

Title: From Black Holes to Quantum Transport

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Abstract: <p>In the last few years there has been significant interest in the possible applications of gauge-gravity duality to condensed matter systems. In this talk I will discuss recent applications of these holographic techniques to strongly correlated systems out of equilibrium. I will argue that insights from general relativity, hydrodynamics and quantum field theory may be combined to yield quantitative predictions for quantum transport.</p>

# From Black Holes to Quantum Transport

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King's College London



Perimeter Institute  
Waterloo, Canada  
22<sup>nd</sup> April 2015







# Strings, Cosmology & Condensed Matter

Benjamin Doyon    Koenraad Schalm    Andy Lucas



Ben Simons    Julian Sonner



Jerome Gauntlett    Toby Wiseman



King's, Leiden, Harvard, Cambridge, Imperial

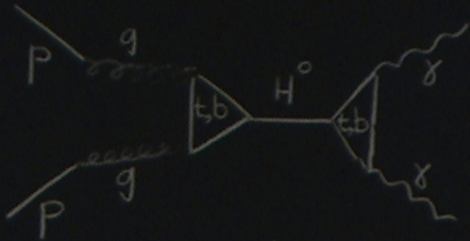


## Outline

- Motivation from condensed matter
- Gauge-gravity duality
- Far from equilibrium dynamics
- Recent work on energy flow
- Einstein equations, hydrodynamics, transport
- Current status and future developments

MJB, Benjamin Doyon, Andrew Lucas, Koenraad Schalm  
*“Far from equilibrium energy flow in quantum critical systems”*  
arXiv:1311.3655

$$m_H \approx 125 \text{ GeV}$$



$$\phi = \frac{1}{\sqrt{2}} \begin{pmatrix} \phi^1 + i\phi^2 \\ \phi^0 + i\phi^3 \end{pmatrix}$$

$$\mathcal{L}_H = \left| \left( \partial_\mu - ig W_\mu^a T^a - \frac{ig'}{2} B_\mu \right) \phi \right|^2 + \mu^2 \phi^\dagger \phi - \lambda (\phi^\dagger \phi)^2$$

$$M_W = \frac{v|g|}{2}$$

$$M_Z = \frac{v\sqrt{g^2 + g'^2}}{2}$$

$$\cos \Theta_W = \frac{M_W}{M_Z} = \frac{|g|}{\sqrt{g^2 + g'^2}}$$

$$\phi_0 = \sqrt{\frac{-\mu^2}{\lambda}} = v$$

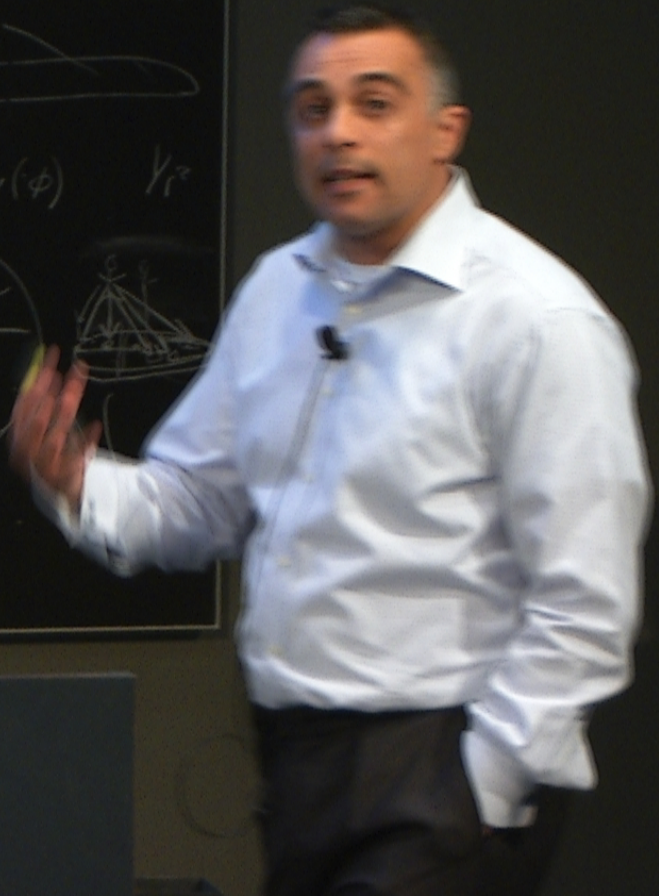
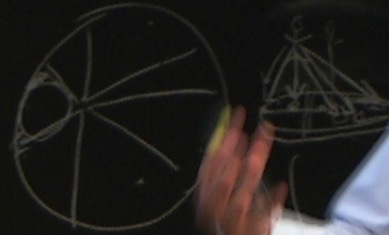
$$\mathcal{L} = \frac{1}{2} (\partial_\mu \eta)^2 - (\lambda v^3) \eta^2 + \frac{1}{2} (\partial_\mu \xi)^2 + \dots \text{ higher order}$$

$$m_\eta = \sqrt{2\lambda v} = \sqrt{2\mu^2} > 0$$

$$\partial_\mu \rightarrow D_\mu = \partial_\mu - ie A_\mu$$

$$A'_\mu = A_\mu + \frac{1}{e} \partial_\mu \alpha$$

$$\mathcal{L} = (D^\mu \phi)^\dagger (D_\mu \phi) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - v(\phi)$$

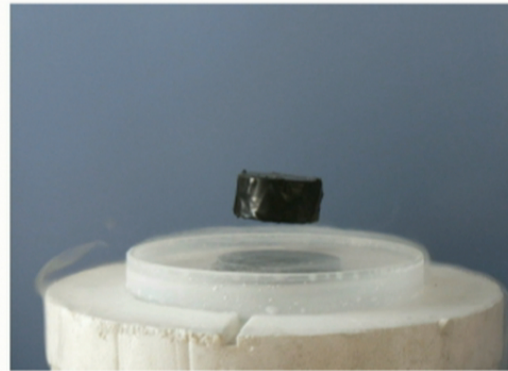
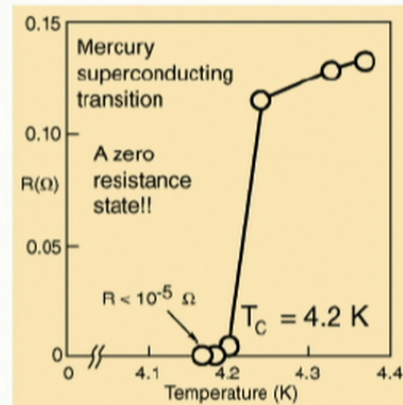




# Condensed Matter Theory

## Novel states of matter

### Superconductor



<http://hyperphysics.phy-astr.gsu.edu/hbase/solids/scex.html#c1>

<http://en.wikipedia.org/wiki/Superconductivity>

**Zero Resistance** Kamerlingh Onnes (1911)  $-269^{\circ}\text{C}$

## Phase Transition



## “More is different”

$10^{23}$

100,000,000,000,000,000,000,000

**Rich cooperative behaviour**

**Emergence**



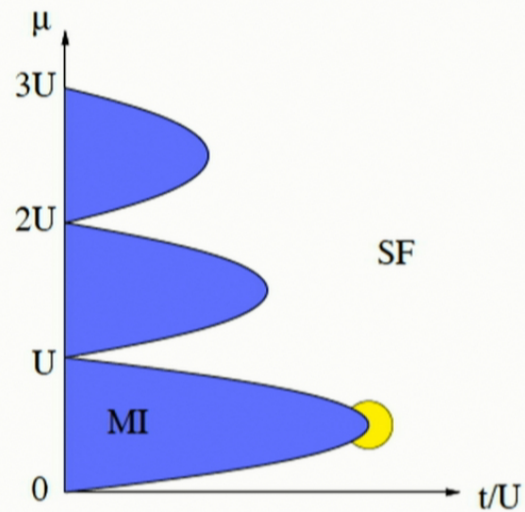
[http://en.wikipedia.org/wiki/Philip\\_Warren\\_Anderson](http://en.wikipedia.org/wiki/Philip_Warren_Anderson)

Cooper Pairs   Fractional Charges   No Quasi-Particles

**Strongly Correlated Systems**

## The Bose–Hubbard Model

$$H = -t \sum_{\langle ij \rangle} (a_i^\dagger a_j + \text{h.c.}) + \frac{U}{2} \sum_i n_i(n_i - 1) - \mu \sum_i n_i$$



Fisher, Weichman, Grinstein and Fisher, PRB 40, 546 (1989)

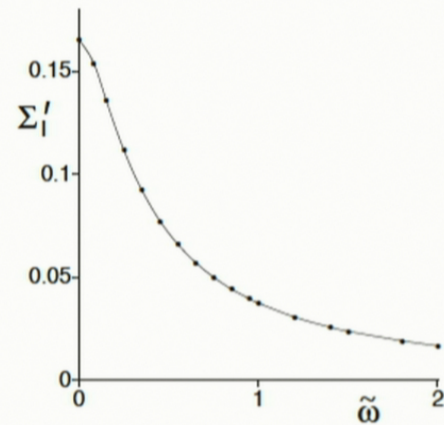
$$L = \int d^D x |\partial_\mu \Phi|^2 - m^2 |\Phi|^2 - \frac{u_0}{3} |\Phi|^4$$



# Universal Transport

Damle and Sachdev, PRB 56, 8714 (1997)

$$\sigma(\omega) = \frac{(2e)^2 T^{d-2}}{\epsilon^2} \Sigma\left(\frac{\omega}{\epsilon^2 T}\right) \quad d = 3 - \epsilon$$



Two spatial dimensions

$$\sigma(0) \approx 1.037 \left( \frac{4e^2}{h} \right)$$

Crossover between hydrodynamic and collisionless regimes

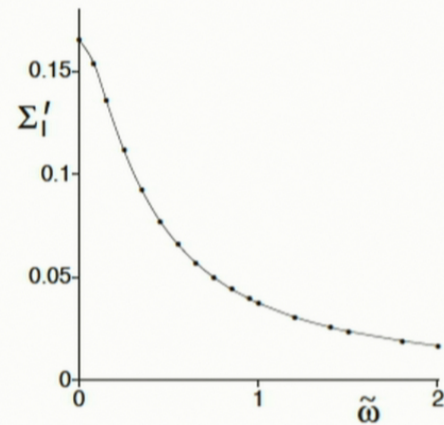




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# Transport near Quantum Critical Points

Linear response for SF-MI transition in Bose–Hubbard

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MJB, Green and Sondhi, PRL **98**, 166801 (2007)

$$\alpha_{xy} = \frac{S}{B} \quad \bar{\kappa}_{xx} \simeq g \epsilon^2 \frac{T^{d+3}}{(2e)^2 B^2} \quad g \approx 5.5$$

Hartnoll, Kovtun, Müller and Sachdev, PRB **76**, 144502 (2007)

$$\bar{\kappa}_{xx}(B) = \frac{TS^2}{B^2 \sigma_{xx}(0)}$$

**Relativistic hydrodynamics & AdS/CFT link all coeffs**

QBE Müller *et al*, PRB (2008) MJB *et al*, PRB (2009)

Viscosity/Entropy

Quark Gluon Plasma

Graphene

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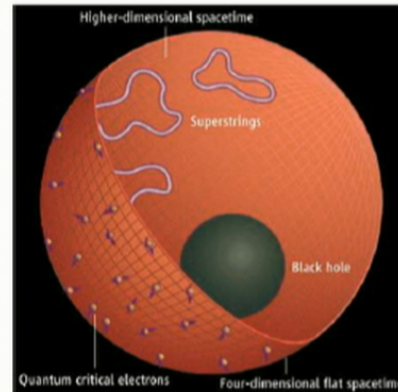
Graphene



# Gauge-Gravity Duality

AdS/CFT correspondence (Anti-de Sitter/Conformal Field Theory)

Maldacena (1997); Gubser, Klebanov, Polyakov (1998); Witten (1998)



S. Hartnoll, *Science* **322**, 1639 (2008)

For an overview see for example John McGreevy, *Holographic duality with a view toward many body physics*, arXiv:0909.0518

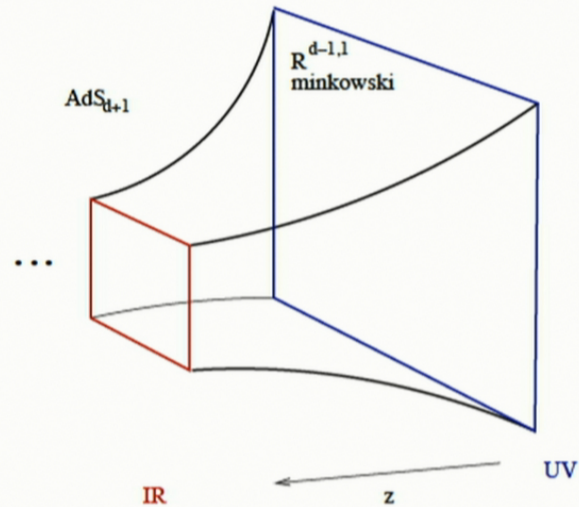
**Strong-Weak Duality**

**A different way of looking at things**





# AdS/CFT Correspondence



Gubser–Klebanov–Polyakov–Witten

$$Z[\phi_0]_{\text{CFT}} \simeq e^{-S_{\text{AdS}}[\phi]} \Big|_{\phi \sim \phi_0 \text{ at } z=0}$$

$$\phi(z) \sim z^{d-\Delta} \phi_0(1 + \dots) + z^\Delta \phi_1(1 + \dots)$$

Fields in AdS  $\leftrightarrow$  operators in dual CFT  $\phi \leftrightarrow \mathcal{O}$





# Progress in AdS/CMT

Anti-de Sitter/Condensed Matter Theory

## Transport Coefficients

Viscosity, Conductivity, Hydrodynamics, Bose–Hubbard, Graphene

## Strange Metals

Non-Fermi liquids, instabilities, cuprates

## Holographic Duals

Superfluids, Fermi Liquid,  $O(N)$ , Luttinger Liquid

Equilibrium or close to equilibrium



# Non-Equilibrium High Energy Physics

## Thermalization in Strongly Coupled Gauge Theories

Chesler & Yaffe (2009), de Boer & Keski Vakkuri (2011), ...

Buchel, Lehner, & Myers, “*Thermal quenches in  $N = 2^*$  plasmas*”,  
JHEP 1208 (2012)

Buchel, Heller, & Myers, “*Equilibration rates in a strongly coupled  
nonconformal quark-gluon plasma*”, arXiv:1503.07114

## Quantum Quenches

Aparício & López (2011), Basu & Das (2012), ...

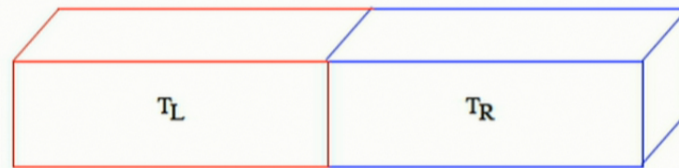
Buchel, Lehner, Myers, & van Niekerk, “*Quantum quenches of  
holographic plasmas*”, arXiv:1302.2924

Das, Galante, & Myers “*Universality in fast quantum quenches*”,  
arXiv:1411.7710

## Hydrodynamics

# Thermalization

Condensed matter and high energy physics



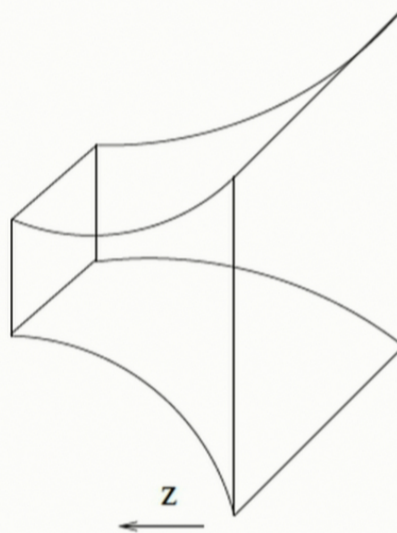
Why not connect two strongly correlated systems together  
and see what happens?



# AdS/CFT

Energy flow may be studied within pure Einstein gravity

$$S = \frac{1}{16\pi G} \int d^{d+2}x \sqrt{-g} (R - 2\Lambda)$$

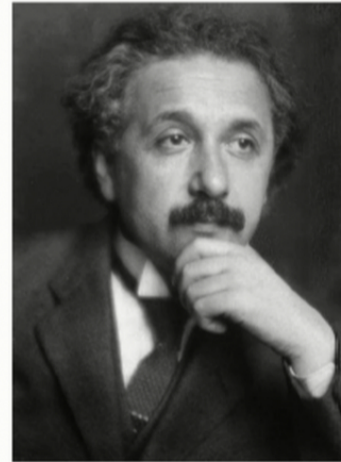
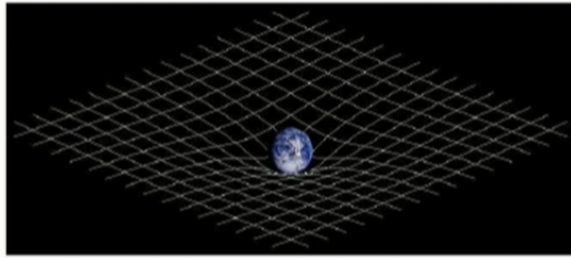


$$g_{\mu\nu} \leftrightarrow T_{\mu\nu}$$



# Einstein Centenary

## The Field Equations of General Relativity (1915)



$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = \frac{8\pi G}{c^4}T_{\mu\nu}$$

[http://en.wikipedia.org/wiki/Einstein\\_field\\_equations](http://en.wikipedia.org/wiki/Einstein_field_equations)

$g_{\mu\nu}$  metric     $R_{\mu\nu}$  Ricci curvature     $R$  scalar curvature

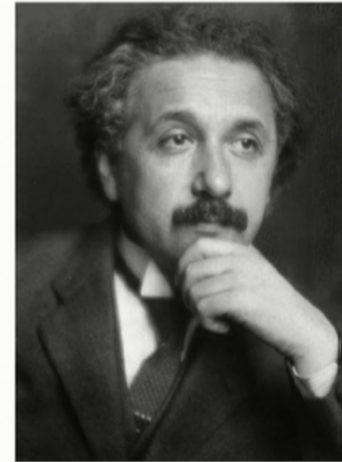
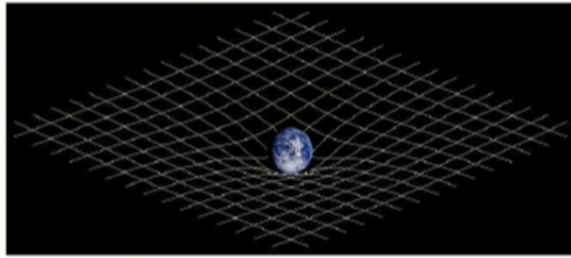
$\Lambda$  cosmological constant     $T_{\mu\nu}$  energy-momentum tensor

### Coupled Nonlinear PDEs

Schwarzschild Solution (1916)     $R_S = \frac{2MG}{c^2}$     Black Holes

# Einstein Centenary

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$g_{\mu\nu}$  metric     $R_{\mu\nu}$  Ricci curvature     $R$  scalar curvature  
 $\Lambda$  cosmological constant     $T_{\mu\nu}$  energy-momentum tensor

### Coupled Nonlinear PDEs

Schwarzschild Solution (1916)     $R_S = \frac{2MG}{c^2}$     Black Holes





# Non-Equilibrium CFT

Bernard & Doyon, *Energy flow in non-equilibrium conformal field theory*, J. Phys. A: Math. Theor. **45**, 362001 (2012)

Two critical 1D systems (central charge  $c$ )  
at temperatures  $T_L$  &  $T_R$



Join the two systems together



Alternatively, take one critical system and impose a step profile

Local Quench

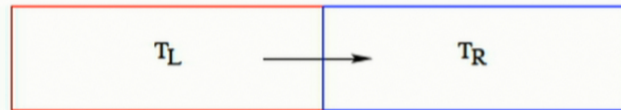


## Steady State Energy Flow

Bernard & Doyon, *Energy flow in non-equilibrium conformal field theory*, J. Phys. A: Math. Theor. 45 362001 (2012)

If systems are very large ( $L \gg vt$ ) they act like heat baths

For times  $t \ll L/v$  a steady energy current flows



**Non-equilibrium steady state**

$$J = \frac{c\pi^2 k_B^2}{6h} (T_L^2 - T_R^2)$$

**Universal result out of equilibrium**

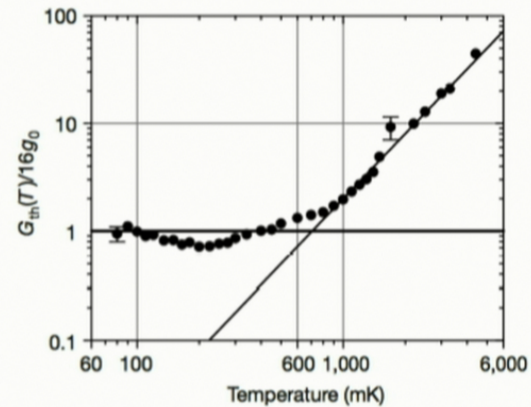
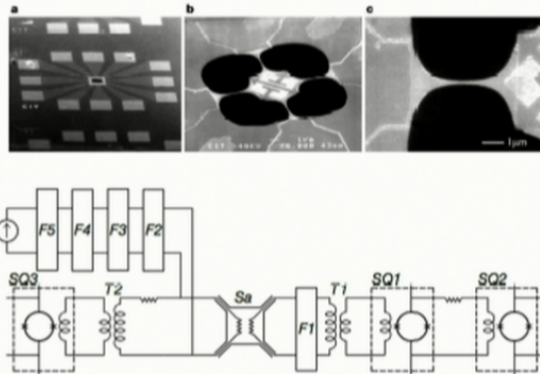
Direct way to measure central charge; velocity doesn't enter

Sotiriadis and Cardy. J. Stat. Mech. (2008) P11003

$c$  doesn't have to be  $\in \mathbb{Z}$  **Emergence** Stefan-Boltzmann

# Experiment

Schwab, Henriksen, Worlock and Roukes, *Measurement of the quantum of thermal conductance*, Nature 404, 974 (2000)



## Quantum of Thermal Conductance

### Numerical Simulations

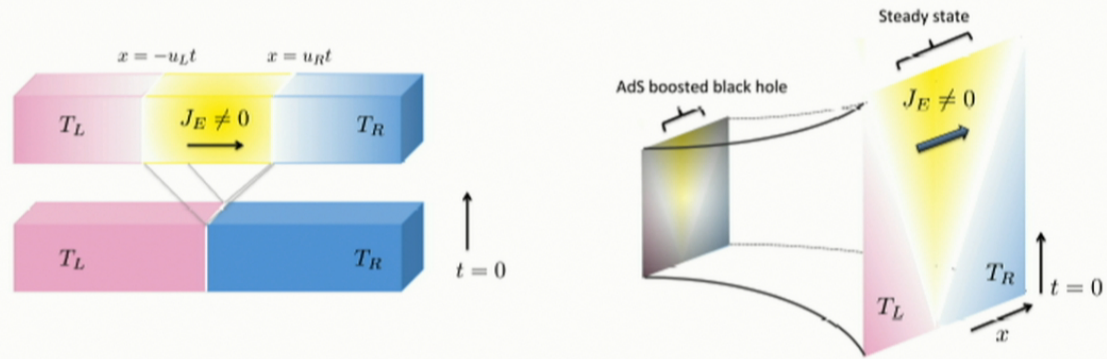
Time-dependent Density Matrix Renormalization Group (DMRG)

Karrasch, Ilan and Moore, Phys. Rev. B 88, 195129 (2013)

Quantum Information and Entanglement



# AdS/CFT



Steady State Region

Spatially Homogeneous

## Boost Solution

Lorentz boosted stress tensor of a finite temperature CFT

Perfect fluid

$$\langle T^{\mu\nu} \rangle_s = a_d T^{d+1} (\eta^{\mu\nu} + (d+1)u^\mu u^\nu)$$

$$\eta^{\mu\nu} = \text{diag}(-1, 1, \dots, 1) \quad u^\mu = (\cosh \theta, \sinh \theta, 0, \dots, 0)$$

One spatial dimension

$$a_1 = \frac{L\pi}{4G} \quad c = \frac{3L}{2G}$$

$$T_L = T e^\theta \quad T_R = T e^{-\theta} \quad \langle T_{tx} \rangle = \frac{c\pi^2 k_B^2}{6h} (T_L^2 - T_R^2)$$

Can also obtain complete steady state density matrix

Describes all the cumulants of the energy transfer process

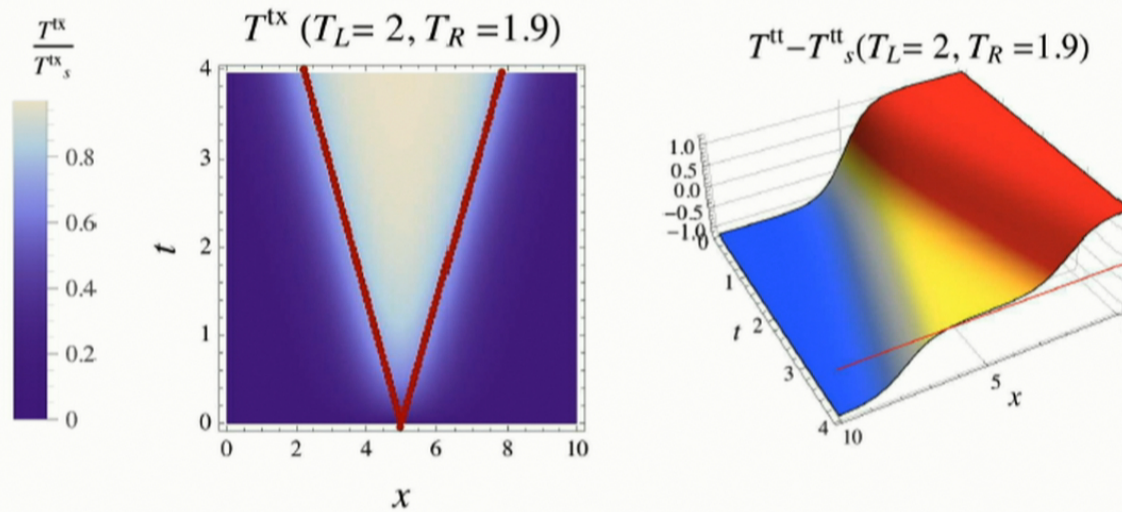
The non-equilibrium steady state (NESS)

is a Lorentz boosted thermal state

Run past a thermal state with temperature  $T = \sqrt{T_L T_R}$



# Numerics I



Excellent agreement with predictions

