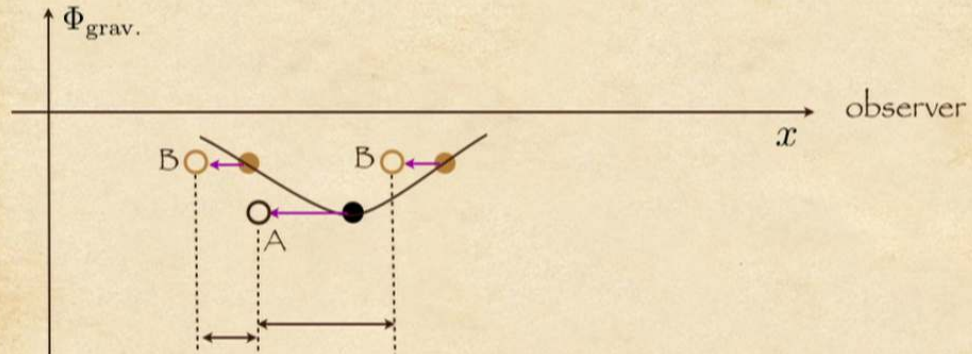
$$\xi_{\text{cross}}(x_A - x_B) \neq \xi_{\text{cross}}(x_B - x_A)$$

i.e. whether B is in front of, or behind A, matters.

Wojtak, Hansen & Hjorth - average by stacking clusters

Also: McDonald; Yoo, Hamaus, Seljak & Zaldarriaga; Zhao, Peacock & Li;
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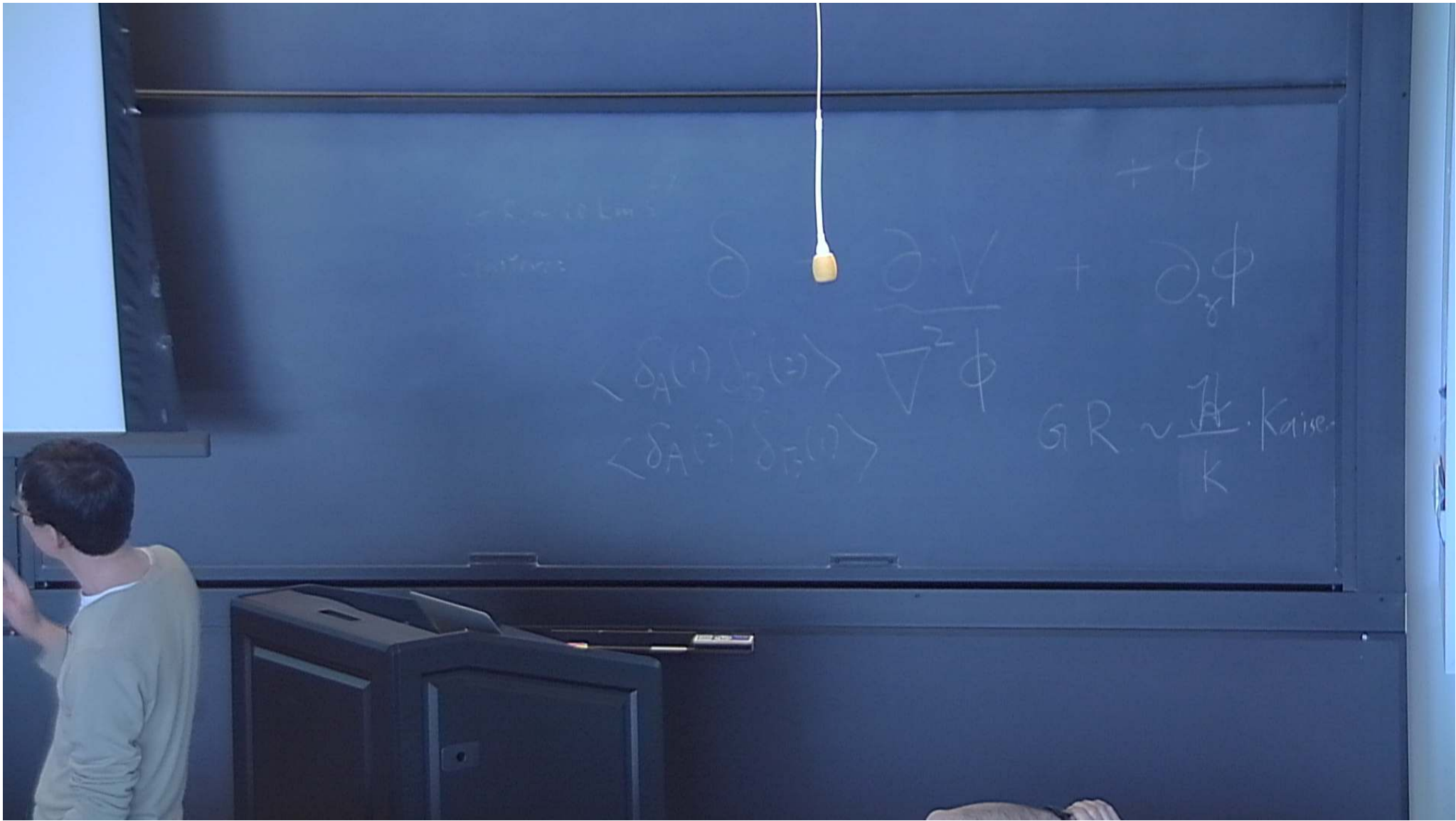


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What are other parity violating effects?

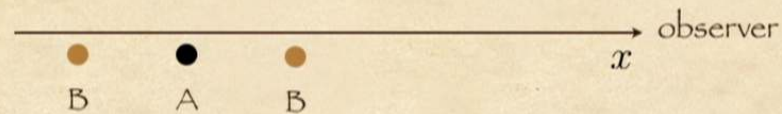
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$$\delta_{\text{obs.}} \sim \delta \left[1 + \frac{\mathcal{H}}{k} + \frac{\mathcal{H}^2}{k^2} \right]$$

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parity violating

Yoo, Fitzpatrick, Zaldarriaga; Challinor, Lewis;
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- More mundane, but present: evolution.



- Can disentangle between the two.

Footnote 1: parity violation only in the z direction.

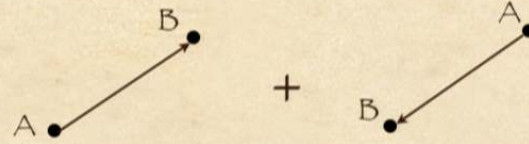
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Gravitational redshift term canceled, assuming geodesic motion.

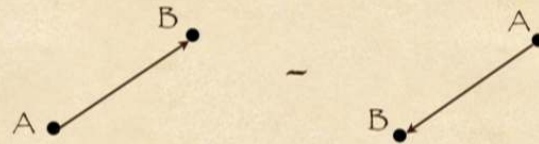
Footnote 3: selection effects.

Lessons for LSS measurement:

- Don't just add:



Subtract too:



Or, more generally: combine different orientations appropriately.

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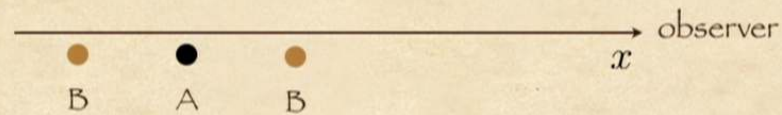
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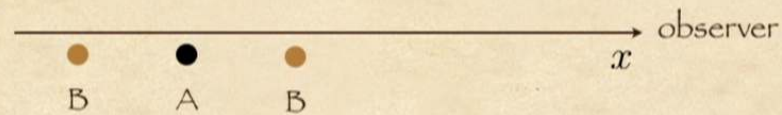
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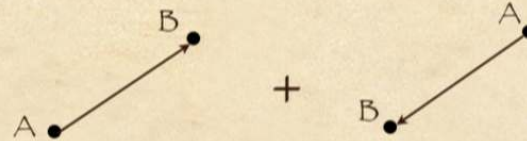
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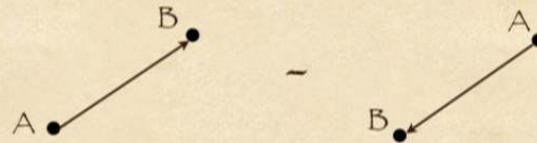
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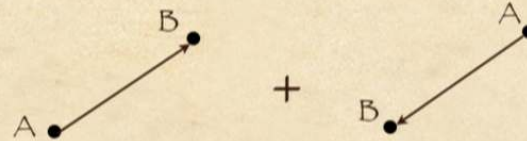


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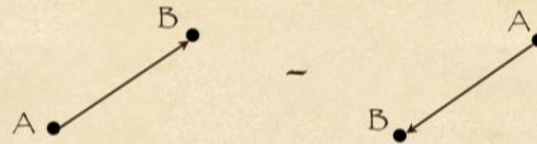
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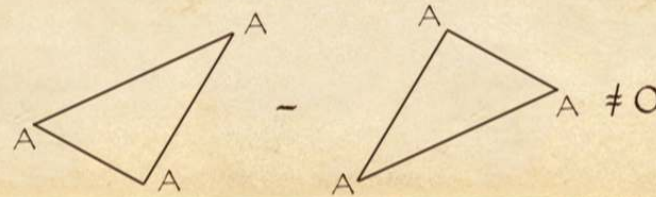
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- Question: do we need to cross-correlate multiple populations to see parity violating effects in higher N-point functions?

Answer: no.



Idea 3: non-perturbative consistency relations in LSS

- 1. Consider a familiar example of symmetry: spatial translation.

$$x \rightarrow x + \Delta x, \quad \text{where } \Delta x = \text{const.}$$

Its consequence for correlation function is well known:

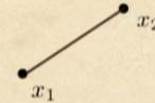
$$\langle \phi(x_1)\phi(x_2)\phi(x_3) \rangle = \langle \phi(x_1 + \Delta x)\phi(x_2 + \Delta x)\phi(x_3 + \Delta x) \rangle$$

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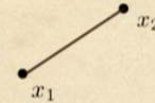
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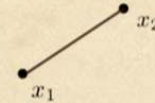
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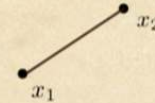
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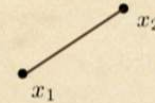
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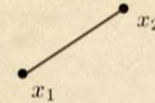
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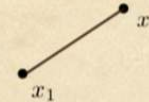
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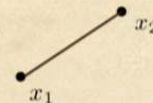
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- What makes the second case so different?



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- 1. Consider a familiar example of symmetry: **spatial translation**.

$$x \rightarrow x + \Delta x, \text{ where } \Delta x = \text{const.}$$

Its consequence for correlation function is well known:

$$\langle \phi(x_1)\phi(x_2)\phi(x_3) \rangle = \langle \phi(x_1 + \Delta x)\phi(x_2 + \Delta x)\phi(x_3 + \Delta x) \rangle$$

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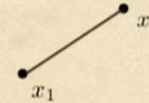
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Conclude: $\langle \phi_1\phi_2\phi_3 \rangle$ is **not** invariant under $\phi \rightarrow \phi + c$

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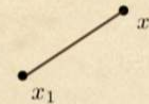
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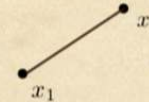
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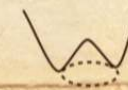
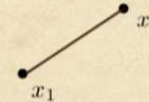
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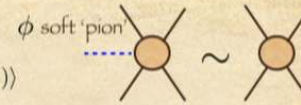
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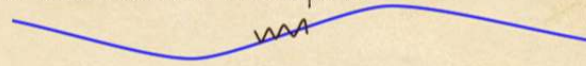


Consistency relations from SSB

- Schematic form: $\lim_{q \rightarrow 0} \frac{1}{P_\phi(q)} \langle \phi(q) \mathcal{O}(k_1) \dots \mathcal{O}(k_N) \rangle \sim \langle \mathcal{O}(k_1) \dots \mathcal{O}(k_N) \rangle$



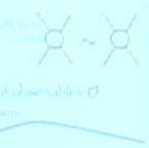
They are (momentum space) statements about how correlations of observables \mathcal{O} behave in the presence of a long wave-mode Goldstone boson/pion.



Consistency relations from SSIS

Schematic form $\lim_{\omega \rightarrow 0} \frac{1}{\omega} \langle \delta \rho(\omega, \vec{k}) \delta \rho(\omega, -\vec{k}) \rangle$

They are "momentum space" statements about how correlations of observables \mathcal{O} behave in the presence of a long-wavelength disturbance ϕ .




Handwritten notes on a chalkboard:

$\delta \rho(\omega, \vec{k})$

$\langle \delta \rho(\omega, \vec{k}) \delta \rho(\omega, -\vec{k}) \rangle$

$\langle \delta A^i(\omega, \vec{k}) \delta \rho(\omega, -\vec{k}) \rangle$

$\delta \rho = \frac{\partial V}{\partial \phi} + \partial_i \phi$

$\nabla^2 \phi$

$GR \sim \frac{H^2}{k^3} K_{IHG}$



Symmetries and consistency relations

comoving gauge $\delta\phi = 0$ $ds_{\text{spatial}}^2 = a^2 e^{2\zeta} [e^\gamma]_{ij} dx^i dx^j$

dilation symm. $x \rightarrow e^{-2\lambda} x$, $\zeta \rightarrow \zeta + \lambda$

$$\lim_{q \rightarrow 0} \frac{1}{P_\zeta(q)} \langle \zeta(q) \zeta_{k_1} \dots \zeta_{k_m} \rangle' \sim k \cdot \partial_k \langle \zeta_{k_1} \dots \zeta_{k_m} \rangle'$$

Maldacena

generalization $x \rightarrow x + M \cdot x^{N+1}$, $\zeta \rightarrow \zeta + M \cdot x^N$, $\gamma \rightarrow \gamma + M \cdot x^N$

$$\lim_{q \rightarrow 0} \partial_q^N \left(\frac{1}{P_\zeta(q)} \langle \zeta(q) \zeta_{k_1} \dots \zeta_{k_m} \rangle' + \frac{1}{P_\gamma(q)} \langle \gamma(q) \zeta_{k_1} \dots \zeta_{k_m} \rangle' \right) \sim k \cdot \partial_k^{N+1} \langle \zeta_{k_1} \dots \zeta_{k_m} \rangle'$$

Note:

1. The symmetries originate as diff. But consistency relations are not empty statements i.e. they can be violated (e.g. curvaton); they are a test of initial conditions (e.g. single clock, etc).
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References:

Maldacena; Creminelli & Zaldarriaga; Creminelli, Norena, Simonovic; Assassi, Baumann & Green; Flauger, Green & Porto; Pajer, Schmidt, Zaldarriaga; Kehagias & Riotto; Peloso & Pietronni; Berezhiani & Khoury; Pimentel; Creminelli, Norena, Simonovic, Vernizzi; Goldberger, LH, Nicolis; Hinterbichler, LH, Khoury; Horn, LH, Xiao.

A Newtonian symmetry:

The Newtonian continuity, Euler and Poisson eqs. are invariant under:

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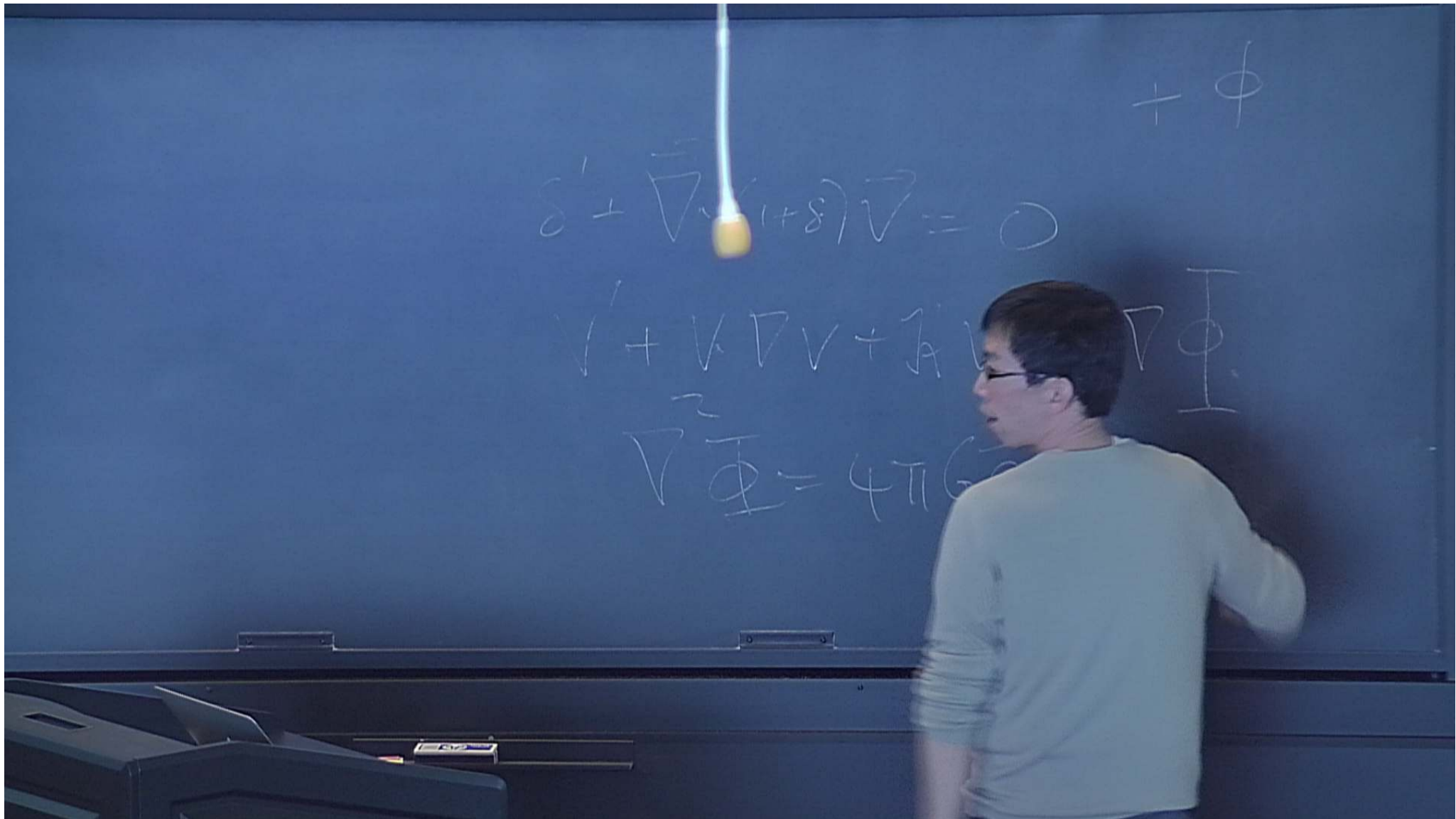
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Comments:

- The high k observables can be highly nonlinear and astrophysically messy.
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- The consistency relation is non-trivial only at unequal times - makes the interesting regime challenging to observe.



$$+ \phi$$

$$\delta' + \vec{\nabla} \cdot (\epsilon \vec{V}) = 0$$

$$\vec{V}' + \epsilon \vec{\nabla} \cdot \vec{V} + \epsilon \vec{\nabla} \phi$$

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giving the consistency relation

$$\lim_{\eta \rightarrow 0} \frac{1}{\eta} (\mathcal{L}' + \mathcal{L}'' - \mathcal{L}) = \sum_{i=1}^N \frac{D_i \mathcal{L}''(t)}{D_i \eta} \mathcal{L}''(t) = \mathcal{L}''(t)$$

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$$\begin{aligned} \partial^2 \Phi - \nabla \cdot \delta \mathbf{V} &= 0 \\ (\partial^2 + \mathcal{H} \partial_t) \Phi + \delta \mathbf{V} \cdot \nabla &= -\mathcal{H} \delta \\ \nabla \cdot \delta \mathbf{V} &= -\mathcal{H} \delta \end{aligned}$$



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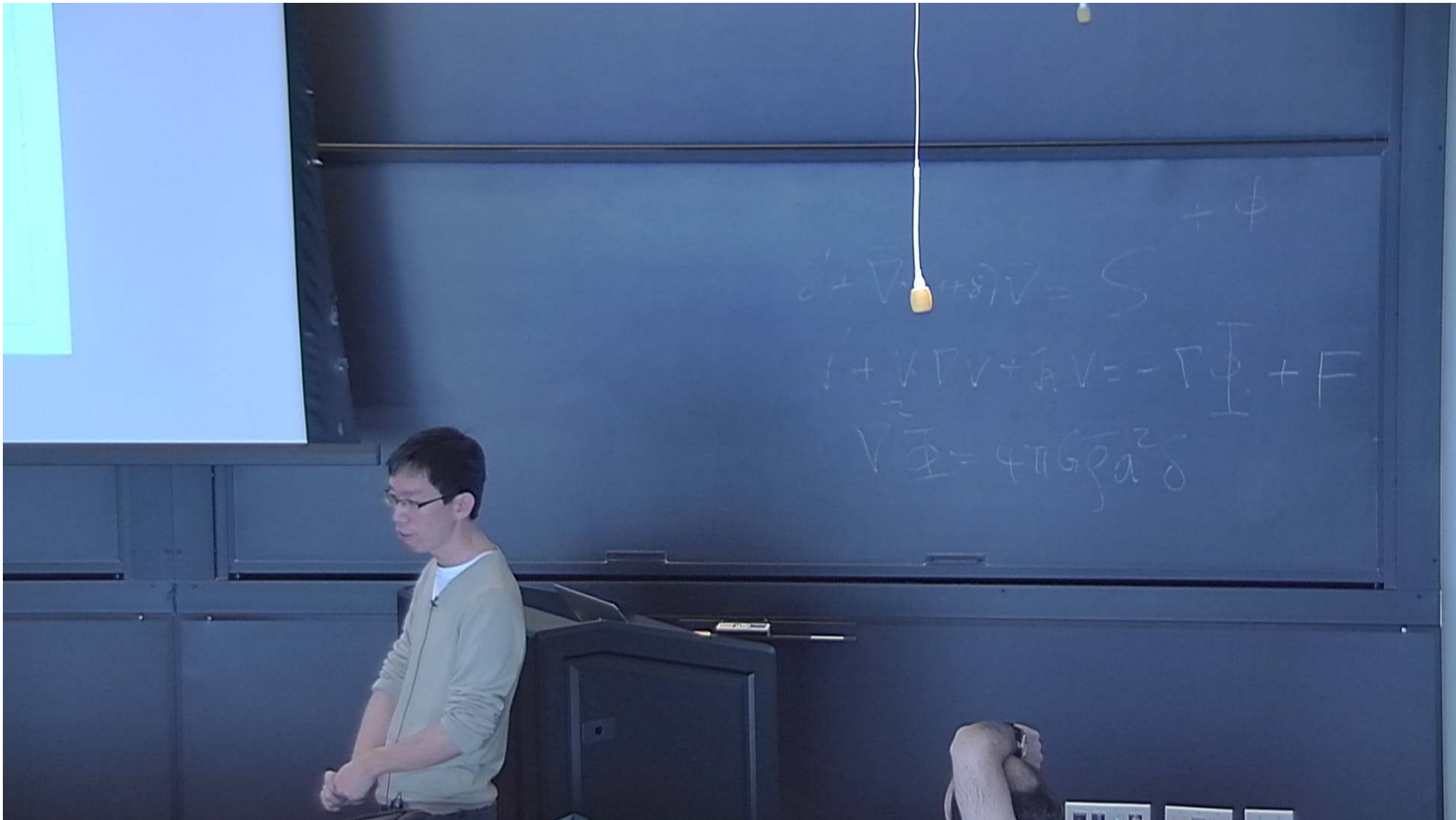
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A side remark:

The Newtonian consistency relation simplifies greatly in Lagrangian space:

$$\lim_{\vec{q} \rightarrow 0} \frac{1}{P_v(q, \eta)} \langle v^i(\vec{q}, \eta) \mathcal{O}(\vec{k}_1, \eta_1) \dots \mathcal{O}(\vec{k}_m, \eta_m) \rangle = 0$$

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Open issues:

- Connection with asymptotic symmetries (e.g. BMS)?

- Why $P_\delta(q) \sim q^n$, $n < -3$?



