

Title: Explorations in Cosmology-4

Date: Apr 09, 2015 10:15 AM

URL: <http://pirsa.org/15040028>

Abstract:

CLARIFY WICK'S THEOREM

$$\hat{X}_1 = \alpha_1 \hat{a} + \beta_1 \hat{a}^\dagger$$

$$\hat{X}_2 = \alpha_2 \hat{a} + \beta_2 \hat{a}^\dagger$$

⋮

$$\langle 0 | \hat{X}_1 \hat{X}_2 | 0 \rangle = \langle 0 | (\alpha_1 \hat{a}) (\beta_2 \hat{a}^\dagger) | 0 \rangle = \alpha_1 \beta_2$$

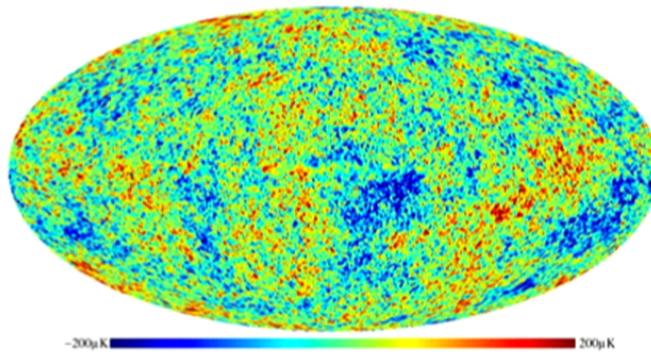
ODD CORRELATOR  $\langle 0 | \hat{X}_1 \dots \hat{X}_{2N-1} | 0 \rangle = 0$

EVEN CORRELATOR

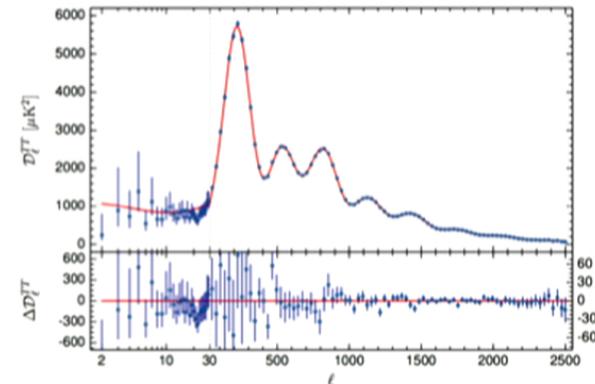
$$\begin{aligned} \langle 0 | \hat{X}_1 \hat{X}_2 \dots \hat{X}_{2N} | 0 \rangle &= \langle 0 | \overbrace{\hat{X}_1 \hat{X}_2} \overbrace{\hat{X}_3 \hat{X}_4} \dots | 0 \rangle + \langle 0 | \overbrace{\hat{X}_1 \hat{X}_2 \hat{X}_3} \overbrace{\hat{X}_4} \dots | 0 \rangle + \dots \\ &= (\langle 0 | \hat{X}_1 \hat{X}_2 | 0 \rangle \langle 0 | \hat{X}_3 \hat{X}_4 | 0 \rangle \dots) + (\langle 0 | \hat{X}_1 \hat{X}_3 | 0 \rangle \langle 0 | \hat{X}_2 \hat{X}_4 | 0 \rangle \dots) + \dots = (\alpha_1 \beta_2 \alpha_3 \beta_4 \dots) \end{aligned}$$

# CMB temperature

Temperature map

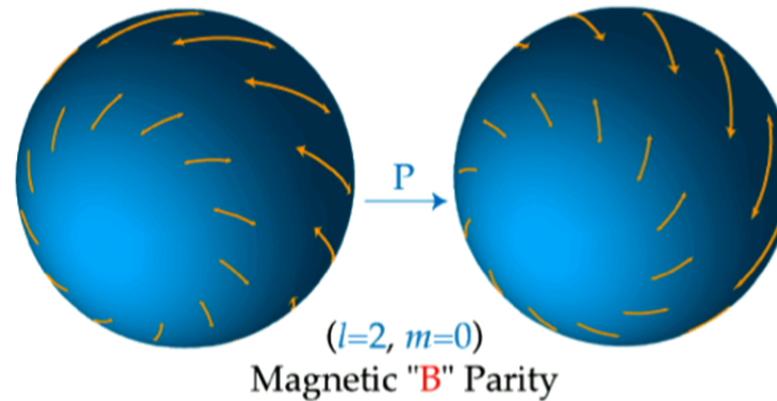
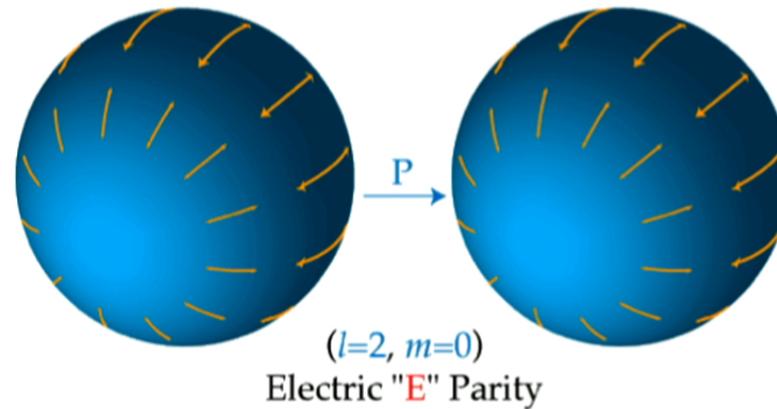


TT power spectrum



Fluctuations are  $\sim 100$  microK  
The CMB is very Gaussian! ( $\sim 0.1\%$ )

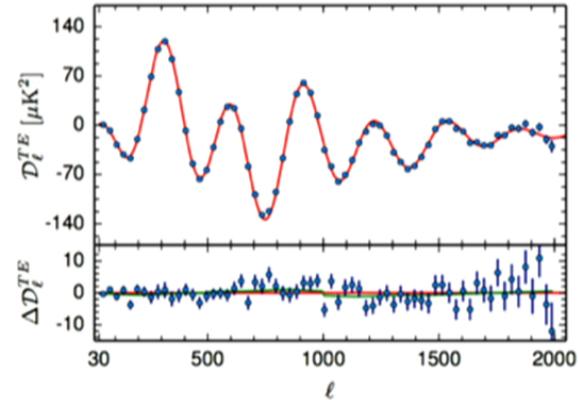
# CMB polarization has an E-B decomposition



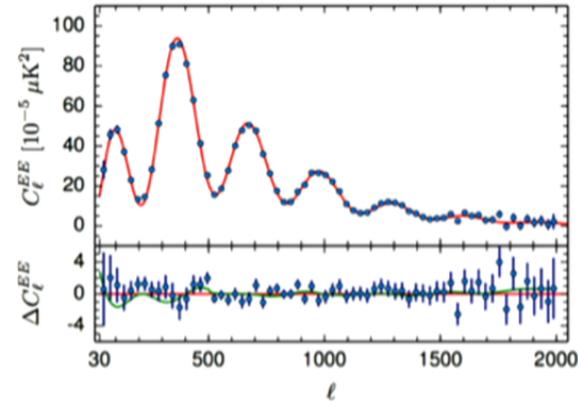
Wayne Hu

# E-modes are qualitatively similar to T

TE power spectrum

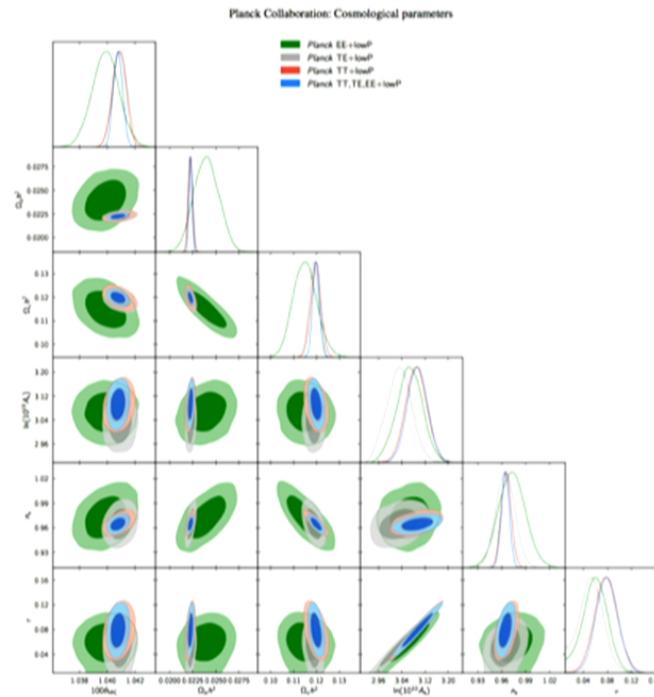


EE power spectrum



E-mode polarization is a few microK

# T+E constrain cosmological parameters



Parameter	TT+lowP 68 % limits	TT+lowP+lensing 68 % limits	TT+lowP+lensing+ext 68 % limits	TT,TE,EE+lowP 68 % limits	TT,TE,EE+lowP+lensing 68 % limits	TT,TE,EE+lowP+lensing+ext 68 % limits
$\Omega_b h^2$ . . . . .	$0.02222 \pm 0.00023$	$0.02226 \pm 0.00023$	$0.02227 \pm 0.00020$	$0.02225 \pm 0.00016$	$0.02226 \pm 0.00016$	$0.02230 \pm 0.00014$
$\Omega_c h^2$ . . . . .	$0.1197 \pm 0.0022$	$0.1186 \pm 0.0020$	$0.1184 \pm 0.0012$	$0.1198 \pm 0.0015$	$0.1193 \pm 0.0014$	$0.1188 \pm 0.0010$
$100\theta_{MC}$ . . . . .	$1.04085 \pm 0.00047$	$1.04103 \pm 0.00046$	$1.04106 \pm 0.00041$	$1.04077 \pm 0.00032$	$1.04087 \pm 0.00032$	$1.04093 \pm 0.00030$
$\tau$ . . . . .	$0.078 \pm 0.019$	$0.066 \pm 0.016$	$0.067 \pm 0.013$	$0.079 \pm 0.017$	$0.063 \pm 0.014$	$0.066 \pm 0.012$
$\ln(10^{10} A_s)$ . . . . .	$3.089 \pm 0.036$	$3.062 \pm 0.029$	$3.064 \pm 0.024$	$3.094 \pm 0.034$	$3.059 \pm 0.025$	$3.064 \pm 0.023$
$n_s$ . . . . .	$0.9655 \pm 0.0062$	$0.9677 \pm 0.0060$	$0.9681 \pm 0.0044$	$0.9645 \pm 0.0049$	$0.9653 \pm 0.0048$	$0.9667 \pm 0.0040$

## B-modes

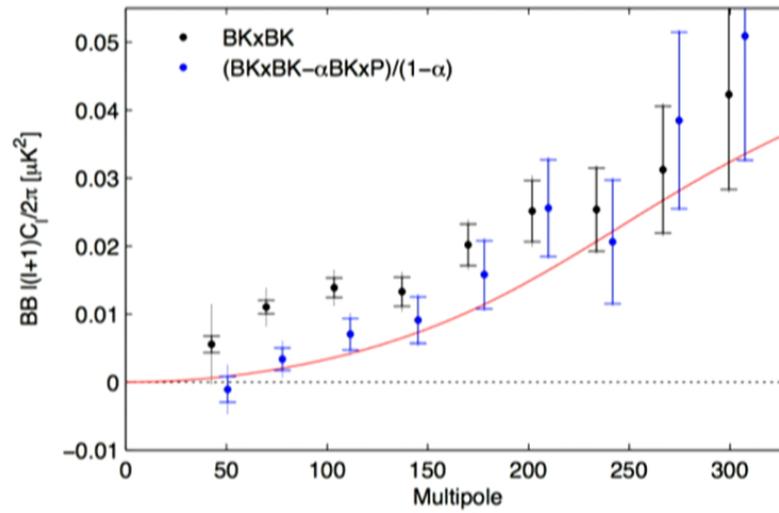
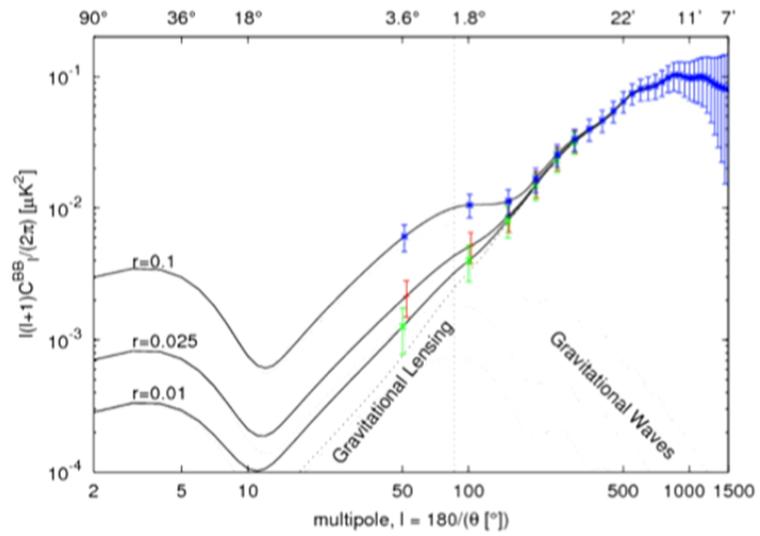
Theorem: scalar sources (e.g. inflaton) do not generate B-mode polarization in linear perturbation theory

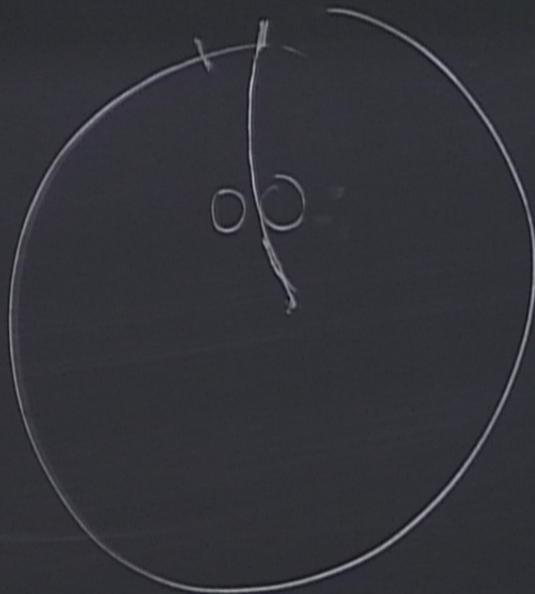
Theorem: parity invariance  $\Rightarrow TB=EB=0$   
i.e. BB is the only B-mode related power spectrum

In slow-roll field inflation, the leading sources of B-modes are

- gravity waves from inflation!! (maybe)
- gravitational lensing (leading second-order effect; foreground for measuring gravity waves)

Because gravitational lensing is small (second order), B-mode polarization is small ( $\sim 100$  nano-K) and challenging to measure, but potentially a very powerful probe of gravity waves





$$|T_l|^2 \sim C_l^{TT}$$

$$|E_l|^2 \sim C_l^{FF}$$

$$|T_l E_l|^2 \sim C_l^{TF}$$

6-PARAMETER MODEL:  $\{\Omega_b, \Omega_c, \Omega_u, \Delta_s^2, n_s, \tau\}$

9 EXPANSION HISTORY:  $\{\Omega_b, \Omega_c, \Omega_u\}$

3 INITIAL FLUCTUATIONS  $\{\Delta_s^2, n_s\}$

"ADIABATIC SCALAR GAUSSIAN POWER-LAW"

"ADIABATIC SCALAR":  $ds^2 = -dt^2 + a(t)^2 e^{2\psi(x)} dx^2$

$$\delta\rho_c^{(i)} = (\#) \xi(x)$$

"ADIABATIC CURVATURE"  
FUNCTION OF  $x$  BUT NOT  $t$

"GAUSSIAN"  $\langle \underbrace{S_{\vec{k}}}_{\text{circled}} S_{\vec{k}'} \rangle = P_S(k) (2\pi)^3 \delta^3(\vec{k} - \vec{k}')$

$$\langle S_{k_1} S_{k_2} S_{k_3} S_{k_4} \rangle = \langle S_{k_1} S_{k_2} \rangle \langle S_{k_3} S_{k_4} \rangle + [2 \text{ PERM.}]$$

"POWER LAW"

$$\underbrace{\frac{k^3}{2\pi^2}}_{\text{circled}} P_S(k) = \Delta_S^2 \left( \frac{k}{k_0} \right)^{n_S - 1}$$

$$n_S = 1 \Rightarrow \text{SCALE INVARIANT}$$

$$\begin{aligned} \langle S^2 \rangle &= \int \frac{d^3k}{(2\pi)^3} P_S(k) \\ &= \int_0^\infty d \log k \underbrace{\frac{k^3}{2\pi^2} P_S(k)} \end{aligned}$$

"GAUSSIAN"  $\langle \vec{S}_{\vec{k}} \vec{S}_{\vec{k}'} \rangle = P_S(k) (2\pi)^3 \delta^3(\vec{k} - \vec{k}')$

$$\langle S_{k_1} S_{k_2} S_{k_3} S_{k_4} \rangle = \langle S_{k_1} S_{k_2} \rangle \langle S_{k_3} S_{k_4} \rangle + [2 \text{ PERM.}] \frac{k^2}{2\pi} C_\ell$$

"POWER LAW"

$$\frac{k^3}{2\pi^2} P_S(k) = \Delta_S^2 \left( \frac{k}{k_0} \right)^{n_S - 1}$$

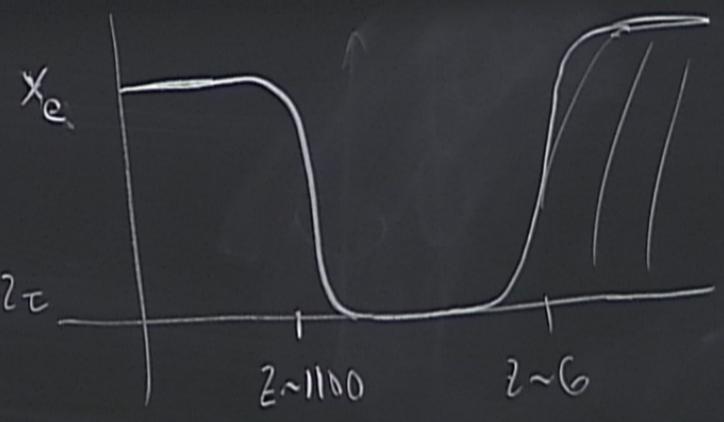
$$n_S = 0.9667 \pm 0.0040$$

$n_S = 1 \Rightarrow$  SCALE INVARIANT

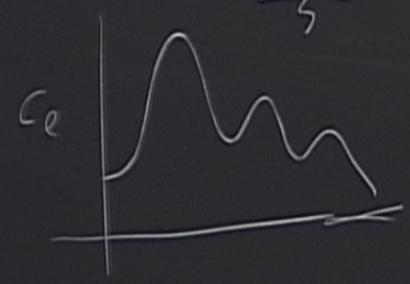
$$\begin{aligned} \langle S^2 \rangle &= \int \frac{d^3k}{(2\pi)^3} P_S(k) \\ &= \int_0^\infty d \log k \underbrace{\frac{k^3}{2\pi^2} P_S(k)} \end{aligned}$$

"OBSERVATION"

OPTICAL DEPTH  $\tau \sim 0.07$



$$\Delta_s^z e^{-2\tau}$$



"OBSERVATIONS" = CMB

CMB = COMPLICATED BUT LINEAR BLACK-BOX

$\int_k \rightarrow (T + \text{POLARIZATION ON SKY})$

SKY IS A SPHERE  $\eta_{ab}$  = EUCLIDAN 2-METRIC ON SPHERE

$I_{ab}(\hat{n}) = \langle E_a(\hat{n}) E_b(\hat{n}) \rangle$  INTENSITY TENSOR

TEMPERATURE  $T(\hat{n}) = \underbrace{f(\omega)}_{\text{BLACKBODY}} \times \eta^{ab} I_{ab}(\hat{n})$

LINEAR POLARIZATION  $P_{ab} = f(\omega) \times \left[ I_{ab} - \frac{1}{2} \eta_{ab} \eta^d I_{ab} \right]$  TRACE

TEMPERATURE IS ANALYZED IN HARMONIC SPACE

IF SKY WERE FLAT:

$$T(\vec{x}) = \int \frac{d^2\ell}{(2\pi)^2} T_{\vec{\ell}} e^{i\vec{\ell} \cdot \vec{x}} \quad \langle T_{\vec{\ell}} T_{\vec{\ell}'}^* \rangle = C_{\ell} (2\pi)^2 \delta^2(\ell - \ell')$$

ON A SPHERICAL SKY

$$T(\hat{n}) = \sum_{l=0}^{\infty} \sum_{m=-l}^l a_{lm}^+ Y_{lm}(\hat{n}) \quad \langle a_{lm}^+ a_{l'm'}^{T*} \rangle = C_{\ell}^{\text{TT}} \delta_{\ell\ell'} \delta_{mm'}$$

LINEAR POLARIZATION  $P_{ab} = f(\omega) \times \left[ I_{ab} - \frac{1}{2} \eta_{ab} \eta^d I_{ab} \right]$

TRACE

TEMPERATURE IS ANALYZED IN HARMONIC SPACE

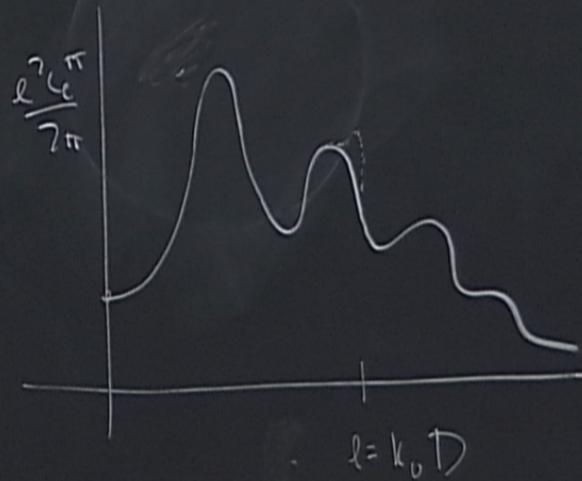
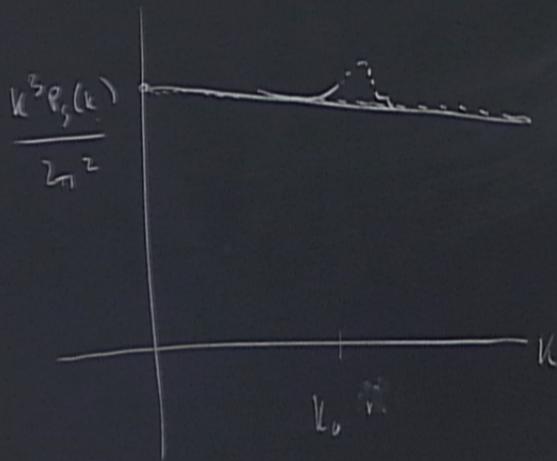
IF SKY WERE FLAT:

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TRACELESS SYMMETRIC TENSOR



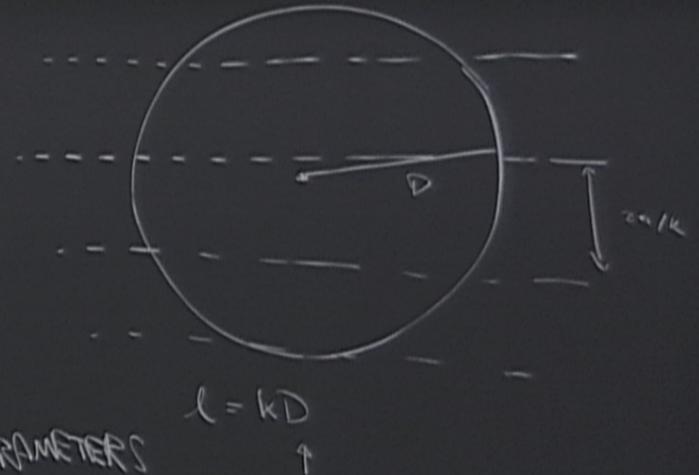
AT A CRUDE LEVEL OF APPROXIMATION

$$C_e^{TT} = P_s(k) \Big|_{k=e/D} \times T(k) \Big|_{k=e/D}$$

D = DISTANCE TO LAST SCATTERING  
= 14000 Mpc

↓  
COSMOLOGICAL PARAMETERS

$\Omega_b, \Omega_c, \Omega_m, \tau$



$$\sum_{l=0}^{\infty} \sum_{m=-l}^l u_{lm} \hat{e}_m^l$$

WHAT IS THE ANALOG OF THE SPHERICAL TRANSFORM FOR A TRACELESS SYMMETRIC TENSOR  $P_{ab}$ ?

FLAT SKY FIRST.

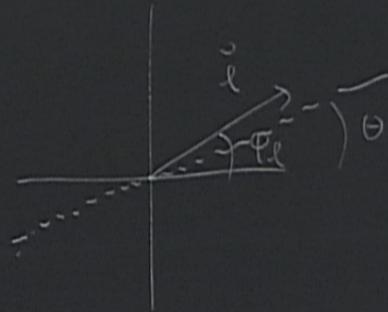
$$P_{ab}(\vec{x}) = Q(x) \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} + U(x) \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

"STOKES PARAMETERS"

TWO NUMBERS FOR EACH  $\vec{\ell}$ :  $Q_{\vec{\ell}}, U_{\vec{\ell}}$

MORE INTERESTING DECOMPOSITION:

$$\begin{pmatrix} E_{\vec{\ell}} \\ B_{\vec{\ell}} \end{pmatrix} = \begin{pmatrix} \cos 2\varphi_{\vec{\ell}} & \sin 2\varphi_{\vec{\ell}} \\ -\sin 2\varphi_{\vec{\ell}} & \cos 2\varphi_{\vec{\ell}} \end{pmatrix} \begin{pmatrix} Q_{\vec{\ell}} \\ U_{\vec{\ell}} \end{pmatrix}$$



ROTATION INVARIANT

$$\varphi' \rightarrow \varphi - \Theta$$

$$\begin{pmatrix} Q' \\ U' \end{pmatrix} \rightarrow \begin{pmatrix} \cos 2\Theta & \sin 2\Theta \\ -\sin 2\Theta & \cos 2\Theta \end{pmatrix} \begin{pmatrix} Q \\ U \end{pmatrix}$$

$$L = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

TRACELESS SYMMETRIC TENSOR  $P_{ab}$ ?

"STOKES PARAMETERS"

PARITY  
 $E \rightarrow E$   
 $B \rightarrow (-B)$

ROTATION INVARIANT

$$\phi' \rightarrow \phi - \theta$$

$$\begin{pmatrix} Q' \\ U' \end{pmatrix} \rightarrow \begin{pmatrix} \cos 2\theta & \sin 2\theta \\ -\sin 2\theta & \cos 2\theta \end{pmatrix} \begin{pmatrix} Q \\ U \end{pmatrix}$$

$$\begin{pmatrix} E' \\ B' \end{pmatrix} = \begin{pmatrix} E \\ B \end{pmatrix}$$